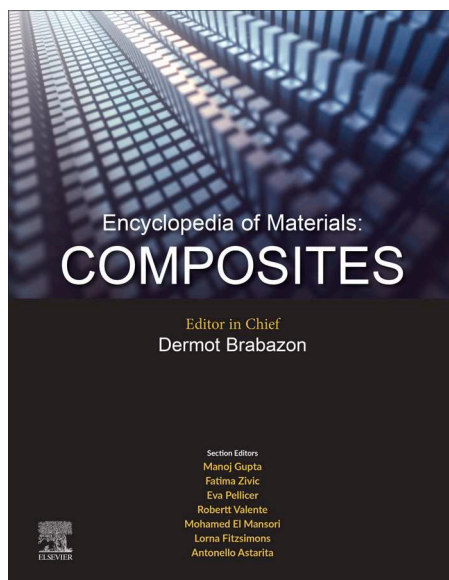


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## Nanoporous Composites With Converse Magnetolectric Effects for Energy-Efficient Applications

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### Nomenclature

ALD Atomic layer deposition

BTO Barium titanate, BaTiO<sub>3</sub>

CME Converse magnetolectric effect

DME Direct magnetolectric effect

EISA Evaporation induced self-assembly

FE Ferroelectric

FM Ferromagnetic

ME Magnetolectric

MOKE Magneto-optic Kerr effect

PC Propylene carbonate

PMN-PT Lead magnesium niobate-lead titanate, Pb(Mg<sub>1/3</sub>Nb<sub>2/3</sub>)O<sub>3</sub>-PbTiO<sub>3</sub>

PZT Lead zirconate titanate, Pb[Zr<sub>(x)</sub>Ti<sub>(1-x)</sub>]O<sub>3</sub>

SQUID Superconducting quantum interference device

VSM Vibrating sample magnetometry

### Glossary

**Converse magnetolectric effect** Modulation of magnetic properties in magnetic materials by means of external electric field.

**Direct magnetolectric effect** Modulation of electric polarization in dielectric materials by external magnetic fields.

**Joule heating effect** Energy dissipation induced by electrical currents flowing through a resistor.

**Magnetolectric composites** Heterostructures which comprise magnetic and dielectric counterparts in which the properties of the magnetic material can be manipulated with voltage, either using solid or liquid electrolytes.

**Magneto-ionics** Magnetolectric mechanism in which the oxidation state of the metal in the magnetic phase is

influenced by the ions (e.g., O<sup>2-</sup>) diffusion back and/or forth from the material of interest toward an ion source/sink (e.g., a high O<sup>2-</sup> mobility thin film, deposited next to the magnetic layer), depending on the voltage polarity.

**Magnetostriction** Ability of most ferromagnetic materials to expand or contract in response to an external magnetic field.

**Single-phase multiferroics** Materials possessing an inherent coupling between magnetic and electric orders.

**Spinels** A class of material, which crystallizes in the cubic crystal system, with general formula AB<sub>2</sub>X<sub>4</sub>, where A and B are cations of (II) or (III) valent element (or the same element, as in Fe<sub>3</sub>O<sub>4</sub>) and X is an anion (typically oxygen or sulfur).

### Introduction

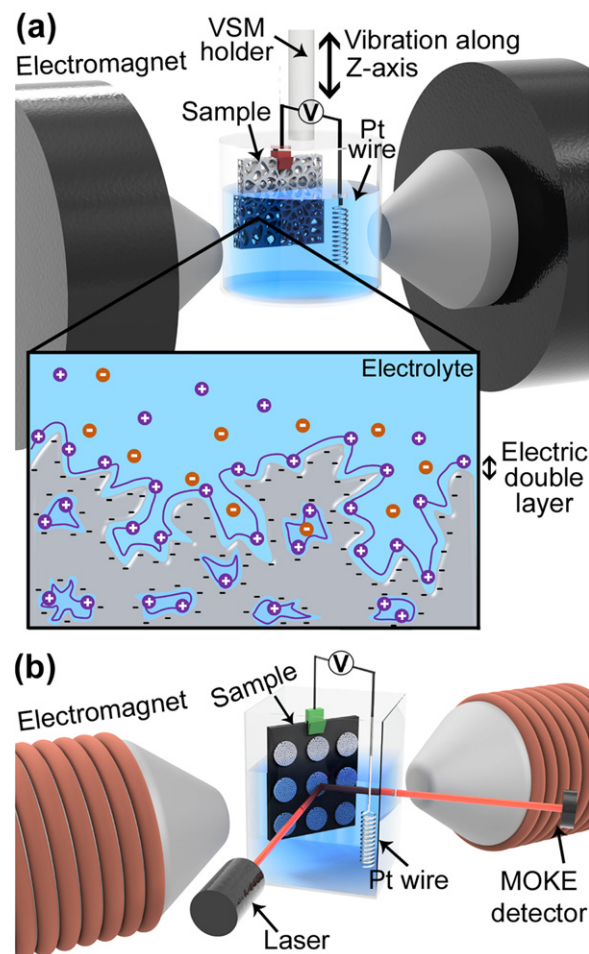
Magnetism and electricity have always had an intimate link. Particularly interesting is the coexistence of magnetic and electric orders in magnetolectric (ME) materials which makes them able to respond, simultaneously, to external magnetic and electric stimuli: (1) Electric polarization can be modulated by external magnetic fields (direct ME effect, DME) and (2) magnetic properties can be largely controlled with an electric field (converse ME effect, CME) (Wang *et al.*, 2010; Hu and Nan, 2019). In conventional ME composites, coupling between piezoelectric/ferroelectric (FE) and ferromagnetic (FM) constituents is mediated by strain and, in some cases, by electric charge effects (Wang *et al.*, 2010; Hu and Nan, 2019; Molinari *et al.*, 2019; Navarro-Senent *et al.*, 2019). In recent years, progress has been made in the so-called magneto-ionic heterostructures, where voltage-driven magnetic changes are mediated by ion migration/intercalation (typically O<sup>2-</sup>, Li<sup>+</sup> (Dasgupta *et al.*, 2014; Zhang *et al.*, 2016) or H<sup>+</sup> species (Tan *et al.*, 2019)) (Bauer *et al.*, 2015; Duschek *et al.*, 2016; Gilbert *et al.*, 2016). DME effects are appealing for healthcare technologies (Chen *et al.*, 2017), water remediation systems (Mushtaq *et al.*, 2019a,b) and sensors/actuators (Wang *et al.*, 2010), whereas CME effects can be exploited in microelectromechanical systems and energy-efficient magnetic memories (Peng *et al.*, 2016).

