IPhT Journal Club – Physics Beyond the Standard Model









Weak Nonlocality in Elementary Particle Physics

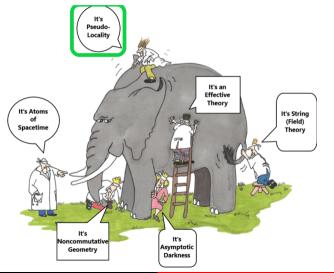
Florian NORTIER

CEA Paris-Saclay, DRF/IPhT, Orme des Merisiers, Gif-sur-Yvette, France florian.nortier@ipht.fr

January 17, 2023

1. Introduction – Blind Physicists & Weak Nonlocality

Weakly Nonlocal Theory \Rightarrow Interpolates btw UV Nonlocal Theory & IR Local QFT



1. Introduction – Weakly Nonlocal QFT

Weakly Nonlocal QFT

Nonlocal Form Factor $\gamma(z)\equiv$ Analytic Function with Regular Taylor Expansion in IR Limit $z=\Box=0$

Selected Reviews on Weak Nonlocality:

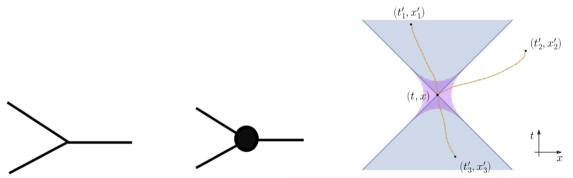
- Toy Scalar QFT's: [Buoninfante, Lambiase, Mazumdar, Nucl.Phys.B 944 (2019) 114646]
- Gravity & Gauge QFT's:
 - Tomboulis, Mod.Phys.Lett.A 30 (2015) 03n04, 1540005
 - Modesto, Rachwal, Nucl. Phys. B 900 (2015) 147-169
 - Modesto, Rachwal, Int.J.Mod.Phys.D 26 (2017) 11, 1730020
 - Bas i Beneito, Calcagni, Rachwal (2022), arXiv:2211.05606
- String Field Theory: [Erbin, Lect.Notes Phys. 980 (2021) 1-421]
- Nonlocal Black Holes: [Buoninfante, Giacchini, de Paula Netto (2022), arXiv:2211.03497]

1. Introduction – Contents

- Introduction
- Why Weak Nonlocality?
- Nonlocal Scalar Field
- Pseudolocality
- Epilogue

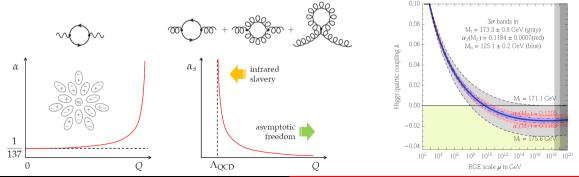
2. Why Weak Nonlocality? – Archaeology

- End 1920s \Rightarrow Birth of Local QFT + Problem of ∞ 's
- ullet Beginning 1930s \Rightarrow 1st Attempts to Regularize UV ∞ 's by Smearing Interactions
- Weak Nonlocality is Almost as Old as Local QFT \rightarrow Book on Old Literature: [Namsrai, Nonlocal QFT and Stochastic Quantum Mechanics, Springer Dordrecht (1986)]
- BUT Clash btw UV Finiteness & Gauge Invariance!
 [Chretien, Peierls, Proc.Roy.Soc.Lond.A 223 (1954) 1155, 468-481]



