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# Challenges for optical nanothermometry in biological environments

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**ABSTRACT** Temperature monitoring is useful in medical diagnosis, and essential during hyperthermia treatments to avoid undesired cytotoxic effects. Aiming to control heating doses, different temperature monitoring strategies have been developed, largely based on luminescent materials, a.k.a. nanothermometers. However, for such nanothermometers to work, both excitation and emission light beams must travel through tissue, making its optical properties a relevant aspect to be considered during the measurements. In complex tissues, heterogeneity, and real-time alterations as a result of therapeutic treatment may have an effect on light-tissue interaction, hindering accuracy in the thermal reading. In this Tutorial Review we discuss various methods in which nanothermometers can be used for temperature sensing within heterogeneous environments. We discuss recent developments in optical (nano)thermometry, focusing on the incorporation of luminescent nanoparticles into complex *in vitro* and *in vivo* models. Methods formulated to avoid thermal misreading are also discussed, considering their respective advantages and drawbacks.

## KEY LEARNING POINTS

1. Basic principles of luminescence nanothermometry.
2. Nanothermometry applied to biological systems.
3. Insights on tissue effects: thermal reading through non-transparent media.
4. Accuracy of spectroscopy-based temperature measurements in complex models.
5. Outlook and perspectives on luminescent nanomaterials for temperature measurements.

1	Introduction	28
2	The concept of hyperthermia for the treatment of disease has	29
3	been reported to exist for over 5000 years. <sup>1</sup> Like many other	30
4	historical medical treatments, with the arrival of modern	31
5	medicine it became clear that applying and monitoring	32
6	temperature changes would need careful refinement.	33
7	Throughout the 20 <sup>th</sup> century, methods to direct heat to	34
8	internal organs were initially achieved via insertable probes	35
9	at the site of application. <sup>2</sup> With the development of	36
10	nanotechnology, new techniques have been proposed to	37
11	achieve localized hyperthermia, based on the ability of some	38
12	nanomaterials to transform magnetic fields or light into heat.	39
13	Indeed, some formulations of nanomaterials are currently	40
14	(January, 2022) at different stages of clinical trials, targeting	41
15	prostate, lung or head and neck cancers. However, optimal	42
16	treatment not only requires spatial control of the heated area,	43
17	but also on the achieved temperature, to avoid overexposure	44
18	and damage of the surrounding tissue. In clinical settings,	45
19	current recommendations to implement thermally monitored	46
20	hyperthermia are mainly related to the treatment of liver	47
21	cancer. Hyperthermia is also of interest in different	48
22	experimental settings, for applications such as rewarming of	49
23	cryopreserved organs, activation of thermo-sensitive	50
24	properties in materials science, ablation of cancer cells in	51
25	oncology studies via specific delivery of strongly absorbing	52
26	nanoparticles, etc.	53
27	Thermal monitoring alone also plays an important role in the	54
	diagnosis of disease state or infections. For example, the	55
	early-stage detection of ischemia, thanks to the so-called	
	Transient Thermometry (TTh) technique, or of breast cancer	
	thanks to a higher local temperature promoted by increased	
	metabolic and vascular changes of the tumor. <sup>3,4</sup> Changes in	
	temperature, whether they are internally or externally	
	provoked, can thus be exploited for both the diagnosis and	
	treatment of many diseases. Notwithstanding, the vast array	
	of tissue heterogeneity as well as the differences in the	
	temperature ranges of interest over space and time for	
	different applications, imply that there is no “one size fits	
	all” rule for thermometry requirements. However, certain	
	requisites such as non-invasiveness, accuracy, and	
	sensitivity are all generally shared among clinical	
	applications.	
	When the focus is on measuring temperature at <i>in vitro</i> and	
	<i>in vivo</i> settings with an ultimate clinical target, an aspect that	
	has not been sufficiently addressed is its accurate	
	determination in real-time with high spatial resolution.	
	Opposite to traditional contact probes, such as liquid-in-	
	glass thermometers, thermocouples or optical fibres, purely	
	non-contact techniques based on spectroscopy have been	
	developed toward this aim. However, commercially	
	available tools (e.g., infrared (IR) thermal cameras) have	
	limited sensitivity and accuracy, and whilst non-invasive,	
	their ability to detect temperature changes is restricted to	
	surface measurements, due to the poor tissue penetration of	
	mid-far IR light, where these cameras are designed to	

1 operate (typically from 6 to 14  $\mu\text{m}$ ). Indeed, there are strong 61  
2 limitations regarding working wavelengths that not only 62  
3 affect the maximum depth of the measurement, but also the 63  
4 accuracy of the technique, both *in vitro* and *in vivo*. Organs 64  
5 comprising different tissues present strong scattering of light 65  
6 because of the various physical interfaces at tissue 66  
7 boundaries.<sup>5</sup> Furthermore, large signal distortions take place 67  
8 in photoluminescence spectra when such interfaces are 68  
9 combined with highly absorbing molecules such as 69  
10 haemoglobin and lipids. To overcome these issues, thermal 70  
11 measurements are generally based on detection mechanisms 71  
12 in the near infrared (NIR) range of the electromagnetic 72  
13 spectrum, which is optimal for biomedical thermography. To 73  
14 date, up to four spectral ranges, also known as “Biological 74  
15 Windows” (BWs), in which both light scattering and 75  
16 absorption from molecules are reduced, have been identified 76  
17 within the NIR. Their exact widths vary depending on the 77  
18 publication in question, but approximate ranges can be 78  
19 identified as 650 - 950 nm (BW-I), 1000 - 1350 nm (BW-II), 79  
20 1450 - 1900 nm (BW-III), and 2100 - 2300 nm (BW-IV).<sup>6,7</sup> 80  
21 Various non-contact thermometry techniques can be adapted 81  
22 to work within these BWs, for example Raman spectroscopy 82  
23 and interferometry thermography. Whereas the former relies 83  
24 on the detection of vibrational modes in the material, the 84  
25 latter deduces temperature values from changes in the 85  
26 refractive index of the sample. Yet, given the heterogeneous 86  
27 nature of biological environments, which include a  
28 multitude of biocomponents with huge variations in 87  
29 refractive indexes, the interpretation of these techniques can 88  
30 be very complex. As an alternative to purely non-invasive 89  
31 spectroscopy and the traditional contact probes, an 90  
32 additional category of thermometry sensors have emerged: 91  
33 luminescent nanomaterials (LNMs). 92

34 We define LNMs as nanoparticles and other nanomaterials 93  
35 such as polymers, nanoclusters, etc., that can be used as 94  
36 thermometers because their luminescence is sensitive to 95  
37 temperature changes. Compared to other temperature 96  
38 sensing techniques, the use of LNMs is generally non- 97  
39 invasive because their luminescence can be excited from 98  
40 sources external to the sample. Additionally, their nanoscale 99  
41 size translates into a relatively high spatial resolution, which 100  
42 in the absence of scattering from tissue, is typically limited 101  
43 by light diffraction. This nanometric scale also facilitates 102  
44 delivery to a wide range of target organelles, cells, tissues, 103  
45 and organs. As opposed to IR cameras, LNMs offer the 104  
46 possibility to detect temperature changes at internal sites, 105  
47 and luminescence offers in principle an easier way to 106  
48 calculate the temperature, as compared to other spectroscopy 107  
49 techniques. As a result, such materials can potentially be 108  
50 used in a variety of therapeutic and diagnostic applications 109  
51 involving internal organs. 110

52 The use of LNM temperature measurements was first 111  
53 described in the 1970's, but it was not until the last decade 112  
54 that they were employed in biological research. During this 113  
55 period, outstanding progress in the synthesis and 114  
56 characterization of LNMs took place, and whereas the initial 115  
57 efforts were directed towards exploring new nanomaterials, 116  
58 further advances were focused on tuning the emission and 117  
59 excitation wavelengths to the (farther) biological windows, 118  
60 hence improving light penetration. Besides, the field has

necessarily grown in parallel with bionanotechnology, including thus concerns on cell internalization, cytotoxicity, and controlled delivery. One of the characteristics that provide versatility to LNM measurements is the large library of available nanoscale materials, the most commonly described LNMs including organic molecules, quantum dots (or other semiconductor nanoparticles), lanthanide-doped nanoparticles, and nanodiamonds.<sup>7-9</sup> The correct selection of nanothermometers, taking into special consideration the available biological windows, has been described in an excellent review by Nexha et al.<sup>7</sup> In this work it becomes apparent that, whereas abundant literature is available on the synthesis of LNMs, their application to *in vitro* and *in vivo* studies is relatively underexplored.

If we ask ourselves why this is the case, the insufficient choice of nanothermometers cannot be the reason, even though a better signal-to-noise ratio would always facilitate sensing and this remains an important drawback of some materials. Instead, what if a poor understanding of the physicochemical properties of LNMs in biological environments were the cause? Or the lack of technological instrumentation and reproducibility for their accurate detection? Indeed, the latest developments in the field of LNMs are directed towards standardization of the probes, and to the development and implementation of measurement protocols to improve accuracy.

In this Tutorial Review, we aim at describing the importance of LNMs for localized nanothermometry, focusing on the experimental conditions to be considered. This analysis includes the description of the biological environment, which determines light penetration and exit, but also the selection of suitable temperature sensing methods. Materials for nanothermometry have been extensively described in recent reviews,<sup>8,10,11</sup> thus, here we describe the pros and cons concerning their performance in heterogeneous environments. We pay special attention to the obstacles that may arise during temperature measurements, and to the solutions that have been proposed. These are key aspects for the current evolution of nanothermometry, as a poor or incomplete analysis will not only result in poor thermometric performance and thus inaccurate temperature measurements, but additionally can cause undesired tissue damage, unforeseen cytotoxicity, cellular changes at a molecular level, or even instability of NPs, all of which should be avoided in any biological situation.

## Relevant Parameters in Luminescence Nanothermometry

Luminescent materials can be used as thermal probes, provided that the light they emit is dependent on temperature in any respect. Indeed, different characteristics of emitted light can be (and have been) exploited to measure temperature (or temperature changes). These parameters include emission intensity and lifetime, spectral width, and spectral shift, as well as intensity ratios and polarization anisotropy (**Figure 1**). The physical processes behind each thermometry technique play a key role in the feasibility of the measurement. First of all, thermal sensitivity, i.e. the

1 extent of change of the measured parameter per degree, is  
 2 determined by the physics of luminescence. Second,  
 3 practical implementation of each specific experimental set-  
 4 up will affect the cost, speed, and ease of data analysis of the  
 5 technique. Finally, both nanomaterials and emitted photons  
 6 will inevitably interact with their environment, resulting in  
 7 alterations of the detected signal, with various degrees of  
 8 relevance depending on the physical origin of the measured  
 9 parameter. Let us briefly discuss the main available options.

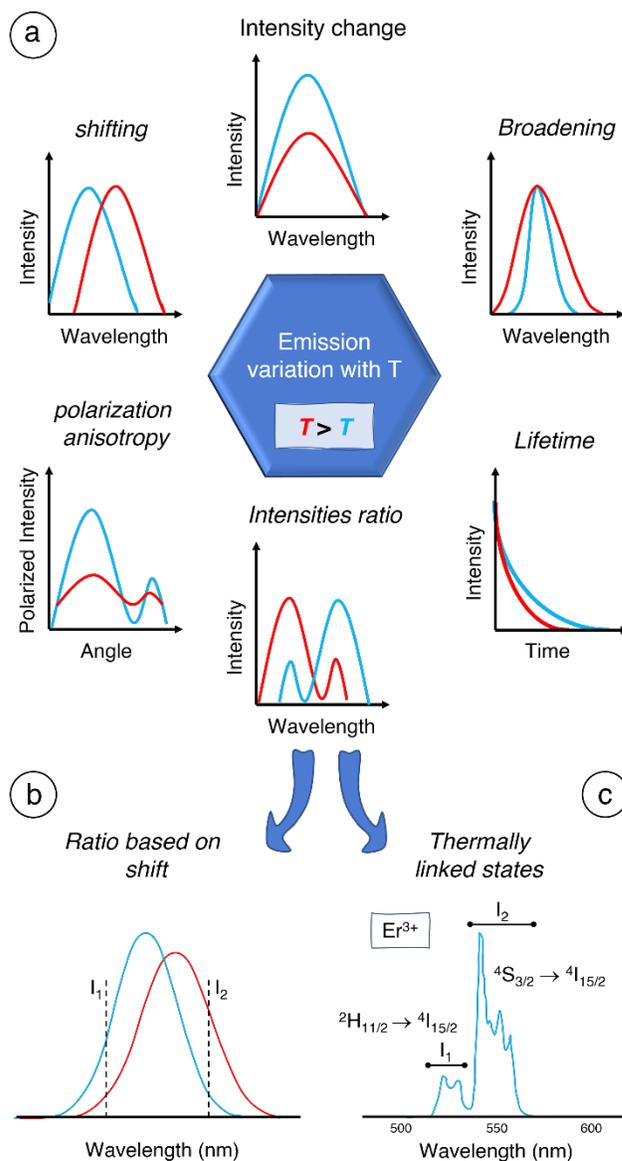
10 The emission probability of an excited fluorophore, i.e., a  
 11 fluorophore that has absorbed sufficient energy to promote  
 12 an electron from the ground state to an excited state, is given  
 13 by the radiative decay rate,  $k_r$ . In the same way, there is a  
 14 non-radiative decay rate,  $k_{nr}$ , which is the probability for the  
 15 fluorophore to transfer its extra energy to the vibrational  
 16 modes of surrounding atoms or molecules. Let us assume the  
 17 simple case in which the fluorophore is located in a  
 18 homogeneous medium, as would be the case for one  
 19 luminescent dopant in a perfect transparent crystal matrix.  
 20 Then,  $k_r$  and  $k_{nr}$  are the only possible relaxation routes and  
 21 thus, there is a competition between them. While  $k_r$  is a  
 22 constant for a specific dopant-host pair,  $k_{nr}$  varies with  
 23 temperature, which means that the observed emission  
 24 intensity will also vary with temperature, thereby allowing  
 25 for thermometry applications. This can be easily seen  
 26 through the definition of photoluminescence quantum yield,  
 27  $\eta$ , which is the ratio between the number of emitted photons  
 28 and absorbed photons:

$$29 \quad \eta = \frac{k_r}{k_r + k_{nr}} \quad (1)$$

30 The other parameter directly related to  $k_r$  and  $k_{nr}$  that can be  
 31 used in thermometry is experimental lifetime,  $\tau_{exp}$ , which  
 32 measures the characteristic time an electron remains in an  
 33 excited state before relaxation:

$$34 \quad \tau_{exp} = \frac{1}{k_r + k_{nr}} \quad (2)$$

35 Most nanothermometry methods are based on either  
 36 emission intensity or lifetime measurements, both of which  
 37 are strongly dependent on temperature and are thus highly  
 38 sensitive. Changes to the band width and wavelength shifts  
 39 are less commonly exploited, as the temperature dependent  
 40 effect is generally weaker and hence less sensitive.  
 41 Nevertheless, the use of materials such as  $\text{Y}_2\text{O}_3:\text{Eu}^{3+}$   
 42 nanoparticles with a strongly temperature-dependent  
 43 bandwidth,<sup>12</sup>  $\text{Ag}_2\text{S}$  nanoparticles featuring emission peak  
 44 redshifts at higher temperatures,<sup>13</sup> or nanodiamonds in  
 45 which the energy separation between two energy levels is  
 46 used as a thermal probe, have been reported.<sup>14</sup> The latter case  
 47 is interesting as it has provided accurate results including the  
 48 study of plasmonic heating and *in vitro* temperature  
 49 reading.<sup>14</sup> However, visible light excitation and microwave  
 50 pulses are required, which hinders their applicability *in vivo*.  
 51 All such alternative options are also interesting in the context  
 52 of the current trend to apply multiparametric analysis based  
 53 on the thermal dependence of more than one luminescence  
 54 parameter.<sup>15</sup> This strategy gives access to improved thermal  
 55 resolution, reaching values as low as 0.05 K, so even if one  
 56 technique alone cannot reach such a high resolution by itself,



**Figure 1.** (a) Luminescence-based strategies applied to measure temperature with NPs. (b,c) Intensity ratios have been used in two modalities: to characterize an emission shift (b), or to characterize the emission from two thermalized energy states (c).

a chance arises if complemented with other thermometer probes.

An important advantage of intensity-based thermometry is the simplicity of the experimental set-up and of data processing. Indeed, if the experiment is properly controlled, an accurate thermal calibration allows for a direct correlation between intensity and temperature. The same reason can however also become a limiting factor, as an accurate agreement between calibration and experimental conditions is required. In other words, excitation and emission light beams must reach the nanoparticles and the detector, respectively, in the purest form for temperature to be determined. This factor complicates temperature measurements in heterogenous biological environments where light scattering and absorption may become dominant. In addition, the concentration of LNMs must remain constant and known, as more LNMs will obviously

1 provide a higher emission intensity. This is a major 60  
2 limitation for *in vivo* situations, as neither the initial density 61  
3 of LNMs in the tissue, nor its variations over time can be 62  
4 easily controlled. 63

5 A solution to these drawbacks has been found by measuring 64  
6 the intensity of light at two different wavelengths and using 65  
7 their ratio as a thermometry probe. In this technique, one 66  
8 intensity can be understood as a reference value or internal 67  
9 calibration for the other. Therefore, for the highest thermal 68  
10 sensitivity it is crucial that their intensity variations due to 69  
11 temperature changes are as dissimilar as possible. If both 70  
12 selected intensities show the same dependence on LNM 71  
13 concentration and on excitation power, then the intensity 72  
14 ratio will not be affected by these parameters. These 73  
15 restrictions mean that not every intensity ratio is valid, 74  
16 especially if coming from two separate fluorophores that 75  
17 may not be colocalized, or if non-linear excitation occurs (as 76  
18 in upconversion techniques), as this can trigger different 77  
19 dependencies between emission intensity and excitation 78  
20 power. In general, there are two main situations related to 79  
21 the emission of one single fluorophore in which intensity 80  
22 ratio thermometry has been safely applied. 81

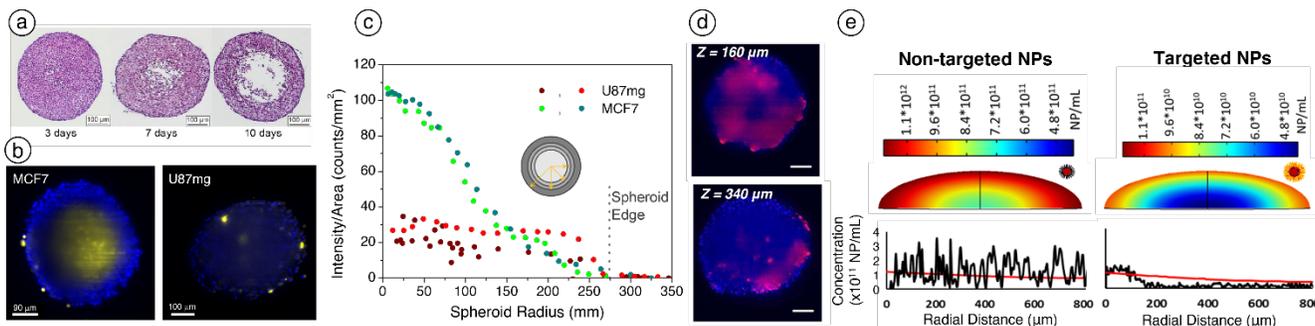
23 The first one involves a shift of the emission spectrum when 82  
24 temperature changes (**Figure 1b**). Instead of measuring the 83  
25 emission peak position, it is then possible to simplify data 84  
26 recording by defining an intensity ratio between two 85  
27 different wavelengths within the same emission band. Such 86  
28 a ratio would also be affected by any intensity changes of the 87  
29 emission, which typically occur simultaneously to the shift, 88  
30 so both effects are included in one single parameter. A good 89  
31 example of the use of intensity ratios in such a case is 90  
32 provided by Ag<sub>2</sub>S semiconductor nanoparticles, which 91  
33 typically feature one emission band within the second 92  
34 biological window, which shifts and loses intensity as 93  
35 temperature increases.<sup>13</sup> It has been reported that, in this 94  
36 strategy, the accuracy and sensitivity of the intensity ratio 95  
37 depends on a careful selection of the two wavelengths at 96  
38 which emission intensity is to be monitored. 97

39 The second situation in which intensity ratios are applied is 98  
40 related to emission bands with contributions from two 99  
41 different excited states (**Figure 1c**). If both energy states are 100  
42 close enough, i.e., typically with a separation smaller than 101  
43 2000 cm<sup>-1</sup>, thermal energy (given by  $k_B \cdot T$ ,  $k_B \approx 0.695$  cm<sup>-1</sup>/K) 102  
44 may be sufficient to promote electrons from the lower 103  
45 energy state to the upper one. Thus, the electronic population 104  
46 of both levels will be thermally linked, and their intensities 105  
47 will constitute a reliable intensity ratio. Most materials 106  
48 displaying this type of energy states are lanthanide-doped 107  
49 nanoparticles. Notwithstanding, some transition metal 108  
50 doped nanoparticles have also been recently reported to 109  
51 display all the required characteristics for this technique.<sup>16</sup> In 110  
52 the case of lanthanides, emission bands are sufficiently 111  
53 narrow so that nearby emission bands can be spectrally 112  
54 resolved and subsequently used for nanothermometry, as 113  
55 shown for I<sub>1</sub> and I<sub>2</sub> in Figure 1c. This is because 4f electrons, 114  
56 responsible for luminescence in these materials, are partially 115  
57 shielded from the electromagnetic field created by the 116  
58 surrounding material, which is known as the crystal field, 117  
59 Furthermore, although the effect of the external crystal field 118

is weak, it can break the degeneracy of the energy states in  
sub-levels separated by only 10<sup>2</sup> cm<sup>-1</sup>. In the emission  
spectrum, this translates into a sub-structure of peaks within  
each emission band (**Figure 1c**, see peaks in each of the two  
main transitions) that will be close enough to being also  
thermally linked. In this case, however, the levels are often  
too close together to be spectrally resolved in measurements,  
but otherwise could also constitute a thermal probe.  
Examples of intensity ratios based on thermally linked states  
are found in most thermometers based on emissions from  
erbium, in the green spectral region (500 nm – 560 nm), as  
shown in **Figure 1c**. Shifting towards the NIR, neodymium  
ions have been also described to display similar properties,  
with emissions located both in the BW-I (*ca.* 800 nm) and in  
the BW-II (*ca.* 1050 nm).<sup>8</sup>

The use of intensity ratios is, in principle, a robust strategy  
to address the issues caused by fluctuations in LNM  
concentration and excitation power. However, as it is based  
on luminescence from lanthanide ions, a detection set-up  
with high spectral resolution and sensitivity is required. This  
second feature is related to the fact that electronic transitions  
between 4f states are forbidden by selection rules, and thus  
only emit because an asymmetric crystal (or molecular) field  
can relax this situation. Still, their absorption cross-sections  
and emission quantum yields are often lower than those of  
other luminescent materials, such as quantum dots or  
fluorescent organic dyes. The need for high spectral  
resolution is due to the narrow emission levels of lanthanide-  
doped materials and the close energies between the  
emissions involved. An alternative energy ratio that would  
solve this issue has been proposed, which uses thermalized  
states at the ground level, i.e. the ground state can thermally  
populate the first excited state. Thus, the same emission band  
is alternatively excited by two different excitation lasers,  
from the ground and the first excited states. This method is  
interesting because it facilitates data recording and analysis,  
but it requires two excitation sources and constant switching  
between them. Such materials have not yet been optimised  
for use within *in vivo* environments because they work  
exclusively in the visible range. In addition, thermalization  
of the first excited state is low at biological temperatures and  
thus, though suitable for measuring temperature changes  
within this range, one of the excitation wavelengths provides  
a low emission intensity that may be hard to distinguish in  
biological environments.<sup>17</sup>

Polarization anisotropy can also be considered as a  
ratiometric technique, because it compares emission  
intensities recorded under two different light polarizations.  
It shares the advantages of intensity ratios and is independent  
of the illumination power and the distribution of LNMs.  
However, this technique has been insufficiently developed  
in the biological context, probably because its temperature  
dependence is complex, being influenced by both fluid  
viscosity and luminescence lifetime. Still, the feasibility of  
the technique has been demonstrated, even in small animals,  
using the Green Fluorescent Protein (GFP) as the LNM.<sup>18</sup>  
Further work has been carried out, aiming to construct  
smaller probes based on proteins, which may soon push  
further development of this strategy.<sup>19</sup>



**Figure 2.** (a) Histology of human head and neck carcinoma cell line FaDu-formed spheroids, showing the characteristic necrotic core devoid of intact cells.<sup>30</sup> (b,c) Fluorescence microscopy images (b) and quantification (c) of propidium iodide penetration into spheroids made of breast (MCF7) and brain (U87.mg) cancer cell lines, showing the increased presence of propidium iodide (in yellow, labelling dead cells) in MCF7 spheroids of the same size and age as U87.mg.<sup>32</sup> (d) Cross section of a spheroid at 160  $\mu\text{m}$  and 340  $\mu\text{m}$  (equivalent to its maximum diameter), where blue colour shows DAPI fluorophore emission labelling cell nuclei and red shows emission from  $\text{Nd}^{3+}$ -based nanothermometers. LNM penetration is inhomogeneous over the spheroid surface. Scale bar: 100  $\mu\text{m}$ .<sup>32</sup> (e) Experimental and computational models of the diffusion of polymeric fluorescent nanoparticles into a spheroid, using non-targeted and targeted NPs. Targeted NPs are sequestered at the spheroid periphery, whereas non-targeted NPs can penetrate (and leave) the spheroid easily.<sup>31</sup> Figures reproduced under the Creative Commons Attribution Licence (CC BY 4.0) from references 30, 31, and 32, with permission from MPDI, Nature Publishing Group, and Ivyspring International Publisher, respectively.

