

Optical Reading of Nanoscale Magnetic Bits in an Integrated Photonic Platform

Hamed Pezeshki, Pingzhi Li, Reinoud Lavrijsen, Jos J. G. M. van der Tol, and Bert Koopmans

Abstract— In this paper, we propose a compact integrated hybrid plasmonic-photonic device for optical reading of nanoscale magnetic bits with perpendicular magnetic anisotropy in a magnetic racetrack on top of a photonic waveguide on the indium phosphide membrane on silicon platform. The hybrid device is constructed by coupling a doublet of V-shaped gold plasmonic nanoantennas on top of the indium phosphide waveguide. By taking advantage of the localized surface plasmons, our hybrid device can enable detection of the magnetization state in magnetic bits beyond the diffraction limit of light and enhance the polar magneto-optical Kerr effect (PMOKE). We further illustrate how combining the hybrid device with a plasmonic polarization rotator provides magneto-optical read-out by transforming the PMOKE-induced polarization change into an intensity variation of the waveguide mode. According to the simulation results based on a three-dimensional finite-difference time-domain method, the hybrid device can detect the magnetization states in targeted bits in a magnetic racetrack medium down to $\sim 100 \times 100 \text{ nm}^2$, regardless of the magnetization state of the rest of the racetrack with a relative intensity contrast of greater than 0.5% for a $\sim 200 \times 100 \text{ nm}^2$ magnetic bit. We believe our hybrid device can be an enabling technology that can connect integrated photonics with nanoscale spintronics, paving the way toward ultrafast and energy efficient advanced on-chip applications.

Index Terms— Photonic integrated circuits, Plasmonics, Spintronics, Indium phosphide, Magneto-plasmonics, Polar magneto-optical Kerr effect

I. INTRODUCTION

WITH the increase in demand for high bit-rate data transfer in the field of telecommunications and quantum information, the need for new technologies for high speed and reliable data reading and writing is foreseen. Advancements in the field of integrated photonics have enabled wide bandwidth optical data transmission on photonic integrated circuits, while spintronics has empowered high bit-rate data storage [1]. As a further upgrade, it is expected that a direct optical reading of spintronic domains, without any intermediate controlling high-frequency electronics, can bring higher bit-rate data transfer as well as enhance the energy efficiency of the communication platform [1], [2].

However, the drastic size mismatch between the diffraction-

limited waveguide mode ($> 400 \text{ nm}$ [3]) and spintronic devices ($< 150 \text{ nm}$ [1]), makes photonic access of spintronics incompatible. Moreover, the limited interaction cross-section results in low magneto-optical (MO) response.

Being able to concentrate the light energy without excessive photonic elements is the key to address the above mentioned problems. For over a couple of decades, it has been demonstrated that nanostructures made of noble metals like gold can confine the light energy in nanoscale spots at the metal-dielectric interface and magnify it as a result of the localized surface plasmon resonance [4], [5]. Due to the localized and enhanced electric field, the light-matter interaction within these nanostructures (plasmonic nanoantennas) gets boosted, which leads to an effectively larger cross-section of scattering and absorption powers of coupled nanoparticles. Nowadays, such unique features of plasmonic nanoantennas allow them to play a key role in addressing nanoscale elements in different areas such as telecommunication [6], [7] and biosensing [8], [9]. Hence, it has been anticipated that we can improve MO activity at the nanoscale based on magneto-plasmonic effects via incorporating plasmonic nanoantennas in our photonic device.

In this paper, we propose an integrated hybrid plasmonic-photonic device for enhanced optical reading of magnetic bits by taking advantage of magneto-plasmonic effects. The targeted functionality is to read out data, which are stored in the up and down magnetization of ferromagnetic bits of a racetrack with perpendicular magnetic anisotropy [10], placed on top of a photonic waveguide on the indium phosphide (InP) membrane on silicon (IMOS) platform. The read-out functionality is based on the polar magneto-optical Kerr effect (PMOKE), which manifests itself as the induced change of the polarization state of light reflected off the surface of a magnetic material depending on the magnetization state. The intrinsically small MO signal is enhanced by coupling a doublet of V-shaped gold plasmonic nanoantennas with the racetrack on top of a photonic waveguide on the IMOS platform. This allows the incident light propagating through the waveguide to be focused efficiently on the targeted nanoscale magnetic bit in the coupled racetrack and to enhance the naturally weak PMOKE response