2. Why Weak Nonlocality? – Asymptotic Safety

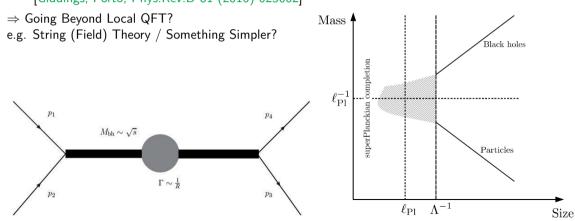
- Local QFT \Rightarrow SM of Particle Physics
 - Lorentz-Poincaré Invariance (Special Relativity) + Unitarity (Quantum Mechanics)
 - Gauge Symmetries + Higgs Mechanism.
 - Perturbative Renormalizability ⇒ Finite Number of Free Parameters
- BUT Landau Poles + Higgs Instability ⇒ Not Asymptotically Safe!
 - ullet \neq Quantum Chromodynamics (QCD) \Rightarrow Asymptotic Freedom
 - Total Asymptotic Freedom [Giudice, Isidori, Salvio, Strumia, JHEP 02 (2015) 137]



2. Why Weak Nonlocality? – Quantum Gravity

Quantum Gravity in Local QFT:

- Drawbacks for Perturbative Renormalizability + Unitarity!
- Black Holes: UV/IR Mixing + Information Paradox ⇒ Nonlocality? [Giddings, Porto, Phys.Rev.D 81 (2010) 025002]



3. Nonlocal Scalar Field

- Introduction
- Why Weak Nonlocality?
- Nonlocal Scalar Field
- Pseudolocality
- 6 Epilogue

3. Nonlocal Scalar Field – Toy Model

Nonlocal Toy Scalar Model:

$$S = \int d^4x \left[\frac{1}{2} \phi(x) \gamma(\Box) \phi(x) - V(\phi) \right]$$

Weakly Nonlocal Form Factor:

$$\gamma(\Box) = \sum_{n=0}^{\infty} c_n \Box^n$$

• Nonlocality Manifest via Delocalization Kernel:

$$\phi(x)\gamma(\Box)\phi(x) = \int d^4y\phi(x)K(y-x)\phi(y)$$

• Euler-Lagrange Equation:

$$\frac{\delta S}{\delta \phi(x)} = 0 \implies \gamma(\Box)\phi(x) - \frac{dV}{d\phi(x)} = 0$$

N.B: Free Field ⇒ Local Theory

3. Nonlocal Scalar Field – Ghost Problem (1/2)

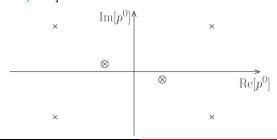
• Truncated Sum ⇒ Higher-Derivative Theory:

$$\mathcal{L}_{\phi} = -\frac{1}{2}\phi \left[\left(\frac{\square}{M^2} \right)^2 + 1 \right] \left(\square + m^2 \right) \phi - \lambda \sum_{n=4}^{N} \frac{c_n}{n!} \phi^n \,, \quad N \in \mathbb{N}$$

Propagator with Complex Conjugate Poles:

$$iD(p^2, m^2, \epsilon) = \frac{iM^4}{(p^2 - m^2 + i\epsilon)((p^2)^2 + M^4)}$$

 Problems with Causality/Unitarity/Stability ⇒ No Consensus on Higher-Derivative Theories [Platania, JHEP 09 (2022) 167]



3. Nonlocal Scalar Field – Ghost Problem (2/2)

• No Ghosts in EFT's!!! Weinberg's Footnote [Weinberg, Phys.Rev.D 77 (2008) 123541]:

$$\mathcal{L} = -\frac{1}{2} [\partial_{\mu} \varphi \partial^{\mu} \varphi + m^{2} \varphi^{2} + M^{-2} (\Box \varphi)^{2}] + J \varphi$$

where $M \gg m$ is some very large mass, and J is a c-number external current. We can easily find the connected part Γ of the vacuum persistance amplitude:

$$\Gamma = i \int d^4k \frac{|J(k)|^2}{k^2 + m^2 + k^4/M^2} \ .$$

If we took this result seriously, then we would conclude that in addition to the usual particle with mass $m + O(m^3/M^2)$, the theory contains an unphysical one particle state with mass $M + O(m^2/M)$. But if we regard \mathcal{L} as just the first two terms in a power series in $1/M^2$, then we must treat the term $M^{-2}(\Box \varphi)^2$ as a first-order perturbation, so that the vacuum persistence amplitude is