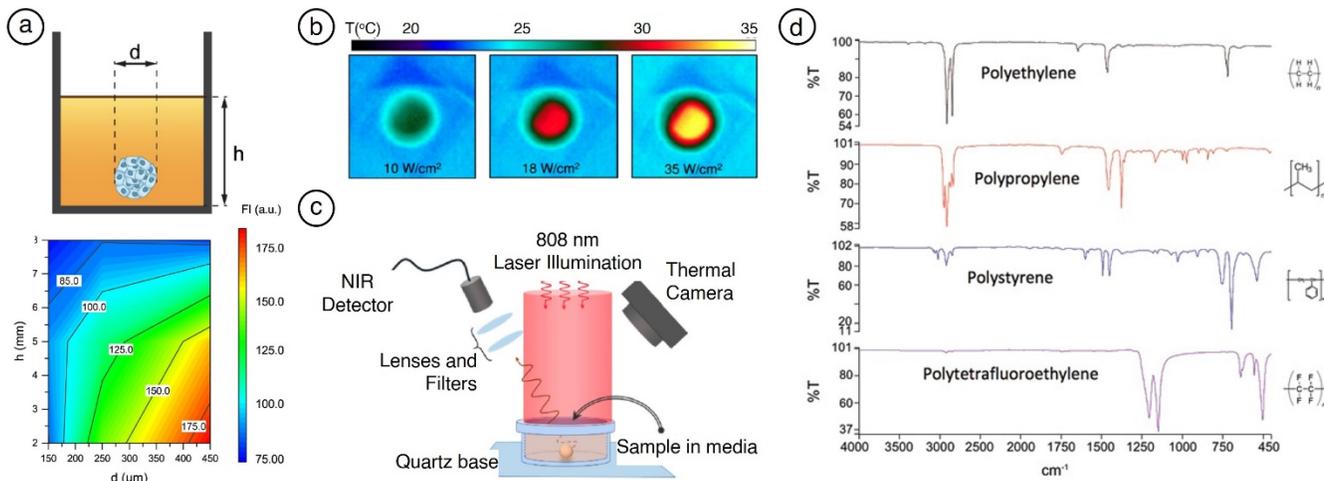
1 Finally, luminescence lifetime, i.e., the characteristic time an  
 2 excited state requires to lose its population (see Equation 2),  
 3 is also a robust parameter to be considered, as it is  
 4 independent of both the concentration of LNMs and the  
 5 illumination intensity. In addition, it requires measurement  
 6 of a single emission wavelength, meaning that spectral  
 7 distortions, often occurring when light is travelling through  
 8 tissue, will not play a role. However, as a principal  
 9 drawback, the experimental set-up is more complex, as it  
 10 requires a pulsed laser and synchronized detection with time  
 11 resolution.<sup>20</sup> Still, this technique has been extensively used  
 12 and constitutes an interesting option, even in biomedical  
 13 applications, first demonstrated with organic dyes<sup>21</sup> and  
 14 more recently with inorganic nanoparticles such as  $\text{Ag}_2\text{S}$ .

## 15 Nanothermometry in Practice

16 One of the first proposed applications of nanothermometry  
 17 has been in the determination of intracellular temperature, to  
 18 study cellular metabolism. Thermogenesis at the cellular  
 19 level refers to chemical processes that cause local changes in  
 20 temperature, as a response to endogenous stimuli. Typically,  
 21 processes involving cellular metabolism are affected, and  
 22 hence the mitochondria and surrounding cytoplasm are the  
 23 most studied organelles.<sup>10</sup> Various sensing strategies have  
 24 been explored to determine endogenous temperature  
 25 changes occurring in live immortalized cells. Whereas most  
 26 techniques address temperature changes of around  $\leq 1$  K for  
 27 endogenous intracellular gradients and those due to external  
 28 stimuli, more striking values close to 3 K (intracellular  
 29 temperature gradient between the nucleus and cytoplasm),  
 30 or 6 to 9 K (increase in temperature around mitochondria)<sup>23</sup>  
 31 have also been reported. Such impressive gradients and  
 32 changes in intracellular temperature understandably lead to  
 33 speculation as to whether the probes, and the employed  
 34 methods, are reliable. To put these results into perspective,  
 35 extracellular temperature changes in cancer tissue are  
 36 generally in the range of 1-2 K, as determined from  
 37 consistent theoretical and experimental data. To better  
 38 understand the putative range of temperature heterogeneity

within a cell, Baffou *et al.* reported theoretical modelling of  
 the expected thermal changes that occur during intracellular  
 thermogenesis. By taking into consideration aspects such as  
 heat diffusion by organelles, the size of mitochondria, their  
 distribution throughout the cytoplasm, thermal conductivity  
 of the plasma membrane, and the mass of glucose that can  
 be transformed into heat, changes in temperature of the order  
 of  $10^{-5}$  K were calculated.<sup>24</sup> Similar theoretical values have  
 been determined by other groups.<sup>25</sup>

This discrepancy between theoretical and experimental  
 values (known as the  $10^5$  gap) makes us wonder about the  
 limitations of each approach. Indeed, in a commentary  
 discussion from 2015,<sup>26</sup> relating to the  $10^5$  gap, it was  
 suggested that a limiting factor in experimental findings was  
 the use, or rather lack, of accurate intracellular  
 thermometers. Experimental limitations are largely based on  
 the sensitivity of the measurement techniques and the  
 adequacy of the implemented protocols, which can be  
 affected by parameters other than temperature. Regarding  
 theoretical modelling, sources of inaccuracy are insufficient  
 data on the validity of the thermodynamic equations that  
 apply in the complex cell environment, and the accuracy of  
 the thermal constants used for cell components (thermal  
 conductivity, mainly).<sup>27</sup> However, as more thermometric  
 materials and sensing strategies have been applied to the  
 measurement of intracellular temperature (fluorescent and  
 non-fluorescent, a good summary can be found in reference  
 10), all providing consistent results, a higher reliability of  
 experimental data is accepted. It should be noted that the  
 maturity of *in vitro* thermometry has recently reached a point  
 at which intracellular thermometric probes are commercially  
 available, as well as detailed measurement protocols.<sup>10</sup>  
 Applications of these probes include, for instance,  
 understanding the mitochondrial thermogenesis in brown  
 adipocytes,<sup>28</sup> which play an important metabolic role in non-  
 shivering thermal regulation of many animals, including  
 humans. From an applied perspective, their contribution to  
 energy consumption makes them interesting targets in  
 obesity treatment studies.



**Figure 3.** (a) Spheroid model of 150 to 450  $\mu\text{m}$  in diameter, shown in a 96 well-plate (top), and the fluorescence intensity of ICG (NIR dye) as a function of the distance from the centre of the spheroid (bottom).<sup>30</sup> (b) IR thermal camera images of irradiated spheroids showing the poor spatial resolution of the technique. These images were taken using the experimental setup in (c), in which hyperthermia was achieved in spheroids treated with plasmonic nanoparticles and nanothermometers. Laser illumination at 808 nm is required for heating and luminescence excitation. Temperature was recorded using both, a thermal camera and nanothermometers emission.<sup>32</sup> (d) IR absorption spectra of the most common plastics used for the manufacture of the cell culture plates.<sup>33</sup> Figures reproduced under the Creative Commons Attribution Licence (CC BY 4.0) from references 30, and 32 with permission from MPDI and Ivyspring International Publisher, respectively.

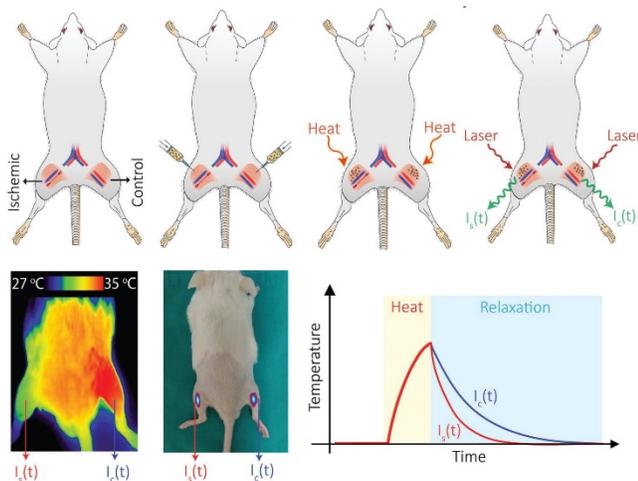
1 Despite early controversy, *in vitro* nanothermometry has 36  
 2 seen a fast evolution, reaching a stage at which it is being 37  
 3 applied in biological research. However, thermal monitoring 38  
 4 of more complex samples involves additional concerns, 39  
 5 largely related to light-matter interactions, but also to the 40  
 6 delivery of LNMs to the area of interest. 41

7 **Nanothermometry in 3D cell models.** It is generally 42  
 8 accepted that 3D *in vitro* cellular models represent a suitable 43  
 9 stepping-stone, or even a valid alternative to *in vivo* models, 44  
 10 as long as physical aspects such as geometry and cellular 45  
 11 heterogeneity are kept as close to the real tissue as possible. 46  
 12 Indeed, advanced 3D cell models may even offer improved 47  
 13 characteristics with respect to *in vivo* models, such as control 48  
 14 over biophysical properties (including oxygen state 49  
 15 and fluid flow) and the ability to genetically modify the 50  
 16 expression of proteins at the single cell level. 51

17 As discussed above, a large body of information is available 53  
 18 on mapping intracellular temperature changes using LNMs. 54  
 19 However, few studies have attempted temperature 55  
 20 measurements in 3D cell models such as spheroids, 56  
 21 organoids, organ-on-a-chip, and 3D-printed tissues, all of 57  
 22 which may provide extensive information as to how 58  
 23 successful the use of LNMs may be *in vivo*. Two of the most 59  
 24 critical aspects in the experimental design are the 60  
 25 spheroid/organoid geometry, as well as the uptake and 61  
 26 spatial distribution of the LNMs in or around the spheroid. 62  
 27 Regarding organoid geometry, it is well known that the type 63  
 28 and number of cells, or relative numbers of the various cell 64  
 29 types when using heterologous cell mixtures, as well as the 65  
 30 growth stage (number of days), are crucial factors that 66  
 31 determine the overall size and shape.<sup>29</sup> Unfortunately, there 67  
 32 is no rulebook that can be used to pre-define organoid size 68  
 33 or shape, and therefore preliminary growth tests must be 69  
 34 carried out. With increasing size and growth time, the 70  
 35 interior core of the spheroid becomes gradually devoid of

nutrients and there is a lack of gas exchange, which is crucial  
 for the delivery and removal of oxygen and carbon dioxide,  
 respectively. The resulting necrotic core, which can be  
 distinguished using either fluorescence staining or histology  
 (Figure 2a-c), features a different light absorption profile  
 compared to the outer spheroid rim, due to the enhanced  
 presence of cell debris and cellular heterogeneity.<sup>30</sup> In  
 practice, few LNMs incubated together with spheroids will  
 ever reach the internal core due to sub-optimal LNM  
 physicochemical properties such as uncontrolled  
 aggregation or sedimentation, both of which can be  
 unfavourable for LNM-cell interactions. Also, the compact  
 nature of the extracellular matrix and cell-to-cell interactions  
 within the external layers of the spheroid may result in  
 inhibited LNM penetration. Nanoparticle trapping within the  
 external layers of spheroids has been described,<sup>31</sup> as has non-  
 homogenous LNM distribution around the spheroidal  
 periphery,<sup>32</sup> most likely due to LNM sedimentation during  
 incubation. Examples of these aspects are shown in Figure  
 2d,e.

With regards to the experimental setup, the overall spheroid  
 size and the spatial resolution of available IR cameras may  
 determine the signal intensity recorded from LNMs. In work  
 by Egloff-Juras et al.,<sup>30</sup> the use of a NIR camera (700 – 850  
 nm detection window) restricted imaging depths to a few  
 mm and required spheroid sizes to be larger than 225  $\mu\text{m}$   
 in diameter (Figure 3a). In contrast, a standard IR thermal  
 camera can only provide information on the surface of  
 spheroids immersed in a medium (i.e. on the liquid), and thus  
 cannot distinguish temperature changes through the spheroid  
 (Figure 3b,c). As opposed to IR cameras, the detection of  
 light emitted by LNMs for thermometry involves the  
 arrangement of laser excitation sources and detectors in a  
 similar location above the sample to be interrogated, which  
 may complicate the setup (see Figure 3c). Thus, any barriers



**Figure 4.** Mouse model in which ischemia was generated in the left leg. Due to the ischemic episode, the temperature of the limb was lower, which was detected using a NIR-based nanothermometer through the transient thermometry (TTh) method. Reproduced from reference 37, with permission from John Wiley and Sons Publishing.

1 that may hinder light transmission (including the cell  
 2 medium) should be considered. In particular, the presence of  
 3 polystyrene (PS) and polypropylene (PP) plastics typically  
 4 used in cell culture, for example in microplate lids, causes a  
 5 significant reduction of light transmission through the IR  
 6 camera detection window (**Figure 3d**).<sup>33</sup> Whilst this may be  
 7 problematic for NIR cameras and hence detection of heat  
 8 release, it also poses a problem for incident NIR light  
 9 sources typically used for LNM measurements. For  
 10 example, a 10% reduction in power density can occur upon  
 11 irradiation with a NIR laser source (808 nm), which is  
 12 significant for experiments in the biological range (*ca.* 30 –  
 13 50 °C), where a temperature increase of a few degrees can  
 14 result in cell death. This issue reiterates the importance of  
 15 conducting appropriate controls and taking into account all  
 16 components that may cause light scattering or absorption.

