# kOmega documentation

## <span id="page-0-0"></span>Giulio Dolcetti<sup>1,a</sup>, Borbála Hortobágyi<sup>2</sup>, Matthew Perks<sup>3</sup>, Simon J. Tait<sup>4</sup>

<sup>1</sup>: Department of Civil, Environmental and Mechanical Engineering, University of Trento, Trento, Italy.

<sup>2</sup>: UMR5600 Environnement Ville Société, Université de Lyon, Lyon, France.

<sup>3</sup>: School of Geography, Politics and Sociology, Newcastle University, Newcastle, UK.

<sup>4</sup>: Department of Civil and Structural Engineering, University of Sheffield, Sheffield, UK.

<sup>a</sup>: corresponding author, giulio.dolcetti@unitn.it

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## **Contents**



# <span id="page-1-0"></span>1 Fundamentals of the Method

kOmega is a Matlab script to calculate the average flow velocity and water depth of rivers and open-channel flows from sequences of images of the water surface recorded with a camera. The analysis is based on the method described in Dolcetti et al., [2022.](#page-14-2) The script computes the frequency-wavenumber spectra of the input set of images, and runs an optimisation routine to compare the measurements with the theoretical dispersion relation of water waves and to identify the set of flow parameters that provide the best fit with the measured data.

The method allows the estimation of the average flow velocity without requiring the presence of artificial tracers. It implements a robust analytical model of the water surface dynamics, therefore the accuracy is not undermined by the presence of gravity waves (including standing/stationary waves). The method is best suited for the analysis of videos or sets of images where surface deformations such as gravity waves are clearly visible, although it can also be applied in the absence of visible waves in the presence of artificial or natural floating tracers with suitable density.

Note: This script is intended for research applications only. The results of the analysis shall be interpreted with caution and always validated against established measurement methods. The Authors decline any responsibility for damage or harm resulting from the inappropriate use of the script or of any part of it.

## <span id="page-1-1"></span>1.1 Types of surface deformations

The water surface of rivers and open-channel flows shows the presence of multiple types of surface deformations (e.g., Muraro et al., [2021\)](#page-14-3). For the kOmega method, it is important to distinguish between gravity-capillary waves and turbulence-generated or turbulence-forced surface deformations. These surface deformations produce distinct features in the frequency-wavenumber spectra of the water surface of river and open-channel flows (Dolcetti et al., [2016\)](#page-14-4).



Figure 1: Examples of water surface dominated by different types of deformations. Left: turbulence-forced deformations (boils). Right: gravity-capillary waves.

Gravity-capillary waves are generally defined as sinusoidal long-crested waves comprising multiple crests and troughs, although the combination of multiple gravitycapillary waves can also produce patterns that are difficult to identify visually. Gravity-capillary waves do not move at the same speed of the underlying flow. This can affect the accuracy of standard optical velocimetry approaches if gravity waves or ripples are used as tracers (e.g., Dolcetti et al., [2020\)](#page-14-5). Instead, gravity-capillary waves propagate relative to the flow with the intrinsic celerity  $c_i$ , which varies with

the wavelength (long waves are faster than short waves) and which also depends on the depth of the flow,  $d$ :

$$
c_i(d) = \sqrt{gd\frac{(1+B)\tanh(kd)}{B}},\tag{1}
$$

where  $B = \rho g k^{-2} \gamma^{-1}$  is the so-called Bond number (which tends to infinity for long gravity-dominated waves),  $\rho$  is the water density, g is the acceleration due to gravity, and  $\gamma$  is the surface tension coefficient. Gravity-capillary waves are also advected by the flow, therefore their velocity in the streamwise direction is increased by an amount proportional to the speed of the flow.

Turbulence-generated surface deformations can appear at the surface of turbulent flows with different forms and shapes such as boils, scars, vortex dimples (e.g., Brocchini and Peregrine, [2001;](#page-14-7) Muraro et al., [2021\)](#page-14-3). Although their dynamics are still largely unknown, they are believed to move downstream approximately at the same speed of the flow. Hence, their speed is independent of their wavenumber.

Note: within kOmega, tracers (either artificial or natural, such as seeding particles, leaves, foam, etc.) are effectively equivalent to turbulence-generated surface deformations, since they also move at the same speed of the flow. Therefore, the method can be used also with artificial or natural tracers.