Unfortunately, in FM films directly grown onto FE substrates, the voltage required to generate ME effects is extremely high (e.g., 4 kV) (Ahmad *et al.*, 2015) due to the large thickness of the substrate (i.e., in capacitors, electric field is inversely proportional to the dielectric thickness). This is not suitable for microelectronics, where much lower voltages (<10 V) are desirable to enhance energy efficiency and not to burn the electronic components. If thin FE/FM bilayers directly grown onto rigid (non-FE) substrates, then the required voltages are lower, but the attainable strain is small due to the clamping with the substrate, thus also limiting the extent of ME effects (Torah *et al.*, 2004; Schmitz-Antoniak *et al.*, 2013; Chien *et al.*, 2016). To overcome these drawbacks, new ME composites based on nanoporous materials filled with either liquid or solid dielectric materials (eventually FE polymers) have been recently developed. These types of materials are overviewed in this article.

During the last two decades, much progress has been made in the synthetic procedures to prepare nanoporous systems with precise control of the porosity degree, pore/ligament sizes and crystallinity of the pore walls. This development has been triggered by the widespread use of these materials in chemistry areas such as catalysis, gas sensors, or supercapacitors, where a high surface-to-volume ratio ( $S/V$ ) favors an enhanced performance. Also remarkable is the literature dealing with the growth and characterization of mesoporous composites containing embedded ferromagnetic particles, to be used in drug delivery applications (Trewyn *et al.*, 2007), ultra-light magnets (Heiligtag *et al.*, 2014) or water remediation (Kharissova *et al.*, 2015). Interestingly, although many of the cutting-edge technological applications in nanomagnetism/spintronics also rely on surface or interface magnetic phenomena (e.g., spring-magnets, exchange bias, skyrmions, etc.) (Bibes *et al.*, 2011), the use of nanoporous materials in this field has been mainly overlooked.

In the particular subject of voltage-driven magnetic actuation, the study of nanoporous materials (i.e., mesoporous metals or oxides filled with a liquid high dielectric constant, or metal/metal oxide porous nanocomposites), all with high  $S/V$  ratio, is relatively new. In such materials, the relatively small effects from seminal works on magnetoelectric phenomena, initially observed in ultra-thin metallic films (Weisheit *et al.*, 2007), are drastically enhanced (Molinari *et al.*, 2018; Navarro-Senent *et al.*, 2019). In fact, it is easy to calculate that the  $S/V$  ratio of a 50-nm-thick porous film covering an area of  $1 \times 1 \mu\text{m}^2$  and being made of a square array of vertically oriented cylindrical pores, with 5 nm pore diameter and 2.5 nm average pore wall (i.e.,  $\sim 50\%$  total porosity) is about 40 times larger than the  $S/V$  ratio of a fully dense layer of 50 nm thickness. Hence, the voltage-driven magnetic interfacial effects that depend on the  $S/V$  ratio (i.e., magnetic anisotropy, coercivity, direction of the magnetization, exchange bias), as well as those proportional to the total volume of affected material (e.g., the net magnetic moment) can be significantly enhanced in the porous architecture, by several orders of magnitude, with respect to non-porous magnetic thin films.

To measure the CME in nanoporous materials, either conventional magnetometry (vibrating sample or SQUID magnetometers) or magneto-optic Kerr effect techniques, in conjunction with custom-made setups to apply voltage, can be utilized (see Fig. 1). Either solid dielectrics (prepared by, e.g., atomic layer deposition) or liquid electrolytes, which form the so-called electric double layer upon voltage application, can be used.



**Fig. 1** Schematic illustration of the experimental setups typically used for magnetolectric measurements in liquid configuration using (a) a vibrating sample magnetometer (VSM) and (b) a magneto-optic Kerr effect setup (MOKE).

The main body of this article is divided into three sections, covering CME in (1) metal-based composites, (2) metal/metal oxide systems and (3) all-oxide heterostructures, all based on nanoporous architectures. Here the use of the “composite” term is extended (as in the review paper by [Molinari et al., 2019](#)) to include those materials comprising a nanoporous framework filled with either a solid or a liquid second phase (often needed to generate the electric field or as buffer materials for magneto-ionics). We begin each section with a short introduction on the growth methods and experimental procedures used to induce and control the nanoporosity. We then focus on the most recent progress on the nanoporous ME composites belonging to each of the three categories, with emphasis on the systems where the ME effects are exacerbated due to the presence of nanoporosity. The article is followed by brief concluding remarks and perspectives.

