

# Tidal interactions between planets and stars

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with: Aurélie Astoul, Craig Duguid, Chris Jones, Jérémie Vidal  
School of Mathematics



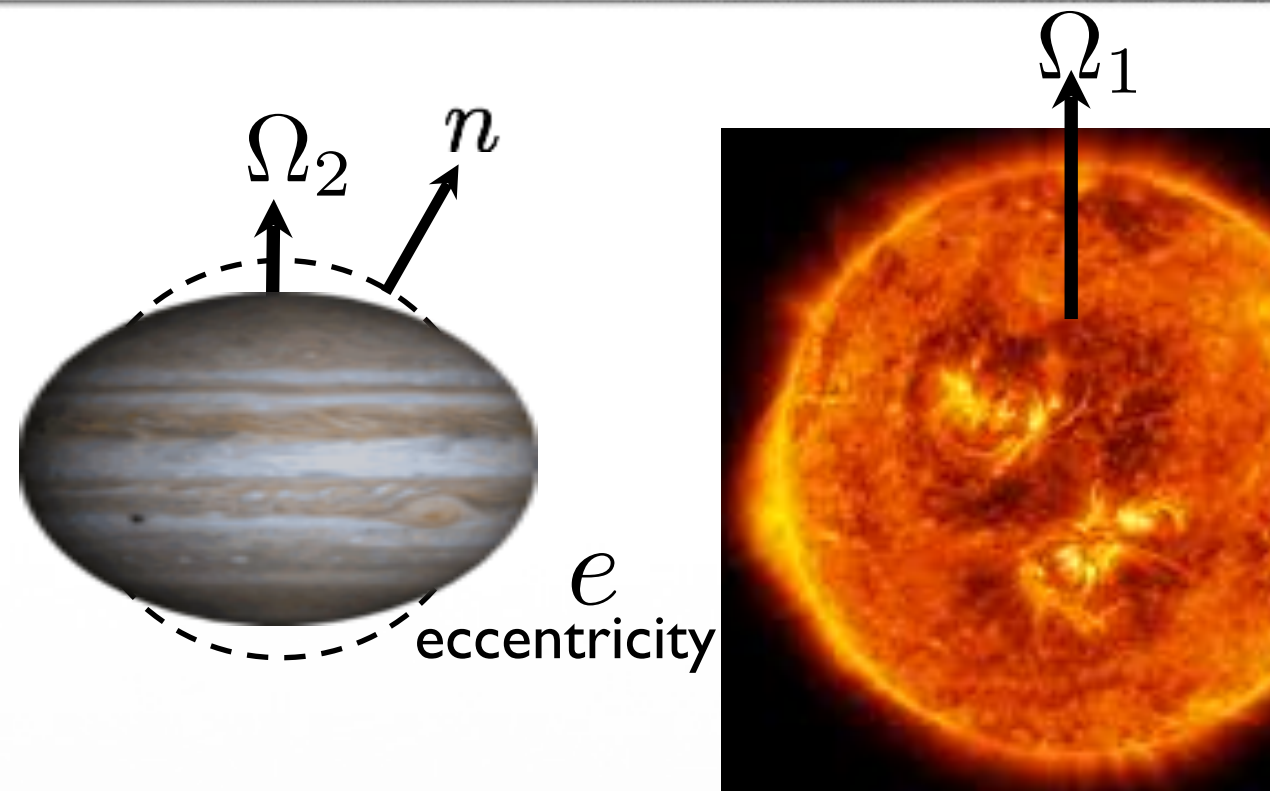
**UNIVERSITY OF LEEDS**

# Tides in stars: orbital decay of hot Jupiters

- Gravitational tidal interactions drive spin and orbital evolution in planetary systems with close-in planets (e.g. hot Jupiters) and binary stars
- First tentative evidence for tidally-driven orbital decay for the hot Jupiter WASP-12 b (Maciejewski et al. 2016; Patra et al. 2017; Yee et al. 2020)

$$\dot{P} = -29 \pm 3 \text{ msyr}^{-1} \Rightarrow P/\dot{P} = 3.2 \text{ Myr}$$

- Evidence against rapid orbital decay for WASP-18 b (Wilkins et al. 2017)
- Tidally-driven orbital decay of hot Jupiters due to dissipation in the star  $\Rightarrow$  *what mechanisms are responsible?*