17 **In vivo and ex vivo tissue.** Attempts to apply  
 18 nanothermometers to real tissues appeared early in the  
 19 development of the field, aiming to demonstrate that  
 20 detection of the signal is possible. To simplify tissue-related  
 21 complications, most studies started by using phantom tissue  
 22 (a preparation with optical absorption and scattering  
 23 resembling that in real tissue) or *ex vivo* animal models (most  
 24 often chicken breast, but not exclusively). Phantom tissues  
 25 have been developed to mimic tissue properties at specific  
 26 wavelengths, also within the BW (mainly BW-I). The  
 27 current catalogue of recipes covers a large number of  
 28 options, including phantom breast, bone, cartilage, skin,  
 29 arterial tissue, etc., so versatility is high.<sup>34</sup> The use of  
 30 phantoms guarantees a high reproducibility in the  
 31 experiments, due to their homogeneity and control over the  
 32 shape and thickness of the sample, which is not always  
 33 possible with *ex vivo* tissue. However, they do not resemble  
 34 the chemical environment of real tissues, which would  
 35 require using *ex vivo* tissues. Chicken breast is a  
 36 homogenous tissue with reduced absorption properties and  
 37 provides a sufficient thickness to study laser penetration  
 38 depth. It should be noted that, commercial chicken breast

39 shows a large variability regarding retention of liquids,  
 40 meaning that the optical properties may vary between  
 41 samples, but also that injected LNMs may spread differently  
 42 depending on the degree of hydration. Obviously, this  
 43 variability may cause inaccuracies in model design and  
 44 temperature readout, as LNM concentration in a certain area  
 45 may change. For the same reason, the way  
 46 nanothermometers are administered is also important;  
 47 injection into chicken breast (essentially an intramuscular  
 48 injection) without backflow is crucial to effectively deliver  
 49 the LNMs (or any other NP, for that matter). Partly due to  
 50 the high current incidence of diabetic population who  
 51 require daily injections, the methods and properties of  
 52 intramuscular injections have been extensively studied.  
 53 Whereas it is recommended to insert needles at a 90° angle  
 54 to the tissue surface, the length of the needle (and hence the  
 55 depth of injection) is largely dependent on the aim and  
 56 precise location of the injection. The “pull-back” technique,  
 57 developed to ensure that needles do not pierce blood vessels,  
 58 involves insertion of the needle followed by aspiration, prior  
 59 to expulsion of the syringe contents. Naturally, in *ex vivo*  
 60 tissues such as commercial chicken breast, no blood would  
 61 be expected, but the process of inserting a needle to create a  
 62 channel, prior to expulsion of the contents, can still be  
 63 useful. Other methods such as incision or biopsy punches are  
 64 interesting alternatives, as they produce cavities of  
 65 controlled dimensions.

These principles also apply to experiments using *in vivo*  
 models. It should however be noted that the administration  
 route in this case primarily depends on the disease model of  
 choice. For the purposes of tumour hyperthermia, *in vivo*  
 models should ideally involve the administration of LNMs  
 intravenously or directly into the organ of interest. In  
 contrast, in studies in which minimum tissue interference  
 and controlled NP concentrations are required, subcutaneous  
 or even intradermal injections may be most appropriate. An  
 important consideration in all cases is the maximum  
 permitted volume, which can vary as much as 10-fold,  
 depending on the administration route and animal model.<sup>35</sup>

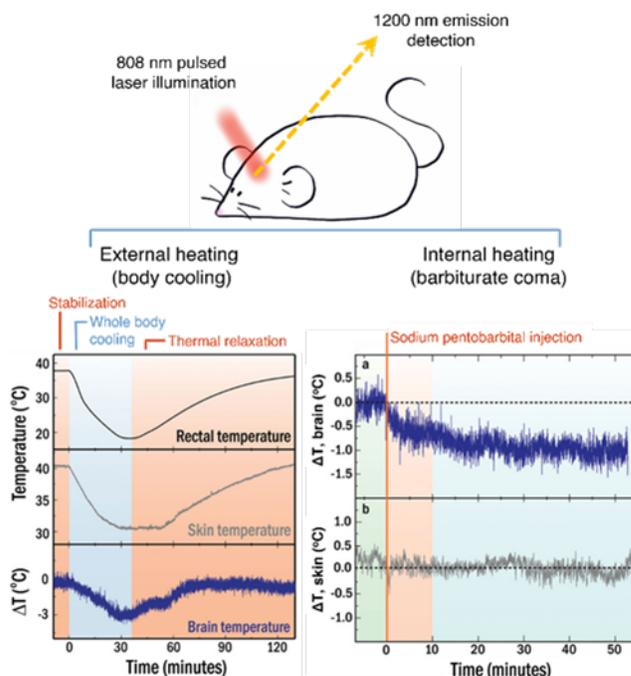
To some extent, the theory applied to nanothermometry *in vitro*  
 can also be translated to *ex vivo* or *in vivo* applications.  
 Notwithstanding, the complexity of tissues and organs in  
 terms of geometry and interaction with light, restricts the use  
 of nanothermometry to proof-of-concept experiments, rather  
 than diagnosis. Yet, some promising results have been  
 reported. One example is the exploration of LNM-quantum  
 dots (QDs) composed of cadmium telluride (CdTe), to detect  
 temperature changes at the nanoscale in *ex vivo Drosophila*  
*melanogaster* skeletal muscle samples. In this case, the  
 nanosized LNM-QDs (2 nm) could bind to myosin and  
 detect heat released during myosin-mediated ATP  
 hydrolysis, via emission intensity measurements in the  
 visible range (*ca.* 520 nm).<sup>36</sup> QDs have also been explored  
 to study tissue relaxation dynamics in an ischemia murine  
 model. Ximendes *et al.*<sup>37</sup> developed a technique to detect  
 either temporal or permanent restrictions of blood supply  
 that can lead to a shortage of oxygen in certain areas of the  
 body, as it occurs in ischemia. As such, the biophysical  
 properties of the relevant tissue are affected, including their  
 thermal relaxation dynamics, which can thus be measured

1 using TTh (**Figure 4**). In this technique, the tissue is first  
2 heated up to subsequently measure how temperature  
3 decreases back to the equilibrium state. In the clinic,  
4 however, transient temperatures are typically monitored  
5 using IR thermal cameras, which are only reliable for  
6 superficial tissues, thereby preventing the extensive  
7 application of TTh.

8 Hence, the authors designed a nanothermometry detection  
9 system in which temperature-sensitive LNM-QDs  
10 (PbS/CdS/ZnS) were injected into the site of interest. As the  
11 emission intensity of these QDs depends on temperature,  
12 irradiation at 800 nm (an efficient excitation wavelength)  
13 and detection over time of LNM-QD emission at *ca.* 1200  
14 nm, allowed the localization of damaged tissue. A similar  
15 TTh strategy has been applied to the early detection of  
16 cancer in mice, in this case using intratumor injection of  
17 Ag<sub>2</sub>S NPs with high thermal sensitivity.<sup>38</sup> These particles  
18 were excited at 808 nm, with emission at 1200 nm, i.e.,  
19 within the first and second biological windows, respectively.  
20 Heating required to induce the TTh effect was provided by  
21 simultaneous laser irradiation at 810 nm, using a higher  
22 power than that for photoluminescence excitation. The  
23 characteristic changes in tumour microvascular thermal  
24 relaxation properties could be used to differentiate it from  
25 healthy tissue. The findings showed an impressive  
26 diagnostic ability, as tumours could be identified up to 7  
27 days prior to optical detection.

28 A final illustrative example of nanothermometry *in vivo* is  
29 provided by the measurement of brain temperature in mice,  
30 upon injection of Ag<sub>2</sub>S NPs. Again, by using excitation and  
31 emission wavelengths of 808 nm and 1200 nm, respectively,  
32 measurements could be recorded in a non-invasive manner  
33 through the skull and scalp. Although thermal calibration  
34 was performed with the particles inside an excised brain, so  
35 as to mimic the experimental environment, conclusions still  
36 relied on thermal variations, as expected for a sensor based  
37 on luminescence intensity. A thermal resolution of  $\pm 0.2$  °C  
38 was estimated. As shown in **Figure 5**, this proof-of-concept  
39 experiment was designed in a way that it nicely underlined  
40 the differences between measuring temperature *in situ* (brain  
41 temperature) or externally through an IR thermal camera  
42 (skin temperature). The results suggested that externally  
43 induced temperature changes (whole body cooling) caused  
44 activation of thermal regulation at the organ level (in the  
45 brain), to reduce the impact of temperature in a homeostatic  
46 method. Thus, thermal changes were up to three-fold larger  
47 on the skin (10 K) than in the brain (3 K). In contrast, when  
48 thermal changes were triggered directly inside the brain  
49 through a barbiturate coma, thermal change was exclusively  
50 observed in the brain, without any external impact on the  
51 skin, so it could only be detected through the nanoparticles.<sup>39</sup>

52 In the above examples, temperature was deduced from the  
53 evolution of the emission intensity from luminescent  
54 nanoparticles. These techniques are valid to determine how  
55 temperature changes over time, but cannot provide absolute  
56 temperature values, which would require knowing exactly to  
57 what extent light gets absorbed by tissue along the way.  
58 Besides, this method assumes that the concentration of  
59 nanoparticles in the spot of interest remains constant during



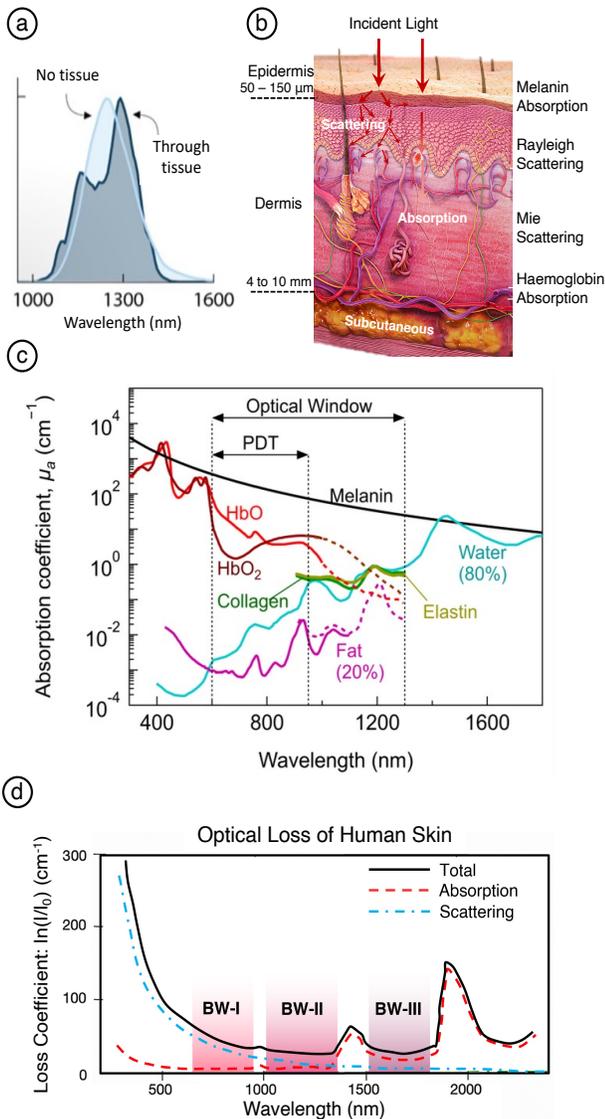
**Figure 5.** Brain thermometry through skull and scalp measured with Ag<sub>2</sub>S nanoparticles emitting in the BW-II (at 1200 nm). Brain temperature is compared with skin temperature measured with a thermal camera, and rectal temperature measured with a thermocouple. Reproduced from reference 39, with permission from John Wiley and Sons Publishing.

the measurement, and requires that the tissue through which light penetrates does not change over time due to the measurement itself. Additionally, the TTh technique also requires control measurements involving healthy tissue. When all of these aspects are taken into consideration, it is clear that nanothermometry measurements *in vivo* are far from simple, and although LNMs can provide advantages such as local delivery and improved sensitivity, ultimately a critical eye must be taken in deciphering the true values.

**Beyond diagnosis: hyperthermia.** Techniques such as radiofrequency or microwave ablation are commonly used to induce tissue necrosis of solid tumors, especially in liver cancers. In both methods, localized heat is induced via insertion of a probe into the diseased tissue, so that sufficiently high local temperatures are achieved. During treatment, a fine control over the applied temperature is crucial to avoid excessive damage to surrounding healthy tissue, and therefore thermocouples or optical fibres are used to determine temperature changes *in situ*. However, these sensors suffer from limited spatial resolution and are intrinsically invasive, two limitations that can in principle be solved with nanothermometry.

Interestingly, various types of nanoparticles can transform incident electromagnetic radiation into heat released to their local environment. This effect is mainly achieved by using plasmonic nanoparticles illuminated at their localized surface plasmon resonance (LSPR) and magnetic nanoparticles subjected to alternating magnetic fields. In the case of plasmonic nanoparticles, heat release is due to the large absorption cross-section at the plasmonic resonance wavelength. At this wavelength, surface electrons

1 participate in a synchronized oscillation, known as a 14  
 2 plasmon, whose energy can be eventually released to the 15  
 3 environment in the form of heat. Regarding magnetic 16  
 4 nanoparticles, alternating magnetic fields are transformed 17  
 5 into heat in different ways, including energy losses during 18  
 6 magnetization-demagnetization cycles (associated to 19  
 7 hysteresis loops), and losses linked to a flipping movement 20  
 8 of the particle (Brown relaxation) or its magnetic moments 21  
 9 (Néel relaxation).