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by improving the MO interactions with the magnetic bits. To detect the magnetic states in a targeted magnetic bit, the resultant polarization change induced by PMOKE is transformed into an intensity variation of the transverse magnetic (TM_0) waveguide mode [11] using our proposed plasmonic polarization rotator. Based on three dimensional finite-difference time-domain simulations, Lumerical, FDTD solver [12], we found that the presented hybrid device can enable the detection of the magnetization state in targeted bits in a magnetic racetrack medium down to a footprint of $\sim 100 \times 100 \text{ nm}^2$, without being disturbed by neighbouring bits. In contrast, for a photonic waveguide without the plasmonic nanoantennas, i.e. the bare waveguide device, this is not possible due to very weak PMOKE, while the extended waveguide mode would mix signals from neighbouring bits, ending up in ambiguous reading. Moreover, by using the

polarization rotator, which works as an integrated quarter-wave plate, we show that the magnetic state in a magnetic bit with a size of $\sim 200 \times 100 \text{ nm}^2$ can be read-out optically with an intensity contrast of more than 0.5% only with using the hybrid device. The proposed device is a generic model which can be implemented in other photonic platforms such as silicon-on-insulator [13], [14] and silicon nitride [15], [16]. Based on our theoretical results, we believe that our proposed device concept can be an enabling technology which offers a method for direct optical reading of magnetic bits without intermediate electronics and it can be as well useful for magnet based sensing technologies.

II. DESIGN STRUCTURE AND FUNCTION PRINCIPLE

The schematic diagram of the proposed integrated device for optical reading of magnetic bits in a racetrack is shown in Fig.