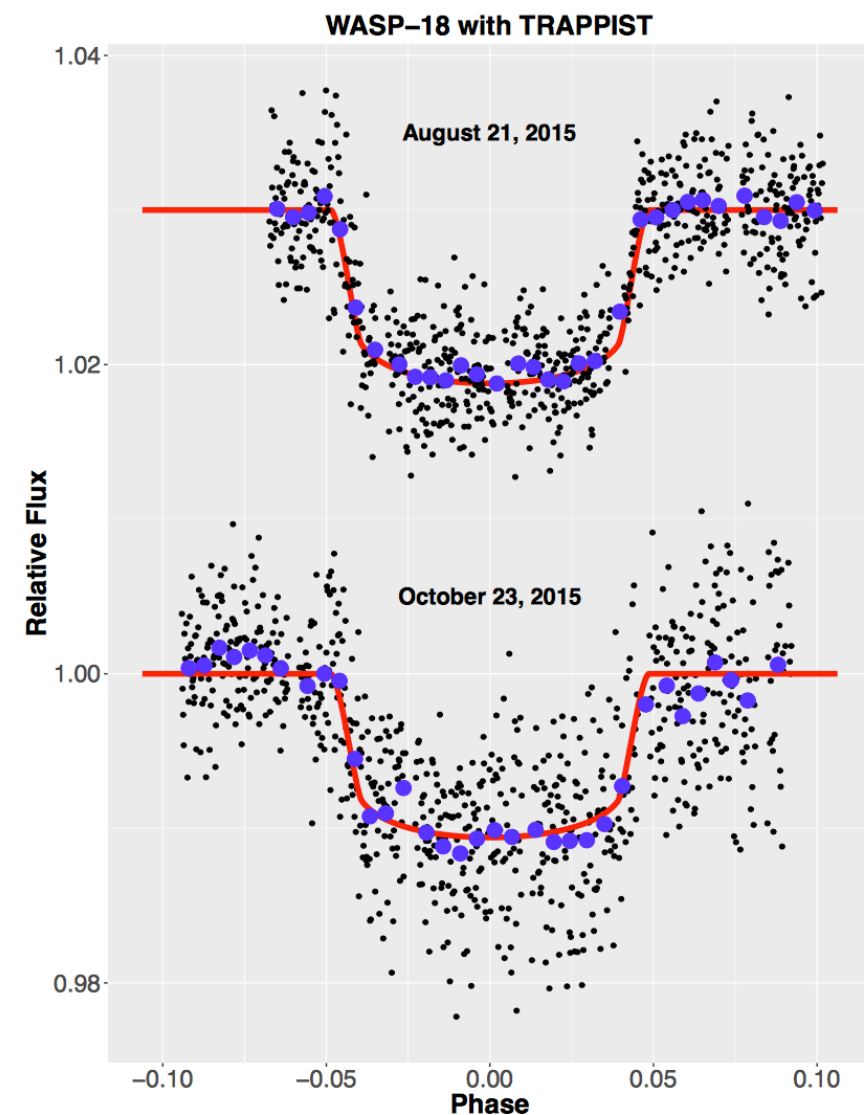
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Wilkins, Delrez, Barker et al. (2017)

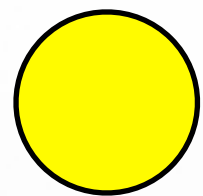


$$T_{shift} = - \left( \frac{27}{8} \right) \left( \frac{M_p}{M_*} \right) \left( \frac{R_*}{a} \right)^5 \left( \frac{2\pi}{P} \right) \left( \frac{1}{Q_*'} \right) T^2$$

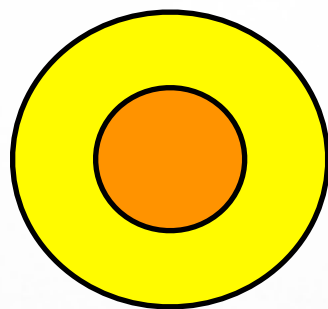
Dimensionless inverse measure of the efficiency of tidal dissipation

# Tides in stars

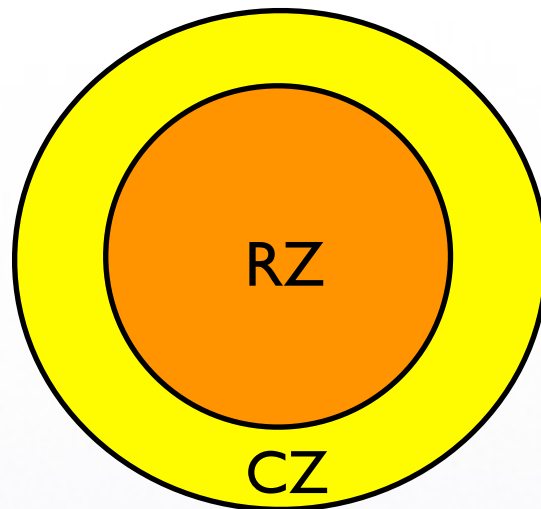
- I will present theoretical calculations of tidal dissipation in stars with masses ranging from 0.1 to 1.6  $M_{\odot}$  throughout their pre-main sequence (PMS) and main sequence (MS) evolution



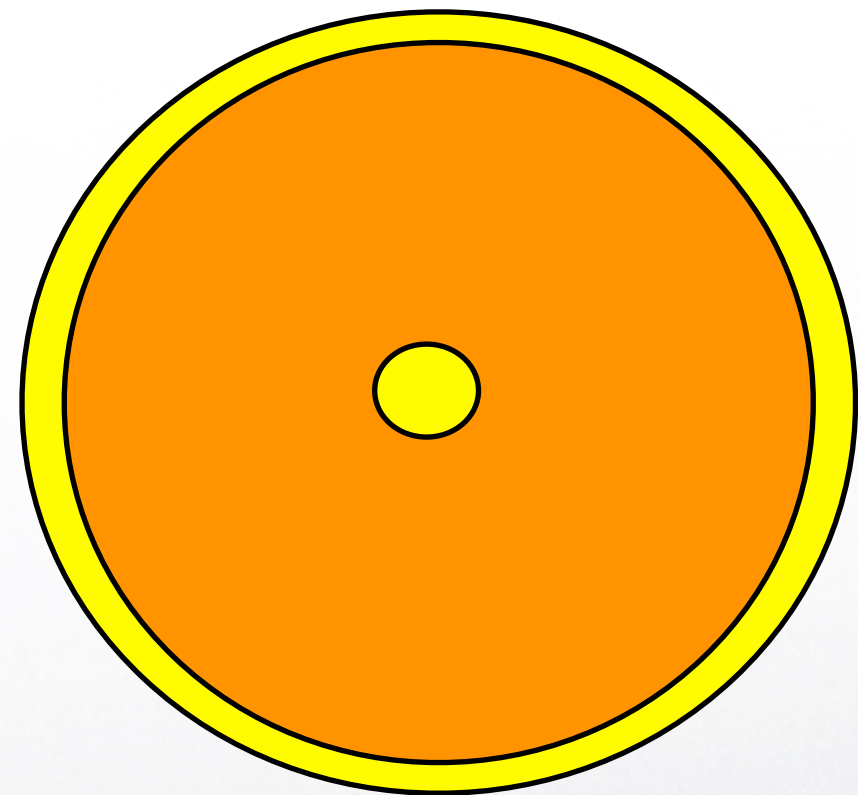
M type  
( $<0.4 M_{\odot}$ )  
Fully  
convective