**Figure 6.** (a) Distorted NIR emission from Ag<sub>2</sub>S nanoparticles measured through tissue, compared with the emission free from tissue interference.<sup>46</sup> (b) Schematic view of the skin, showing the main interactions with electromagnetic radiation at different depths, and pointing out the most important causes of optical losses.<sup>47</sup> (c) Main biological molecules and their absorbances in the Vis-NIR electromagnetic range.<sup>48</sup> (d) Optical losses in the NIR for a standard model tissue, showing the biological windows (BW) and differentiating absorption and scattering losses.<sup>53</sup> Figures reproduced under the Creative Commons Attribution License (CC BY 4.0) from references 46, 47, 48, and 53, with permission from John Wiley and Sons Publishing, Wikipedia, MDPI, and The Royal Society of Chemistry, respectively.

in tissues, as well as for other heating applications (see section on cryopreservation below), because they would minimize the invasiveness of localized hyperthermia. In principle, it is also possible to generate heat from some LNMs, which would thereby act as both nanothermometers and nanoheaters, simultaneously. However, this process would involve simultaneous radiative and non-radiative energy processes, the latter requiring significantly higher power densities to achieve sufficient heating. Alternatively, co-administration of nanothermometers together with nanoheaters offers an exceptional flexibility in the design of the system (including excitation range, NP size, surface charge, colloidal stability, etc.). Although simply mixing nanoheaters and nanothermometers could indeed be sufficient to detect local heating upon externally applied electromagnetic excitation, biological systems are affected by aspects such as protein corona formation, EPR effect, or size-dependent cellular uptake. As a result, co-administration is unlikely to result in co-localization of nanoheaters and nanothermometers. Consequently, a more interesting alternative that offers considerable flexibility in the design of the nanoparticle system, is the mutual binding of nanoheaters to nanothermometers via either covalent bonds or electrostatic forces, or even using encapsulation techniques. Approaches toward covalent binding must take into consideration the length of linker molecules or the growth of an intermediate shell, because photoluminescence quenching by nanoheaters may occur. A potentially simpler approach comprises the encapsulation of both nanoparticle types within a biologically compatible shell, such as a polymeric coating. Indeed, the dual encapsulation approach further protects LNMs from any undesired protein corona formation. The protein corona, increasingly termed biocorona because it includes other molecules such as lipids and nucleic acids, has likely a more important effect in altering the amount of LNMs that are delivered to the site of interest, as opposed to directly altering the luminescence readout. Still, it may have some influence, as will be discussed in the section devoted to the particle microenvironment.

**Beyond diagnosis: Cryopreservation.** The work by the Bischof group aimed at bringing the heating properties of highly absorbing NPs to the forefront of biomedicine, by applying them in the rewarming of cryopreserved tissue.<sup>40-42</sup> Specifically, the technique involves the perfusion of tissues and organs, prior to cryopreservation, with a colloidal dispersion of gold NPs featuring a LSPR in a BW. In the presence of chemical cryoprotectors, the tissue sample is vitrified following a critical cooling rate (CCR), specific to the amount and type of cryoprotector. To rewarm or devitrify the sample, rapid warming must take place to avoid any cell damage due to the formation of ice crystals during the devitrification stage. Such a rapid rewarming can be achieved via resonant excitation of a plasmonic mode in gold NPs, reporting rewarming rates above 10<sup>7</sup> K/min, which are considered to be the gold-standard to avoid crystallization.<sup>40</sup> However, it should be noted that the experimental setup relies on a thermocouple that is inserted into a droplet of less than 1 μL, allowing biological sample diameters up to 1 mm in diameter only. Whereas the thermocouple offers an

10

11  
12  
13

Such nanoparticles, which we term nanoheaters, are interesting choices for the induction of coagulative necrosis

1 inexpensive and fast method to detect temperature changes, 59  
2 its level of accuracy is low ( $\approx \pm 2.5$  K RMS).<sup>43</sup> Indeed, such 60  
3 a degree of accuracy is irrelevant in the context of rewarming 61  
4 cryopreserved samples where temperature changes are 62  
5 nearly 100-fold higher. However, in the case of 63  
6 photoablation where a few degrees can determine whether 64  
7 cellular necrosis occurs or not, or for temperature-based 65  
8 measurements of cell malignancy, inaccuracies of  $\pm 2.5$  K 66  
9 are disproportionate. Again, the use of LNMs as a suitable 67  
10 alternative has gained increasing interest due to its accuracy 68  
11 and biocompatibility. Specifically focusing on *in vitro* 69  
12 biological models being exposed to external electromagnetic 70  
13 radiation, the amount of available data is surprisingly small. 71  
14 Conceptually, the simultaneous application of nanoheaters 72  
15 and nanothermometers is not complicated, yet most studies 73  
16 in which photoablation is monitored *in vitro* rely on 74  
17 temperature measurements using an infrared camera. Hence, 75  
18 we should ask the question of why nanothermometers are not 76  
19 used more often. The answer may come from the difficulties 77  
20 in delivering sufficient LNMs to the site of interest, so that 78  
21 the LNM signal is strong enough. A second possibility 79  
22 relates to the structural and biological complexity of the 80  
23 tissue itself, which has an important effect on the intensity 81  
24 of both incident and reflected light, hence questioning the 82  
25 reproducibility of LNMs (disused in further detail below). 83

## 26 Performing the Experiment: How Accurate 85 27 Is the Measured Temperature? 86

28 Most studies related to 2D intracellular *in vitro* temperature 88  
29 measurements are performed using NP probes, some of 89  
30 which are commercially available.<sup>10</sup> With the exception of 90  
31 transfection agents, most temperature nanoprobe enter cells 91  
32 via endocytosis. Once located intracellularly, temperature 92  
33 measurements can be made using a fluorescence 93  
34 microscope, a plate reader, or a spectrofluorometer, for 94  
35 fluorescence emission readout. Such readings must be 95  
36 compared with a previously obtained calibration curve. 96  
37 Thus, the light detection system and the medium used for 97  
38 both the calibration curve and the *in vitro* measurements 98  
39 must be identical, which often involve highly complex 99  
40 media containing high quantities of glucose, amino acids 100  
41 such as glutamine, and albumins or other proteins. Whilst 101  
42 mammalian cells can be studied in transparent buffers with 102  
43 adjusted ionic strength and pH, these are significantly 103  
44 different to their usual environment and therefore might 104  
45 result in undesired stress, which would manifest itself as a 105  
46 temperature change. Therefore, any LNM must ideally be 106  
47 stable in complex media and compatible with visible-light 107  
48 detection techniques, which would drastically facilitate their 108  
49 use in standard commercial set-ups available in most biology 109  
50 laboratories. 110

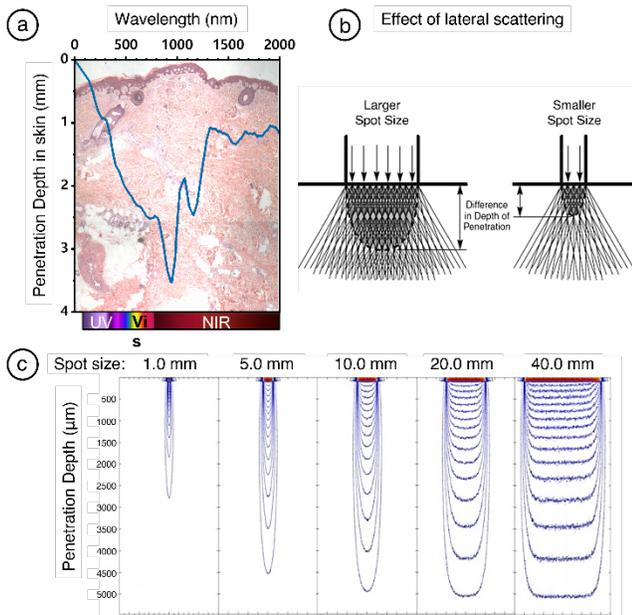
51 On the other hand, application to *in vivo* studies requires that 11  
52 the stability, size, and surface chemistry of 12  
53 nanothermometers be carefully controlled, to avoid 13  
54 unwanted non-specific removal from the circulation via the 14  
55 reticuloendothelial system (RES), and to accumulate at the 15  
56 site of interest. Such an accumulation may require the use of 16  
57 targeting molecules such as antibodies or aptamers, or 17  
58 simply be based on the so-called enhanced permeability and

retention effect (EPR), which is known to occur in  
tumours.<sup>44</sup> From the spectroscopy perspective, conditions  
are also demanding, as the LNMs must be as bright as  
possible and their excitation and emission wavelengths  
match one of the biological windows, for best penetration  
combined with minimal harm. Furthermore, when heating is  
aimed at producing hyperthermia, a common problem is the  
dissipation of heat due to blood flow. Obviously, this is  
highly dependent on the tissue of interest, highly oxygenated  
organs being more affected. The opposite may also happen,  
as poorly oxygenated tissues, such as dense tumours, cause  
a heat-sink and inhibit heat dissipation. Therefore, accurate  
temperature measurements become even more relevant.  
However, can we trust the thermal readout obtained from  
nanothermometers? Although thermal monitoring has been  
demonstrated in small animals, it is still a matter of debate  
how nanothermometers, or more specifically, their emitted  
light, should be recorded and processed to provide accurate  
thermal information. Indeed, several factors limiting the  
accuracy of the measurements have been identified, which  
may hinder the adequacy of the calibration curve, thereby  
compromising the thermal readout.

### Limitations associated to the light path: the tissue issue.

One of the most important aspects of *in vivo* experiments is  
related to tissue scattering and absorption,<sup>45</sup> as on one hand,  
these reduce the fluence of incident light reaching the  
nanothermometer and, eventually, the nanoheater; and on  
the other hand they can change the spectral shape of  
luminescence (**Figure 6a**).<sup>46</sup> The same can be said for the  
luminescent signal emitted by the nanothermometers in  
question, which must be recorded with a sufficient signal-to-  
noise ratio, to guarantee a meaningful thermal readout. In  
general, absorbance is caused by soluble components  
present in the intra- and extra-cellular environments,  
whereas scattering occurs due to the physical organization  
of the tissue, which leads to refractive index contrast (**Figure  
6b**).<sup>47</sup> Both effects are remarkably dependent on wavelength.  
In the case of absorption, the spectral range is defined by the  
allowed electronic transitions in molecules, and thus  
strongly depends on chemical composition. The absorption  
spectrum of a tissue will thus depend on the concentration of  
each component, haemoglobin and water typically making  
the main contributions (**Figure 6c**). However, for some  
tissues, melanin and yellow pigments ( $\beta$ -carotenes and  
bilirubin) play a major role in the visible range, while fat and  
adipose tissue contribute in the NIR.<sup>48,49</sup>

Absorption is important, not only because of the optical  
losses in the light beam, but also because the absorbed  
optical energy is then transformed into a different type of  
energy, most often into heat, thereby leading to an increase  
in local temperature. A maximum illumination dosage  
(power and time) to avoid tissue overheating has thus been  
defined (guidelines are regularly published and updated by  
the International Commission on Non-Ionizing Radiation  
Protection, ICNIRP)<sup>50</sup>. Alternatively, absorbed energy can  
be released radiatively, in which case the tissue itself  
presents light emission, constituting what is known as  
autofluorescence. Just as light emission does,  
autofluorescence may overlap in the detector with the  
emission from LNMs and cause problems, including



**Figure 7.** (a) Light penetration depth in human skin, overlaid on a histological cut stained with eosin and haematoxylin, showing the biological structures that light can reach depending on its wavelength. (b, c) Dependence of the penetration depth of light, ranging from 525 to 1100 nm, as a function of the incident beam diameter. Reproduced under the Creative Commons Attribution Licence (CC BY 4.0) from reference 56, with permission from Springer Nature.