$$\Gamma = i \int d^4k |J(k)|^2 \left[\frac{1}{k^2 + m^2} - \frac{k^4}{M^2(k^2 + m^2)^2} + \dots \right] ,$$

and the only pole is at $k^2 = -m^2$. This is just the same result for Γ that we would find if we were to eliminate the second time derivatives in the $O(M^{-2})$ term in \mathcal{L} by using the field equation derived from the leading term in the Lagrangian

$$\Box \varphi = m^2 \varphi - J \ .$$

In this case the effective Lagrangian becomes

$$\mathcal{L} = -\frac{1}{2} [\partial_{\mu} \varphi \partial^{\mu} \varphi + m^2 \varphi^2 + m^4 M^{-2} \varphi^2] + (1 + m^2/M^2) J \varphi - J^2/2M^2 \; .$$

Original Argument: [Simon, Phys.Rev.D 41 (1990) 3720]

3. Nonlocal Scalar Field – Ghost-Free Form Factors (1/5)

Ghost-Free Infinite-Derivative Form Factor (Nonlocal Scale M):

$$\gamma(\Box) = (\Box - m^2)e^{H(\Box)}$$

 $H(\Box)$ is an Entire Function with H(0) = 1; Euclidean Bare Propagator:

$$-\frac{e^{-H(-k^2)}}{k^2+m^2}$$

• Wataghin Form Factor (String-Inspired) [Wataghin, Z.Phys. 88 (1934) 92-98]:

$$H_{\mathsf{Wat}} = -rac{\square}{M^2}$$

• Krasnikov Form Factor [Krasnikov, Theor.Math.Phys. 73 (1987) 1184-1190]:

$$H_{\mathsf{Kras}} = \frac{\square^2}{M^4}$$

3. Nonlocal Scalar Field – Ghost-Free Form Factors (2/5)

• Kuz'min Form Factor [Kuz'min, Sov.J.Nucl.Phys. 50 (1989) 1011-1014]:

$$H_{\mathsf{Kuz}}(\square) = \log m(\square) + \Gamma\left[0, m(\square)
ight] + \gamma_{\mathsf{E}}\,, \quad m(\square) = a\left(rac{\square}{M^2}
ight)^{n_{\mathsf{deg}}}$$

- $m(\Box) \equiv \mathsf{Real} \; \mathsf{Monomial} \; (a \in \mathbb{R}^*) \; \mathsf{of} \; \mathsf{Degree} \; n_{deg} \in \mathbb{N}^*$
- Upper Incomplete Gamma Function / Exponential Integral:

$$\Gamma(0,z) = \int_{z}^{\infty} dx \frac{e^{-x}}{x}$$

• Asymptotically Monomial in Conical Region:

$$e^{\mathrm{H}(z)} \stackrel{|z| \to \infty}{\simeq} |z|^{n_{\mathrm{deg}}} \qquad z \in \mathcal{C}, \quad -\theta < \arg(z) < \theta$$

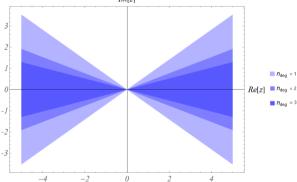
3. Nonlocal Scalar Field – Ghost-Free Form Factors (3/5)

• Tomboulis Form Factor: [Tomboulis (1997), arXiv:hep-th/9702146]

$$H_{\mathsf{Tom}}(\square) = rac{1}{2} \left[\log p^2(\square) + \Gamma \left[0, p^2(\square)
ight] + \gamma_E
ight] \; ; \; \mathsf{Polynomial} \; p(z) \; \mathsf{of} \; \mathsf{Degree} \; n_{deg} \in \mathbb{N}^*$$

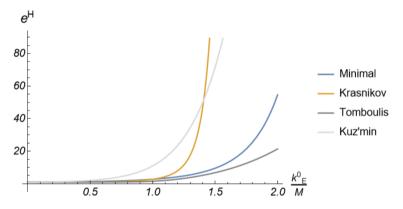
Asymptotically Polynomial in Conical Region:

$$\mathcal{C} = \{z | -\theta < \arg(z) < \theta \quad \cup \quad \pi - \theta < \arg(z) < \pi + \theta\}, \qquad \theta = \frac{\pi}{4n_{\text{deg}}}$$