#### <span id="page-2-0"></span>1.2 Frequency-wavenumber spectra

kOmega exploits the relationship between the speed of the different types of waves and the flow parameters (velocity and water depth) to estimate the flow conditions. It does so by means of a Fourier-based approach applied to sequences of images of the water surface, which produce the so-called frequency-wavenumber spectrum. In fact, the speed or celerity of a wave is proportional to the ratio between its frequency f and wavenumber  $k$ :

$$
c = \frac{2\pi f}{k}.\tag{2}
$$

Frequency and wavenumber describe the temporal and spatial scales of a sinusoidal fluctuation. The frequency f (Hz) is the inverse of the period P (s),  $f = 1/P$ , and it indicates the number of cycles per second measured at a fixed location in space. The wavenumber  $\bf{k}$  (rad/m) is the spatial equivalent of frequency, it is proportional to the inverse of the wavelength  $\lambda$  (m),  $|\mathbf{k}| = 2\pi/\lambda$ , and it indicates how many full wavelengths are found in a 1 m length. The wavenumber is a vector with modulus  $k = |{\bf k}|$  directed in the direction of propagation of the wave. The two components of the wavenumber vector along the x and y directions are denoted as  $k_x$  and  $k_y$ , respectively.

The relationship between frequency and wavenumber (the so-called dispersion relation) can be approximated with the following equations (see Dolcetti et al., [2022\)](#page-14-2), for turbulence-generated deformations:

<span id="page-2-1"></span>
$$
f = \frac{1}{2\pi} \mathbf{k} \cdot \mathbf{U_s},\tag{3}
$$

and for gravity-capillary waves:

<span id="page-2-2"></span>
$$
f = \frac{(1 - \beta)}{2\pi} \mathbf{k} \cdot \mathbf{U_s} \pm \frac{1}{2\pi} \sqrt{(\beta \mathbf{k} \cdot \mathbf{U_s})^2 + (kc_i)^2},\tag{4}
$$

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Figure 2: Dispersion relation of gravity-capillary waves in two-spatial dimensions.  $\omega = 2\pi f$ . Taken from Dolcetti et al., [2022.](#page-14-2)

where  $U_s$  is the surface flow velocity,

$$
\beta = (1 - \alpha) \frac{\tanh(kd)}{kd},\tag{5}
$$

and  $\alpha = \bar{U}/U_s$  is the so-called velocity index, i.e., the ratio between the depthaveraged flow velocity  $\bar{U}$  and the surface velocity  $U_s$ . Note that the dispersion relation of turbulence-generated surface fluctuations (eq. [\(3\)](#page-2-1)) depends linearly on the surface flow velocity  $U_s$ , while the dispersion relation of gravity-capillary waves (eq. [\(4\)](#page-2-2)) shows a more complex dependence on both the surface flow velocity  $U_s$  and on the water depth d.

The discrete frequency-wavenumber spectrum is employed for decomposing a space-time signal such as a sequence of images into a discrete set of sinusoidal components (waves), each with wavenumber components  $k_{x,q}$  and  $k_{y,p}$  and with frequeny  $f_n$ . This is achieved by means of a Fourier transform in three dimensions  $(x,$ y, and time). It is assumed that a set of digital images is represented by a 3D matrix  $Z_{N_y \times N_x \times N_t}$ , where  $N_y$  and  $N_x$  are the number of pixels in the y and x-direction, respectively, and  $N_t$  is the number of frames.  $Z_{\eta,\xi,\tau}$  indicates the instantaneous intensity of the pixel corresponding to the spatial co-ordinates  $y = \eta L_y/N_y$  and  $x =$  $\xi L_x/N_x$  at the time  $t = \tau T/N_t$ , where  $L_y$  and  $L_x$  are the physical dimensions covered by the image (in m), and  $T$  is the duration of the set of images or video. Then, the discrete frequency-wavenumber spectrum  $I_{N_y \times N_x \times N_t}$  is a 3D matrix obtained by means of three discrete Fourier transforms applied along all three dimensions of the data, i.e.,

$$
I_{p,q,n} = \frac{1}{N_x N_y N_t} \left| \sum_{\eta=0}^{N_y - 1} \sum_{\xi=0}^{N_x - 1} \sum_{\tau=0}^{N_t - 1} Z_{\eta,\xi,\tau} \exp\left[-i2\pi \left(q\eta/N_y + p\xi/N_x - n\tau/N_t\right)\right] \right|^2. \tag{6}
$$