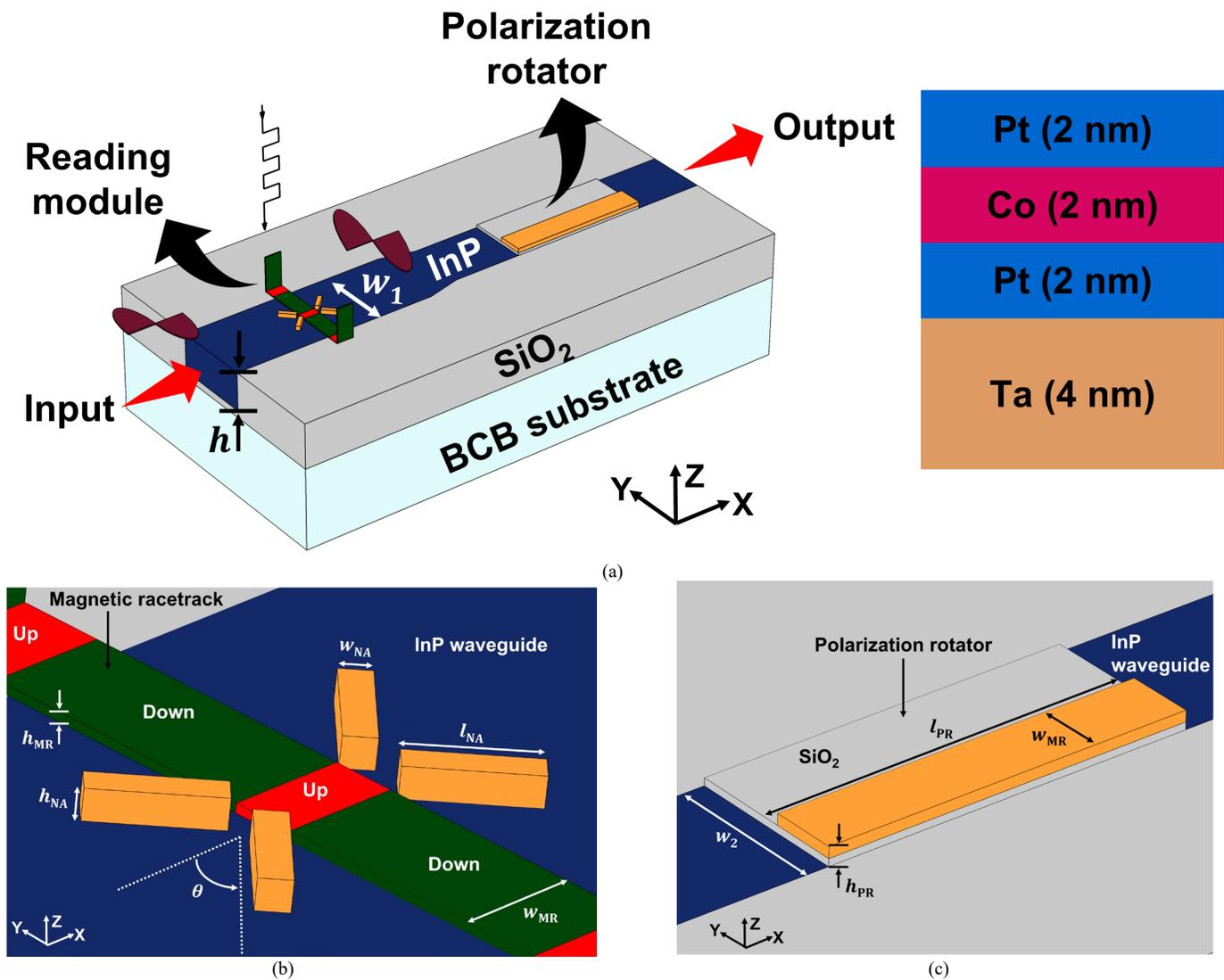


Fig. 1. The concept of the integrated device for all-optical reading of magnetic bits. (a) Left: a schematic diagram of the device illustrating the operation principle; right: the inset showing the transversal schematic diagram of the multilayer stack, comprising the magnetic racetrack. (b) Magnified view of the reading module showing the dimensions of the plasmonic nanoantenna elements and magnetic racetrack. The waveguide's width and height are $w_1 = 570 \text{ nm}$ and $h = 280 \text{ nm}$ (see (a)), the racetrack width and height are $w_{MR} = 120 \text{ nm}$ and $h_{MR} = 10 \text{ nm}$, respectively; the length, width and height of the plasmonic nanoantenna elements are $l_{NA} = 120 \text{ nm}$, $w_{NA} = h_{NA} = 30 \text{ nm}$, and they are oriented at $\theta = 45^\circ$ with reference to the waveguide direction. (c) Magnified view of the polarization rotator, where the width of the waveguide is $w_2 = 440 \text{ nm}$, respectively; the length, width, and height of the polarization rotator are $l_{PR} = 1860$, $w_{PR} = 180 \text{ nm}$, and $h_{PR} = 50 \text{ nm}$, respectively. 'Ta', 'Pt', 'Co', 'NA', 'MR', and 'PR' stand for tantalum, platinum, cobalt, nanoantenna, magnetic racetrack, and polarization rotator.

1(a). As depicted, the device consists of two sections: the reading module and the polarization rotator section. The whole device is based on the IMOS technology [17]. The reading module is composed of a doublet of V-shaped gold plasmonic nanoantennas coupled with a magnetic racetrack as a top-cladding on the InP waveguide. A magnetic racetrack enables densely storing of information as up and down magnetization states [18], which can be moved along the racetrack by electrical current [19]–[21]. The racetrack is modelled as a multilayer stack (from bottom to top) of 4 nm heavy metal tantalum seed layer, 2 nm platinumium layer, 2 nm ferromagnetic cobalt layer with a MO Voigt constant of $Q = 0.154 - 0.100i$ [22] (responsible for PMOKE in our simulation model), and a 2 nm platinumium capping layer. A continuous wave laser light source with a free space wavelength of $\lambda_0 = 1550$ nm is coupled to the transverse electric (TE_0) waveguide mode. Upon interaction between the TE_0 waveguide mode and the plasmonic nanoantenna under the resonance condition, the localized surface plasmons of the plasmonic nanoantenna get excited and enhance the electric field at its nanoscale hot spot, where the magnetic racetrack is coupled. The concentrated electric field leads to an enhanced PMOKE. As the pure TE_0 mode interacts with the magnetic cladding, partly a TM_0 mode with a small magnitude is created due to PMOKE, whose phase is magnetization dependent, i.e. changes by 180° when the magnetization reverses. As a result of the birefringence in the waveguide, the TE_0 mode and PMOKE-induced TM_0 mode beat along the propagation distance. Therefore, the rotation and phase of this beating is also magnetization dependent. When the propagating light passes through the polarization rotator, this phase difference is transformed to an intensity variation of the emerged TM_0 mode due to the partial conversion of the TE_0 mode to TM_0 mode [11], by which the change in the magnetization state can be detected using a photodetector.

The design parameters are presented in Fig. 1. The width and height of the waveguide are $w_1 = 570$ nm and $h = 280$ nm. The plasmonic nanoantenna is based on a doublet of V-shaped gold nanoantennas, each of which is oriented at $\theta = 45^\circ$ with reference to the waveguide direction and has a length, width and height of $l_{NA} = 120$ nm, $w_{NA} = h_{NA} = 30$ nm, where the subscript ‘NA’ is the acronym for nanoantenna. The geometry of the plasmonic nanoantenna is optimized to have a resonance peak at $\lambda_0 = 1550$ nm. The racetrack has a width and height of $w_{MR} = 120$ nm and $h_{MR} = 10$ nm, where the subscript ‘MR’ stands for magnetic racetrack. Figure 1(c) shows a schematic of the polarization rotator. A gold metal film with a thickness of 30 nm is stacked on top of the waveguide, where to minimize the absorption loss by the metal film, a silica (SiO_2) spacer layer with a thickness of 20 nm is placed in between the metal film and the waveguide. The length, width, and height of the polarization rotator are $l_{PR} = 1860$ nm, $w_{PR} = 180$ nm, and $h_{PR} = 50$ nm (where ‘PR’ stands for polarization rotator), respectively, in order to match the quarter of beating length and rotate the eigenmode by 45° . In other words, we optimized our polarization rotator to perform as an integrated quarter-wave plate. The width of the waveguide in the polarization rotator section is reduced to $w_2 = 440$ nm to maximize the TE_0 to TM_0