3. Nonlocal Scalar Field – Ghost-Free Form Factors (4/5)

Form Factors vs Euclidean Energies k_E^0 :

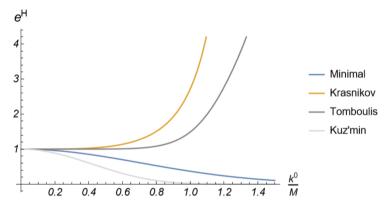


Minimal \equiv Wataghin; Kuz'min ($n_{deg} = 3$); Tomboulis ($n_{deg} = 2$)

⇒ All Form Factors Dampen Propagators in UV!

3. Nonlocal Scalar Field – Ghost-Free Form Factors (5/5)

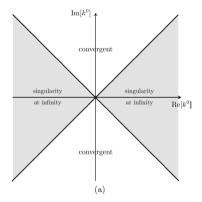
Form Factors vs Minkowskian Energies k^0 :



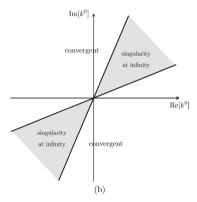
- \Rightarrow Wataghin (Minimal) & Kuz'min ($n_{deg} = 3$) Form Factors Make Propagators Blow Up in UV!
- \Rightarrow Only Krasnikov & Tomboulis ($n_{deg} = 2$) Form Factors Dampen Propagators in UV!

3. Nonlocal Scalar Field – Analyticity & Unitarity (1/2)

- Review: [Buoninfante, Phys.Rev.D 106 (2022) 12, 126028]
- Essential Singularity at Complex Infinity ⇒ Wick Rotation Not Defined!
- Form Factors Blow Up for Large Real Energies ⇒ Perturbative Unitarity Lost [Koshelev, Tokareva, Phys.Rev.D 104 (2021) 2, 025016]



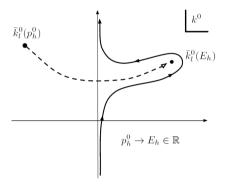
(a) Wataghin Form Factor $e^{-\Box}$

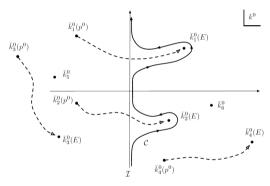


(b) Krasnikov Form Factor e^{\Box^2}

3. Nonlocal Scalar Field – Analyticity & Unitarity (2/2)

Efimov Analytic Continuation ⇒ Only External Momenta (Euclidean Starting point)
 [Efimov, Sov.J.Nucl.Phys. 4 (1967) 2, 309-315]





• Cutkosky Rules [Briscese, Modesto, Phys.Rev.D 99 (2019) 10, 104043]:

$$\frac{i e^{-H(p^2)}}{p^2 - m^2 + i\epsilon} \longrightarrow (-2\pi i) \sigma(p^0) \delta(p^2 - m^2 + i\epsilon)$$

3. Nonlocal Scalar Field – Higgs Mechanism & Naturalness (1/2)

- Naturalness of Higgs Boson Mass [Biswas, Okada, Nucl.Phys.B 898 (2015) 113-131]
- Toy ϕ^4 Theory (Klein-Gordon Mass):

$$S = \int d^4x \left[-\frac{1}{2} \phi e^{\frac{\Box + m^2}{M^2}} (\Box + m^2) \phi - \frac{\lambda}{4!} \phi^4 \right]$$

• Radiative Corrections to ϕ -mass (External Momenta $p^2=-m^2$):

$$\delta m^2 = i\Gamma_2 = \frac{\lambda}{32\pi^2} \left[e^{-\frac{m^2}{M^2}} + \left(\frac{m^2}{M^2}\right) Ei\left(-\frac{m^2}{M^2}\right) \right] M^2$$

$$M^2 \gg m^2 \quad \Rightarrow \quad \delta m^2 \simeq \frac{\lambda}{32\pi^2} M^2$$

• Terascale Fuzziness $\Rightarrow M = \mathcal{O}(1)$ TeV \Rightarrow Electroweak Scale Stabilized?