## Magnetoelectric Composites Based on Nanoporous Metals or Metal Alloys

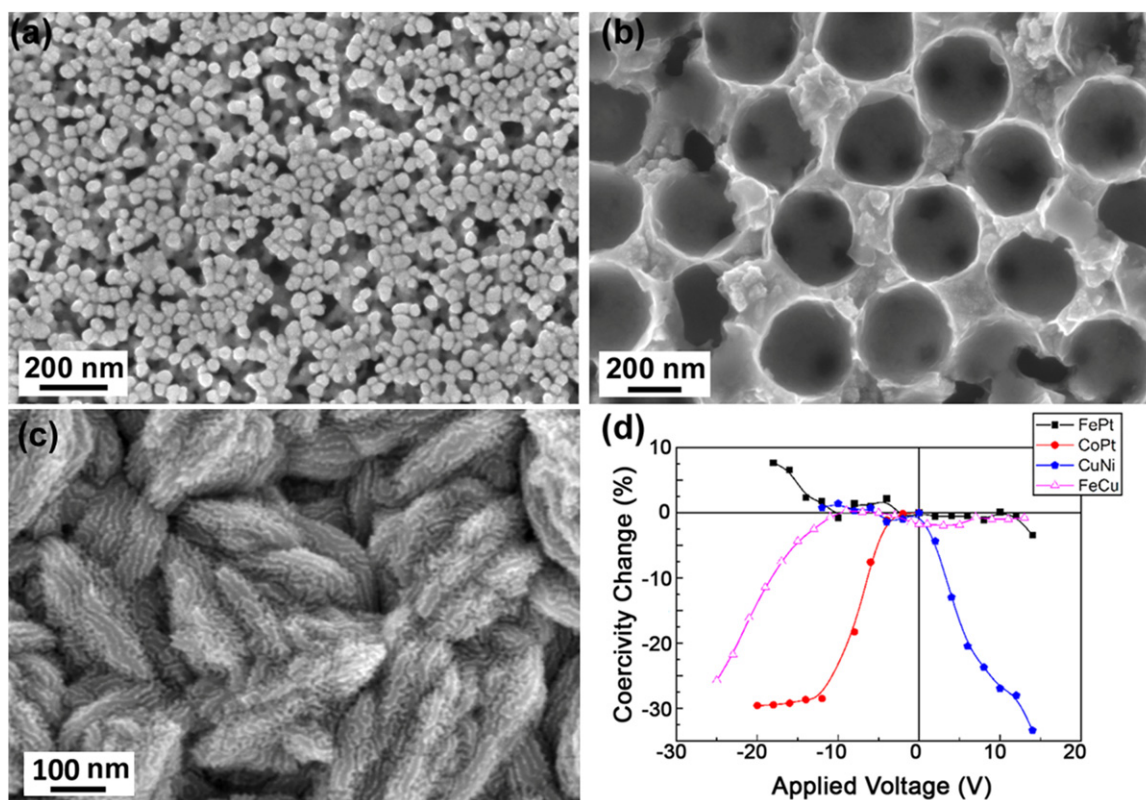
Based on previous works on magnetism tuned by carrier density modulation in thin film magnetic semiconductors ([Ohno et al., 2000](#); [Chiba et al., 2003](#)), as well as preliminary theoretical works that predicted a strong dependence of the magnetic anisotropy energy on the number of valence band electrons in alloys such as Fe–Pt or Co–Pt ([Daalderop et al., 1991](#)), [Weisheit et al.](#) demonstrated in 2007 that the coercivity of ultra-thin Fe–Pt and Fe–Pd films could be modified by up to 4.5% under the action of external voltage during electrolyte gating ([Weisheit et al., 2007](#)). Since electric field in metals is confined within the Thomas–Fermi screening length (which is of the order of 0.5 nm), the changes in coercivity driven by voltage in metallic alloys were initially only observed in films of 2–3 nm in thickness.

In parallel to these works on ultra-thin metallic films, [Weissmüller et al.](#) showed that reversible changes in strain could be induced in nanoporous Pt due to electric surface charge accumulation ([Weissmüller et al., 2003](#)). Further investigations in that direction revealed that the magnetic susceptibility of nanoporous Pd ([Drings et al., 2006](#)) and the ferromagnetic properties of Pd–Co ([Ghosh et al., 2006](#)), Pd–Ni ([Ghosh, 2011, 2013](#)) and Au–Fe ([Mishra et al., 2010](#)) nanoporous alloys could be also manipulated with voltage due to these electrically induced variations of strain. In these works, the typical size of the crystallites comprised in the porous frameworks ranged between 8 and 20 nm. The charge-induced pressure on the nanocrystallites comprised in the nanoporous frameworks and the concomitant changes in sample dimensions during electrochemical charging and discharging processes were found to be responsible for these observations. In spite of their fundamental interest, the typical changes in magnetization due to strain effects are usually rather small (up to 0.5% at room temperature and 3% near the Curie temperature; [Ghosh et al., 2006](#)). An interesting theoretical work by [Subkow and Fähnle \(2011\)](#) analyzed in detail the effects of (1) filling of the electronic  $d$  band upon electric charging and (2) the magnetoelastic response to charge-induced strain variations in the nanoparticles forming nanoporous systems, both of which play a role during magnetoelectric actuation of these metallic alloys using liquid electrolytes. The authors concluded that the magnetic anisotropy energy of this type of systems can be modified with voltage only at the surface of nanoparticles (constituting the nanoporous framework) that are not in grain–boundary contact, which are the ones that can be effectively charged in the electrolyte due to the formation of the electric double layer when voltage is applied.

Subsequent works have been carried out to try to enhance the effects of voltage in the magnetic properties of nanoporous metallic alloys. Different experimental techniques have been used to improve the quality of the synthesized materials (i.e., to reduce the pore wall width or ligament size, or to make the porosity pseudo-ordered). Some of the most standard methods include micelle-assisted electrodeposition, electrochemical dealloying, or electrodeposition combined with colloidal templating (see [Fig. 2](#)).

Electrochemical dealloying is a method that allows fabricating porous films from previously deposited fully dense counterparts. In this case, the porous metal or metal alloy is obtained by selectively dissolving the most electrochemically active element comprised in an alloy. This technique has been shown successful to generate porosity in bulk metals, ribbons, as well as thin films (the latter grown by, e.g., electrodeposition or sputtering, as in the example shown in [Fig. 2\(a\)](#)) ([Sun et al., 2004](#); [Robbenolt et al., 2018](#)). Colloidal templating is a facile and cost-effective method to create well-defined, pseudo-ordered 3D porous structures by using 3D assembled nanospheres as a soft mask ([Dislaki et al., 2018](#)). Colloid size and self-assembly are factors that determine the nanopore dimensions, pore wall width and overall layer thickness (see [Fig. 2\(b\)](#)). Various methods to assemble monodisperse micro- and nanospheres have been put forward in the literature, such as dip coating, spin-coating, self-assembly, solvent evaporation, or electrophoresis. Among these methods, electrophoretic deposition is amongst the most attractive options for achieving homogeneous coverage of close-packed arrangements of spheres. Finally, micelle-assisted electrodeposition relies on the formation of micelles in the electrolytic bath which act as soft templates (structure-directing agents) during the growth of the metallic layers of interest. When a block co-polymer (e.g., P123) is added to the electrolyte at a concentration above “critical micelle concentration,” micelles start to form spontaneously in the solution, getting progressively in contact and tending to self-assemble at the solid–liquid interface ([Quintana et al., 2017](#); [Isarain-Chávez et al., 2018](#); [Li et al., 2018](#)). When voltage is applied to the cathode and reduction of the metal cations takes place, these micelles interfere and guide the electrodeposition process, leading to the growth of mesoporous metallic films (see [Fig. 2\(c\)](#)). The pore size (and pore wall size) is typically very small (even sub-5 nm). If the surfactant concentration is high enough to form a lyotropic liquid crystal, regular arrangements of pores can be obtained.