conversion efficiency. Note that the refractive indices of the materials used in the model are taken from the built-in library of the Lumerical, FDTD solver [12]. The optical mode profiles at the cross section with the magnetic racetrack, i.e. in the YZ plane at $X=0$, for the bare waveguide and hybrid devices are illustrated in Figs. 2 (a) and (b), respectively. By comparing these two cases, we can clearly see an enhanced optical confinement in the magnetic racetrack with the help of the plasmonic nanoantenna. The spatial field mode distributions cut-through the middle of the reading module’s waveguide in the XY plane for the bare waveguide and hybrid devices are presented in Figs. 2(c) and (d), respectively. Based on Fig. 2 (c), we can see that as light travels toward the polarization rotator, the change in the magnitude of the electric field is insignificant, due to a weak interaction between the magnetic racetrack and waveguide mode. In contrast, in the hybrid device, due to the enhanced light-matter interaction by the plasmonic nanoantenna, we can see a drop in the magnitude of the electric field of the propagating light toward the polarization rotator, as shown in Fig. 2(d). These changes in the mode profile and magnitude of the electric field are consequences of the enhanced interaction between the waveguide mode and plasmonic nanoantenna, which ultimately results in enhanced MO response of the device.

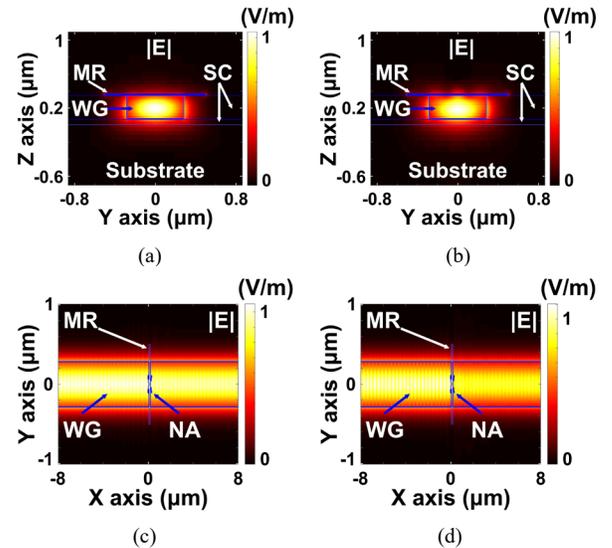


Fig. 2. (a, b) The optical mode profiles at the cross section with the magnetic racetrack, i.e. in the YZ plane at $X=0$, for the bare waveguide and hybrid devices, respectively. (c, d) The spatial field mode distribution cut-through the middle of the waveguide in the XY plane for the bare waveguide and hybrid devices. ‘MR’, ‘NA’, ‘SC’, and ‘WG’ stand for magnetic racetrack, nanoantenna, SiO_2 cladding, and waveguide, respectively.

III. RESULTS

In this section, we present the steps involved in the process of reading the magnetization state optically. The first step is based on PMOKE which is an induced change in the polarization state of the TE_0 mode due to the MO activity.

In the next step, we will convert the PMOKE-induced polarization change to an intensity variation of the TM_0 mode which can be easily detected using an on-chip photodetector.

A. Optical Reading using PMOKE

By optically reading the magnetization using PMOKE, data stored in up and down magnetization states can be transferred to the photonic state. We assessed the performance of the optical reading using PMOKE based on two major factors. The first factor is the magnitude of the resultant PMOKE-induced polarization rotation (Kerr rotation) defined as the real part of the ratio of the magnitudes of TM_0 to TE_0 , i.e. $\Re(TM_0/TE_0)$. The second one is the minimum footprint of the targeted magnetic bit in which the magnetization state can be identified regardless of all adjacent magnetic bits in the rest of the racetrack.