M and K type  
Convective  
envelope and  
radiative core



G type (solar-type)  
Convective  
envelope and  
radiative core



F type  
 $>1.1 M_{\odot}$  ( $<1.6$ )  
Convective core and  
envelope



# Tides in stars: mechanisms

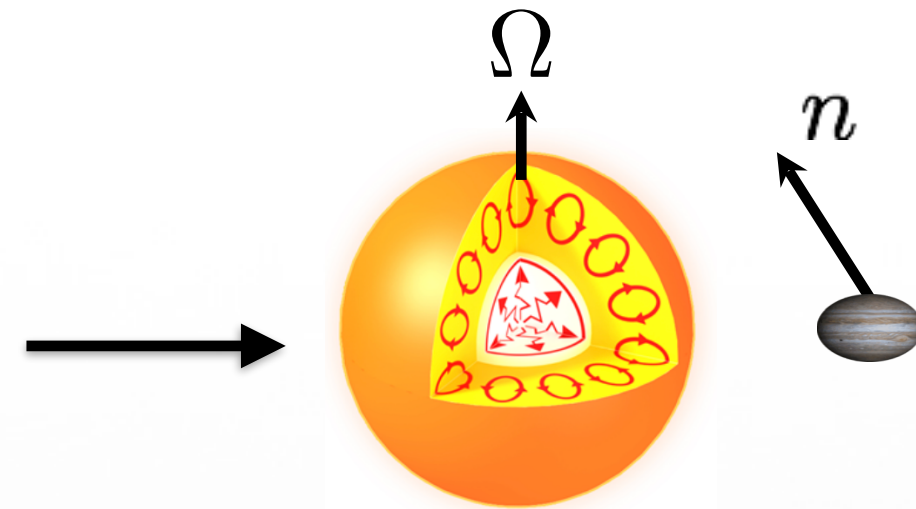
- The tidal response is usually decomposed into two components: an **equilibrium** and **dynamical** tide.

1. **Equilibrium (non-wavelike) tides:** large-scale quasi-hydrostatic deformation of the body, with associated flow. This flow is dissipated by:

- a) interaction with turbulent convection, which can act as an effective **viscosity** (e.g. Zahn 1966/1989; Goldreich & Nicholson 1977; Goodman & Oh 1997; Ogilvie & Lesur 2012; Braviner 2015; Duguid et al. 2020a,b; Vidal & Barker 2020a,b; Terquem 2021)
- b) “nonlinear tidal effects” such as the elliptical instability (e.g. Kerswell 2002; Rieutord 2004; Le Bars et al. 2010; Barker & Lithwick 2013; Barker 2016; Le Reun et al. 2017...)

2. **Dynamical (wave-like) tides:** waves restored by Coriolis or buoyancy (or Lorentz?) forces. Dissipated by viscosity/diffusion/turbulent viscosity (?) /nonlinear interactions.

- a) Internal gravity waves in radiative regions (e.g. Cowling 1941; Zahn 1977; Goodman & Dickson 1998; Ogilvie & Lin 2007; Barker & Ogilvie 2010; Barker 2011; Weinberg et al. 2012; Chernov et al. 2017...)
- b) Inertial waves in convective regions (e.g. Wu 2005; Ogilvie & Lin 2007; Ivanov & Papaloizou 2007; Goodman & Lackner 2009; Rieutord & Valdetaro 2010; Favier, Barker et al. 2014; Gallet et al. 2017; Astoul & Barker, in prep...)





# Tides in stars: mechanisms

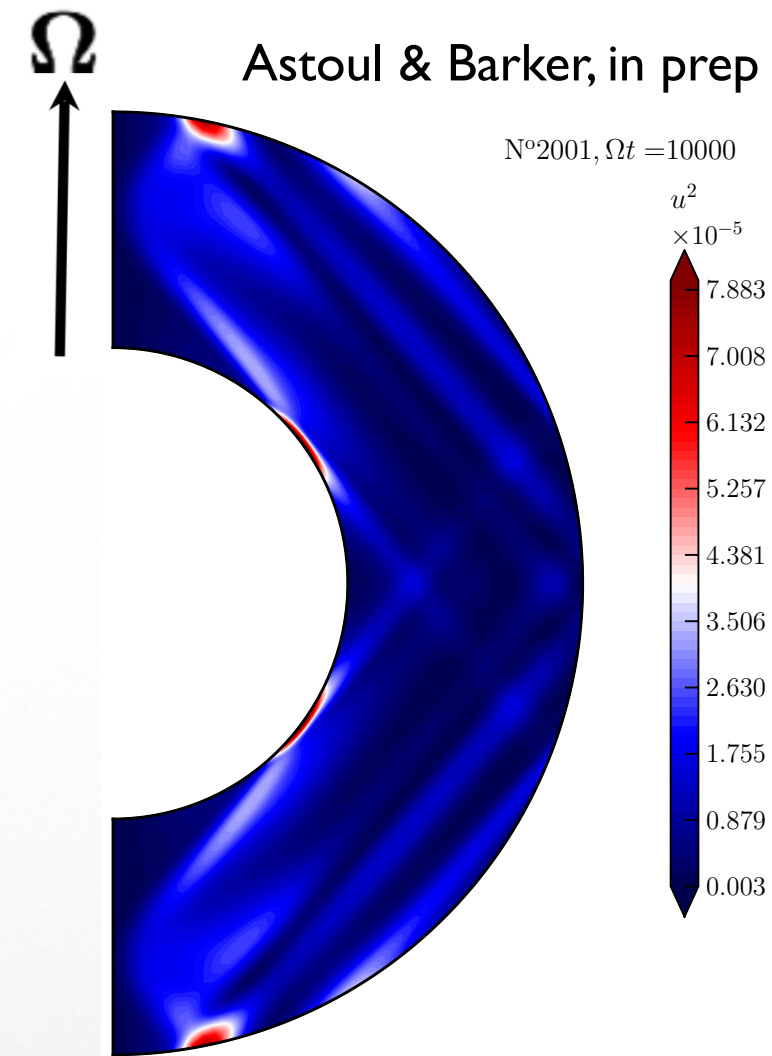
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Probably important in rapidly rotating young stars, but most HJ hosts rotate too slowly for this to operate *at the present day*. Probably dominant for binary star tidal evolution.