Remarkably, using this structure-optimized nanoporous alloys, changes up to 32% in coercivity (with virtually no significant variations in the saturation magnetization) have been reported in recent years in relatively thick films (0.5–1  $\mu\text{m}$ ) of various alloy compositions (nanoporous Cu–Ni, Co–Pt, Fe–Cu), mainly due to changes in the magnetic anisotropy energy stemming from electric surface charge accumulation, as evidenced by ab-initio calculations ([Quintana et al., 2017](#)). Depending on the alloy system,



**Fig. 2** Representative scanning electron microscopy (SEM) on-top images of (a) dealloyed Fe-Cu sputtered films, (b) electrodeposited Fe-Cu porous films grown onto colloidal templated substrates, (c) nanoporous Co-Pt films prepared by micelle-assisted electrodeposition. Panel (d) shows the relative variations of coercivity, as function of voltage, in various porous alloys electrolyte-gated using an anhydrous electrolyte (propylene carbonate). In all these cases, the voltage effects are mainly due to surface charge accumulation. Reprinted from (d) Navarro-Senent, C., Quintana, A., Menéndez, E., Pellicer, E., Sort, J., 2019. Electrolyte-gated magnetolectric actuation: Phenomenology, materials, mechanisms, and prospective applications. *APL Materials* 7 (3), 030701. doi:10.1063/1.5080284.

the coercivity can either increase or decrease for either positive or negative applied voltages (Fig. 2(d)). Also, the changes in coercivity as a function of voltage polarity are often not symmetric. This can be due to the way the magnetic anisotropy energy varies with the electric field or because of the dissimilar thickness (and nature) of the electric double layer formed at the interfaces between the pore walls and the electrolyte during positive/negative voltage application. The experimental demonstration that the coercivity of thick nanoporous magnetic alloy films can be drastically reduced by simply subjecting this type of materials to the action of an electric field (i.e., applying a moderate DC voltage to them) is very appealing for energy-efficient magnetic actuation. If the coercivity is reduced, lower magnetic fields are needed to induce magnetization reversal (i.e., to write the magnetic bits of information in memory devices). As a consequence, less electric current is required, and this concurrently lowers undesirable Joule heating effects. The large S/V ratio and the ultra-narrow pore walls of these systems play a crucial role in the observed effects. Since electric field in metals is confined at their surface, the nanoporous morphology of the investigated materials allows for much larger accumulation of electrostatic charges compared to fully-dense films. That is, the whole nanoporous structure is affected by the electric field and not only the outer topmost surface, thus resulting in a very significant voltage-induced reduction of coercivity. This waives the stringent “ultrathin-film requirement” from previous studies, wherein smaller voltage-driven coercivity variations were reported (Weisheit *et al.*, 2007). Finally, besides coercivity, tuning of the superparamagnetic state has been achieved in nanoporous Co-Pd alloys electrically gated in a 1 M aqueous KOH electrolyte (Göbller *et al.*, 2019). The effect seems to be related to electrochemical hydrogen sorption. More specifically, a strong magneto-ionic ( $H^+$ ) effect arises from coupling of the magnetic clusters via a Ruderman-Kittel-Kasuya-Yoshida-type interaction in the Pd matrix which is enhanced by the hydrogen sorption.

An overview of the different types of electrolytes, materials and main magnetolectric effects observed in nanoporous metallic alloys, metal/metal oxide composites and pure oxide systems is given in Table 1.

### Magnetolectric Composites Based on Nanoporous Metal/Metal Oxide Heterostructures

In most of the examples discussed in the previous section, an anhydrous electrolyte (e.g., propylene carbonate) was utilized in order to induce electric charge accumulation at the surface of nanoporous metallic alloys (through the formation of the

**Table 1** Representative materials, typical electrolytes and most prominent magnetolectric effects induced in nanoporous materials with voltage (electrolyte-gating)

Electrolyte	Material	Variation	Reversibility	References
LiClO <sub>4</sub> in EC	Nanoporous Au–Fe alloy	M = 0.2%	Yes	Mishra <i>et al.</i> (2010)
LiClO <sub>4</sub> in EC	Nanoporous Pd–Ni alloy	M ≈ 0.5%	Partially ( $\Delta V$ , ~1 h) Yes ( $\Delta V$ , ~1 h)	Ghosh (2011)
PC	Nanoporous Cu–Ni film	H <sub>C</sub> = –32%	–	Quintana <i>et al.</i> (2017)
PC	Pseudo-ordered porous Fe–Cu film	H <sub>C</sub> = –25%	Partially (0 V, 5 min)	Dislaki <i>et al.</i> (2018)
LiClO <sub>4</sub> in EC	Nanoporous Pd–Co alloy	M = 3%	Partially ( $\Delta V$ , ~2 h)	Ghosh <i>et al.</i> (2006)
1 M KOH	Nanoporous Pd–Ni alloy	M = 24.6%	Yes ( $\Delta V$ , ~2 h)	Ghosh (2013)
1 M KOH	Nanoporous Co–Pd alloy	M = 100% (ON–OFF)	Yes ( $\Delta V$ , ~ h)	Göbler <i>et al.</i> (2019)
1 M NaOH	Nanoporous Cu–Ni	M <sub>S</sub> = +33%	Yes ( $\Delta V$ , 10 min)	Quintana <i>et al.</i> (2018a)
PC	Nanoporous Co–Pt lithographed disks	$\Delta M_S$ = 66% $\Delta H_C$ = –88%	Partially ( $\Delta V$ , ~ min)	Navarro-Senent <i>et al.</i> (2018)
PC	Nanoporous Co–Pt + HfO <sub>2</sub> /Al <sub>2</sub> O <sub>3</sub> films	$\Delta m_S$ = 76% $\Delta H_C$ = –58%	Partially ( $\Delta V$ , ~2 h)	Navarro-Senent <i>et al.</i> (2020)
1 M KOH	Porous $\gamma$ -Fe <sub>2</sub> O <sub>3</sub> /Pt-nanocomposite	$\Delta m$ = 10.4%	Yes ( $\Delta V$ , ~5 min)	Topolovec <i>et al.</i> (2013)
LITFSI in EMIM-TFSI	Nanoporous Co <sub>0.5</sub> Ni <sub>0.5</sub> Fe <sub>2</sub> O <sub>4</sub> and CoFe <sub>2</sub> O <sub>4</sub> films	$\Delta M$ ≈ 2–5%	Yes ( $\Delta V$ , ~ min)	Dubraja <i>et al.</i> (2018)
LITFSI in EMIM-TFSI	Nanoporous LiFe <sub>5</sub> O <sub>8</sub> films	$\Delta M$ ≈ 4%	Yes ( $\Delta V$ , ~ h)	Reitz <i>et al.</i> (2016)
PC	Nanoporous Fe–Cu films	$\Delta M_S$ = 20% $\Delta H_C$ = 100%	Partially ( $\Delta V$ , 40 min)	Robbenolt <i>et al.</i> (2018)
PC	Nanoporous CoFe <sub>2</sub> O <sub>4</sub> films	$\Delta M_S$ = 15% $\Delta H_C$ = +28%	Yes ( $\Delta V$ , ~2 h)	Robbenolt <i>et al.</i> (2019a)
PC	Nanoporous CoFe <sub>2</sub> O <sub>4</sub> + HfO <sub>2</sub> films	$\Delta M_S$ = +56% $\Delta H_C$ = +69%	Yes ( $\Delta V$ , ~1 h)	Robbenolt <i>et al.</i> (2020)
PC	Nanoporous FeO <sub>x</sub> films	$\Delta M_S$ = 1310% $\Delta H_C$ = +100%	Partially ( $\Delta V$ , ~1 h)	Robbenolt <i>et al.</i> (2019b)