To investigate the minimum size of the targeted bit whose magnetization state can be determined, we perform a study of the evolution of the polarization rotation as a function of the size of the targeted magnetic bit, along the light propagation direction inside the waveguide. This analysis is done in the presence of the oppositely magnetized rest of the racetrack (see the inset in Fig. 3(a)). Here, we study the polarization rotation magnitude and phase as a function of the targeted bit size, by

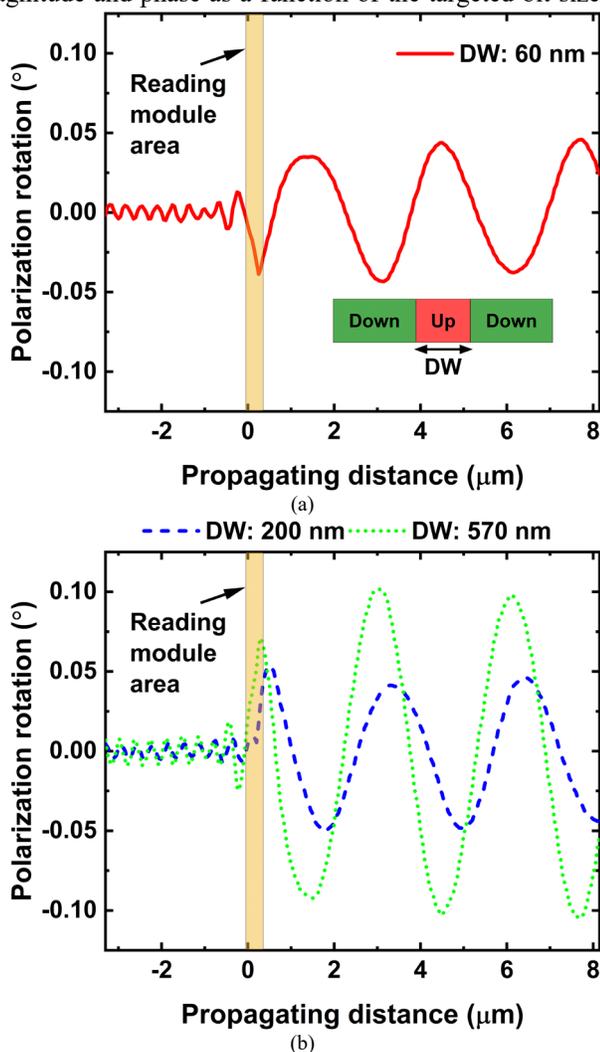


Fig. 3. The evolution of the PMOKE response for the hybrid device in terms of the different targeted domain width (DW) of: (a) 60 nm, and (b) 200 and 570 nm, respectively. The inset in (a) shows the targeted bit in red (Up magnetization) is surrounded by oppositely magnetized bits in green (Down magnetization).

which the smallest readable bit size can be identified.

Figure 3(a) illustrates the evolution of the polarization rotation amplitude as the TE_0 light mode propagates through the waveguide of the hybrid device, where the targeted bit has a domain width (DW) of 60 nm and is surrounded by the oppositely magnetized background (red domain in green background in the inset in Fig. 3(a)). As light interacts with the magnetic racetrack (orange region), we can clearly see a rise in the rotation of the polarization as a result of PMOKE. Here, we also note that the small polarization oscillation before the magnetic section is due to the light reflection from the plasmonic nanoantenna and the magnetic racetrack back to the input. Due to the birefringence in the waveguide, we have a beating between the TE_0 mode and the PMOKE-induced TM_0 mode which results in the oscillation of the polarization rotation magnitude as light propagates through the waveguide.

By increasing the width of the targeted bit (see Fig. 3(b)) and comparing the results, we can observe that not only the magnitude of the polarization rotation changes but also the phase varies. This change is such that the Kerr rotation at $DW = 60$ nm is reversed by approximately 180° compared to DW widths of 200 and 570 nm. To explain the reason behind this reversal, it is important to note that the detected polarization rotation is the superposition of the contributions of both the targeted red domain and oppositely magnetized green domains. When the targeted bit width is very small, e.g. $DW = 60$ nm, the sum of the polarization rotations from the green domains dominate that of the targeted bit due to the small MO contribution from the targeted region. In this case, we cannot detect the magnetization state in the targeted bit unambiguously, since the outcome would depend on the content of neighboring bits. In contrast, when DW is either 200 nm or 570 nm (or any value in between), the polarization rotation of the targeted bit becomes larger than the sum of the polarization rotations of the oppositely magnetized regions. In this case, the magnetization state of the targeted magnetic bit can be uniquely identified.

Hence, we plotted the magnitude and phase of the polarization rotation for both the bare waveguide and hybrid devices as a function of DW in Fig. 4. As shown, a sudden transition in the phase of the polarization rotation is observed at DW for which the magnitude of the polarization rotation has a minimum. The minimum Kerr rotation happens at $DW = 120$ nm (200 nm) for the hybrid (bare waveguide) device, which is accompanied by a jump in the phase of the polarization rotation. For very small DW s, i.e. $DW < 120$ nm (200 nm) in the hybrid (bare waveguide) device, the PMOKE response of the targeted bit is weaker than the superposition of the rest of the bits in the racetrack due to the limited MO contribution from the targeted region. Thus, the oppositely magnetized bits in green determine the magnitude and phase of the polarization rotation. On the other hand, for $DW > 120$ nm (200 nm), the PMOKE response from this bit (red domain) has become dominant. Thus, the magnetization state in the targeted magnetic bit can be explicitly identified above this value of DW , regardless of the magnetization state in the rest of the racetrack.