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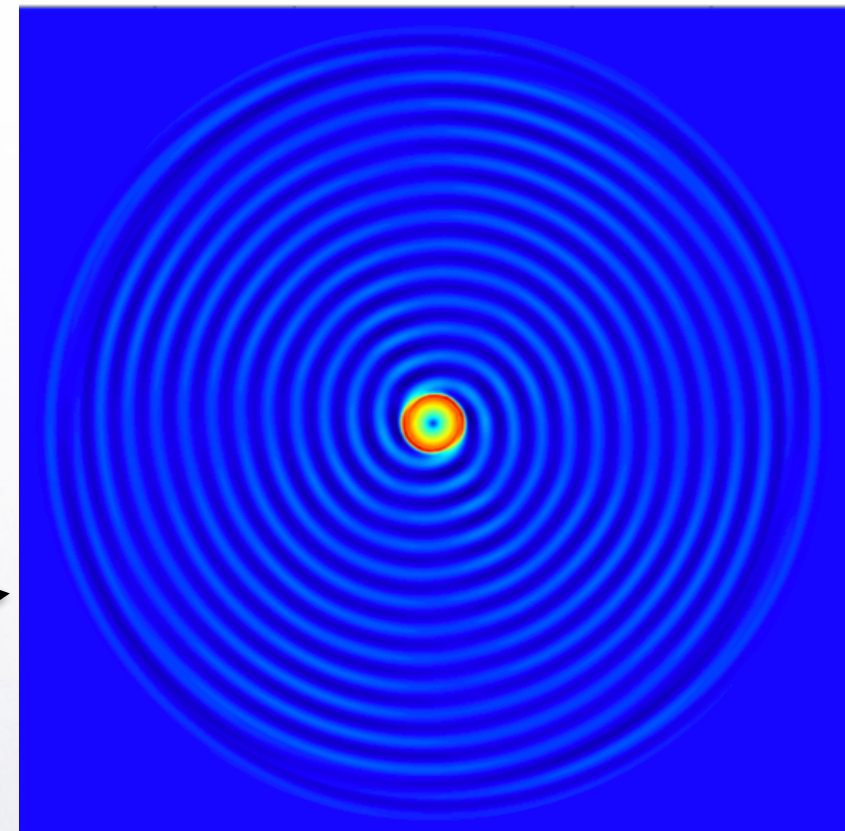
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Barker & Ogilvie 2010