3. Nonlocal Scalar Field – Higgs Mechanism & Naturalness (2/2)

- Ghost-Free Higgs Mechanism Is Tricky [Hashi, Isono, Noumi, Shiu, Soler, JHEP 08 (2018) 064]
- Tachyon Field ⇒ Ghost-Free in Tachyonic Vacuum (Unstable):

$$S = \int d^4x \left[\frac{1}{2} \phi \, e^{-\frac{\Box}{M^2}} (\Box + \mu^2) \phi - \frac{1}{3} g \phi^3 \right] \quad \text{with} \quad \mu^2 > 0$$

• Spinless Boson ⇒ Not Ghost-Free in Physical Vacuum:

$$S = \int d^4x \left[\frac{1}{2} \varphi e^{-\frac{\Box}{M^2}} (\Box + \mu^2) \varphi - \mu^2 \varphi^2 - \frac{1}{3} g \varphi^3 \right].$$

The propagator of φ reads

$$\Pi_{\varphi}(k) = \left[e^{\frac{k^2}{M^2}} (k^2 - \mu^2) + 2\mu^2 \right]^{-1} .$$

⇒ No Ghost-Free Spectrum Anymore!

4. Pseudolocality

- Introduction
- Why Weak Nonlocality?
- Nonlocal Scalar Field
- Pseudolocality
- 6 Epilogue

4. Pseudolocality – Local Higher-Derivative Theories (1/4)

• Stelle Quadratic Gravity [Stelle, Phys.Rev.D 16 (1977) 953-969]:

$$S = \frac{1}{2\kappa^2} \int d^4x \sqrt{-g} \left(R + \alpha_1 R^2 + \alpha_2 R_{\mu\nu} R^{\mu\nu} + \alpha_3 R_{\mu\nu\rho\sigma} R^{\mu\nu\rho\sigma} \right)$$

- \Rightarrow Renormalizable in D = 4 (Euclidean) BUT 1 Ghost
- Higher-Derivative Gravity [Asorey, Lopez, Shapiro, Int.J.Mod.Phys.A 12 (1997) 5711-5734]:

$$S_{local} = \frac{1}{2\kappa^2} \int d^4x \sqrt{-g} \left[R + \sum_{n=0}^{n_{\text{deg}}-1} Ra_n \Box^n R + \sum_{n=0}^{n_{\text{deg}}-1} R_{\mu\nu} b_n \Box^n R^{\mu\nu} \right]$$

Superficial Degree of Divergence:

$$\omega(\mathcal{F}) = 4 - 2(n_{\text{deg}} - 1)(L - 1)$$

 $n_{deg} > 3 \Rightarrow$ Only 1-loop UV $\infty \Rightarrow$ Super-Renormalizable BUT Ghosts.

• Possibility to get a UV Fixed Point [Fradkin, Tseytlin, Phys.Lett.B 104 (1981) 377-381]

4. Pseudolocality – Local Higher-Derivative Theories (2/4)

Local QFT of Gravity ⇒ Ghost Problem!



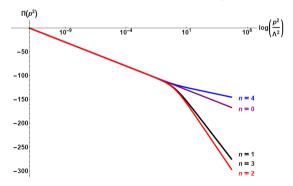
4. Pseudolocality – Local Higher-Derivative Theories (3/4)

• Eulidean Higher-Derivative Yang-Mills [Asorey, Falceto, Rachwal, JHEP 05 (2021) 075]:

$$S = \frac{1}{4g^2} \int d^4x \, F^a_{\mu\nu} F^{\mu\nu a} + \frac{1}{4g^2 \Lambda^{2n}} \int d^4x \, F^a_{\mu\nu} \, \Delta^n \, F^{\mu\nu a}$$

$$\Delta^{\nu b}_{\mu \, a} = -D^2 \delta^{\nu}_{\mu} \delta^b_a + 2 f^b_{\ ca} \, F^{\nu c}_{\mu}$$

• $2 \le n \le 4 \Rightarrow$ Superrenormalizable + Asymptotic Freedom (Only 1-loop UV ∞)



4. Pseudolocality – Local Higher-Derivative Theories (4/4)

How to Get Rid of those Ghosts?