 $I_{p,q,n} = I(k_{y,p}, k_{x,q}, f_n)$  indicates the contribution in terms of energy that can be attributed to a wave with wavenumber components  $k_{y,p} = p2\pi/L_y$  and  $k_{x,q} = q2\pi/L_x$ and with frequency  $f_n = n/T$ . If gravity-capillary waves or turbulence-generated waves are present at the water surface, the frequency-wavenumber spectrum will have a larger amplitude at the frequency-wavenumber combinations that satisfy equations  $(3)$  and  $(4)$ .



Figure 3: Effect of a variation of the flow velocity (a, b) or of the water depth (c, d) on the frequency-wavenumber spectra of the water surface fluctuations. In (a,b) the depth is kept constant and the flow velocity is varied. In  $(c,d)$  the flow velocity is kept constant and the depth is varied.  $(a, c)$ : cross-section along the streamwise direction. (b,d): cross-section along the lateral direction. Taken from Dolcetti et al., [2022.](#page-14-2)

#### <span id="page-4-0"></span>1.3 Estimation of the flow velocity and depth

The approach of kOmega consists in fitting equations [\(3\)](#page-2-1) and [\(4\)](#page-2-2) to the high-energy peaks of the measured frequency-wavenumber spectrum in order to identify the optimal values of  $U_x$  and  $U_y$  (the two components of  $\mathbf{U}_s$ ) and eventually of d that better approximate the data. The fitting is performed by means of an optimiser that searches for the maximum of the Normalised Scalar Product:

<span id="page-4-1"></span>NSP = 
$$
\left[\sum_{p,q,n} Z_{p,q,n} M_{p,q,n} \right] \left[\sum_{p,q,n} Z_{p,q,n} \sum_{p,q,n} M_{p,q,n} \right]^{-1},
$$
 (7)

where  $M_{N_y \times N_x \times N_t}$  is a synthetic Gaussian-weighed frequency-wavenumber spectrum that follows equations [\(3\)](#page-2-1) and [\(4\)](#page-2-2).

Note: Unlike Dolcetti et al., [2022,](#page-14-2) who used a Self-Adaptive Differential Evolution (SADE) algorithm (Qin and Suganthan, [2005\)](#page-14-8) for the optimisation, kOmega employs the nonlinear constrained multivariable solver fmincon in Matlab, which is part of the Optimization Toolbox. Please ensure that the toolbox is installed in order to run the code.



Figure 4: Example of the fitting procedure. (a) and (b) show the streamwise (a) and lateral (b) cross-section through the 3D measured frequency wavenumber spectrum. (c) and (d) show the theoretical Gaussian-weighed spectrum M after the optimisation. The green lines are the theoretical dispersion relations, eq. [\(3\)](#page-2-1) and eq. [\(4\)](#page-2-2). Taken from Dolcetti et al., [2022.](#page-14-2)

### <span id="page-5-0"></span>1.4 Uncertainties and guidelines

The main sources of uncertainty of the method have been discussed in detail by Dolcetti et al., [2022.](#page-14-2) As is customary with other optical velocimetry approaches such as LSPIV (Muste et al., [2008\)](#page-14-9), the whole size of an image is usually split into portions (Areas of Interest, AOI) where conditions are assumed to be homogeneous, and for which a single value of velocity and/or depth is estimated. A smaller AOI improves the resolution of the measured distribution of the flow parameters. However, Dolcetti et al., [2022](#page-14-2) demonstrated that the uncertainty of the Fourier-based approach implemented in kOmega is strongly dominated by the spectral resolution. This defines the minimum detectable change of the dispersion relation and therefore of the flow parameters, and it is inversely proportional to the size of the AOI and to the duration of the sequence of images.

For accurate results, the area of interest (AOI, the portion of image that is actually used for the analysis) should be large enough to include multiple wavelengths (ideally 6-10).