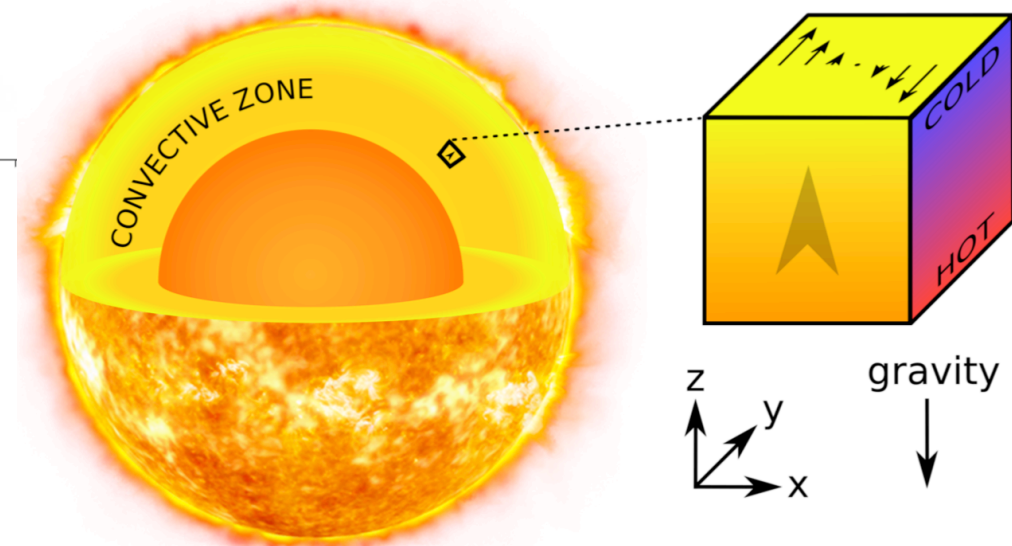
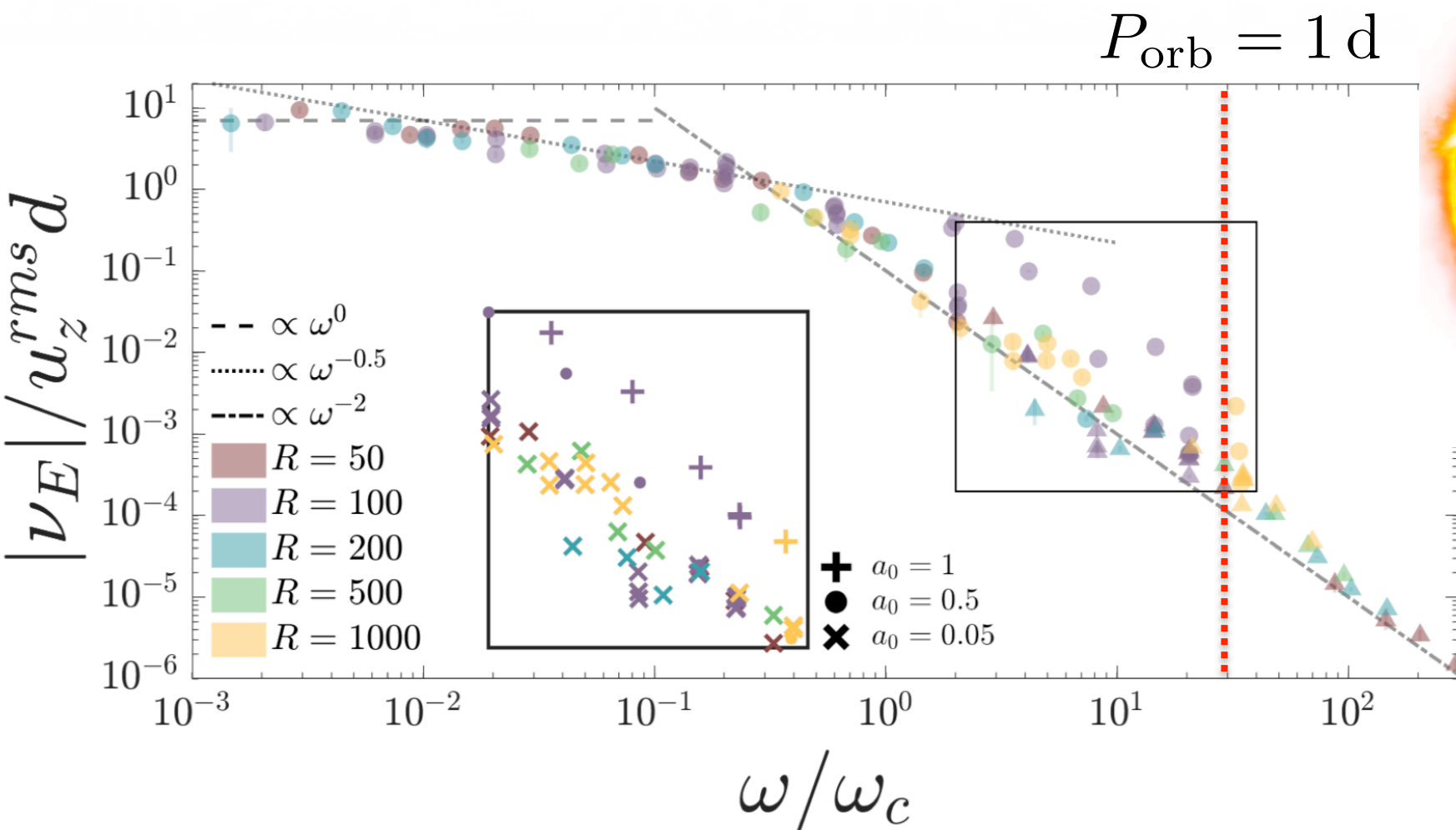


# I. Equilibrium tides in CZs

- Equilibrium tide damping is theoretically uncertain, but recent simulations have shed light on this mechanism.

To summarize, the main weaknesses of the tidal theory, when applied to stars with a convective envelope, reside in our limited knowledge of the dynamics of the convective motions and of their interaction with the tidal flow. (Zahn 1989)

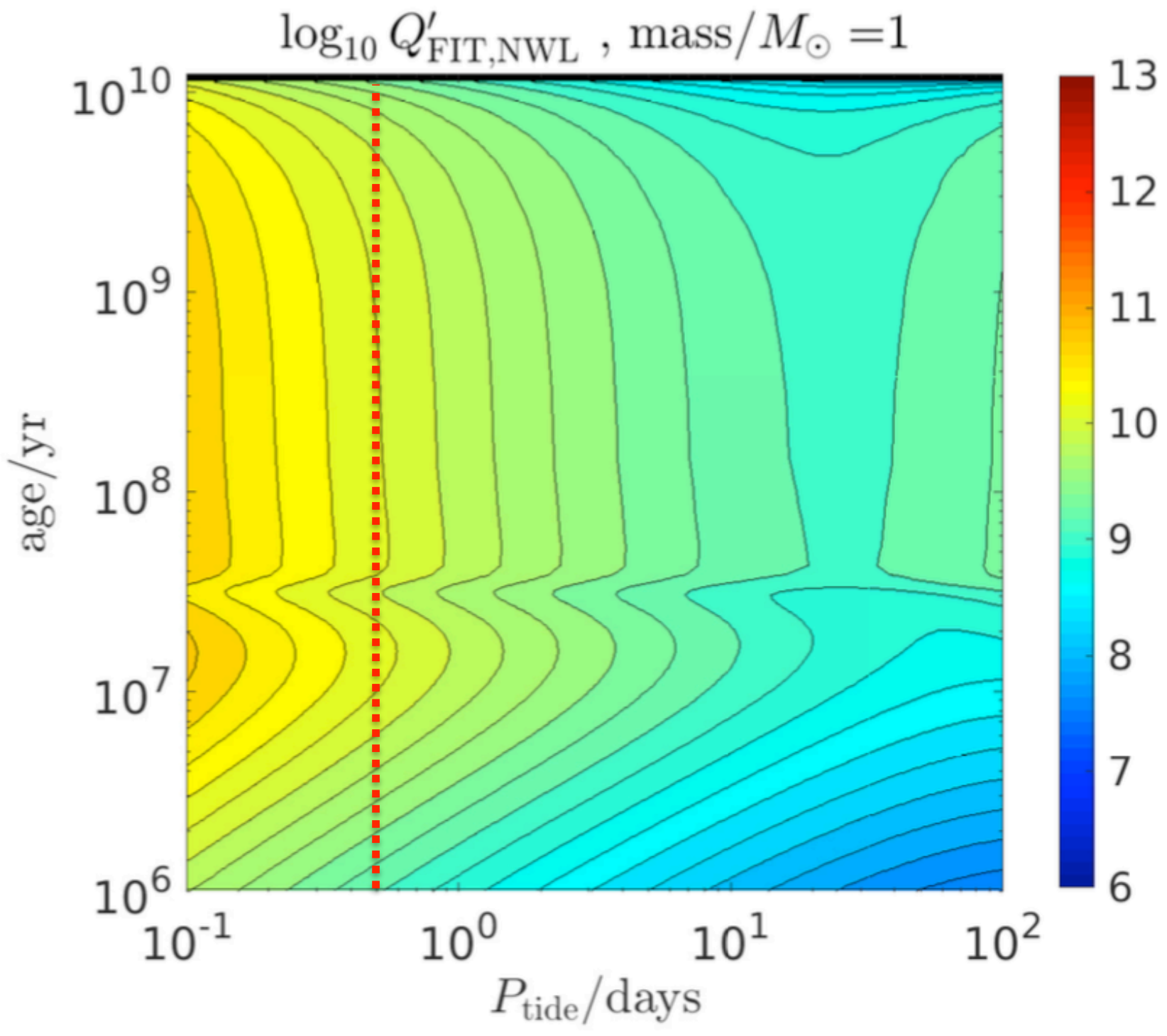
- Turbulent convection can act as an effective viscosity in damping equilibrium tides...  
Less efficient for rapidly orbiting planets (large ratios of tidal/convective frequencies)...