Based on the results, we can see that the minimum footprint

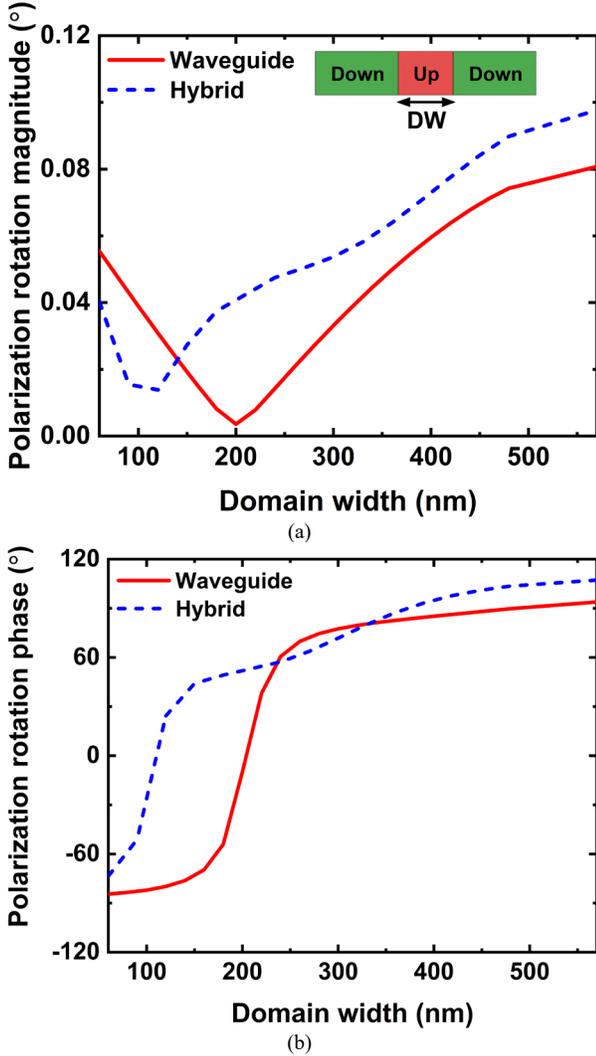


Fig. 4. The polarization rotation magnitude and phase in terms of DW for the bare waveguide and hybrid devices, respectively.

for determining the magnetization state in the targeted bit is DW = 120 nm (200 nm) for the hybrid (bare waveguide) device, which is a measure for the resolution of the device. Comparison of the performance of the bare waveguide and hybrid devices indicates that the hybrid device enhances the resolution of magnetization read-out due to magneto-plasmonic effects beyond the diffraction limit. In general, this section illustrated the possibility of optically reading of the magnetization states in magnetic bits with subwavelength sizes of down to $\sim 100 \times 100 \text{ nm}^2$ ($\sim \text{DW} \times w_{\text{MR}}$), regardless of the magnetization state in the rest of the racetrack.

B. Detecting the Change in Magnetization State

As mentioned earlier, change in the magnetization state of the targeted bit induces a phase difference between the TE_0 and emergent TM_0 modes due to PMOKE. To be able to detect such a phase difference, we introduce a method for converting the phase difference to an intensity variation of the TM_0 mode [11] with the aid of a plasmonic polarization rotator [23]–[26] (see Fig. 1(c)). With using the polarization rotator, as light passes through it, due to the partial conversion of the TE_0 mode to the TM_0 mode, the phase difference between the TE_0 and TM_0

modes gets transformed into an intensity variation of the TM_0 mode. In this way, the magnetization state is encoded in the TM_0 intensity. As stated in section II, the polarization rotator section is comprised of a bilayer of SiO_2/gold (from bottom to top) which is asymmetrically positioned on top of the InP waveguide. The design parameters and schematic of the device are shown in Fig. 1(c).