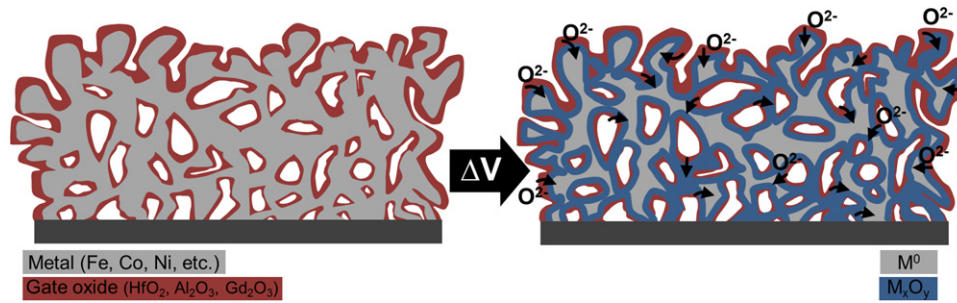
Abbreviation: EC, ethylene carbonate; EMIM-TFSI, 1-ethyl-3-methylimidazolium bis(trifluoromethylsulfonyl)imide; H<sub>C</sub>, coercivity; LITFSI, lithium bis(trifluoromethanesulfonyl)imide; m, magnetic moment; M, magnetization; M<sub>S</sub>, saturation magnetization; PC, propylene carbonate.

Note: Table adapted from Navarro-Senent, C., Quintana, A., Menéndez, E., Pellicer, E., Sort, J., 2019. Electrolyte-gated magnetolectric actuation: Phenomenology, materials, mechanisms, and prospective applications. *APL Materials* 7 (3), 030701. doi:10.1063/1.5080284.

electric double layer). As aforementioned, this causes electric field induced changes in the magnetic properties (converse magnetolectric effect), which are often reversible. In some cases, though, an aqueous electrolyte (e.g., 1 M KOH or NaOH) has been purposely utilized in the literature to trigger reduction–oxidation electrochemical reactions at the surface of the metallic porous frameworks or in metal/metal oxide islands when voltage is applied (Duscek *et al.*, 2018; Quintana *et al.*, 2018a). Sometimes, unexpected results have been encountered. For example, when positive voltage is applied to electrodeposited nanoporous Cu–Ni alloys and the films get oxidized, an increase (rather than a decrease) of the saturation magnetization is observed (Quintana *et al.*, 2018a). This counter-intuitive result can be understood because of the preference for Cu (rather than Ni) to get selectively oxidized. As a consequence, upon application of a positive voltage, the Cu–Ni alloy gets progressively enriched in Ni (as Cu is oxidized) and the resulting magnetic moment increases. The process can be fully reversed by application of a negative voltage. Oxygen migration can also cause changes in the coercivity of the nanoporous alloy, as well as in the perpendicular magnetic anisotropy (Ibrahim *et al.*, 2018). However, in nanoporous systems, since shape anisotropy is not clearly defined (i.e., these materials are typically polycrystalline and the ligaments and pore walls are randomly oriented), no pronounced anisotropy effects are typically reported.

Some nanocomposite porous materials contain oxide phases in the as-prepared states (as depicted in Fig. 3) (Navarro-Senent *et al.*, 2018; Robbenolt *et al.*, 2018). Some examples are nanoporous Co–Pt + CoO<sub>x</sub> lithographed disks prepared by electrodeposition or nanoporous Fe–Cu + FeO<sub>x</sub> + CuO<sub>y</sub> films prepared by electrochemical dealloying. In these cases, migration of the structural oxygen (i.e., oxygen present in the as-prepared material) can occur when voltage is applied, even when utilizing anhydrous electrolytes, thereby rendering truly magneto-ionic effects (Navarro-Senent *et al.*, 2018). When this occurs, drastic changes in the coercivity and the magnetic moment at saturation take place. In some cases (e.g., nanoporous Fe–Cu solid solutions), this oxygen migration causes structural transformations in the metallic counterpart (e.g., from Cu-rich face centered cubic to Fe-rich body centered cubic phases) (Robbenolt *et al.*, 2018). Eventually, ON-OFF transitions (reversible transformations from ferromagnetic to paramagnetic or antiferromagnetic states) might be induced in nanoporous metal + metal oxide thin films, similar to what has been observed in nanostructured oxide thin films (Duscek *et al.*, 2018; Quintana *et al.*, 2018b).

Another effect that could be of interest in this type of nanocomposites would be the study of the eventual exchange bias phenomenon (i.e., shift of the hysteresis loop, along the magnetic field axis, due to the coupling between ferromagnetic and the electrochemically generated antiferromagnetic or ferrimagnetic phases). While this has been studied in flat, non-porous, multi-layered films (Gilbert *et al.*, 2016; Zhou *et al.*, 2016), comprehensive studies in nanoporous composites have not yet been reported.



**Fig. 3** Schematic drawing of voltage-induced oxygen migration in metal (M)/metal oxide ( $M_xO_y$ ) magnetolectric composite under the action of an applied voltage,  $\Delta V$ .

