

# Practical parametric amplification with tunable working frequency

**Abstract:** Superconducting traveling wave parametric amplifiers (TWPA) based on kinetic or Josephson nonlinear inductance are known to be broadband and low noise. A high gain of TWPA can be achieved for long and ideally matched system. In practice, the unmatched nonlinear impedance of the transmission line with finite size results in resonant character of its response, and thus limits the bandwidth of the amplifier. Nevertheless, many applications do not require wide bandwidth, which increases the total noise. Instead, wide tunability of frequency range would be preferred. We present design of 14 mm long superconducting coplanar waveguide with 2080 Josephson junctions arranged as a metamaterial with proper dispersion. The highest amplification achieved was 15 dB within 35 MHz bandwidth around central frequency  $f_0 = 6.2$  GHz was achieved with the dynamic range of TWPA around 10dB. The  $f_0$  can be widely tuned in frequency ranges 3.2 – 4.3 GHz and 5.1 – 6.4 GHz by pump frequency and pump power. In the first frequency region the average amplification is around 4 dB with 8 dB peak and in the second region the gain oscillates between 6 dB and 15 dB. Similar behaviour was obtained theoretically within the coupled mode theory. The properties of the stepped impedance resonator and the role of the metamaterial are studied by numerical solution of nonlinear telegrapher's wave equations utilizing the finite element method.

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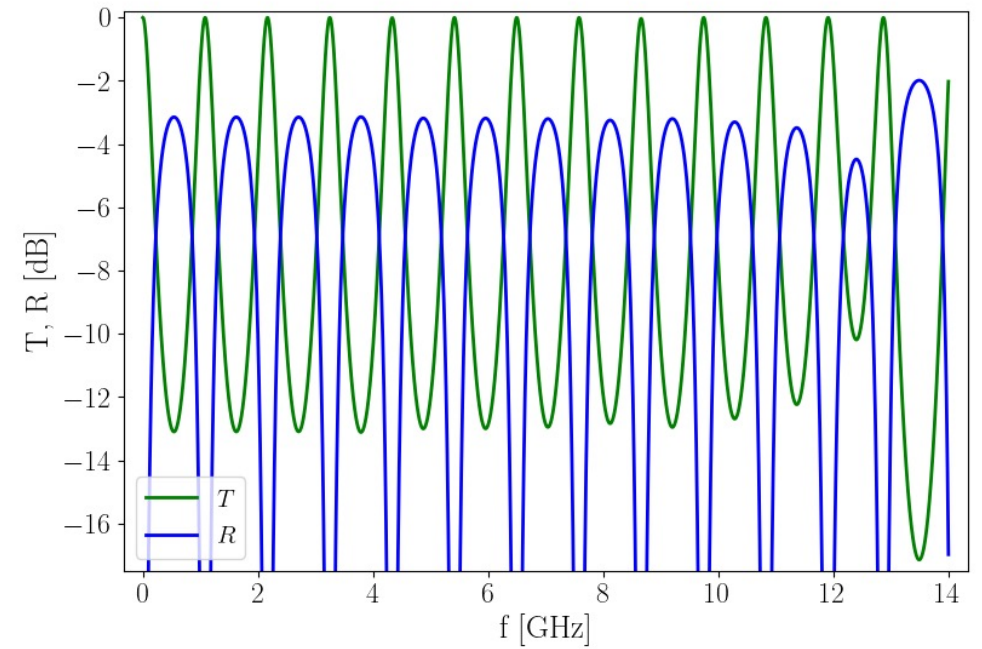
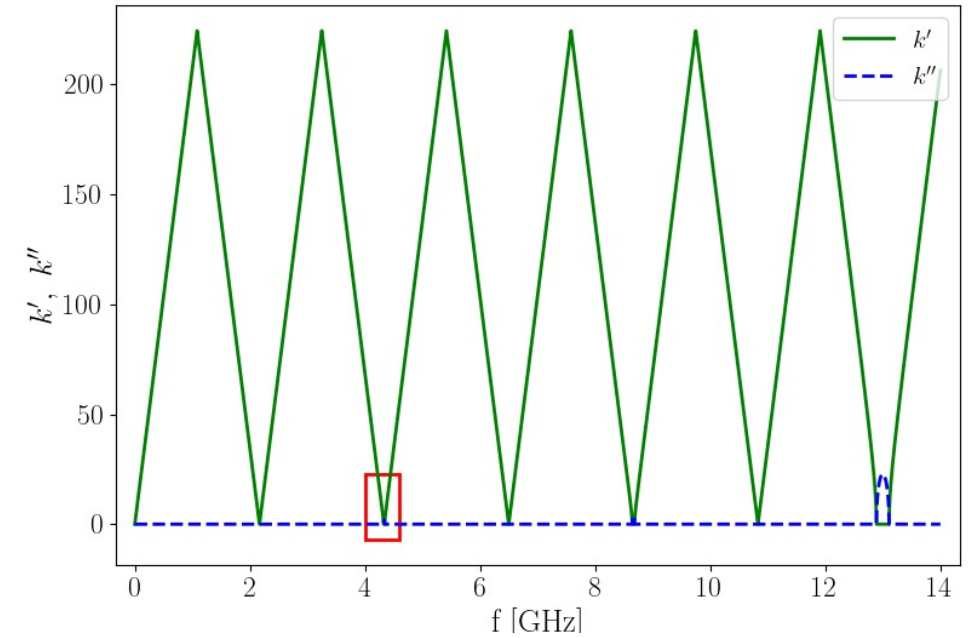
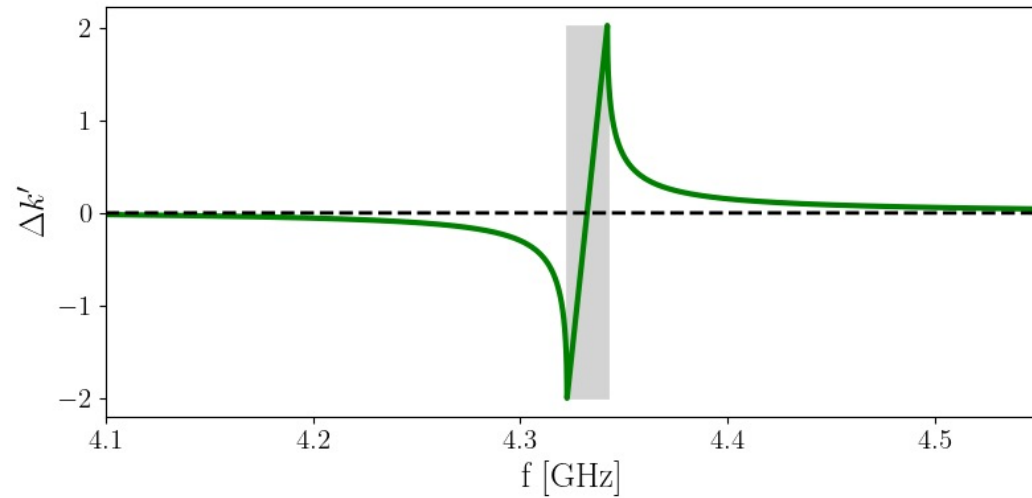
# Metamaterial

$$M = \prod_j \begin{pmatrix} e^{ik_j \Delta l_j} & 0 \\ 0 & e^{-ik_j \Delta l_j} \end{pmatrix} M_{Z_0, j Z_0} \begin{pmatrix} e^{ik_0 l} & 0 \\ 0 & e^{-ik_0 l} \end{pmatrix} M_{Z_0, j Z_0}^{-1}$$

$$k = \frac{1}{2L} \arccos(\text{Tr}M)$$

Phase matching  $\beta = 0 \Rightarrow$ :

$$\Delta k = \frac{2k_p \gamma}{(1 + 2\gamma)}$$



# TWPA vs. Resonator

$$\beta = \Delta k(1 + 2\gamma) - 2k(\omega_p)\gamma$$

$$K = \sqrt{k(\omega_s)k(\omega_i)\gamma^2 - \frac{\beta^2}{4}}$$

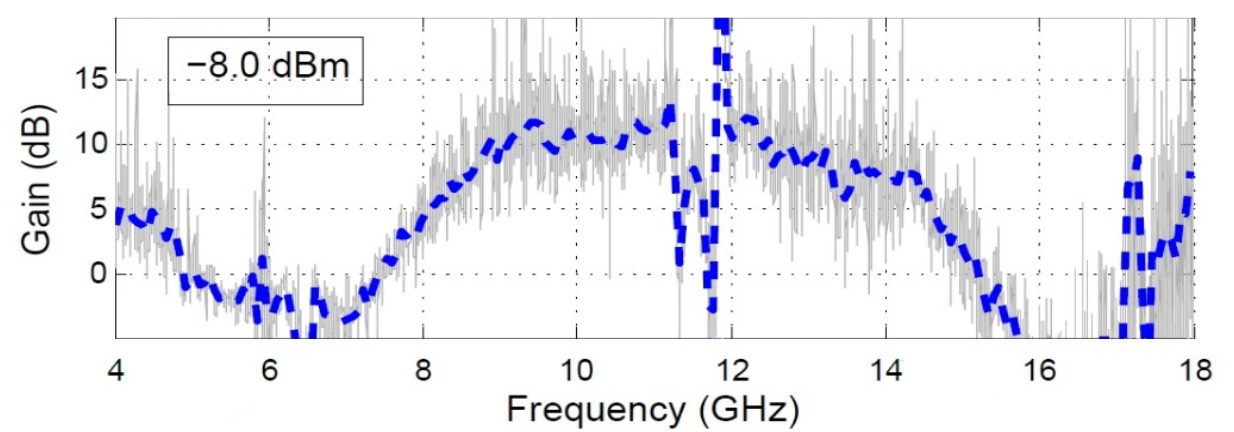
$$G \equiv \frac{P_s(z)}{P_s(0)} = \left( \cosh^2(Kz) + \left( \frac{\beta}{2K} \right)^2 \sinh^2(Kz) \right)$$

$$\gamma = \left( \frac{I_p}{4I_c} \right)^2$$

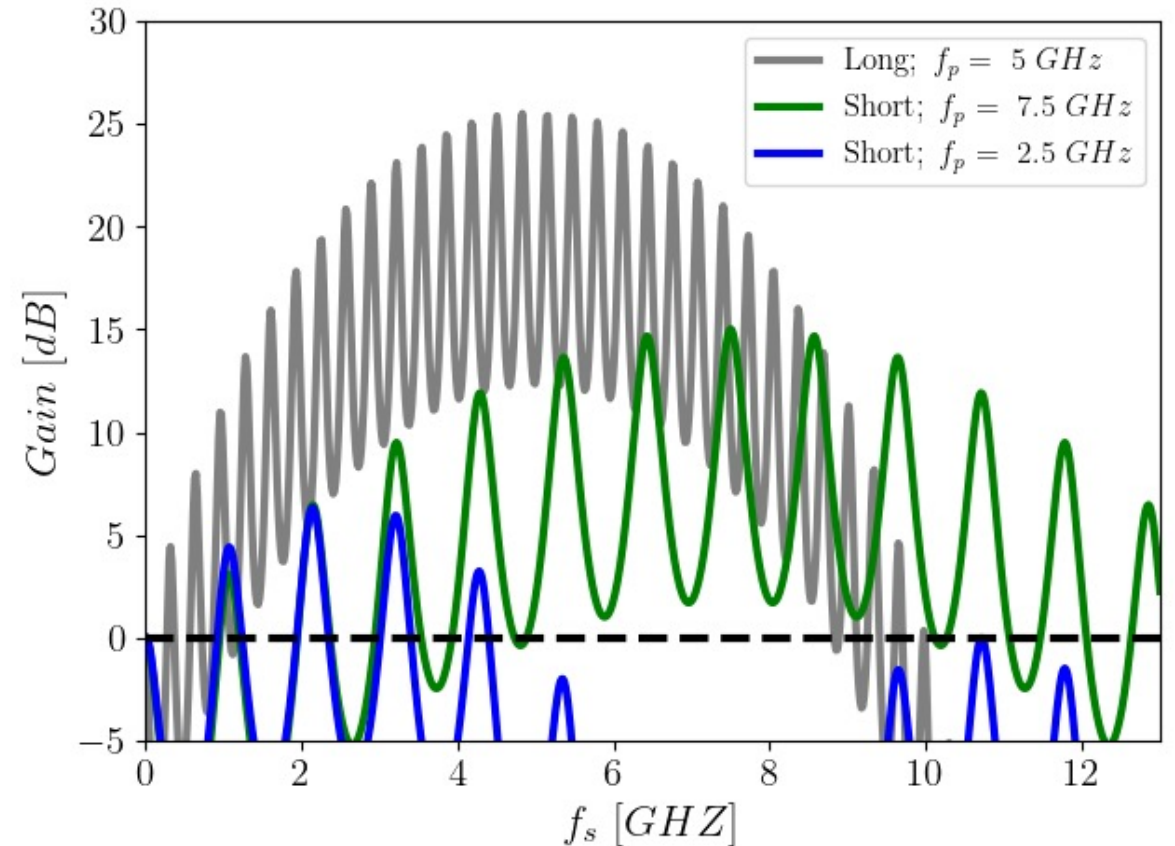
$$I_p^2 = 2QT(f) \frac{P}{Z_0}$$

$$T(f) = \frac{(1 - \Gamma^2(I_p))^2}{(1 - \Gamma^2(I_p)\cos(2\frac{2\pi f}{v_p}l))^2 + (\Gamma^2(I_p)\sin(2\frac{2\pi f}{v_p}l))^2}$$

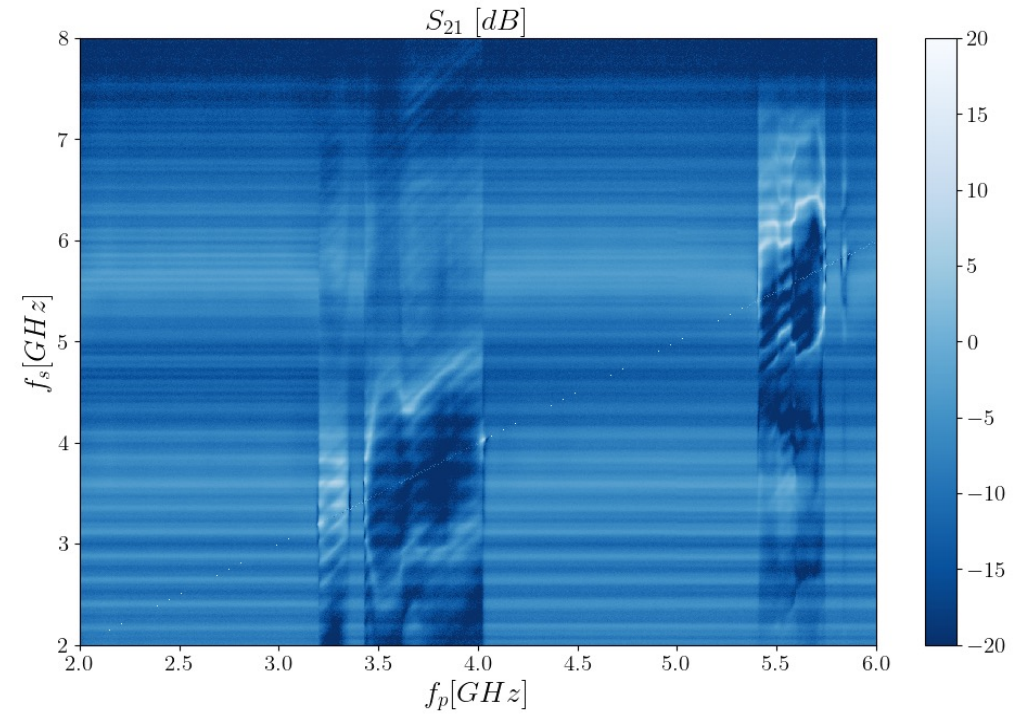
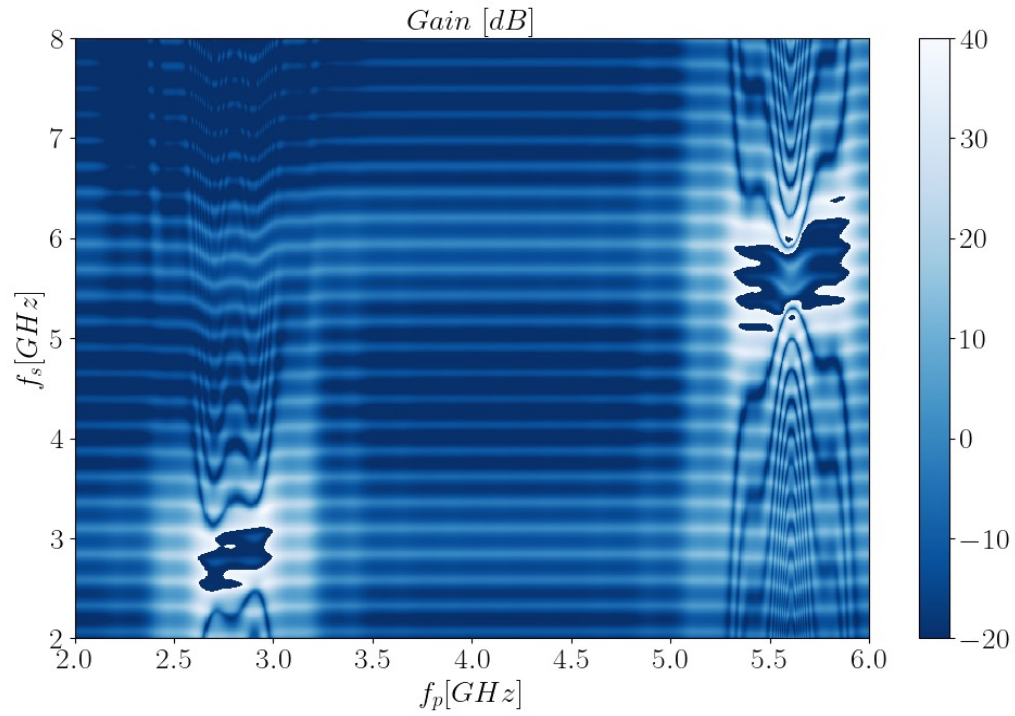
$$Q = \frac{2\pi fL}{v_p(I_p)(1 - \Gamma^2(I_p))}$$



B. H. Eom, P. K. Day, H. G. LeDuc, & J. Zmuidzinas. Nature Physics, **8**(8), 623-627.(2012).



# Experiment



$$\Gamma(I_p) = \frac{Z_0(I_p) - Z_l}{Z_0(I_p) + Z_l}; \quad Z_0(I_p) = \sqrt{\frac{L(1 + I_p^2/I_c^2)}{C}};$$

$$v_p(I_p) = \frac{1}{CL(1 + I_p^2/I_c^2)}$$

Nonlinear response solved iteratively