To get insight about transforming a phase change to an intensity variation of the TM_0 mode using the polarization rotator, we use the polarization ellipse to illustrate the polarization states of the light mode with and without the polarization rotator. More explanation can be found in the previous work of our group [11]. The polarization state shown by the polarization ellipse can be quantified using the Stokes parameters S_1 to S_3 as follows [27]:

$$S_1 = \cos 2\varepsilon \cos 2\theta, \quad (1)$$

$$S_2 = \cos 2\varepsilon \sin 2\theta, \quad (2)$$

$$S_3 = \sin 2\varepsilon, \quad (3)$$

where θ and ε are the polarization rotation (Kerr rotation) and ellipticity angle (Kerr ellipticity), respectively. The Stokes parameter S_1 to S_3 show that whether the light mode is a pure TE_0 or TM_0 mode (S_1), an elliptically polarized mode (S_2) or a circularly polarized mode (S_3). Table I shows the values of polarization rotation (θ) and ellipticity (ε) at the output of the bare waveguide and hybrid devices in terms of the magnetization states at DW = 200 nm, with and without the polarization rotator. Based on Table I, the values of θ and ε do not vary with the change in the magnetization state for the bare waveguide device with and without the polarization rotator. The reason is that at DW = 200 nm for the bare waveguide, the MO contribution is almost vanishing as shown in Fig. 4(a) which leads to $S_1 \approx 1$ and $S_{2,3} \approx 0$ for the cases of with and without the polarization rotator based on Eqs. (1) – (3). Figure 5(a) shows the polarization state of light (with some exaggeration for the sake of clarification) at the output of the hybrid device without the polarization rotator section for both up (red solid-line curve) and down (blue dashed-line curve) magnetization states, where two curves are overlapped. Based on the values of θ and ε in Table I and Eqs. (1) - (3), the Stokes parameters are $S_1 \approx 1$ and $S_{2,3} \approx 0$ for both up and down magnetization states. In this case, the effect of PMOKE-induced polarization rotation is so small that it cannot alter the polarization state of the input mode significantly, and consequently we have a TE_0 light mode without the polarization rotator section. On the other hand, Fig. 5(b) shows the polarization ellipse for the hybrid device in the presence of the polarization rotator section which shows the two curves do not completely overlap. According to Table I, the values of θ and ε differ by 0.2° and 0.1° , respectively, when the magnetization state changes. Based on the values of θ and ε in this case, $S_2 \approx 0.8$ and $S_3 \approx -0.5$ for both the magnetization states. In contrast, we have two slightly different values of $S_1^\uparrow \approx -0.251$ and $S_1^\downarrow \approx -0.256$ for up and down magnetization states, respectively. This difference in the values of S_1 originates from a variation in the intensity of the PMOKE-induced TM_0 mode, which explains the phase to intensity transformation that is used

TABLE I
POLARIZATION ROTATION (θ) AND ELLIPTICITY (ε), FOR BARE
WAVEGUIDE AND HYBRID DEVICES IN TERMS OF MAGNETIZATION STATES

| Kerr components | Bare waveguide | | | | Hybrid | | | |
|----------------------------|-------------------|-----------------|-----------------------|------|-------------------|------|-----------------------|------|
| | Up (\uparrow) | | Down (\downarrow) | | Up (\uparrow) | | Down (\downarrow) | |
| | WoP ^a | WP ^b | WoP | WP | WoP | WP | WoP | WP |
| θ ($^\circ$) | 0.003 | 53.4 | -0.003 | 53.4 | 0.04 | 53.3 | -0.04 | 53.5 |
| ε ($^\circ$) | 0.04 | 14.6 | 0.04 | 14.6 | 0.09 | 14.5 | 0.09 | 14.6 |

The values of θ and ε are for the target domain width (DW) of 200 nm.

^aWoP: Without the polarization rotator

^bWP: With the polarization rotator

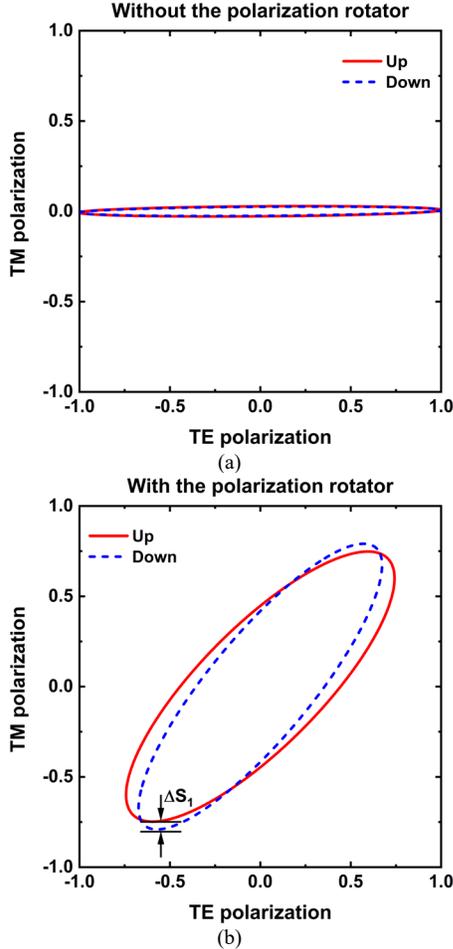


Fig. 5. Demonstration of transformation of the phase change to the intensity variation of a light mode propagating through the waveguide of the hybrid device using the polarization ellipse without and with the polarization rotator. The legends refer to up and down magnetization states. For the sake of clarity, the value of the magneto-optical Voigt constant, Q , is multiplied by 20.

