

Speculation on Temperature in Pure Quantum Bound States

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In (1) it is suggested that Newton's second law for constant acceleration, i.e. $F=ma$, may be derived using the first law of thermodynamics (with $dE=0$) and the special relativistic idea of a constant acceleration being linked to a temperature as shown by Unruh i.e. $T=a C1$ (where $C1$ is a constant given in terms of \hbar , c etc).

In this note, we consider the two Lorentz invariants $E^2 = p^2 + m^2 c^4$ (and its generalization $(E-V(x))^2 = p^2 + m^2 c^4$) and $-Et+px$. The former becomes Newton's energy conservation law in the nonrelativistic limit i.e. $E = p^2/2m + V(x)$ which is equivalent to Newton's second law (upon taking d/dx), but contains a variety of accelerations.

The Lorentz invariant $-Et+px$ suggests t and x are independent and that t is linked with an eternal clock with frequency \hbar/E and x with an internal ruler with spacing \hbar/p . Thus even though a particle with constant speed moves as $x = p/m t$ (as measured using an external clock and ruler) there is an internal wavelength \hbar/p in which there is a probability for the particle to be at various x points (as seen with an external ruler) and no internal time.

Given that one may consider externally $x = p/m t$, one may also consider externally the situation of acceleration. Each constant p , however, is associated with $\exp(ipx)$ (an eigenfunction of the translation operator $-i\hbar d/dx$) so at a given x , one has a superposition of $\exp(ipx)$'s. This suggests a different "probability" distribution for p at each x or in other words a temperature linked with $\{\sum_p a(p) p^2/2m \exp(ipx)\} / \{\sum_p a(p) \exp(ipx)\}$. The temperature changes from one x point to another. If one considers that E is constant at each x . then applying the first law of thermodynamics yields $F dx = dx d/dx (-1/2m d/dx dW/dx)$ which should be equivalent to $T dS$. Using Unruh's results, this leads to the expression for entropy given in (1).

This, however, seems to be an entropy which is different from one constructed using Shannon's entropy equation. We consider some of these ideas in this note.

Newton's Law Using Unruh Temperature

In (1) it is noted that using Unruh's temperature and the first law of thermodynamics with $dE=0$ i.e. $F dx = T dS$ ((1)) and dS assumed to be proportional to dx , one may obtain Newton's second law $F=ma$ for a constant acceleration.

Unruh's temperature/constant acceleration relationship follows from special relativity applied to a constant acceleration (2):

$$kT = \hbar a / (2\pi \cdot 3.14 \cdot c) \quad ((2))$$

Thus a constant acceleration is associated with the notion of a temperature.

Using ((1)) with $F dx = T dS$ and assuming

$$dS = 2\pi \cdot 3.14 \cdot k^* m \hbar / c dx \quad ((3)) \quad \text{yields} \quad F=ma \quad ((4)) \quad \text{i.e. Newton's second law for a constant}$$

acceleration.

We try to argue that the idea of a “temperature” associated with acceleration appears in quantum mechanics which we argue also follows from special relativity.

Quantum Mechanics from Special Relativity

We have argued in previous notes that quantum mechanics follows from Lorentz invariants in special relativity.

$$\text{First we use: } E^2 = p^2 + m^2 \quad (c=1) \quad ((5))$$

Here $p = m_0 v / \sqrt{1 - v^2/c^2}$ ((6a)) where $x = vt$ for a constant velocity. Thus there is an external clock and ruler which may be used to observe the particle move according to $x = vt$.

Next, one introduces the notion of v , as measured from an external clock and ruler, changing in each “ dx ” interval, where dx is chosen to be small i.e. $dx \rightarrow 0$. Using a function $V(x)$ (called a potential) one may write:

$$(E - V(x))^2 = p^2 + m^2 \quad (c=1) \quad ((7))$$

Taking the nonrelativistic limit i.e. p and $V(x) \ll m$ ($c=1$) yields the Newtonian conservation of energy equation:

$$E = p^2/2m + V(x) \quad ((8)) \quad \text{with } p = m_0 v \quad \text{and } F = -dV/dx \quad ((9))$$

((8)) is associated with a different acceleration dv/dt at each x if $V(x)$ is not a constant.