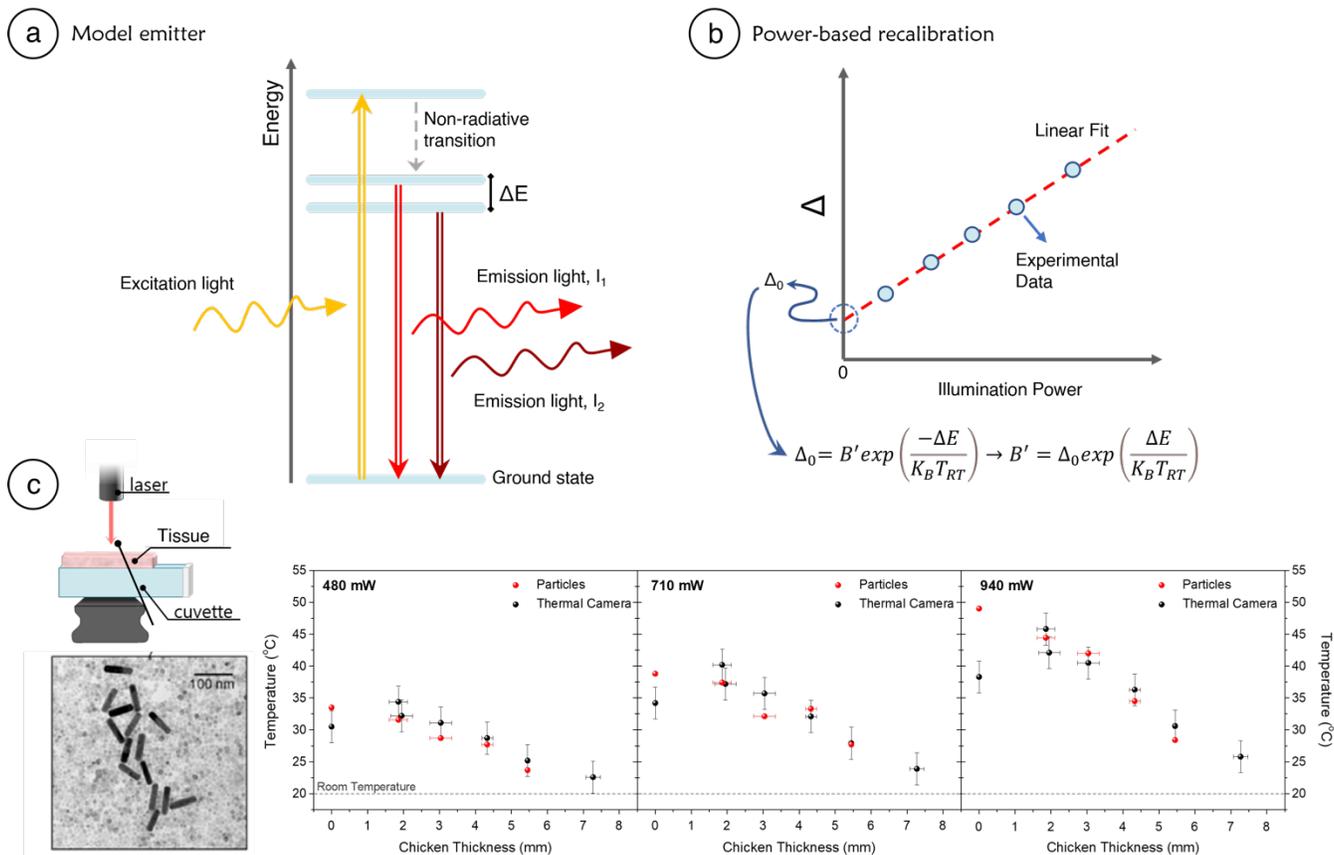
1 uncontrolled background signal or additional emission  
 2 bands that may hinder those used for thermometry.  
 3 Fortunately, most endogenous components of *in vivo* and *in*  
 4 *vitro* biological samples display autofluorescence in the  
 5 blue-green optical range and thus it does not represent a  
 6 significant issue for measurements within the biological  
 7 windows. However, autofluorescence in the NIR has also  
 8 been reported, e.g. associated to melanin or chlorophylls,  
 9 which are present in commercial food for experimentation  
 10 animals,<sup>50</sup> oftentimes even if labelled as “fluorescence free”.  
 11 It is thus important to be aware of the expected composition  
 12 of the tissue under study and its potential autofluorescence,  
 13 to avoid misinterpretation of thermal data.<sup>51</sup>

14 For scattering, on the other hand, the degree of interaction  
 15 depends on the relative size of tissue fibrils, cells, or other  
 16 obstacles for light, compared to the wavelength of incident  
 17 light. Traditionally, two scattering modes are described:  
 18 Rayleigh scattering, when the obstacle is significantly  
 19 smaller than light wavelength; and Mie scattering, when  
 20 both dimensions are in the same range. For longer  
 21 wavelengths, scattering becomes negligible, which for  
 22 tissues applies in the NIR (**Figure 6d**).<sup>53</sup> In summary, both  
 23 Rayleigh and Mie scattering contributions should be  
 24 considered to optically describe a tissue, even though often  
 25 one of them dominates over the other. For instance, fatty  
 26 tissues or brain tissue are usually dominated by Mie  
 27 contributions, whereas the skin and other fibrous tissues  
 28 display predominantly Rayleigh scattering.<sup>52</sup>

29 The exact light attenuation (a term that includes both  
 30 scattering and absorption) will thus vary for different tissues.  
 31 However, a generic spectrum can be approximated as shown

32 in **Figure 6d**, where the biological windows can be clearly  
 33 identified. It is well documented that successful penetration  
 34 of light through *in vivo* environments, with multiple tissue  
 35 boundaries, requires the selection of wavelengths within one  
 36 of the BWs, in the NIR. By focusing light within these BWs,  
 37 the absorbing properties of water, haemoglobin, lipids, etc.  
 38 are partially avoided,<sup>54</sup> thereby allowing the maximum  
 39 energy to be transferred to the site of interest. The  
 40 characteristic length that light of a certain wavelength can  
 41 travel through tissue is known as penetration depth. Since  
 42 light intensity can be assumed to decay exponentially in the  
 43 propagation direction, the penetration depth is typically  
 44 defined as the characteristic distance of the exponential  
 45 decrease (i.e., the distance at which the normalized intensity  
 46 decays to 1/e times its initial value). Penetration depth varies  
 47 with wavelength,<sup>55</sup> for the reasons discussed above, and is  
 48 obviously longer for light within the biological windows  
 49 (**Figure 7a**). However, it also depends on illumination  
 50 conditions such as spot size, because lateral scattering of the  
 51 light beam is less relevant for bigger spot sizes (**Figure 7b**).  
 52 As an example, Monte Carlo simulations applied to light-  
 53 tissue interactions through skin show that maximum  
 54 penetration is reached for a 10 mm beam width, while a  
 55 further increase of the beam size does not have any  
 56 additional effect on penetration (**Figure 7c**).<sup>56</sup> Considering  
 57 all of the above discussion, at the time of designing an  
 58 experiment it is important to select an appropriate  
 59 wavelength within the biological windows, but also to  
 60 consider the maximum illumination dose that can be used to  
 61 avoid tissue damage, and even to select the appropriate size  
 62 of the illuminated area and improve penetration depth.

The extent of scattering and absorption varies in tissues,  
 between healthy and diseased states. Of significant concern  
 are changes in scattering and absorption that may occur *in*  
*situ*, due to an applied therapeutic treatment.<sup>57</sup> In thermal  
 ablation techniques, it is well-known that real-time  
 monitoring is required to control excessive heating, which  
 may easily cause undue damage to surrounding tissues. A  
 wide range of cellular and molecular changes can take place,  
 including protein denaturation, cell membrane damage,  
 production of apoptotic bodies, changes in osmolarity, etc.,  
 all of which can in turn enhance scattering and absorption of  
 both excitation and emission light.<sup>58</sup> This may also happen,  
 for instance, during the measurement of transient  
 temperatures in TTh, as a preliminary heating step is  
 required prior to recording the cooling curve that provides  
 tissue information. To study the extent of such a change,  
 Lifante et al. used Ag<sub>2</sub>S semiconductor nanoparticles (13 nm  
 diameter), specifically prepared to display a lower  
 photoluminescence quantum yield, but a high heating  
 efficiency instead. Upon injection in mice and illumination  
 at different powers, a range of target temperatures to  
 measure TTh curves was obtained. From the results, it was  
 concluded that the characteristic relaxation time varied with  
 temperature, which could be correlated to a different heat  
 diffusion within the tissue. Based on numerical modelling, it  
 was concluded that changes were produced in the tissue at  
 the higher illumination powers, which could be ascribed to  
 changes in blood perfusion through tissue.<sup>59</sup>



**Figure 8.** (a) Example of a diagram of energy states with thermalized levels. (b) Scheme for the recalibration of  $B$ , based on the dependence of the intensity ratio with illumination power. (c) Model *ex vivo* experiment, comparing temperature measured through tissue and externally with a thermal camera. Reproduced with permission from reference 60, American Chemical Society.

1 **The need to recalibrate.** The absorption and scattering of 30  
 2 light by tissue, even within the biological windows, 31  
 3 represent a major source of inaccuracy in optical 32  
 4 nanothermometry. This is most clear in the case of intensity- 33  
 5 based thermometers not calibrated *in situ*, which can only be 34  
 6 used to obtain transient temperatures, as long as neither the 35  
 7 tissue properties nor the NP distribution changes during the 36  
 8 experiment. If we shift from one-intensity to intensity-ratio 37  
 9 thermometry, inaccurate control over excitation light or NP 38  
 10 distribution would no longer be an issue. However, both  
 11 scattering, and absorption depend on wavelength, thus the 39  
 12 shape of the emission spectrum may change along the path  
 13 of light through tissue (**Figure 6a**). The spectral distortion  
 14 becomes more significant in thicker tissue, but it will also 40  
 15 depend on other tissue features (blood flow, percentage of 41  
 16 fat, melanin, etc.). Consequently, working at two different 42  
 17 wavelengths generates the drawback that each of them will 43  
 18 interact differently with the tissue. As a result, if the 44  
 19 calibration of the intensity ratio has not been performed in 45  
 20 *situ*, it will most likely not be valid in the final experiment. 46  
 21 Unfortunately, calibration requires measuring the emission 47  
 22 spectrum along the whole thermal range of interest 48  
 23 (commonly from room temperature up to 50 or 60 °C, at 49  
 24 least), which is not viable during diagnosis or therapy. 50  
 25 Besides, conducting multiple calibrations is a time- 51  
 26 consuming task, which is not the most convenient situation 52  
 27 in biomedical experiments. 53  
 28 Solutions to the above issue have been proposed by using 54  
 29 intensity ratios in thermally linked states of lanthanide-

doped nanoparticles. The procedure involves a preliminary  
 calibration in a controlled environment, which will be  
 subsequently adapted to the final experiment *in situ*,  
 following simple steps that do not require heating the sample  
 to specific set points. This is possible because, in thermally  
 linked ratios, two states share the electronic population, with  
 a distribution that is described by a Boltzmann function.  
 Therefore, their emission intensities are also linked through  
 the same law, and the intensity ratio can be described as:

$$\Delta = \frac{I_1}{I_2} = B \cdot e^{\left(\frac{-\Delta E}{k_B T}\right)} \quad (3)$$

Here,  $\Delta$  is the thermal parameter to be calibrated, i.e. the  
 ratio between intensities  $I_1$  and  $I_2$  (**Figure 1c**),  $I_1$  being the  
 intensity of the most energetic state and  $I_2$  the intensity of  
 the least energetic one (**Figure 8**);  $K_B$  is Boltzmann's  
 constant,  $T$  the temperature and  $\Delta E$  is the energy separation  
 between the barycentres of the two involved energy levels  
 (**Figure 8a**).<sup>60</sup> Finally,  $B$  is a value that depends on  
 spectroscopy parameters such as the degeneracies of the  
 involved states, their spontaneous emission rates, and their  
 average spectral frequency. In the first reports using this  
 formula,<sup>61</sup> the definition of  $B$  also included the wavelength  
 dependence of the detection system response, which is fair  
 because it also affects the recorded spectrum. However, this  
 consideration was made assuming a homogeneous,  
 transparent medium between the emitting sample and the  
 detector. If, instead, light travels through a non-transparent

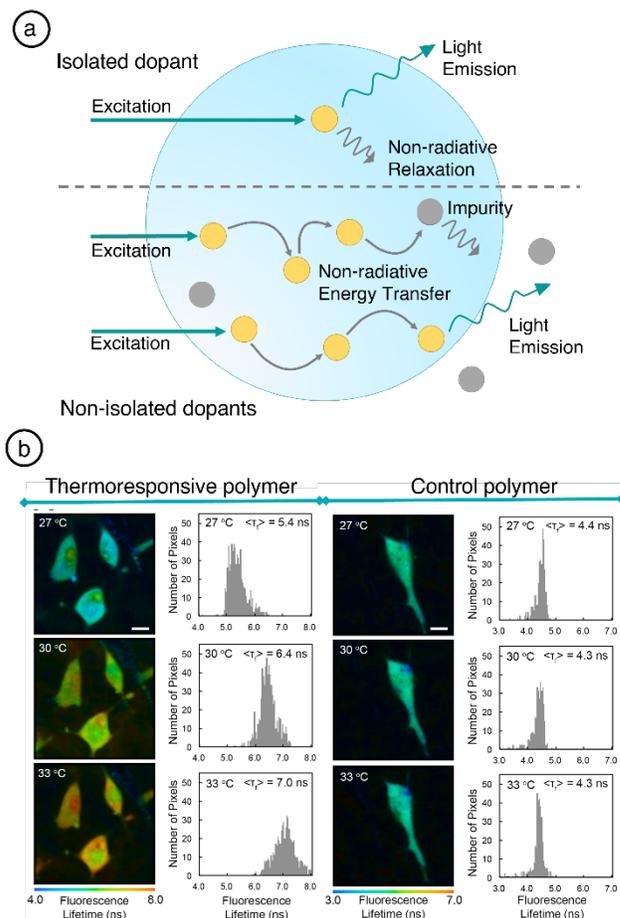
1 medium, the spectrum will be distorted as discussed above,  
 2 in turn leading to a change in the B value. Thus, B is the  
 3 parameter that must be recalibrated in situ, whereas  $\Delta E$  can,  
 4 in principle, be considered constant (note that, in lanthanide-  
 5 doped crystals, this consideration is accurate because of the  
 6 shielding of the  $4f$  orbital).