to read out the magnetic bits. Hence, we defined a figure of merit ΔS_1 , i.e. the relative contrast of the TM_0 mode, as follows: $\Delta S_1(\%) = |S_1^\uparrow - S_1^\downarrow| \times 100$. (4) Based on Eqs. (1) and (4), for the bare waveguide $\Delta S_1 \approx 0$. In contrast, our hybrid device offers ΔS_1 of greater than 0.5% for the same targeted magnetic bit. One can see that using the bare waveguide, we cannot detect the change in the magnetization state in such a sub diffraction limit domain. However, our hybrid device with the help of magneto-plasmonic effects (offered by the proposed plasmonic nanoantenna) can overcome the diffraction limit and detect the magnetization change in a targeted bit size down to $\sim 100 \times 100 \text{ nm}^2$ in the presence of oppositely magnetized neighboring bits. Note that

the absorption losses by the plasmonic nanoantenna, magnetic racetrack, and polarization rotator are 8.68%, 9.45%, and 11.56%, respectively. The total reflection in the hybrid device is 4.84% and the transmission coefficient is 65.47%, which means that the insertion loss of the complete hybrid device is only -1.84 dB. As the SNR is inversely proportional to the square root of the insertion loss, we can see that this amount of insertion loss has negligible impact on the performance of our device.

IV. DISCUSSION

We numerically illustrated an ultracompact hybrid plasmonic-photonic device for optical read-out of the magnetization state in nanoscale magnetic bits. Using the hybrid device and a racetrack with a width of $\sim 100 \text{ nm}$, the read-out of magnetic bits with a nanoscale size down to $\sim 100 \times 100 \text{ nm}^2$ is possible using PMOKE, irrespective of the magnetization state in the rest of the racetrack. To complete the reading function, i.e. detecting the magnetization change on-chip, we proposed a method based on the polarization rotation principle. Our plasmonic polarization rotator transforms the PMOKE-induced polarization change in the TE_0 mode to an intensity variation of the TM_0 mode. Based on the simulation results, we showed that the hybrid device can detect the change in the magnetization state for a $\sim 200 \times 100 \text{ nm}^2$ targeted magnetic bit with a relative contrast of greater than 0.5%, while the bare waveguide device is not able to detect.

During recent years, few groups have reported on enhancing PMOKE using plasmonic effects. Maccaferri *et al.* presented a periodic array of nickel nano-ellipsoids to enhance PMOKE using plasmonic resonance of this element [28]. It is noteworthy to mention that unlike gold, nickel cannot offer a strong plasmonic resonance as it suffers from a strong interband absorption at resonance condition [29]. Luong *et al.* demonstrated that in contrast to a titanium-cobalt composite nanohole array, a silver-cobalt composite nanohole array can provide larger PMOKE enhancement due to plasmonic effects offered by silver [30]. Freire *et al.* exhibited a PMOKE enhancement with growing a periodic 2D array of [cobalt/platinum]₁₀ (10 is the number of the repetition of the bilayer) on a gold metal layer [31]. In spite of reporting on improved PMOKE by these groups, the approaches introduced in these works are not ideally suitable for optically reading of nanoscale magnetic bits. In addition, from fabrication point of view, the integration of such structures to photonic integrated circuits is difficult. Very recently, our group demonstrated the first experimental report on the on-chip MO reading of a diffraction-limited magnetic bit with a cross section of $600 \times 400 \text{ nm}^2$ and without the presence of oppositely magnetized nearby domains [11]. In contrast, our proposed device concept, by offering MO read-out beyond the diffraction limit, can play a key role in the realization of the future technology of hybrid spintronic-photonic memories with energy efficient switching and reading with high bit-rate data transfer and data storage capacity.

V. CONCLUSION

We introduced an ultracompact integrated hybrid plasmonic-photonic device for optical reading of nanoscale ferromagnet bits having perpendicular magnetic anisotropy on the IMOS platform. The hybrid device, which is based on coupling a doublet of V-shaped gold plasmonic nanoantennas on top of the waveguide, strengthen the MO interaction beyond the diffraction limit of light with the help of magneto-plasmonic effects. According to the simulation results, the hybrid device can make possible the identification of the magnetization state for $\sim 200 \times 100 \text{ nm}^2$ magnetic bits with a relative contrast of greater than 0.5%, but in general, targeted bits down to $\sim 100 \times 100 \text{ nm}^2$ can be unambiguously detected irrespective of the magnetization state in the rest of the racetrack. We believe this device can have a potential impact on direct optical read-out and can encode information in the optical state.

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