### All-Oxide Nanoporous Magnetolectric Composites

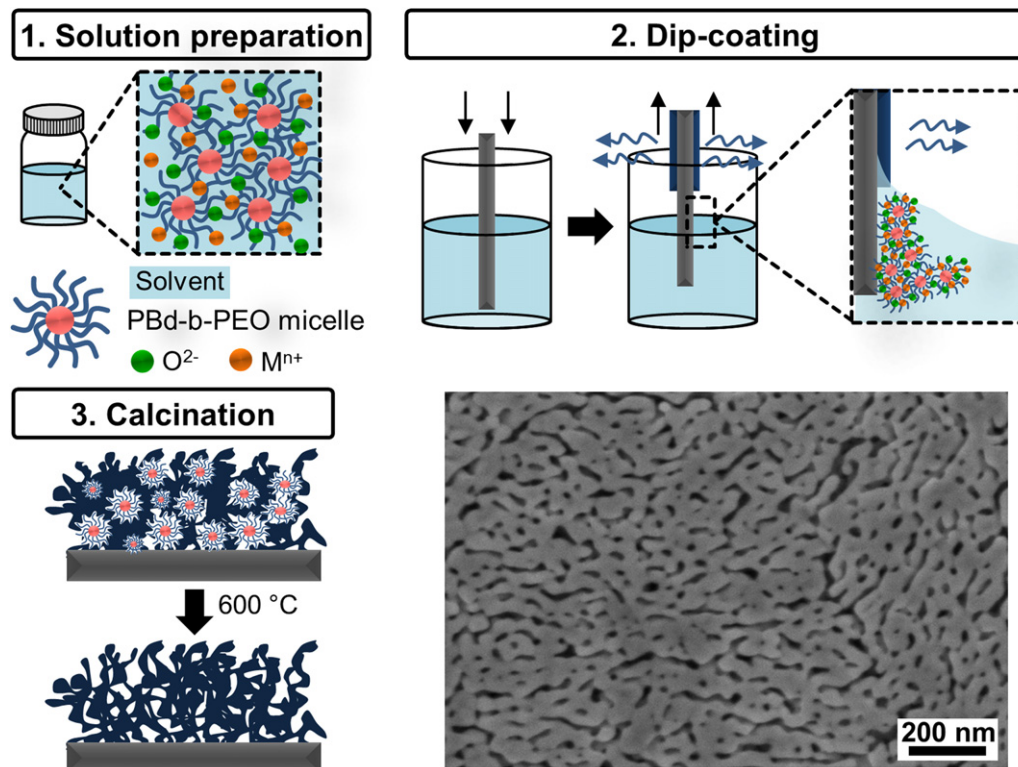
Up to date, the vast majority of magnetolectric studies (in liquid and solid state) have been performed in oxide systems. As a matter of fact, all intrinsic multiferroics are oxides, for example, spinels ( $Fe_3O_4$ ,  $CoFe_2O_4$ ), perovskites ( $BiFeO_3$ ,  $BaTiO_3$ ) (Liu and Yang, 2017) and hexagonal manganites ( $RMnO_3$ , where R is a rare-earth element) (Vaz *et al.*, 2010; Rao *et al.*, 2012). Furthermore, many oxides exhibit high magnetostriction coefficients ( $CoFe_2O_4$ ,  $La(Sr)MnO_3$ ) or high piezoelectric constants (PZT, PMN-PT, BTO) that makes them appealing for the design of various composite architectures where the magnetolectric effects are mediated by strain. Strain coupling requires an optimum mechanical matching between the two constituent phases, which is sometimes very difficult to achieve in ceramics. For this reason, the effects are only observed in epitaxial composite films or patterned structures where the crystallographic orientation and interfacial roughness are accurately controlled. Nevertheless, even in epitaxially grown composite films, the changes in magnetization induced by voltage are typically rather low (< 5%), due to clamping effects with the substrate, which strongly limit the available strains.

Apart from the strain-mediated mechanism, an applied electric field can modulate the charge carrier density in oxide composite heterostructures, inducing noticeable changes in magnetization and coercivity. However, a direct demonstration of modulation of the magnetic properties in oxides heterostructures occurring through this mechanism remains relatively less explored (Schmitz-Antoniak *et al.*, 2013; Chu *et al.*, 2018). Furthermore, the presence of oxygen vacancies in the oxides has been shown to contribute to the magneto-ionic mechanism, where, as aforementioned, the voltage-driven oxygen ions migration can cause significant changes in coercivity, magnetization and magnetic anisotropy (Chien *et al.*, 2016). Various systems have been comprehensively analyzed and reviewed in recent articles (Song *et al.*, 2017; Chu *et al.*, 2018; Molinari *et al.*, 2019; Navarro-Senent *et al.*, 2019). In this section, we will focus specifically on the several oxide systems where the occurrence of nanoporosity has brought a significant contribution to the observed ME effects: Multiferroic  $BiFeO_3$ , solid state  $CoFe_2O_4/PZT$ , solid/liquid  $CoFe_2O_4/PC$ , and  $FeO_x/PC$  composites, and lithiated spinel ferrites.

The synthesis of the nanoporous oxides traditionally relies on the evaporation induced self-assembly (EISA) of the sol-gel type inorganic precursors and the amphiphilic diblock co-polymer templates, as depicted in Fig. 4. In aqueous or organic solvents, the precursors (e.g., metal chloride or nitrate) are hydrolysed and condensed to form inorganic polymers composed of M–O–M bonds (M is the metal). At the same time, the polymer forms micelles which attach to the metal ions (Step 1). During the dip-coating process, the polymer and sol-gel precursors are progressively concentrated by evaporation, leading to aggregation, gelation, and final drying to form a type of dry gel layer (Step 2). As a final step, the layers are heated in a controlled atmosphere in order to crystallize the oxide and burn out the polymer template which leaves behind the porous structure (Step 3). Remarkably, the nanoporous films can be also obtained directly by dip-coating, without the use of the polymer template (Robbenolt *et al.*, 2019b). In this case the degree of nanoporosity is governed by the solvent evaporation rate. A representative image of the  $FeO_x$  film produced by dip-coating in the absence of polymer template is also given in Fig. 4. This is a simple and relatively fast method to grow nanoporous oxide films with the controlled thickness, porosity and crystalline structure. More details on the methodology and the effect of the synthesis parameters can be found elsewhere (Schneller *et al.*, 2013; Zhang *et al.*, 2019; Galy *et al.*, 2020).