4. Pseudolocality – Nonlocal Exorcism (1/3)



4. Pseudolocality – Nonlocal Exorcism (2/3)

Weakly Nonlocal QFT's with Ghost-Free Form Factor:

$$\mathcal{L}_{YM} = -\frac{1}{4g_{YM}^2} \left[\operatorname{tr} F e^{H(-\mathcal{D}_{\Lambda}^2)} F + \mathcal{V}_{YM} \right] \quad (GAUGE - THEORY)$$

$$\mathcal{L}_{gr} = -2\kappa_D^{-2} \sqrt{|g|} \left[R - \frac{1}{2} R \frac{e^{H(-\Box_{\Lambda})} - 1}{\Box} R + R_{\mu\nu} \frac{e^{H(-\Box_{\Lambda})} - 1}{\Box} R^{\mu\nu} + \mathcal{V}_{gr} \right]$$
(GRAVITY)

- Original Articles:
 - Kuz'min, Sov.J.Nucl.Phys. 50 (1989) 1011-1014
 - Tomboulis (1997), arXiv:hep-th/9702146
 - Modesto, Phys.Rev.D 86 (2012) 044005
 - Modesto, Rachwal, Nucl. Phys. B 889 (2014) 228-248
 - Modesto, Piva, Rachwal, Phys.Rev.D 94 (2016) 2, 025021
- Exponential Form Factors (Wataghin & Krasnikov) ⇒ Violate Power Counting Theorem!
- $\bullet \ \, \text{Asymptotically Polynomial Form Factors (Kuz'min \& Tomboulis)} \Rightarrow \text{Superrenormalizability!}$
- Theory Interpolates btw 2 local QFT's (deep IR & deep UV) via a Nonlocal Window

4. Pseudolocality – Nonlocal Exorcism (3/3)

Renormalization of Weakly Nonlocal Yang-Mills QFT [Tomboulis (1997), arXiv:hep-th/9702146]:

$$\mathcal{L}_{R} = -\frac{1}{2g^{2}} \operatorname{tr} \mathbf{F}_{\mu\nu} \mathbf{F}^{\mu\nu} - \frac{\alpha}{2} \operatorname{tr} \mathbf{F}_{\mu\nu} h(-\frac{\mathbf{D}^{2}}{\Lambda^{2}}) \mathbf{F}^{\mu\nu} - \frac{1}{2g^{2}} (Z_{3} - 1) \operatorname{tr} \mathbf{F}_{\mu\nu} \mathbf{F}^{\mu\nu}$$

• Ghost-Free Form Factor (α Is Not Renormalized):

$$\overline{h}(z) \equiv 1 + g^2 \alpha h(z)$$

- No Ghosts \Rightarrow Renormalization Condition at Scale $\mu_0 \sim \Lambda$ with $\alpha = 1/g(\mu_0)^2$.
- Dressed Propagator $\Rightarrow \infty$ -Tower of CC Poles? [Shapiro, Phys.Lett.B 744 (2015) 67-73]:

$$-i\frac{e^{-H(k^2)}}{k^2\left(1+\beta_{\alpha}\,e^{-H(k^2)}\log(k^2/\mu_0^2)\right)}$$