Duguid, Barker & Jones, 2020, MNRAS, 491, 923  
Duguid, Barker & Jones, 2020, MNRAS, 497, 3400  
Vidal & Barker, 2020, MNRAS, 497, 4472  
Vidal & Barker, 2020, ApJL, 888, L31



# Equilibrium tides in CZs



$$Q'_{\text{eq}} \sim 10^{10}$$

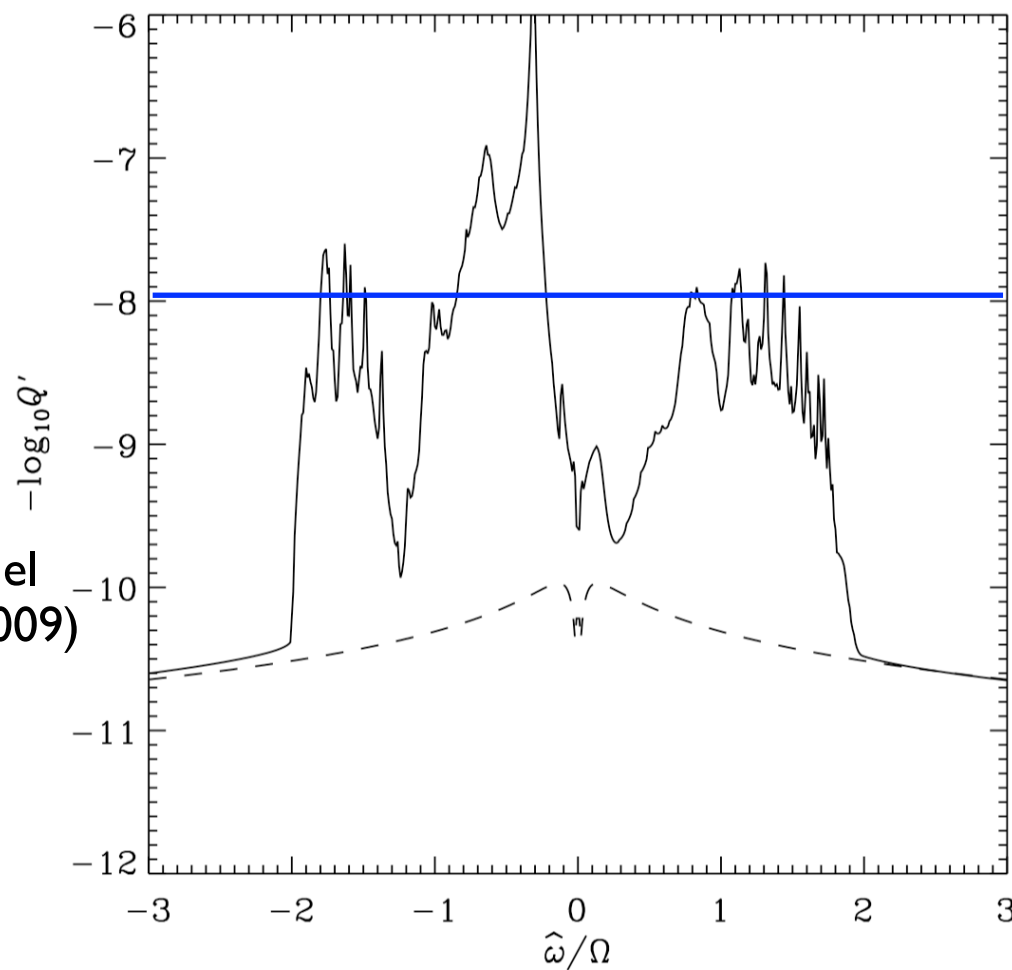
for planets on  $\sim$  day orbital periods  
 $\Rightarrow$  orbital decay timescales much longer than MS lifetime ( $> 100$  Gyr)

Equilibrium tide dissipation in CZs is probably unimportant for tidal dissipation in PMS and MS stars but is probably dominant for giant stars