The velocity estimations rely mostly on short ripples that can be detected accurately even with a small AOI. Therefore, measurements of the velocity alone can be performed with high accuracy even with a small AOI, thus enabling a high spatial resolution of the velocity distributions.

On the other hand, the sensitivity of the spectra to depth variations is small and controlled by waves with a wavelength larger than the water depth itself. As a result, depth estimations require a large AOI, at least 6-10 times larger than the water depth. Moreover, depth estimations are highly uncertain when long waves are absent. This could be the case of relatively calm, slow and/or deep flows with relatively small bed roughness. In these cases an accurate estimation of the water depth could be impossible. Additionally, the depth can only be estimated in the presence of gravity-capillary waves, while velocity estimations are possible even with turbulence-generated surface deformations alone.

When this is possible without affecting the relevance of the results (e.g., in the case of a relatively flat bathymetry) it is suggested to perform the analysis in two steps. In the first step, it is recommended to use a very large AOI covering almost the entire width of the river/channel to estimate an average water depth and velocity. In the second step, a refined spatial distribution of velocity and eventually depth can be calculated by means of multiple smaller AOI's opportunely distributed in space. During this second step, it is possible to either use the value of depth estimated during the first step, or to select relatively narrow boundaries to at least constrain the depth within reasonable limits. An exemplification of this two-step approach is given in the examples included with the code.

## <span id="page-6-0"></span>1.5 Camera placement

The relationship between the pixel intensity of an image of the water surface and the local surface deformation depends on multiple factors such as the camera angle and sensitivity to light and the illumination conditions. These factors can affect the detectability of the surface deformations. As discussed by Dolcetti et al., [2022,](#page-14-2) better visibility of the waves is obtained with oblique-viewing cameras in spite of nadir-looking ones (e.g., from drones), although too low angles may result in increased ortho-rectification errors. Direct sun reflections and sun glint should be avoided as well as shadows from river banks, bridges, trees, etc., whenever possible. The camera should be fixed and stable, or the images stabilised.

Note: kOmega does not include the algorithms to perform the image orthorectification or stabilisation. These steps should be performed externally with another software or code prior to the analysis.



Figure 5: Left: successful optimisation. Right: unsuccessful optimisation caused by a local minimum. Note the discrepancy between the measured spectrum and the theoretical relations in the latter case. Top: streamwise cross-section of the spectrum. Bottom: lateral cross-section of the spectrum.

## <span id="page-7-0"></span>2 Matlab Script

### <span id="page-7-1"></span>2.1 Requirements

kOmega uses the nonlinear constrained multivariable solver fmincon in Matlab, which is part of the Optimization Toolbox. The toolbox must be installed in order to run the code.

## <span id="page-7-2"></span>2.2 Algorithm Steps

- 1. Input data and parameters
- 2. Images pre-processing
- 3. Calculation of the frequency-wavenumber spectrum
- 4. Spectrum pre-processing
- 5. Optimisation
- 6. Output

The main input data consists in sets of orthorectified grayscale images provided in the form of a 3-dimensional array, where the 3 dimensions correspond to the spatial y-axis, to the spatial x-axis, and to the time (frame) axis, respectively. By default, kOmega analyses the whole data provided as input. If downsampling in time and/or space is required, or if the analysis should be limited to a smaller Area of Interest, then these should be identified and selected before running the code.



<span id="page-8-0"></span>Figure 6: Main steps of kOmega algorithm.



Figure 7: The system of reference used in kOmega.

<span id="page-9-3"></span>By default, kOmega runs the optimisation only once, using the average of the lower and upper boundaries for each parameter as starting point (see [2.3\)](#page-9-0). In the case of very noisy data, for small AOI size, or when the starting point is distant from the actual flow conditions, the optimisation algorithm may fail to identify the global maximum of the SNP and give erroneous results. Often these errors are easily identified by comparing the measured spectra with the theoretical relations of eq. [\(3\)](#page-2-1) and [\(4\)](#page-2-2) (see Fig. [6\)](#page-8-0). For these cases, kOmega provides an option to run the optimisation multiple times with randomly selected starting points. The options is activated by setting the parameter options.initialisation runs equal to the desired number of runs. The duration of the computation will increase accordingly. The optimal parameters are then identified as those with the highest SNP value across all runs.