The mechanical flexibility brought by nanoporosity can provide a powerful tool to strain engineer nanostructured materials (Quickel *et al.*, 2010). The studies performed on internal multiferroic  $BiFeO_3$  (Quickel *et al.*, 2015) revealed that the voltage-induced change in  $d$  spacing was 10 times greater than in the dense film of the same composition. However, it is not only the extent of the lattice strain that matters, but also the nature of strain. As the material is clamped to a rigid substrate, there is a gradient of strain (anisotropic strain) developed during the application of electric field. The strain symmetry lowering appears to be the key to the observed amplified voltage-induced magnetization of  $BiFeO_3$ , which exceeds the one of epitaxial  $BiFeO_3$  systems. Indeed, the numerical results are impressive: upon electric field application the material shows a large change in saturation magnetization, from 0.04 to 0.84  $\mu_B$  per Fe, which is >20 times higher than in the dense sample.

This seminal work by Quickel *et al.*, paved a way towards the use of flexible nanoporous architectures for strain engineered ME composite materials. The later reports, where the pores are employed as a host to accommodate dielectric compounds, have proved the importance of nanoporosity for strain-engineered oxide systems. Nanoporous  $CoFe_2O_4$  was produced using polymer templating of sol-gel derived thin films, followed by conformal filling the pores with a piezoelectric  $Pb(Zr,Ti)O_3$  using atomic



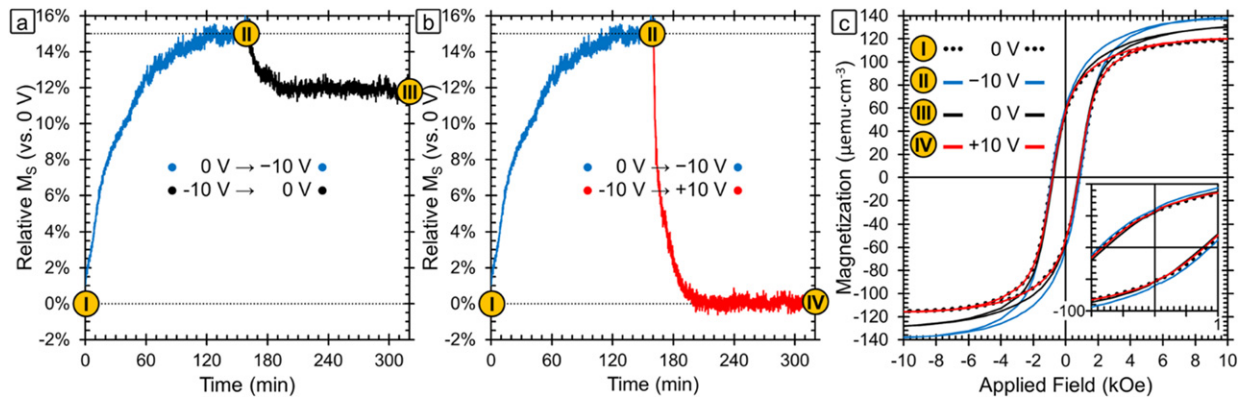
**Fig. 4** Schematic pictures of the template-assisted sol-gel process used to induce nanoporosity in oxide films and a representative SEM image of nanoporous FeO<sub>x</sub> layer obtained after calcination.

layer deposition, which resulted in a solid state composite ME material with a very high interfacial area (Chien *et al.*, 2016). Both materials are well documented in the literature and have generated significant interest due to their ferroic properties (Schmitz-Antoniak *et al.*, 2013; Rodzinski *et al.*, 2016; Paudel *et al.*, 2019). However, previous studies were focused on dense films where the magnetolectric effects are significantly limited by substrate clamping and the effects of nanoporosity were mainly overlooked. In porous CoFe<sub>3</sub>O<sub>4</sub> films grown on rigid substrates, the bottom of the layer remains clamped and therefore the in-plane strain is hindered. Nevertheless, the nanoporosity allows for out-of-plane flexing and thus, for more bond distortion in both Pb(Zr,Ti)O<sub>3</sub> and in the coupled CoFe<sub>3</sub>O<sub>4</sub> as compared to dense films. Indeed, with the application of positive voltage, the most pronounced change in saturation magnetization was measured along the perpendicular-to-plane direction, that is,  $M_S$  increases by 15.4% (compared to 3.6% for in-plane measurements). Remarkably, the observed ME effects are reduced when the thickness of the Pb(Zr, Ti)O<sub>3</sub> increases. This again underscores the importance of the mechanical flexibility of the nanoporous frameworks, as thicker piezoelectric layer seals the pores, thus lowering the mechanical distortion and the strain transfer. However, the voltage-induced strain in Pb(Zr,Ti)O<sub>3</sub> is lost when the driving electric field is removed, that is, the observed effects are rather volatile.

Non-volatile, fully reversible electric-field control of magnetism in nanoporous CoFe<sub>3</sub>O<sub>4</sub> was achieved by magneto-ionic means, where a liquid electrolyte (propylene carbonate treated with metallic Na) was used to access the open porosity of the material (Robbenolt *et al.*, 2019a). Applying moderate negative voltage to the sample, between  $-10$  and  $-50$  V, causes a partial reduction of metal ions to zero valent state ( $\text{Co}^{2+} \rightarrow \text{Co}^0$  and  $\text{Fe}^{3+} \rightarrow \text{Fe}^{2+} \rightarrow \text{Fe}^0$ ) which is accompanied with the oxygen ions migration towards the surface or/and out of the sample. This leads to a maximum increase of  $M_S$  by 15% (see Fig. 5) and a decrease in  $H_C$  by 28% in the porous sample after applying  $-50$  V. Analogous samples with lower porosity degree (synthesized by dip coating but without the diblock co-polymer) exhibit more modest magneto-ionic effects, that is, 2% increase in  $M_S$  and 4% decrease in  $H_C$ . This underscores the importance of nanoporosity in this system. The oxygen ion exchange between the sample and the liquid electrolyte occurs mainly at the solid/liquid interface, thus there are two main contributions from nanoporosity: (1) an enhanced overall interface area and (2) larger electric fields due to the reduced size of the ligaments (reduction of the effective thickness of the film since the electrolyte can penetrate towards the interior of the pores).