## 2. Dynamical tides: inertial waves in CZs

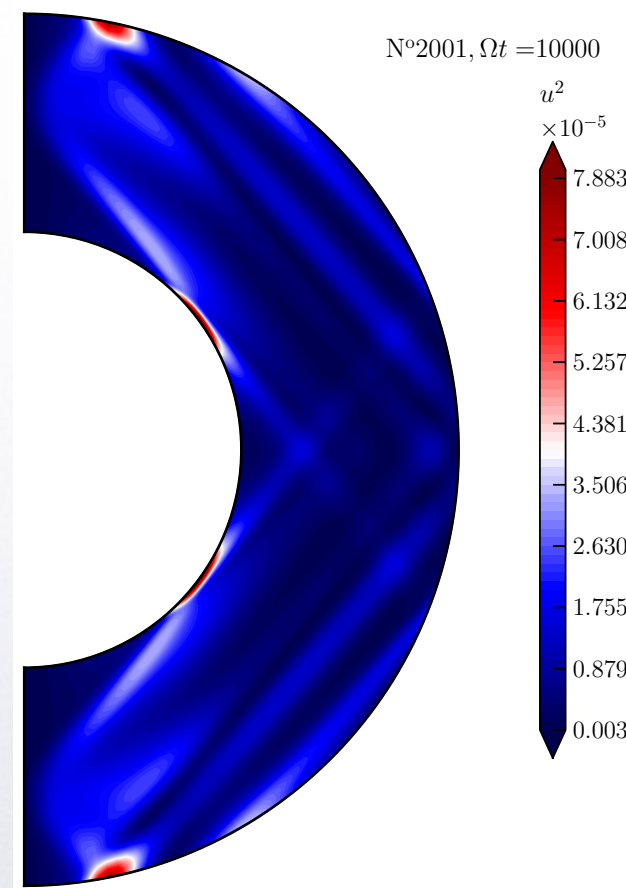
- Tidal forcing excites inertial waves in CZs of rotating stars if the tidal frequency is less than twice the stellar spin frequency i.e.  $P_{\text{rot}} < 2P_{\text{tide}}$  (not satisfied for most HJs currently!)
- I compute dissipation due to inertial waves using a frequency-averaged formalism building upon Ogilvie (2013) & Mathis (2015), Gallet et al. 2017+ *but accounting for the realistic stellar structure*

1.2  $M_{\odot}$  F star model  
Barker & Ogilvie (2009)



$\langle Q'_{\text{IW}} \rangle$

Astoul & Barker, in prep

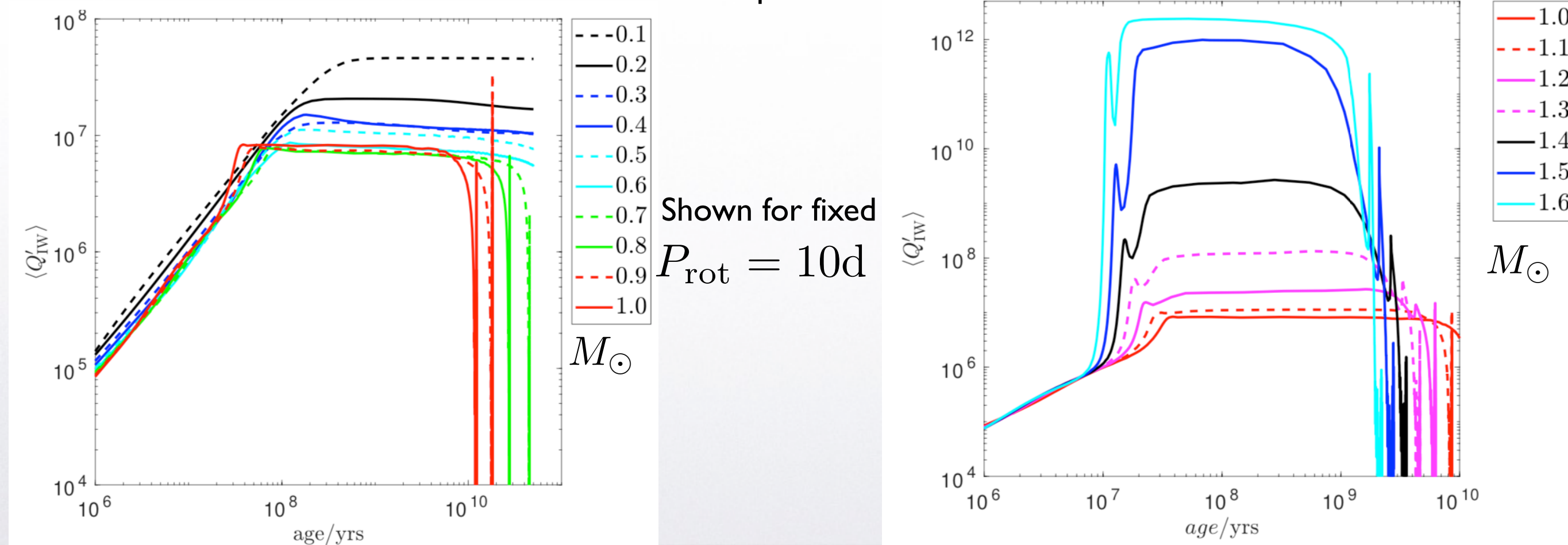




2.

# Dynamical tides: inertial waves in CZs

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- I compute dissipation due to inertial waves using a frequency-averaged formalism building upon Ogilvie (2013) & Mathis (2015) *but accounting for the realistic stellar structure*
- Solar-like stars have  $\langle Q'_{\text{IW}} \rangle \approx 10^7 (P_{\text{rot}}/10\text{d})^2$  on the MS *when these waves are excited*. Consistent with statistical analysis of observations by Collier Cameron & Jardine (2018)
- Dominant tidal mechanism in PMS stars... Less dissipative in F stars.