An important issue, however, arises if one considers the Lorentz invariant $A = -Et + px$. This invariant suggests that t and x are independent and represent an internal time and internal ruler with:

$$\text{Frequency} = \hbar/E \quad ((10a)) \quad \text{and internal ruler spacing} = \hbar/p \quad ((10b))$$

How can this be consistent with $x = vt$ which suggests infinite resolution? We suggest that $x = vt$ only holds on average and is measured using an internal clock and ruler with infinite resolution. There are, however, physical repeated half wavelength units of length $\hbar/2p$ which contain no internal time, but rather a probability for the particle to be at x given by $\cos(px)$. $\exp(ipx)$ is a sum of $\cos(px)$ and $i \sin(px)$ i.e. shifted distributions to indicate the direction of motion because internal time is removed. It may also be noted that $\exp(ipx)$ is an eigenfunction of $-i\hbar/dx$ and so has repeated wavelength units to demonstrate a kind of spatial invariance.

Thus there is no notion of $dx \rightarrow 0$, but one may still achieve and average acceleration using ((8)) i.e.

$$E = \left\{ \sum_{\text{over } p} a(p) \frac{p^2}{2m} \exp(ipx) \right\} / \left\{ \sum_{\text{over } p} a(p) \exp(ipx) \right\} + V(x) \quad ((11))$$

The first term on the RHS is equivalent to $-1/2m \frac{d}{dx} \frac{dW/dx}{W}$ where $W = \sum_p a(p) \exp(ipx)$.

Thus many constant p values create the average kinetic energy at x using interference. This $KE(x)$ is equivalent to the classical value $E - V(x)$ and:

$$\frac{d}{dx} KE(x) = m a(x) \quad ((12)) \quad \text{where } a(x) \text{ is acceleration at } x.$$

In classical statistical mechanics kT is equivalent to average kinetic energy. In ((11)) one has an average kinetic energy consisting of a distribution of p values at x , each with weight $a(p) \cos(px)$ if one considers the symmetry $a(p) = a(-p)$.

Thus this approach also leads to the notion that a given acceleration is equivalent to a temperature i.e. a $pp/2m$ distribution and also follows directly from special relativity (which we argue is the source of the $\exp(ipx)$ formalism of quantum mechanics).

Entropy Considerations

If one uses the first law of thermodynamics at each x point with $E = \text{constant}$ so $0 = TdS - Fdx$ then:

$$TdS = Fdx = (dx) \frac{d}{dx} KE(x) \quad ((13))$$

Thus dS is proportional to dx as in (1). Then according to ((2)) $T = C_1 a$ so ((13)) implies that:

$$C_1 a dS = dx m a(x) \quad \text{or} \quad dS = m/C_1 dx \quad \text{where } C_1 = \hbar / (2 \cdot 3.14 \cdot c \cdot k) \quad ((14))$$

((14)), however, is equivalent to ((3)) presented in (1).

One may note that this $S(x)$ is a constant and not the same entropy density one would obtain using Shannon's entropy formula

Conclusion

In conclusion, we argue that using Lorentz invariants from special relativity i.e. $E^2 = p^2 + m^2$ (with $c=1$) generalized to $(E - V(x))^2 = p^2 + m^2$ and $A = -Et + px$ and taking the nonrelativistic limit one obtains $E = pp/2m + V(x)$ and $\exp(ipx)$ representing a quantum particle. Thus there is no $dx \rightarrow 0$ limit and one must use a superposition of $\exp(ipx)$ s to create an average $v_{rms}(x)$ and the notion of acceleration. Given that a distribution of p 's is used with weights $a(p) \cos(px)$ at each x for symmetry $a(p) = a(-p)$, one has a temperature $T(x)$ associated with $a(x)$. Using the first law of thermodynamics at x with $dE=0$ one has $Fdx = TdS$ and may show that $TdS = dx \frac{d}{dx} KE(x)$. Using Unruh's relationship for T and a constant acceleration, one finds

that $dS = 2 \cdot 3.14 \cdot k \cdot m \cdot hbar / c \cdot dx$ which matches the result in (1) obtained using different arguments. We note that this entropy is not the same as one obtained using Shannon's entropy formula.

References

1. Verlinde, E. On the Origins of Gravity and the Laws of Newton (2010)
<https://arxiv.org/pdf/1001.0785.pdf>
2. https://en.wikipedia.org/wiki/Unruh_effect