7 This technique was first applied by Balabhadra et al., who  
 8 showed the relevance of this effect on the calibration curve  
 9 of  $\text{SrF}_2:\text{Er}^{3+}, \text{Yb}^{3+}$  nanothermometers in different solvents.<sup>62</sup>  
 10 The proposed method comprised a first thermal calibration  
 11 of the emission, to determine  $\Delta E$ . Upon transfer of the  
 12 particles into a different solvent, B was re-evaluated through  
 13 a study of emission as a function of power, which is simpler  
 14 because it does not require accurate thermal control. As any  
 15 solvent will typically absorb some light, it will heat up  
 16 mildly, in a way that  $\Delta$  follows a linear dependence with the  
 17 intensity ratio.<sup>63</sup> Thus, a linear fit of  $\Delta$  versus illumination  
 18 power allows determining  $\Delta_0$ , i.e. the intensity ratio at room  
 19 temperature,  $T_{RT}$  (**Figure 8b**). This value of  $\Delta_0$  can then be  
 20 used to define the required new B, which we denote as B'  
 21 (**Figure 8b**) as:

$$B' = \Delta_0 \cdot e^{\left(\frac{-\Delta E}{k_B \cdot T_{RT}}\right)} \quad (4)$$

22  
 23 With this new B' parameter it is now possible to accurately  
 24 measure temperature in the new solvent, with no need to  
 25 change the temperature externally. This technique has also  
 26 been tested for both excitation and emission light beams  
 27 travelling through *ex vivo* tissue (chicken breast, see **Figure**  
 28 **8c**).<sup>62</sup> In this experiment, an aqueous colloidal dispersion,  
 29 containing a mixture of  $\text{CaF}_2:\text{Nd}^{3+}, \text{Y}^{3+}$  nanothermometers  
 30 and gold nanorods, was prepared to simulate a hyperthermia  
 31 experiment. The mixed dispersion was poured into an  
 32 optical-glass cuvette, which could be readily monitored with  
 33 a thermal camera. Both  $\text{Nd}^{3+}$  ions and gold nanorods could  
 34 then be excited with a laser at 808 nm, within the BW-I. The  
 35  $\text{Nd}^{3+}$  emission at 1050 nm (BW-II) was recorded during  
 36 irradiation, to measure temperature changes. The results  
 37 showed a good agreement between thermal readouts from  
 38 the nanoparticles and from the IR camera. However, both  
 39 values only matched if B was readjusted following the  
 40 protocol explained above. Otherwise, differences as large as  
 41 tens of degrees would have been detected when light was  
 42 travelling through only 5 to 6 mm of tissue.

43 Aiming for a one-step process when nanothermometers are  
 44 administered into a biological sample, the same group tested  
 45 an alternative option for readjustment of B. The experiments  
 46 were performed *in vitro*, using 3D tumour spheroids  
 47 immersed in cell media, which had been previously treated  
 48 with a hybrid probe comprising polystyrene colloidal  
 49 spheres covered by both  $\text{CaF}_2:\text{Nd}^{3+}, \text{Y}^{3+}$  nanoparticles  
 50 (nanothermometers) and gold nanostars (nanoheaters).<sup>32</sup> The  
 51 measurement protocol in this case was based on monitoring  
 52 the evolution of temperature upon turning the illumination  
 53 on, which was found to follow an initial fast increase before  
 54 slowly reaching a plateau. The first ramp could be fitted to a  
 55 straight line, which at time zero would provide the value of  
 56  $\Delta_0$  needed to calculate B'. The main advantage of this  
 57 method is that it avoids the need for a preliminary power  
 58 calibration, which is an asset in hyperthermia experiments.



**Figure 9.** (a) Different mechanisms affecting photoluminescence intensity and lifetime, as described in the text. Note that all available processes compete. (b) Thermal maps calculated from LNM fluorescence lifetime. Reproduced with permission from reference 21.

However, it does require that the system can record several data points during the first minute of illumination, which is not always possible if the luminescence signal is too weak.

In fact, although these techniques solve the problem caused by the interaction of emitted light with tissue, the fact that they are based on lanthanide-doped nanoparticles, whose luminescence is often weaker than luminescence derived from other types of LNMs, precision and penetration depth are hindered. For instance, **Figure 8c** does not show thermal data from LNMs emitting from under a tissue thicker than 5.5 mm, as the measured signal in such experiments showed low signal-to-noise ratios. Besides, equation (3), which allows the described recalibration protocol, also involves a limitation regarding the sensitivity of the technique, which is defined by  $\Delta E$  and B.<sup>20</sup> Thus, whereas many materials might not have enough sensitivity for intracellular temperature readings, it is generally enough for therapeutic and diagnostic applications.

We have seen in this section that thermometry methods based on either intensity or parameters that vary with wavelength (bandwidth and peak position) should be applied with care, taking into consideration the varying response of biological tissue at different wavelengths. This is especially important for calibrations being performed out of the final

1 experimental situation, which may not sufficiently represent 60  
2 the real environment. Additionally, if harsh conditions are 61  
3 applied to the tissue, either due to a disease condition or 62  
4 because of a therapy, optical and heat diffusion properties of 63  
5 the tissue may change, even during the treatment itself. It 64  
6 should be noted that emission lifetime is in principle the one 65  
7 property that is independent of tissue attenuation, because 66  
8 lifetimes are measured at a specific wavelength and thus not 67  
9 affected by spectral variations. However, the spectral 68  
10 distortion of light travelling through tissue is not the only 69  
11 reason why calibrations may lose significance. Indeed, the 70  
12 actual microenvironment where the LNMs are to perform 71  
13 may also influence their optical properties, lifetime in 72  
14 particular. 73

15 **Limitations associated to the microenvironment.** When 75  
16 we introduced above the main emission parameters in 76  
17 luminescence thermometry, we explained that an excited 77  
18 state can release its energy both radiatively and non- 78  
19 radiatively, and that the probability of each process is given 79  
20 by the rates  $k_r$  and  $k_{nr}$ , respectively. Both  $k_r$  and  $k_{nr}$  define the 80  
21 photoluminescence quantum yield of the fluorophore,  $\eta$ , and 81  
22 its experimental lifetime,  $\tau_{exp}$ , as given by equations (1) and 82  
23 (2). As mentioned above, in the simplest scenario of an 83  
24 isolated luminescent atom in a perfect crystal matrix, the 84  
25 direct link between  $k_{nr}$  and temperature triggers a thermal 85  
26 dependence of  $\eta$  and  $\tau_{exp}$ . However, further than temperature, 86  
27 are these parameters related to other environmental features? 87

28 In most materials, the situation is not as simple as the one 88  
29 described above. Continuing with the example of the dopant 89  
30 in a crystal matrix, if we increase the concentration of 90  
31 dopants to the extent that we can no longer consider them as 91  
32 isolated, new relaxation paths arise (**Figure 9a**).<sup>64</sup> In this 92  
33 case, energy can travel from one dopant to another through 93  
34 non-radiative energy transfer processes. Alternatively, some 94  
35 crystal impurities with additional vibrational modes can 95  
36 enhance  $k_{nr}$  if they are present in the surroundings of the 96  
37 dopant. Both non-radiative energy transfer and additional 97  
38 vibrational modes will reduce then the emission quantum 98  
39 yield and are thus called quenching routes. These additional 99  
40 relaxation routes are more likely to occur as the distance 100  
41 between the quencher and the luminescent dopant is 101  
42 reduced. Although the threshold proximity is different for 102  
43 each interaction, as a reference we can say that it is typically 103  
44 in the order of few nm. A final effect to be considered is the 104  
45 potential modification of transition rates due to alterations of 105  
46 the local electromagnetic field, including e.g., defects in the 106  
47 crystal matrix, such as strain, which trigger variations of the 107  
48 crystal field. However, the presence of external fields should 108  
49 also be considered, especially when plasmonic nanoparticles 109  
50 are used to enhance the emission intensity of nearby 110  
51 fluorophores.

52 Keeping this general overview in mind, it is straightforward 111  
53 to see why the environment around LNMs can strongly 112  
54 affect their luminescence properties. Sticking to the dopant 113  
55 example, the small size of NPs imposes interactions of the 114  
56 luminescent dopant, not only with other atoms in the particle 115  
57 itself, but also with atoms or molecules in the surrounding 116  
58 medium (**Figure 9a**). In terms of optical sensing, this means 117  
59 that the properties of the sensor, including the thermal 118

calibration for temperature determination, may change if the 119  
environment does. In applications where the environment is 120  
homogeneous this is a negligible problem, as the effect will 121  
be the same during the experiment and the calibration. 122  
However, heterogeneous and complex environments, such 123  
as those found *in vitro* and *in vivo*, represent a major 124  
challenge because modifications to  $k_r$  and  $k_{nr}$  are hard to 125  
predict. This includes the fact that, upon introduction in the 126  
biological environment, a biocorona directly surrounding the 127  
particle will be formed, with a composition that will depend 128  
on the availability of molecules at the specific location of the 129  
particles, but also on the NP surface charge. Considering that 130  
the use of LNMs in biological environments is relatively 131  
recent, little data are available on the direct effect of the 132  
biocorona on optical sensors or at different temperatures. 133  
One example described the effect of the hard and soft 134  
coronas on the emission intensity of magnetic NPs 135  
functionalized with an  $\text{Eu}^{3+}$  complex.  $\text{Eu}^{3+}$  ions present 136  
certain electric dipole transitions that are sensitive to the 137  
electromagnetic local symmetry, and thus their light 138  
emission can be used to extract information about it, 139  
including atom bond distances.<sup>65</sup> These findings suggest a 140  
distortion in the local point symmetry upon addition of 141  
different concentrations of blood plasma to the sample, 142  
supporting an increased interaction with plasma proteins. 143  
Also important is the fact that the biocorona may yield 144  
autofluorescence and light absorption bands that potentially 145  
interact with luminescent probes, causing energy transfer 146  
processes especially in the visible range.

Another excellent example is the effect of the environmental 147  
refractive index on the lifetime of optical nanosensors. 148  
Experimental lifetime,  $\tau_{exp}$ , which is the parameter measured 149  
in lifetime-based thermometry, as given by equation (2), 150  
includes radiative and a non-radiative contributions. 151  
Focusing on the radiative lifetime ( $\tau_r=1/k_r$ ), described as 152  
early as 1926, Perrin wrote for the first time an expression 153  
for  $\tau_r$  stating that it depends on the refractive index,<sup>66</sup> as does 154  
 $k_r$ . The question then arises, how can the refractive index be 155  
defined? If we take again the example of dopants in a crystal 156  
host, the refractive index of a bulk material would be that of 157  
the crystal host. However, when the material is reduced in 158  
size and becomes smaller than the wavelength of light, this 159  
assumption is no longer accurate. Indeed, a definition has 160  
been proposed for an effective refractive index that takes 161  
into consideration the indexes of both the host and the 162  
solvent.<sup>67, 68</sup> This effective index depends on a filling factor, 163  
 $x$ , which represents how big the nanoparticle is, compared to 164  
the wavelength of light:

$$n_{eff} = x \cdot n_h + (1 - x) \cdot n_{sol} \quad (5)$$

where  $n_{eff}$ ,  $n_h$  and  $n_{sol}$  are the effective refractive index, the 165  
index of the host, and the index of the solvent, respectively. 166  
Illustrative examples have been provided by Meijerink and 167  
co-workers, using 4 nm diameter  $\text{LaPO}_4:\text{Ce}^{3+}$  nanocrystals 168  
to demonstrate that a variation of the solvent refractive index 169  
from 1.36 to 1.48 triggers a  $\approx 26\%$  change in  $\text{Ce}^{3+}$  lifetime 170  
(from 35 to 26 ns).<sup>69</sup> Similar results were obtained in 4 nm 171  
diameter  $\text{LaPO}_4:\text{Tb}^{3+}$  nanoparticles coated by a 1 nm thick 172  
undoped shell. The shell was shown to have a negligible 173  
effect on the refractive index, because its thickness was too 174

1 small compared to the wavelength. It was however 61  
2 beneficial to increase the luminescence intensity, as it 62  
3 provided a separation between emitting ions in the particle 63  
4 ( $Tb^{3+}$ ) and molecules in the solvent that could otherwise 64  
5 non-radiatively act as quenchers through their vibrational 65  
6 modes (vibrational modes of C-H and N-H bonds, mainly). 66  
7 Indeed, the luminescence quantum yield was shown to 67  
8 improve in  $\approx 13\%$  in the presence of the coating shell. 68  
9 In a different example, based on the same type of NPs co- 69  
10 doped with  $Ce^{3+}$  and  $Tb^{3+}$ , the same group demonstrated that 70  
11 the efficiency of energy transfer between both ions 71  
12 decreased as the external refractive index increased. In this 72  
13 material,  $Ce^{3+}$  could absorb the excitation energy, and then 73  
14 transfer it to  $Tb^{3+}$ , which acted as the emitter. When 74  
15 changing the solvent refractive index, the emission rate of 75  
16  $Ce^{3+}$  also changed, affecting the energy transfer efficiency 76  
17 and ultimately the emission from  $Tb^{3+}$ .<sup>70</sup> In these examples, 77  
18 we have seen the consequences of environmental changes on 78  
19 the luminescence emission from metal ions within a crystal. 79  
20 These conclusions can also be extended to other types of 80  
21 fluorophores, but the extent of the resulting effect will be 81  
22 different for each specific case. For instance, in some 82  
23 organic molecules such as GFP, lifetime is so sensitive to the 83  
24 surrounding refractive index that it becomes a suitable 84  
25 technique to measure its value in organelles inside cells.<sup>71,72</sup> 85  
26 Using this and other techniques, it has been determined that 86  
27 the refractive index within a cell varies with protein 87  
28 concentration, salt concentration and, to a lower extent, with 88  
29 temperature. Still, typical refractive index values range from 89  
30 1.355-1.365 for the nucleus, 1.400-1.420 for the 90  
31 mitochondria, 1.46 for the plasma membrane, and around 91  
32 1.60 for lysosomes.<sup>71-74</sup> If we instead consider *in vivo* 92  
33 situations in which whole tissues or organs are to be studied, 93  
34 the same range of refractive indexes can be identified when 94  
35 comparing different locations in bones, skin, eyes, etc.<sup>75</sup> 95  
36 According to the above referred studies under controlled 96  
37 conditions, it is clear that the optical properties of a thermal 97  
38 LNM probe may vary (even significantly) at different 98  
39 locations within a cell or tissue. The medium refractive 99  
40 index should be known prior to devising a specific 100  
41 thermometry experiment, to avoid inaccurate thermal 101  
42 readings. Indeed, the three parameters described above 102  
43 regarding the effect of refractive index (lifetime, intensity, 103  
44 and energy transfer) also depend on temperature, and have 104  
45 been proposed for nanothermometry. 105