• Naive Guess for Matter & Higgs Fields [Modesto, Rachwal, Nucl.Phys.B 900 (2015) 147-169]:

$$\sum_{a}^{N_f} \bar{\psi}_a i \mathcal{D}_a e^{H(-\mathcal{D}_{a,\Lambda}^2)} \psi_a + (\mathcal{D}_\mu \Phi)^{\dagger} e^{H(-\mathcal{D}_{\Lambda}^2)} (\mathcal{D}^{\mu} \Phi) - \mu^2 \Phi^{\dagger} e^{H(-\mathcal{D}_{\Lambda}^2)} \Phi - \lambda (\Phi^{\dagger} \Phi)^2$$

BUT Not Ghost-Free in Physical Vacuum! [Hashi, Isono, Noumi, Shiu, Soler, JHEP 08 (2018) 064]

4. Pseudolocality – Matter & Higgs Sectors

- Pseudolocal Recipe from a Mother Local QFT (Minimal Nonlocality):
 - Modesto (2021), arXiv:2103.04936
 - Modesto, Calcagni, JHEP 10 (2021) 169

$$\begin{split} S(\Phi_i) &= \int d^D x \sqrt{|g|} \left[\mathcal{L}_{\text{loc}} + E_i \, F_{ij}(\Delta) \, E_j + \mathcal{V}(E) \right], \\ S_{\text{loc}} &= \int d^D x \sqrt{|g|} \, \mathcal{L}_{\text{loc}} \,, \qquad \mathcal{L}_{\text{loc}} = \frac{1}{2\kappa^2} R + \mathcal{L}_{\text{m}} \,, \\ E_i(x) &:= \frac{\delta S_{\text{loc}}}{\delta \Phi_i(x)} \,, \\ \Delta_{ki} &:= \frac{\delta E_i}{\delta \Phi_k} = \frac{\delta^2 S_{\text{loc}}}{\delta \Phi_k \delta \Phi_i} \,, \\ 2\Delta_{ik} F_{kj}(\Delta) &\equiv \left[e^{\text{H}(\Delta)} - 1 \right]_{ij} \,, & \mathcal{E}_k = \left[e^{\text{H}(\Delta)} \right]_{kj} \, E_j + O(E^2) = 0 \end{split}$$

- Classical Solutions of Local Theory ⇒ Solutions of Weakly Nonlocal Theory!
- $\bullet \ \mathsf{Tree\text{-}Duality} \Rightarrow \mathsf{Perturbative} \ \mathsf{Spectrum} + \mathsf{Tree\text{-}Level} \ \mathsf{Scat}. \ \mathsf{Amplitudes} \equiv \mathsf{Mother} \ \mathsf{Local} \ \mathsf{QFT!}$
- Ghost-Free Higgs Mechanism: [Modesto, JHEP 06 (2021) 049]

5. Epilogue

- Introduction
- Why Weak Nonlocality?
- Nonlocal Scalar Field
- Pseudolocality
- 6 Epilogue

5. Epilogue – Summary & Outlook

Summary:

- Intrinsic to Gravity + Inspired by String (Field) Theory
- Weak Nonlocality ⇒ UV Softening Without Ghosts
- ullet Euclidean o Lorentzian Signature via Efimov Analytic Continuation
- ullet Gauge Interactions + Gravity \Rightarrow Asymptotically Polynomial Form Factors
- Matter + Higgs \Rightarrow Tree-Duality BUT Detailed Analysis Still Missing...

Outlook:

- Which Notion of Causality in a Pseudolocal QFT?
- Analyticity/Unitarity of Scattering Amplitudes at Loop-Level?
- Consistency ⇒ Isolate a Unique Form Factor?
- Other Consistent Weakly Nonlocal Theories? \Rightarrow Worldline Inversion Symmetry? $t\mapsto 1/t$ [Abel, Dondi, JHEP 07 (2019) 090] \Rightarrow UV/IR Duality?

$$\Pi(p^2) = \int_0^\infty dt e^{-T(t)(p^2+m^2)}$$

5. Epilogue – UV/IR Mixing in String Theory

Modular Invariance in String Theory [Abel, Dienes, Phys.Rev.D 104 (2021) 12, 126032]:

