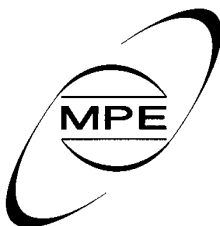


PSD ANALYSIS OF OPTICAL QSO LIGHT CURVES

Torben Simm

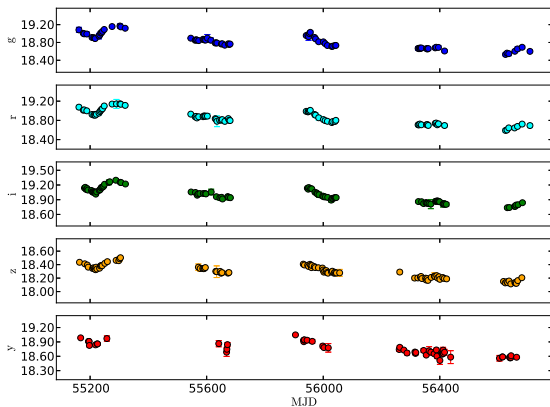
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OPTICAL VARIABILITY OF QSOs

- flux variations are aperiodic and of stochastic nature
- $\sim 10\%$ – 20% fractional amplitude on timescales of months to years
- variability highly correlated across optical bands



QSO multi-band magnitude light curves (Simm et al. 2015)

Power spectral density (PSD) estimators:

- Fourier techniques: periodogram (even sampling), Lomb–Scargle periodogram (uneven sampling) (Priestly 1981, Scargle 1982, Vio et al. 2010)
- normalized excess variance (Nandra et al. 1997, Ponti et al. 2012) and structure functions (MacLeod et al. 2010, Kozłowski 2016)
- continuous-time modeling via stochastic differential equations:
 - put unknown physics (complex processes with large number of degrees of freedom, e.g. magnetic turbulence) into Gaussian white noise process (Vio et al. 2005)
 - Damped Random Walk (DRW) process (Kelly et al. 2009)
 - Continuous-time AutoRegressive Moving Average (CARMA) process (Kelly et al. 2014)

PSD ESTIMATION VIA CARMA PROCESS

For a time series $y(t)$ the variability is assumed to be driven by a continuous-time white noise process $\epsilon(t)$ with zero mean and variance σ^2

DRW = CARMA(1,0) process:

$$\frac{dy(t)}{dt} + \alpha_0 y(t) = \epsilon(t) \quad \leftrightarrow \quad \text{PSD}(\nu) = \sigma^2 \frac{1}{\alpha_0^2 + (2\pi\nu)^2}$$

General CARMA(p,q) process:

$$\frac{d^p y(t)}{dt^p} + \alpha_{p-1} \frac{d^{p-1} y(t)}{dt^{p-1}} + \dots + \alpha_0 y(t) = \beta_q \frac{d^q \epsilon(t)}{dt^q} + \beta_{q-1} \frac{d^{q-1} \epsilon(t)}{dt^{q-1}} + \dots + \epsilon(t)$$

with autoregressive coefficients α_p and moving average coefficients β_q

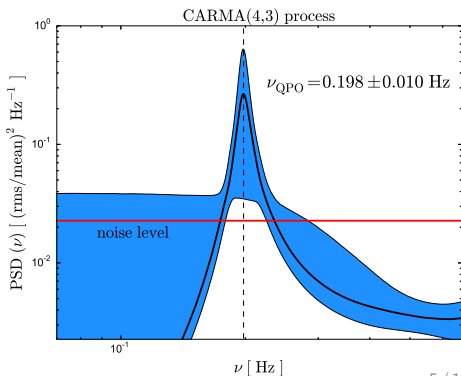
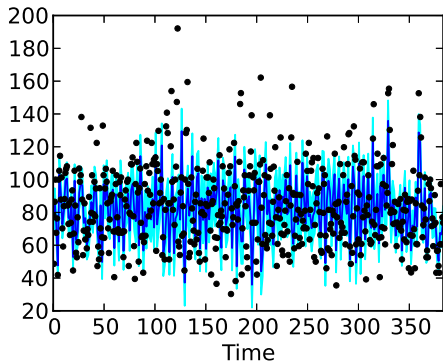
Stationary CARMA(p,q) process with $q < p$ has the PSD:

$$\text{PSD}(\nu) = \sigma^2 \frac{|\sum_{j=0}^q \beta_j (2\pi i\nu)^j|^2}{|\sum_{k=0}^p \alpha_k (2\pi i\nu)^k|^2}$$

CONTINUOUS-TIME LIGHT CURVE MODELING

Example: high frequency PSD of X-ray binary H1743-322
(137 ks count rate light curve from XMM-Newton)

- has well defined QPO at $\nu_{\text{QPO}} \sim 0.20$ Hz (obtained from Fourier spectrum analysis by B. De Marco)
- original light curve has $\Delta t = 0.02$ s \rightarrow use downsampled light curve with $\Delta t = 0.75$ s $\rightarrow \nu_{\text{Nyquist}} = 3.3 \times \nu_{\text{QPO}}$

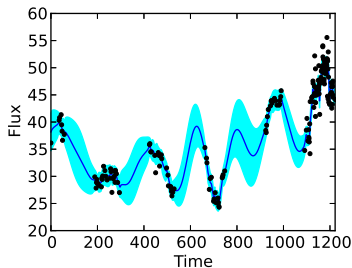


XMM–COSMOS catalogue (Brusa et al. 2010):

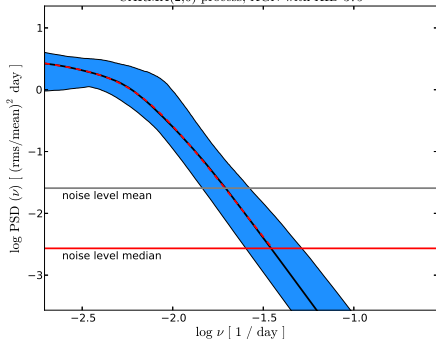
- X–ray selected AGN with $0.3 < z < 2.5$ (100% redshift complete)
- consider point–like and isolated QSOs with known L_{bol} and M_{BH} (Lusso et al. 2012, Trakhtenbrot and Netzer 2012, Rosario et al. 2013)
- 5 band (g, r, i, z, y) multi-epoch light curves from the Pan-STARRS1 (PS1) Medium Deep Field survey
- sample of 187 (g), 184 (r), 165 (i), 135 (z), 76 (y) variable QSOs
- CARMA PSD analysis done in g, r, i, and z bands (Simm et al. 2016)

PSD SHAPE – BROKEN POWER LAW

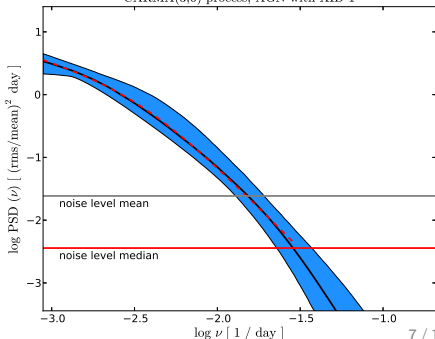
$$\text{PSD}(\nu) = \frac{A(\nu/\nu_{\text{br}})^\alpha}{1 + (\nu/\nu_{\text{br}})^{\beta-\alpha}}$$



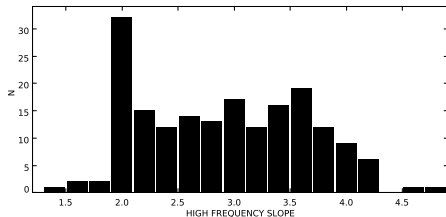
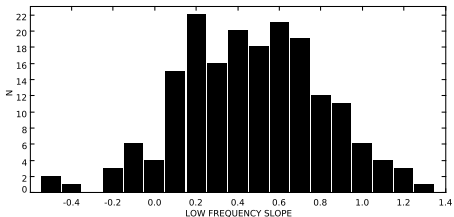
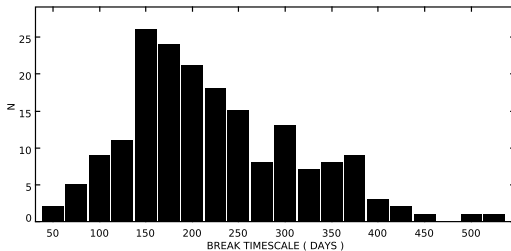
CARMA(2,0) process, AGN with XID 375



CARMA(3,0) process, AGN with XID 1

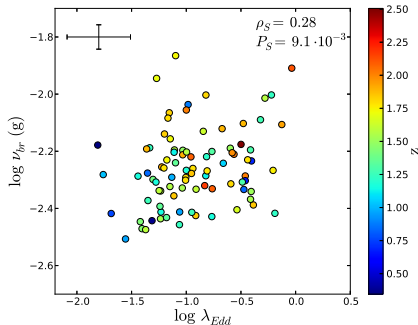
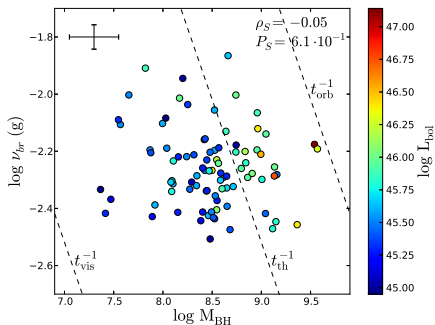


PSD SHAPE – DEVIATIONS FROM DRW MODEL



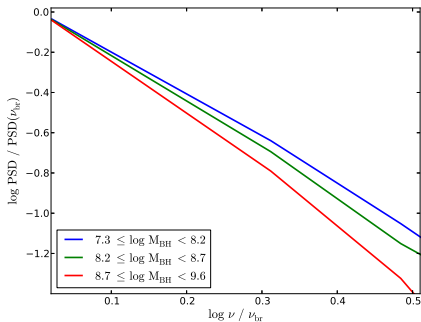
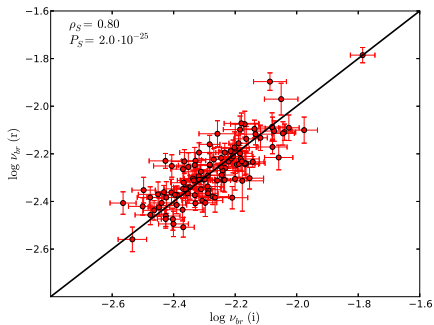
SCALING OF THE OPTICAL BREAK FREQUENCY

X-ray variability studies found: $\nu_{\text{br}} \propto M_{\text{BH}}^{-1}$ (e.g. McHardy et al. 2006)



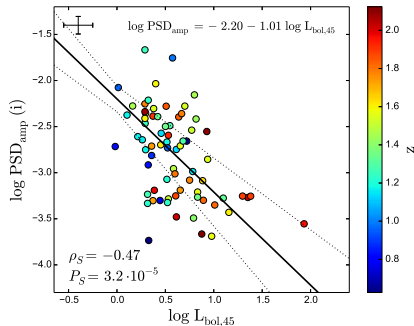
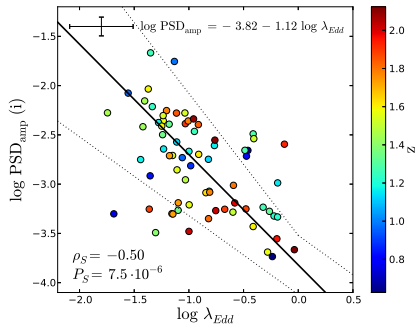
→ the optical break frequency does not scale with M_{BH} , L_{bol} and λ_{Edd} !?

SCALING OF THE HIGH FREQUENCY PSD SLOPE



- the break frequency does not depend on radiation wavelength λ_{rad} !?
- the high frequency PSD slope seems to be anti-correlated with M_{BH}

SCALING OF THE PSD AMPLITUDE

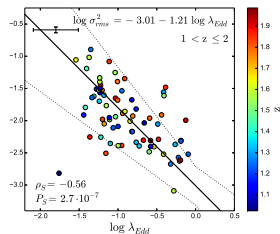
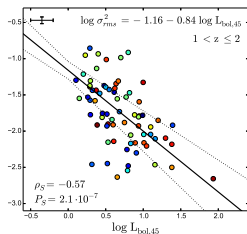
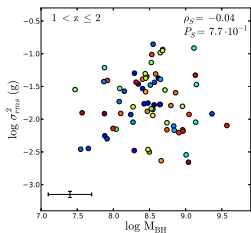


→ $\text{PSD}_{\text{amp}} = A\nu_{\text{br}}$ scales inversely with Eddington ratio λ_{Edd} and luminosity with slope ~ -1

SANITY CHECK: EXCESS VARIANCE ANALYSIS

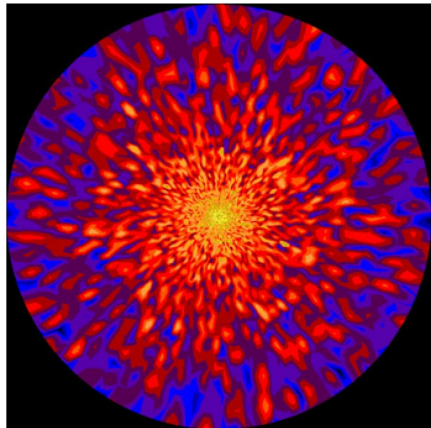
$$\text{PSD} = \begin{cases} A \left(\frac{\nu}{\nu_{br}} \right)^{-1} & \nu \leq \nu_{br} \\ A \left(\frac{\nu}{\nu_{br}} \right)^{-2} & \nu > \nu_{br} \end{cases} \leftrightarrow \sigma_{rms}^2 = \frac{s^2 - \langle \sigma_{err}^2 \rangle}{\langle f \rangle^2} = \int_{\nu_{min}}^{\nu_{max}} \text{PSD} d\nu$$

$$\rightarrow \sigma_{rms}^2 = \begin{cases} \text{PSD}_{amp} (\nu_{min}^{-1} - \nu_{max}^{-1}) \nu_{br} & \nu_{min} > \nu_{br} \\ \text{PSD}_{amp} \left[\ln \left(\frac{\nu_{br}}{\nu_{min}} \right) - \frac{\nu_{br}}{\nu_{max}} + 1 \right] & \nu_{min} < \nu_{br} < \nu_{max} \end{cases}$$



STRONGLY INHOMOGENEOUS ACCRETION DISC?

Temperature map of accretion disc



Toy model of Dexter and Agol 2011
(revised model see Cai et al. 2016):

- disc divided into N independent zones of fluctuating temperature
- variability amplitude $\propto 1/N$
- larger L_{bol} \rightarrow radiation pressure instability enhanced \rightarrow larger number of zones N ?
- M_{BH} -timescale relation smeared out by many zones with same temperature at different radii?

CONCLUSIONS

- for the first time we derived optical PSDs for a large QSO sample in a wide redshift range
 - the optical PSD resembles a broken power law with a low frequency slope of -0.5 and a high frequency slope ranging between -2 and -4
→ significant deviations from simple DRW model
 - the PSD amplitude scales inversely with L_{bol} and λ_{Edd} with the same logarithmic slope of ~ -1
 - the magnitude of the break timescale ($T_{\text{br}} \sim 200$ days) is consistent with the thermal timescale, but seems to be uncorrelated with the AGN parameters M_{BH} , L_{bol} , and λ_{rad} ?
- with more sophisticated accretion disc models and big observational programs such as eROSITA/LSST we may be able to understand the physical origin of these variability correlations