

A comparative study of dark matter flow & hydrodynamic turbulence and its applications

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Dark matter, if exists, accounts for five times as much as ordinary baryonic matter. Therefore, dark matter flow might possess the widest presence in our universe. The other form of flow, hydrodynamic turbulence in air and water, is without doubt the most familiar flow in our daily life. During the pandemic, we have found time to think about and put together a systematic comparison for the connections and differences between two types of flow, both of which are typical non-equilibrium systems.

The goal of this presentation is to leverage this comparison for a better understanding of the nature of dark matter and its flow behavior on all scales. Science should be open. All comments are welcome.

Thank you!

(all slides available at zenodo.org by searching "dark matter flow")

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Data repository and relevant publications Structural (halo-based) approach: Statistics (correlation-based) approach:

0. Data <https://dx.doi.org/10.5281/zenodo.6569898>

ark matter flow for velocity, density,

rXiv.2202.00910

ark matter flow and high order lations for velocity and density correlations <https://doi.org/10.48550/arXiv.2202.02991>

riation of density and velocity If flow and two-thirds law for pairwise velocity<https://doi.org/10.48550/arXiv.2202.06515>

and properties from two-thirds law **and right** matter flow rXiv.2202.07240

leration and deep-MOND from ad energy cascade in dark matter 50/arXiv.2203.05606

is relation from mass and energy

rXiv.2203.06899

d density slope for dark matter and energy cascade rXiv.2209.033

0. Data<https://dx.doi.org/10.5281/zenodo.6541230>

- 1. Inverse mass cascade in dark matter flow and effects on halo mass functions <https://doi.org/10.48550/arXiv.2109.09985>
- 2. Inverse mass cascade in dark matter flow and effects on halo deformation, energy, size, and density profiles<https://doi.org/10.48550/arXiv.2109.12244>
- 3. Inverse energy cascade in self-gravitating collisionless dark matter flow and effects of halo shape <https://doi.org/10.48550/arXiv.2110.13885>
- 4. The mean flow, velocity dispersion, energy transfer and evolution of rotating and growing dark matter halos <https://doi.org/10.48550/arXiv.2201.12665>
- 5. Two-body collapse model for gravitational collapse of dark matter and generalized stable clustering hypothesis for pairwise velocity <https://doi.org/10.48550/arXiv.2110.05784>
- 6. Evolution of energy, momentum, and spin parameter in dark matter flow and integral constants of motion<https://doi.org/10.48550/arXiv.2202.04054>
- The maximum entropy distributions of velocity, speed, and energy from statistical mechanics of dark matter flow <https://doi.org/10.48550/arXiv.2110.03126>
- 8. Halo mass functions from maximum entropy distributions in collisionless dark matter flow<https://doi.org/10.48550/arXiv.2110.09676>

- [Some fundamentals of dark matter research](#page--1-0)
- [Basic concepts in hydrodynamic turbulence](#page--1-0)
- [Dark matter flow \(SG-CFD\) vs. hydrodynamic turbulence](#page--1-0)
- [Theory of dark matter flow](#page--1-0)
	- [Structural \(halo-based\) approach](#page--1-0)
	- [Statistical \(correlation-based\) approach](#page--1-0)
- [Applications of dark matter flow](#page--1-0)
	- **[Predicting dark matter particle properties](#page--1-0)**
	- [Understanding the origin of MOND](#page--1-0)
	- [The baryonic-halo mass ratio and total baron faction](#page--1-0)
	- [Universal scaling laws and halo density slope](#page-5-0)

Applications of dark matter flow

Universal scaling laws and scales for dark matter halos from rotation curves and energy cascade

[Xu, Zhijie arXiv:2209.03313 \[astro-ph.GA\]](https://doi.org/10.48550/arXiv.2209.03313) <https://doi.org/10.48550/arXiv.2209.03313>

- Is there an asymptotic density slope for halos?
- **Why exists a nearly universal density profile?**
- Why different inner slopes γ exist in simulations?
- Core-cusp solutions
	- Within CDM framework
		- **Baryonic feedback processes**
	- Beyond CDM
		- **Self-interacting dark matter**

Pacific Northwest Introduction

- Standard CDM Models' successes for large scale structure formation and evolution.
- Small scale challenges suggest missing pieces:
	- Galaxy scale (<1Mpc)
	- Core-cusp problem
	- Missing satellite
	- Too-big-to-fail
	- Baryonic Tully-Fisher and MOND
- Core-cusp problem
	- **Dark matter halo density** $\rho(r) \propto r^{\gamma}$
	- Cored density (γ =0) from observational data
	- Cuspy density (γ ~-1) from N-body simulations
		- No consensus
		- γ =-1.0 in NFW profile
		- γ =-1.2 (Diemand & Moore 2011)
		- $\gamma = -1.3$ (Governato et al. 2010)
		- $y = -1.3$ (McKeown et al. 2022)
- What are the critical length or density scales for dark matter if exist?
- What is the effect of self-interaction on these scales?
- What are the fundamental properties (mass, crosssection etc.) of dark matter?

No matter collisionless or self-interacting

Energy cascade in hydrodynamic turbulence

 $(-\varepsilon_{_u})$

- There exist an inertial range with a scaleindependent rate of energy cascade (ε does not depend on eddy size *l*) for eddy size η< *l <L.* η is a dissipative scale determined by viscosity ν and ε.
- In inertial range, inertial force is dominant over viscous force. A general scaling for velocity structure functions ${\sf S}_{\sf m}$ (r) for pairwise velocity $\Delta {\sf u}_{\sf L}$ can be identified:

 $m \left(\begin{array}{ccc} & f & f \\ f & & \end{array} \right)$

 $S_m(r) \propto (\varepsilon_u)^{m/3} r^{m/3}$ **b** $S_2 \propto (-\varepsilon_u)^{2/3} r^{2/3}$

 $m=2$

 $S_m(r, a) = \langle (\Delta u_L)^m \rangle = \langle (u_L - u_L) \rangle$

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Two-thirds law

 $\mathbf x$

Big whirls have little whirls, That feed on their velocity; And little whirls have lesser whirls, And so on to viscosity.

Mass/Energy cascade in dark matter flow (SG-CFD)

Little halos have big halos, That feed on their mass; And big halos have greater halos, And so on to growth

- Long-range gravity requires a broad spectrum of halos to be formed to maximize system entropy.
- A continuous cascade of [mass/energy](#page--1-0) from smaller to larger mass scales with a scale-independent rate of mass transfer ε_m and ε_u in a certain range of mass scales (propagation range).
- The mass/energy cascade is an intermediate statistically steady state for non-equilibrium systems to continuously maximize system entropy.
- The maximum entropy distribution of dark matter flow. $\ddot{\mathbf{\tau}}$

Collisionless nature and long-range interaction.

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Mass/Energy cascade in dark matter flow (SG-CFD) Northwest

- Collisionless, no dissipation range in SG-CFD.
- The smallest length scale of inertial range is not limited by viscosity.
- This enable us to extend [the scale-independent](#page--1-0) $\epsilon_{\rm u}$ down to the smallest scale, where quantum effects become important
- Dark matter flow exhibits [scale-dependent flow behaviors](#page--1-0) for peculiar velocity, i.e. a constant divergence flow on small scales and an irrotational flow on large scales.
- The constant divergence flow shares the same even order kinematic relations with those of [incompressible \(divergence free\) flow. This hints to similar scaling laws holds for dark matte](#page--1-0)r.

Constant (time and scale independent) rate of Pacific Northwest energy cascade

Power-law time evolution for energy in terms of rate of energy cascade $ε_{0}$:

7 5 $P_y = \frac{7}{5} \varepsilon_u t$

 $K_p = -\varepsilon_{\rm u}t$

Power-law for **Peculiar** kinetic energy

Also see detail analysis for inverse kinetic energy cascade.

Power-law for potential energy

$$
\varepsilon_{u} = -\frac{K_{p}}{t} = -\frac{3}{2}\frac{u_{0}^{2}}{t_{0}} \approx -4.6 \times 10^{-7} \frac{m^{2}}{s^{3}}
$$

The two-thirds law on small scales from N-body Pacific Northwest Simulations

$$
S_2^{lp}(r) - 2u^2 = S_{2r}^{lp} = \beta_2^* (r/r_s)^{2/3} \propto r^{2/3}
$$

$$
S_{2n}^{lp}(r) - 2^{n} u^{2n} K_{2n}(\Delta u_L, 0) = \beta_{2n}^{*} (r/r_s)^{2/3} \propto r^{2/3}
$$

Second order [\(two-thirds law](#page--1-0)):

$$
S_{2n+1}^{lp}(r) = (2n+1) S_1^{lp}(r) S_{2n}^{lp}(r) \propto r^1
$$

Even order ([two-thirds law](#page--1-0)):

Odd order moment (generalized stable clustering hypothesis):

$$
v_r^2 = S_{2r}^{lp} (r) / (2^{2/3} \beta_2^* a^{3/2})
$$

\n
$$
(-\varepsilon_u) = \frac{2v_r^2}{r} v_r = \frac{2v_r^2}{r/v_r} = \frac{2v_r^3}{r}
$$

\n
$$
\frac{1}{r}
$$

\n1
\n1
\n1
\nTurnaround
\ntime

Introduce a velocity scale:

Pacific Constant rate of cascade from galaxy rotation curves

 $v_s^2 = Gm_r (r_s)/r_s$ $(-\varepsilon_u)$ 3 *s u s s v r* ε γ $-\varepsilon_{u}^{-}$) =

 $\gamma_s \approx 6.83$

 $10⁴$

- rate of cascade;
- Dispersion from spatial intermittence of energy cascade;
- Halos in different local environment may have different $ε_u$
- Dwarf galaxies tends to have smaller ε_{u} due to tidal stripping.

SPARC NFW fitting SPARC pISO fitting DMS NFW fitting DMS pISO fitting SOFUE NFW fitting \blacksquare -E $_u = 4.6 \times 10^{-7} m^2/s^3$.

 $10³$

 $10²$

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Relevant scales for self-interacting dark matter Northwest

Relevant scales for collisionless dark matter Northwest

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Uncertainty principles & energy cascade

For fully collisionless dark matter:

- 1) A unique "symmetry" between **x** and **v** in phase space:
- At given **x** , particles can have multiple **v** (multi-stream)
- With given **v,** particles can be at different **x**
- NOT possible for non-relativistic baryons
- 2) Due to the long-rang gravitational interaction,
- Fluctuations (uncertainty) in **x**
- **Fluctuations (uncertainty) in v**
- Fluctuations (uncertainty) in **a**
- 3) Two pairs of conjugate variables:
- Position **x** and momentum **p**
- Momentum **p** and acceleration **a**

Position (**x**), Velocity (**v** = d**x**/dt), Acceleration (**a** = d**v**/dt)

Uncertainty principles: $\sigma_x \sigma_p \geq \hbar/2$

- Wave function for position: $\psi(x)$ Wave function for momentum: $\varphi(p)$ Wave function for acceleration: $\mu(a)$
	- $2.44 \times 10^{-22} \text{ kg} \cdot \text{m}^2/\text{s}^3$

 $\psi(x) = \frac{1}{\sqrt{2\pi}} \int \varphi(p) e^{ipx/\hbar} dp$ ∞ −∞ $=\frac{1}{\sqrt{2\pi\hbar}}\int \varphi(p)e^{ipx/\hbar}$ \hbar ipa/μ_X *X* $\varphi(p) = \frac{1}{\sqrt{2\pi}} \int \mu(a) e^{ipa/\mu_x} da$ ∞ −∞ $\sigma_{p} \sigma_{q} \geq \mu_{X}/2$ $\varepsilon_u = \mu_X/m_X = a_X v_X$

Postulated uncertainty principle for *a* and *p* leads to the constant rate of energy cascade:

$$
\mu_X = -m_X \varepsilon_u = 7
$$

Northwest Universal halo density from energy cascade

 r *c*(*t*)*r*

Reduced spatial/ temporal coordinate: $x(r,t) = \frac{r}{r_s(t)}$

$$
m_r(r,t) = m_h(t) \frac{F(x)}{F(c)}
$$

 $=\frac{1}{\sqrt{2}}$

 $NFW:$ $F(x) = \ln(1+x) - x/(1+x) \propto x^2$ Einasto: $F(x) = \Gamma(3/\alpha) - \Gamma(3/\alpha , 2x^{\alpha}/\alpha) \propto x^3$ Isothermal: $F(x) = x/c$

Function *F(x)* for enclosed mass at given *r*:

 $h(\lambda) = u_r$

 $\gamma = \frac{C \ln \rho}{2I}$

 $u_{h} (x) = u_{r} - 1 x$

Halo density:

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 $(r, t) = \frac{1}{1} \frac{\partial m_r(r, t)}{\partial} = \frac{m_h(t)}{1} \frac{c^3 F'(x)}{r^3}$ (c) 3 \boldsymbol{E}^\prime 2 λ_n $4\pi r^3$ x^2 1 $\partial m_r(r,$ $f(t) = \frac{1}{4\pi r^2} \frac{cm_r (r,t)}{\partial r} = \frac{m_h}{4\pi r^2}$ *h h* $m_r(r,t)$ *m*_h (t) $c^3F(x)$ $\rho_h(r,t) = \frac{1}{4\pi r^2} \frac{1}{\epsilon_0} \frac{r}{r} \frac{r}{r} \frac{r}{r} = \frac{m_h(r)}{4\pi r_h^3} \frac{1}{x^2 F(r)}$ ∂ $=\frac{1}{1}$ $\frac{cm_r(r, v)}{2}$ = ∂

 $x(r,t$

 $f(t) = \frac{r}{r_s(t)} = \frac{\epsilon(t)}{r_h(t)}$ $r_{s}(t)$ $r_{h}(t)$ $F(x)$

 (t)

$$
\frac{\partial \rho_h(r,t)}{\partial t} + \frac{1}{r^2} \frac{\partial \left[r^2 \rho_h(r,t) u_r(r,t)\right]}{\partial r} = 0
$$

 $\frac{\ln \rho_h}{\ln t} = \frac{\partial x}{\partial \ln t} \frac{\partial \ln t}{\partial \ln x} - 2$

 ∂u_{ν} $\partial \ln m_{r} (r_{s}, t)$ (∂

 $f(x) = u_r \frac{t}{t} = \frac{\partial \ln r_s}{\partial x} + \frac{\partial \ln F(c)}{\partial x} - \frac{\partial \ln m_h}{\partial x} + \frac{F(x)}{F(x)}$

 $h \perp$ $\frac{C \ln m_r (r_s, t)}{s}$

 $x \sim \frac{\partial \ln t}{\partial x}$

 u_h $\partial \ln m_r(r_s,t)$ $\partial \ln r_s$

Density slope:

Radial continuity equation:

s

 \int $\partial \ln r_{s}$ \int $\partial \ln F(c)$ $\partial \ln m_{h}$ $\sqrt{F(x)}$ $= u_r \frac{v}{r_s} = \left[x \frac{\partial \ln r_s}{\partial \ln t} + \left(\frac{\partial \ln r}{\partial \ln t} - \frac{\partial \ln m_h}{\partial \ln t} \right) \frac{r(x)}{F'(x)} \right]$ $m_{_r}\left(r_{_S},t\right) \propto\varepsilon _{u}^{2/3}G^{-1}r_{_S}^{^{5/3}}$ cascade

Radial flow equation:

h

 $\ln x$ $\partial \ln x$

 $x \longrightarrow \partial \ln r_s u$

 $= \frac{\partial \ln \rho_h}{\partial \ln x} = \frac{\partial x}{\partial \ln r_s} = \frac{\partial \ln t}{\partial \ln t} = \frac{u_h}{x}$

 $\partial \ln \rho_h \left(\frac{\partial u_h}{\partial x} \right) + \frac{\partial \ln m_r(v_s, v)}{\partial \ln t} - \left(\frac{\partial u_h}{\partial x} \right)$

ln *t x*

 s_{\perp} $\frac{u_h}{h}$

 $\ln r_{s}$ $\int \partial \ln F(c)$ $\partial \ln$

 $t \quad \Big| \quad \partial \ln r_{s} \quad \Big(\partial \ln F(c) \quad \partial \ln m_{h} \Big/ F(x)$

 $\ln t \quad \vert \quad \partial \ln t \quad \partial \ln$

 r_{s} | $\partial \ln t$ | $\partial \ln t$ $\partial \ln t$ $\partial \ln t$ $\int F(x)$

 $\ln m_r(r_s,t)$ $\big\hat{\partial}\ln$

 s_{\perp} $\frac{\text{cm}_{\perp} \times \text{m}_{\perp}}{2}$ $\frac{\text{cm}_{\parallel} m_{h}}{2}$

Density slope is dependent on

Gradient of radial flow (spatial variation)

$$
\frac{(r,t)}{r} = 0
$$

 $\dot{ }$ (x)

$$
\sqrt{\frac{1}{r}} = -4/3 \quad \rho_h(r) \propto r^{-4/3}
$$

 Nearly universal density profile for fully virialized halos with $u_h=0$;

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Evolution of halo density profile

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- Asymptotic density slope $γ=4/3$;
- Simulated halos have different slope due to u_h and mass accretion
- Halos in different local environment may have different γ
- A modified Einasto profile is proposed to reflect different slopes γ

- Smalls scale challenges suggest missing pieces in our current understandings of dark matter
- Review [inverse mass and energy cascade](#page--1-0) in dark matter flow with a constant rate $\varepsilon_{\rm u}$
- Energy cascade leads to a 2/3 law for kinetic energy $\propto r^{2/3}$ on scale r, as confirmed by N-body simulation
- Energy cascade leads to a -4/3 law for halo density $\propto r^{-4/3}$ on scale r, as confirmed by rotation curves.
- **The largest halo scale is determined by** u_0^3 **/** ε u
- The smallest halo scale is dependent on the nature of dark matter:
	- Collisionless dark matter: $r_n \propto (Gh \epsilon_u)^{1/3}$, where h is Planck constant
	- Self-interacting dark matter: $r_n \propto G^{-3} \epsilon_u^2$ (σ/m)³, where (σ/m) is cross-section
- Asymptotic density slope -4/3 for virialized halos with vanishing radial flow
- Simulated halos have different slope due to radial flow and mass accretion

Pacific
Northwest Summary and keywords

Backup Slides

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	-
-

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PROFILE: Zhijie (Jay) Xu

- **Computational Scientist**
- Team lead

EXPERIENCE:

INTERESTS:

- **Fluid dynamics**
- **Cosmological flow**
- Cosmological flow

Multiscale Modeling

The Multiscale Modeling

- **Travel**
- **Hiking**
- **Biking**

Grand canyon Yellowstone