46 With this information at hand, the practical use of LNMs in 106  
47 heterogeneous environments may seem unlikely. However, 107  
48 understanding the problem is the basis on top of which 108  
49 accurate nanothermometers can be developed. Indeed, some 109  
50 good examples have been reported that pave the way 110  
51 towards practical sensors. Naturally, the proposed sensors 111  
52 should first be tested, to check the extent to which 112  
53 environmental effects may interfere with thermal readings. 113  
54 Indeed, in some reported examples the effect of the 114  
55 microenvironment appears to be negligible (or at least within 115  
56 measurement error). This was shown by Hayashi et al.,<sup>21</sup> 116  
57 using a thermoresponsive polymer whose emission lifetime 117  
58 changed with temperature. A series of careful control 118  
59 experiments showed that the LNM lifetime did not depend 119  
60 on pH or viscosity, but slightly changed with ionic strength 120

As the nanothermometers were designed for intracellular thermometry, ionic strength might compromise the accuracy of the thermal reading. To solve this problem, a control experiment was performed using another polymer with negligible thermal sensitivity, but with the same dependence of lifetime on ionic strength as that of the LNM. Introducing both polymers into HeLa cells, it was possible to map the fluorescence lifetime internally (**Figure 9b**) and, from a comparison of both samples, thermal information was derived, ensuring that ionic strength was not affecting the readout. In this way, a  $\approx 1$  °C thermal difference between nucleus and cytoplasm was determined.

**Towards accurate thermal measurements.** When a new material is conceptually proposed for nanothermometry, a complete thermal calibration must be performed under a well-defined state, typically as a dry powder or as a colloidal dispersion. However, when the targeted applications involve complex biological environments, more extensive calibration may be required. Aiming to clarify what a complete characterization would be, a set of steps for a standard protocol have been proposed,<sup>11</sup> specifically focusing on biomedical applications. When dealing with LNMs destined for *in vitro* studies, the two main identified parameters are pH and dynamic viscosity, in particular focusing on a pH range of 6.0 - 7.5, and viscosity range of  $1 \times 10^{-3}$  - 20 Pa.s. The situation is more complex for *in vivo* studies because LNM concentration and aggregation are difficult to control. It is therefore advised to perform an *ex vivo* calibration in extracted organs relevant to the expected biodistribution.

Although the tested nanothermometers may appear to be accurate, a good strategy to ensure correct thermal readings comprises including an additional nanothermometer control material within the same experiment. This test is particularly interesting if the selected thermal probe is based on a different luminescence temperature measurement strategy, so that environmental changes are unlikely to equally interfere with both mechanisms. Interference of environmental conditions other than temperature would then lead to a disagreement between both probes, which can be readily detected. This strategy has been used in intracellular studies analysing thermogenesis in brown adipocytes.<sup>28</sup> In this work, one of the probes was based on a thermoresponsive polymer and its fluorescence lifetime,<sup>21</sup> whereas a second thermometer was based on the emission intensity of a hydrophilic gel that would collapse upon heating, expelling water molecules and thereby enhancing the emission of a fluorophore inside.<sup>76</sup> The combination of these two temperature measurement strategies is particularly relevant for *in vivo* experiments where many environmental factors may interfere with the thermal reading.<sup>11</sup> One practical requirement is that both luminescence signals can be easily differentiated. In the previous example, both polymers can be excited at around 450 nm and present an emission band ranging from 500 to 700 nm, with a maximum at *ca.* 560 nm. Since there is a significant overlap in the emission properties of both polymers, ideally two separate experiments should be performed to achieve maximum thermal accuracy. Whilst feasible *in vitro*, using standard fluorescence microscopes programmed to do sequential

1 imaging in a short time frame, this is not practical *in vivo*, 61  
2 under diagnostic and/or therapeutic situations. Hence, both 62  
3 thermal probes should display clearly differentiated 63  
4 luminescent properties that allow simultaneous 64  
5 measurements of both LNM emissions.<sup>77</sup> 65

6 It is thus clear that both careful experimental design and a 66  
7 complete calibration are needed, and rationally adapted to 67  
8 each proposed application. However, these tasks may 68  
9 become simpler if nanothermometers are specifically 69  
10 designed to minimize interferences from surrounding 70  
11 molecules with their optical properties. For instance, the 71  
12 distance between luminescent ions and a possible quencher 72  
13 is key to determine how likely the quenching effect would 73  
14 be noticed. Taking lanthanide-based luminescence as an 74  
15 example, the influence of surrounding molecules on doped 75  
16 nanoparticles is typically lower than it would be on 76  
17 lanthanide-based organic complexes, in which the 77  
18 luminescent ion is less isolated.<sup>68</sup> For the same reason, NP 78  
19 size also plays a role because in larger particles the dopants 79  
20 are farther away from the surrounding environment. Another 80  
21 size-related effect can be derived from the definition of the 81  
22 effective refractive index, as shown in equation (5), as larger 82  
23 particles will be closer in size to the illumination wavelength 83  
24 and thus scatter light more efficiently. A suitable alternative 84  
25 to making larger nanoparticles may be the encapsulation 85  
26 within (polymeric) micelles or growing a non-luminescent 86  
27 inorganic shell. In the case of luminescent NPs such as QDs, 87  
28 growth of a non-luminescent inorganic shell is the preferred 88  
29 option because core particle size largely determines the 89  
30 emission wavelength of the QD. In the case of lanthanide- 90  
31 doped NPs, a coating shell would also be efficient, as 91  
32 quenching paths to external molecules are interrupted, an 92  
33 approach that might be beneficial to keep size to a minimum. 93  
94

## 34 Conclusion and perspectives 95

35 Precise monitoring of temperature is critical toward 96  
36 understanding many naturally occurring biological 97  
37 processes, ultimately aiming at the diagnosis and treatment 98  
38 of various pathophysiological disease states, and even to 99  
39 improve current cryopreservation techniques. Regardless of 100  
40 the final application for temperature determination at the 101  
41 cell, tissue, or organ level, the use of non-invasive methods 102  
42 is of crucial importance. Luminescent nanomaterials are 103  
43 proposed as the optimum choice. These materials feature 104  
44 temperature-dependent photoluminescence, with tailored 105  
45 excitation and emission in the biological windows. As such, 106  
46 the interactions between the photoluminescent properties of 107  
47 LNMs and the absorbing and scattering properties of 108  
48 biological tissue are minimized. However, both the 109  
49 excitation beam and emitted light interact with tissue in ways 110  
50 that must be understood and considered to accurately 111  
51 measure temperature. 112

52 If we focus on the effects of a biological environment on the 113  
53 excitation beam, optimization of beam width can be an asset 114  
54 because lateral scattering may contribute to reach deeper 115  
55 tissue. Also, optimization of the final thermal reading 116  
56 requires a good knowledge of autofluorescence within the 117  
57 biological windows. With the development of biomedical 118  
58 techniques exploiting these wavelength ranges, it has been 119  
59 discovered that some components such as melanin or animal 120  
60 food are luminescent in the NIR. It would not be surprising, 121

that, during coming years, other biological components are  
also found to emit, but it can be expected that this  
autofluorescence will be lower than that in the visible.

If we focus instead on the interaction of emitted light with  
tissue, the major problem to resolve is the deformation of  
luminescence spectra due to the absorption by tissue, which  
is a relevant source of error when absolute temperature is to  
be measured. For this reason, most examples in which  
temperature has been measured *in vivo* evaluate transient  
rather than absolute temperature. Some recalibration  
solutions have been proposed to circumvent this problem,  
which occurs mainly in LNMs based on rare-earth doped  
nanoparticles with thermalized energy states. Additional  
solutions will arise through the development of materials  
with a thermalized ground level in which recalibration is not  
needed because the intensity ratio is built from the same  
emission band, excited at two different wavelengths.  
However, these options are based on lanthanide-doped  
materials, typically with a low emission intensity compared  
to other fluorophores. This is a major limiting factor  
regarding penetration depth and thermal resolution, and thus  
more efficient emitters are desired for in depth applications.  
Regarding lanthanides, strategies based on the use of optical  
antennas to better absorb the excitation light have been  
proposed, but to our knowledge not yet implemented in  
(bio)applications. An alternative path may involve LNMs  
doped with transition metals, but so far they have not  
reached the performance of lanthanide-doped materials,  
especially in the thermal range of interest in biology.

Given the increasing awareness of the problems arising from  
the interaction of light with tissue, new measurement  
strategies will likely arise, aiming for better thermometry  
performance with simpler experimental setups that can be  
implemented in biology labs or in clinical settings.  
Improvement of penetration depth or thermal resolution to  
measure absolute temperature may be achieved with  
lifetime-based thermometers, as they are not affected by  
spectral distortions. However, a complete characterization of  
the material and its luminescence in different environments  
will be key to ensure sufficient accuracy.

The relative sensitivity of most nanothermometers is often  
insufficient, at least when the variation of only one  
parameter is taken into account. Hence, methods have been  
developed to measure temperature via two luminescence  
emission characteristics simultaneously, e.g., peak position,  
band width, polarization anisotropy, etc., and then apply  
multiparameter linear regression for temperature  
determination. This can be carried out by using either two  
different probes or just one with multiple readout options,  
which is possible with Ag<sub>2</sub>S nanoparticles. Given its high  
versatility and brightness, Ag<sub>2</sub>S is one of the most promising  
materials towards practical *in vivo* applications, and  
particularly interesting synthesis strategies have been  
patented. Whereas multiparameter and multiprobe schemes  
improve the accuracy and reproducibility of LNMs for  
temperature sensing, it is paramount that the interactions  
between excitation and emission photons with the biological  
milieu be kept to a minimum, to avoid loss of signal  
intensity, and hence sensitivity, during the measurement.  
This multiparameter technique is also relevant if we consider  
the effect of the LNM's microenvironment on their

1 spectroscopic properties. To avoid measurement errors 58  
2 based on this interaction, standardization protocols aim to 59  
3 unveil any uncontrolled change on the luminescence that 60  
4 may trigger inaccuracies during thermal reading. However, 61  
5 since different luminescence characteristics depend 62  
6 differently on the microenvironment, the use of more than 63  
7 one emission parameter, or more than one LNM probe, 64  
8 appears as a promising strategy to identify unexpected 65  
9 effects. Alternatively, the design of thermal probes protected 66  
10 from the microenvironment, e.g. by encapsulation, may be 67  
11 an optimal strategy in some scenarios. However, for 68  
12 hyperthermia treatments in particular, the stability of such 69  
13 encapsulations at different temperatures must be tested, in 70  
14 addition to the effect of the biocorona.  
15 In non-clinical settings, one method to reduce the interfering 71  
16 properties of complex biological tissue is to simplify the 72  
17 biological model, i.e., moving from *in vivo* studies involving 73  
18 multi-organ systems, to *in vitro* organ-on-a-chip devices or 74  
19 spheroids. Our understanding of the natural or controlled 75  
20 formation (i.e., additive manufacturing) of *in vitro* tissue and 76  
21 organ models has vastly improved, allowing reproducible 77  
22 and realistic pre-*in vivo* studies. Thanks to their controllable 78  
23 physiological and biochemical traits, 3D biological models, 79  
24 combined with LNMs, will play an important role in 80  
25 furthering our understanding of hyperthermia. That said, 81  
26 when writing this tutorial review, we discovered a 82  
27 considerable lack of data relating to the use of LNMs for 83  
28 temperature sensing in 3D cell models such as spheroids, 84  
29 organoids, and organ-on-a-chip devices. We propose that 85  
30 these models will be of special importance to better 86  
31 understand the temperature sensing abilities of LNMs, prior 87  
32 *in vivo* investigation. Issues such as the uncontrolled LNM 88  
33 aggregation as a result of *i.e.* biocorona formation, or a 89  
34 reduced luminescence readout due to light attenuation across 90  
35 tissue interfaces, can be studied with *in vitro* 3D models, 91  
36 thereby adhering to the 3R's of Replacement, Reduction, 92  
37 and Refinement. 93

## 39 Conflicts of interest

40 There are no conflicts to declare. 94

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