The way the oxygen ions are stored appears to be the key factor determining the reversibility of the magnetolectric process. It has been demonstrated that the O<sup>2-</sup> can be either dissolved in the electrolyte (Quintana *et al.*, 2018b; Robbenolt *et al.*, 2018) or aggregated at the grain boundaries forming oxides/hydroxides. In the case described above, that is, CoFe<sub>3</sub>O<sub>4</sub>/PC, the migratory oxygen was likely stored at the surface or at the grain boundaries which made it easier to reincorporate it back into the cobalt ferrite structure with the voltage of opposite polarity (Robbenolt *et al.*, 2019a). Indeed, a similar study performed on nanoporous FeO<sub>x</sub> immersed in a liquid electrolyte demonstrated that the voltage-induced changes in  $M_S$  can be only partially reversed in the nanoporous sample when the oxygen ions are stored in the solution (Robbenolt *et al.*, 2019b). The O<sup>2-</sup> ions can be stabilized by





**Fig. 5** (a) Relative change in saturation magnetization ( $M_S$ ) vs. the initial sample magnetization as a function of time and applied voltage for nanoporou  $\text{CoFe}_3\text{O}_4$  films. Point I is the initial sample as-synthesized. The blue line shows the evolution of the magnetization during the application of  $-10$  V for 160 min and point II is the maximum magnetization reached. After the application of  $-10$  V, the voltage was turned to 0 V (black line) in one case and  $+10$  V (red line) in another, as it is shown in panel (b). Points III and IV are the sample state after the sample relaxed at 0 V and recovered under  $+10$  V respectively. (c) Room-temperature magnetic hysteresis loops corresponding to points I-IV. The figure is reprinted from Robbenolt, S., Menéndez, E., Quintana, A., *et al.*, 2019a. Reversible, electric-field induced magneto-ionic control of magnetism in mesoporou cobalt ferrite thin films. *Scientific Reports* 9, 10804. doi:10.1038/s41598-019-46618-6.

polar solvent molecules and, depending on the voltage intensity, these ions remain in the electrolyte or become neutralized at the counter-electrode. In the latter case, the overall system loses part of the oxygen in form of bubbles which can make the magnetolectric changes fully irreversible. In this case, the ability of the system to be re-oxidized will be dictated by the diffusion kinetics of the process.

Recently, magneto-ionic effects in spinel oxides have attracted the attention of scientists working on lithium ion batteries. Such materials ( $\text{LiFe}_5\text{O}_8$ ,  $\text{CoFe}_2\text{O}_4$ ,  $\text{NiFe}_2\text{O}_4$ , etc.) show ability to reversibly incorporate  $\text{Li}^+$  ions from a liquid electrolyte and subsequently release these ions in a specific range of applied voltage (Reitz *et al.*, 2016; Dubraja *et al.*, 2018). Lithium intercalation causes a valence change and a partial redistribution of metal cations in the spinel structure, hence resulting in a large increase in magnetization at room temperature, that can be as high as 30% in some cases (Dasgupta *et al.*, 2014). Remarkably, magnetization reversal dictated by lithium intercalation mechanism is highly reversible (endurance of thousands of cycles), although only to a certain degree of lithiation, beyond which the original spinel structure cannot be fully recovered anymore. This is in line with the other works on oxides discussed previously in this article. Furthermore, the highly porous structure of these oxides allows for a reduced diffusion distance, while enhancing, to some extent, the mechanical resistivity of the material against failure due to volume changes (electrode breathing) during lithiation/de-lithiation.

## Challenges and Perspectives

In recent years, exciting new results in the field of magnetolectric actuation have been obtained in composites comprising metallic, semiconducting and dielectric nanoporous structures. The presence of nanoporosity allows for significant changes in the magnetic properties under application of moderate voltages. Furthermore, such systems are able to outperform at room temperature, compared to fully dense films with the same composition. This eventually leads to reduced energy cost and facile integration of these materials into devices. Indeed, there is a huge potential for these systems to be utilized in a myriad of energy-efficient applications: spintronic devices, energy harvesters, radio-frequency/microwave devices, non-invasive biomedical technologies, and neuromorphic computing platforms (Hu and Nan, 2019; Mishra *et al.*, 2019). Nevertheless, in most cases, the use of a liquid electrolyte is required to generate ultra-high electric fields (thanks to the formation of ultra-narrow electric double layers) and to access the large surface area of the complex 3D nanoporous networks (Navarro-Senent *et al.*, 2019). For further technological exploitation of these enhanced magnetolectric effects it is important to fully adapt ME composites to operate in solid state. In fact, the pores can accommodate various host FE materials that can conformally coat the overall surface of highly porous FM phases, thus resulting in a new type of composite materials able to operate in all-solid state via magneto-ionic, strain-mediated or charge accumulation mechanisms, depending on the nature of the FM/FE materials. Particularly, there is large potential for nanoporous materials to be integrated in strain-mediated composites, for example, filling the pores with FE polymers to take advantage of the synergies between magnetostriction and piezoelectricity (e.g., PDVF (Martins *et al.*, 2015; Poddar *et al.*, 2018)), analogous to bioinspired three-dimensional magnetoactive scaffolds (Fernandes *et al.*, 2019). In addition, magneto-ionic control of magnetism can be achieved in nanoporous ME composites, where the guest material can be a solid electrolyte (ionic conductor) grown throughout the surface of the porous framework using, for example, ALD such as a high  $\text{O}^{2-}$  mobility thin film like  $\text{Gd}_2\text{O}_3$

or  $\text{HfO}_2$  (Detavernier *et al.*, 2011). In fact, it has been recently demonstrated that larger magneto-ionic effects can be obtained in nanoporous cobalt ferrite film conformally coated with  $\text{HfO}_2$ , as compared to a film having the same degree of porosity but simply immersed in an anhydrous electrolyte (Robbenolt *et al.*, 2020). Nevertheless, in this case, a liquid electrolyte was still employed during voltage application to avoid the problem of electrical pinholes in  $\text{HfO}_2$ . Furthermore, the magneto-ionic speed needs to be further increased, while keeping the reversibility and cyclability of the effects at room temperature. This challenge is being tackled by using hydrogen magneto-ionics instead of oxygen (Tan *et al.*, 2019), or optimizing the electric contacts and sample designs (de Rojas *et al.*, 2020). However, in spite of all these works, the mechanisms responsible for magnetolectric actuation are, in general, still rather poorly understood.

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