### <span id="page-9-0"></span>2.3 Input data and parameters

The code requires three inputs: the array of images that need to be analysed  $(\text{input\_images});$  a structure array containing the input parameters  $(\text{input\_params});$ and a structure array containing options for the optimisation (options).

### <span id="page-9-1"></span>2.3.1 input\_images

The input images should be a set of ortho-rectified grayscale images provided in the form of a 3-dimensional array, where the 3 dimensions correspond to the spatial yaxis, to the spatial x-axis, and to the time (frame) axis, respectively. For example,  $input\_images(:,:,10)$  is the 10-th frame. kOmega accepts input as grayscale images or double.

Note: In kOmega, the  $x$  and  $y$  axis correspond to the first and second dimensions of the input images file, respectively. The positive x and  $y$  axis point in the direction of increasing indices. Therefore, the  $x$ -axis is the horizontal axis pointing towards the right. The y-axis is the vertical axis pointing downwards (see Fig. [7\)](#page-9-3).

#### <span id="page-9-2"></span>2.3.2 input\_params

These parameters include metadata for the set of images, expected boundaries of the flow parameters, and pre-processing parameters. Some parameters are required,

while others are optional. input params must be provided as a structure array with the following fields:

input_params.fps	Frame rate of the input images.	$s^{-1}$
input_params.pxl_size	Size of a pixel in the physical space.	m
input_params.velocity_indx	Velocity index, i.e., ratio between the depth-	
	averaged velocity $\bar{U}$ and the surface velocity	
	$U_s$ . See Hauet et al., 2018 for guidelines on	
	the choice of this parameter <sup>2</sup> . Typical values	
	are between 0.7 and 0.9, usually 0.83 to 0.85.	
input_params.segment_duration	Segments duration. The whole dataset is seg-	$\mathbf S$
	mented into shorter segments and the spectra	
	of each segment are then averaged. Shorter	
	segment durations help the noise converge	
	but also reduce the spectral resolution.	
input_params.overlap	Default 0. Overlap between consecutive seg-	
(optional)	input_params.overlap must be a nu- ments.	
	merical value between $0$ and $1$ , where $0$ means	
	no overlap and $0.3$ means $30\%$ overlap. An	
	overlap can be useful in the case of a small	
	number of frames to improve the spectrum	
	convergence.	
input_params.boundaries.velx	[b,ub] where lb and ub are the minimum and	m
	maximum allowed values for the target pa-	$s^{-1}$
	rameter velx. Setting $\mathsf{lb} = \mathsf{ub}$ effectively fixes	
	the parameter $\mathsf{velx} = \mathsf{lb} = \mathsf{ub}.$	
input_params.boundaries.vely	[b,ub] where lb and ub are the minimum and maximum allowed values for the target pa-	m $s^{-1}$
	rameter vely. Setting $\mathsf{lb} = \mathsf{ub}$ effectively fixes	
	the parameter $\text{vely} = \text{lb} = \text{ub}$ .	
input_params.boundaries.depth	[b,ub] where lb and ub are the minimum and	m
	maximum allowed values for the target pa-	
	rameter depth. Setting $\mathsf{lb} = \mathsf{ub}$ effectively	
	fixes the parameter depth $=$ lb $=$ ub.	
input_params.depth	Default []. Fixed water depth value, so that	m
(optional)	output_params.depth $=$ input_params.depth. If	
	input_params.depth is not declared or if in-	
	put_params.depth $= [ ]$ , then the estimation	
	of the water depth is attempted.	
input_params.GravityWaves	Default 'on'. Set input_params.GravityWaves	
(optional)	$=$ 'off' to ignore the spectrum of gravity-	
	capillary waves, if needed.	
input_params. TurbulenceWaves	Default 'on'. Set input_params. Turbulence	
(optional)	Waves $=$ 'off' to ignore the spectrum of	
	turbulence-generated waves, if needed.	

<sup>2</sup>Note that unlike typical non-contact velocimetry methods where the velocity index is only used for estimating the depth-averaged velocity and thus calculating the flow rate, here the velocity index also affects the wave dynamics (see eq. [\(4\)](#page-2-2)). Therefore a different value of input params.velocity indx will yield different estimates of the flow velocity and depth.

#### <span id="page-11-0"></span>2.3.3 options

These parameters include various options for the optimisation. In addition to a few dedicated parameters listed below, any option for the fmincon optimiser described in the fmincon documentation can be passed within the field options.optimoptions using the dedicated Matlab function optimoptions.



<sup>3</sup>The examples included with the script show a successful application of the method obtained with the following options: options.optimoptions = optimoptions('fmincon','Algorithm','interiorpoint','Display','notify','FunValCheck','on','OptimalityTolerance',1e-12,'PlotFcn','optimplotx'); These examples could serve as a starting point for the identification of the optimal parameters, which will ultimately depend on the input data.

## <span id="page-12-0"></span>2.4 Outputs

By default, kOmega attempts to estimate the three following parameters:



The value of each parameter can be fixed if needed by setting identical lowerand upper-boundary values for that parameter (see Sect. [2.3\)](#page-9-0). An alternative way to fix the water depth is to set the input parameter input params.depth equal to the desired value (see Sect. [2.3\)](#page-9-0). Both methods are equivalent.

kOmega also outputs the following optimisation diagnostics:



Additionally, kOmega outputs the 2D cross-sections of the measured frequency-wavenumber spectra along the x and y-directions, and the data to plot eq.[\(3\)](#page-2-1) and [\(4\)](#page-2-2) according to the estimated velocity and depth:

komega_spectrum.Spectrum_kx	Pre-processed frequency wavenumber spec-	
	trum <sup>1</sup> . Cross-section along the x-direction	
	(with $k_y = 0$ ).	
komega_spectrum.Spectrum_ky	Pre-processed frequency wavenumber spec-	
	trum <sup>1</sup> . Cross-section along the y-direction	
	(with $k_x = 0$ ).	
komega_spectrum.fTurbx	of turbulence- Theoretical frequency	$s^{-1}$
	generated waves, eq. (3), as calculated based	
	on the estimated flow velocity and water	
	depth. Cross-section along the $x$ -direction,	
	$f(k_x)$ .	
komega_spectrum.fTurby	frequency of turbulence- Theoretical	$s^{-1}$
	generated waves, eq. (3), as calculated based	
	on the estimated flow velocity and water	
	depth. Cross-section along the $y$ -direction,	
	$f(k_y)$ .	
komega_spectrum.fGWx_plus	Theoretical frequency of gravity-capillary	$\rm s^{-1}$
	waves, eq. $(4)$ , as calculated based on the	
	estimated flow velocity and water depth.	
	Cross-section along the x-direction, $f(k_x)$ .	
komega_spectrum.fGWy_plus	Theoretical frequency of gravity-capillary	$\overline{s^{-1}}$
	waves, eq. $(4)$ , as calculated based on the	
	estimated flow velocity and water depth.	
	Cross-section along the y-direction, $f(k_y)$ .	

<sup>&</sup>lt;sup>1</sup>Note that this is the non-dimensional normalised spectrum used for the optimisation, which has been pre-processed in order to facilitate the fitting, and which may differ significantly from the raw frequency-wavenumber spectrum.



If the full 3D arrays of the frequency-wavenumber spectra are needed, these can be obtained in alternative to the 2D cross-section by setting options.output  $=$  'Spectra projections' in the input options structure. In that case, the komega spectrum structure will be as follows:



# <span id="page-13-0"></span>3 Reusing and Sharing

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Dolcetti, G., Hortobágyi, B., Perks, M., Tait, S. J. & Dervilis, N. (2022). Using non-contact measurement of water surface dynamics to estimate water discharge. Water Resources Research, 58(9), e2022WR032829. [https://doi.org/10.1029/2022WR032829.](https://doi.org/10.1029/2022WR032829)

<sup>&</sup>lt;sup>1</sup>Note that this is the non-dimensional normalised spectrum used for the optimisation, which has been pre-processed in order to facilitate the fitting, and which may differ significantly from the raw frequency-wavenumber spectrum.

## <span id="page-14-6"></span><span id="page-14-0"></span>4 Support

For support, questions, and to flag eventual bugs, please contact Dr Giulio Dolcetti: giulio.dolcetti@unitn.it

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