# 3. Dynamical tides: internal gravity waves in RZs

- Tidal forcing excites internal gravity waves in the RZ that propagate towards the centres of solar-type stars
- If waves are weakly damped, obtain global standing modes (g-modes), with efficient dissipation only for narrow (resonant) ranges of tidal frequencies.
- If waves are fully damped (e.g. by wave breaking), we obtain efficient dissipation:

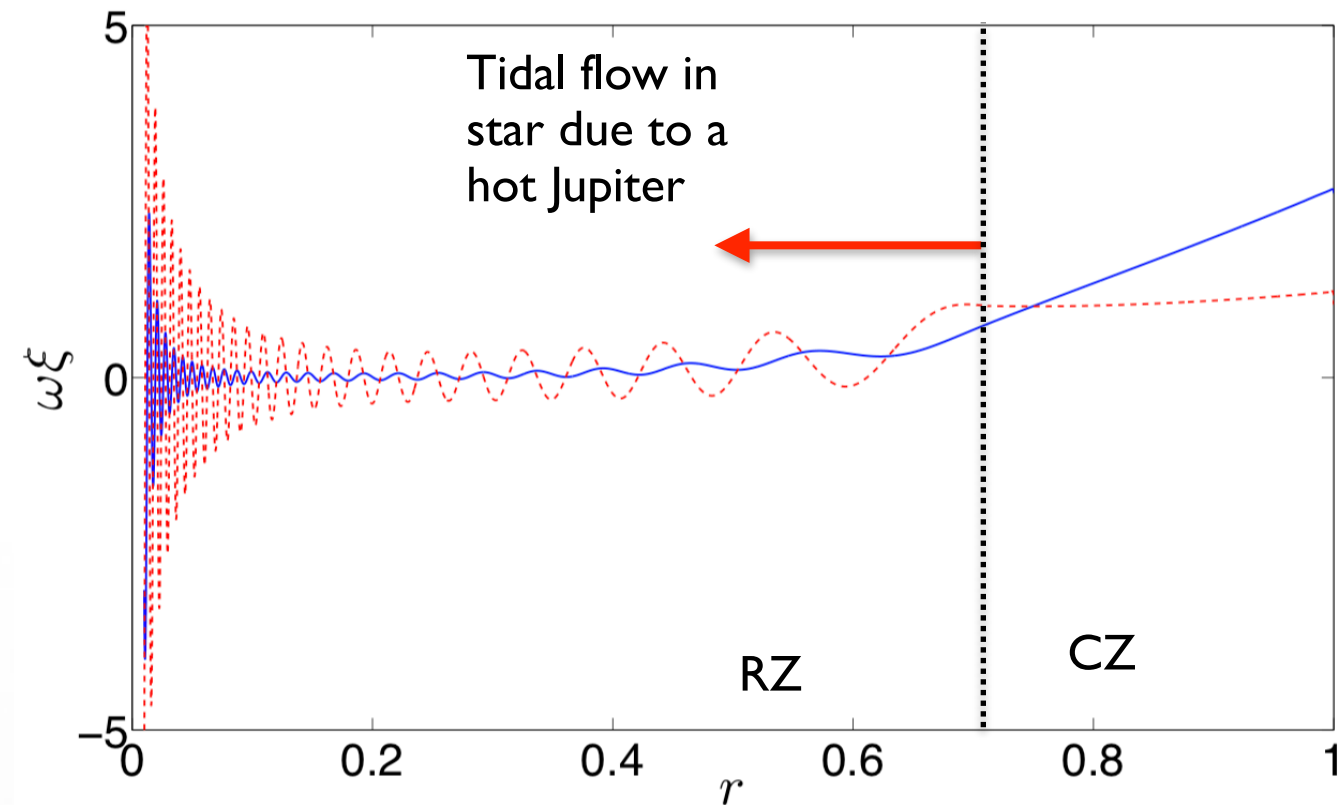
$$Q'_{\text{IGW}} \approx 10^5 \left( \frac{P_{\text{orb}}}{1\text{d}} \right)^{3/10}$$

- Drives hot Jupiter orbital decay on the timescale

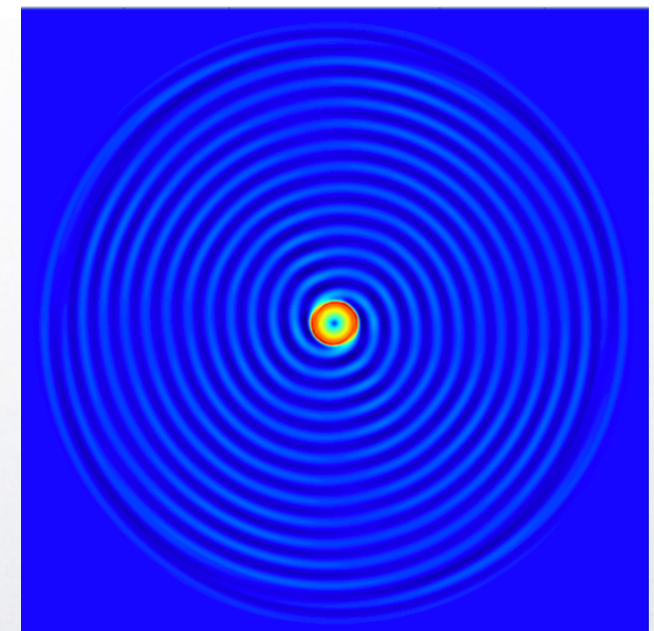
$$\tau_a \approx 2 \text{ Myr} \left( \frac{M_J}{M_p} \right) \left( \frac{P}{1\text{d}} \right)^7$$

- This mechanism may explain the observed orbital decay of e.g. WASP-12 b and the absence of decay in WASP-18 b...

However: *should the waves be fully damped in WASP-12 b? Depends on stellar model...*



Wave breaking  
can lead to  
efficient wave  
damping

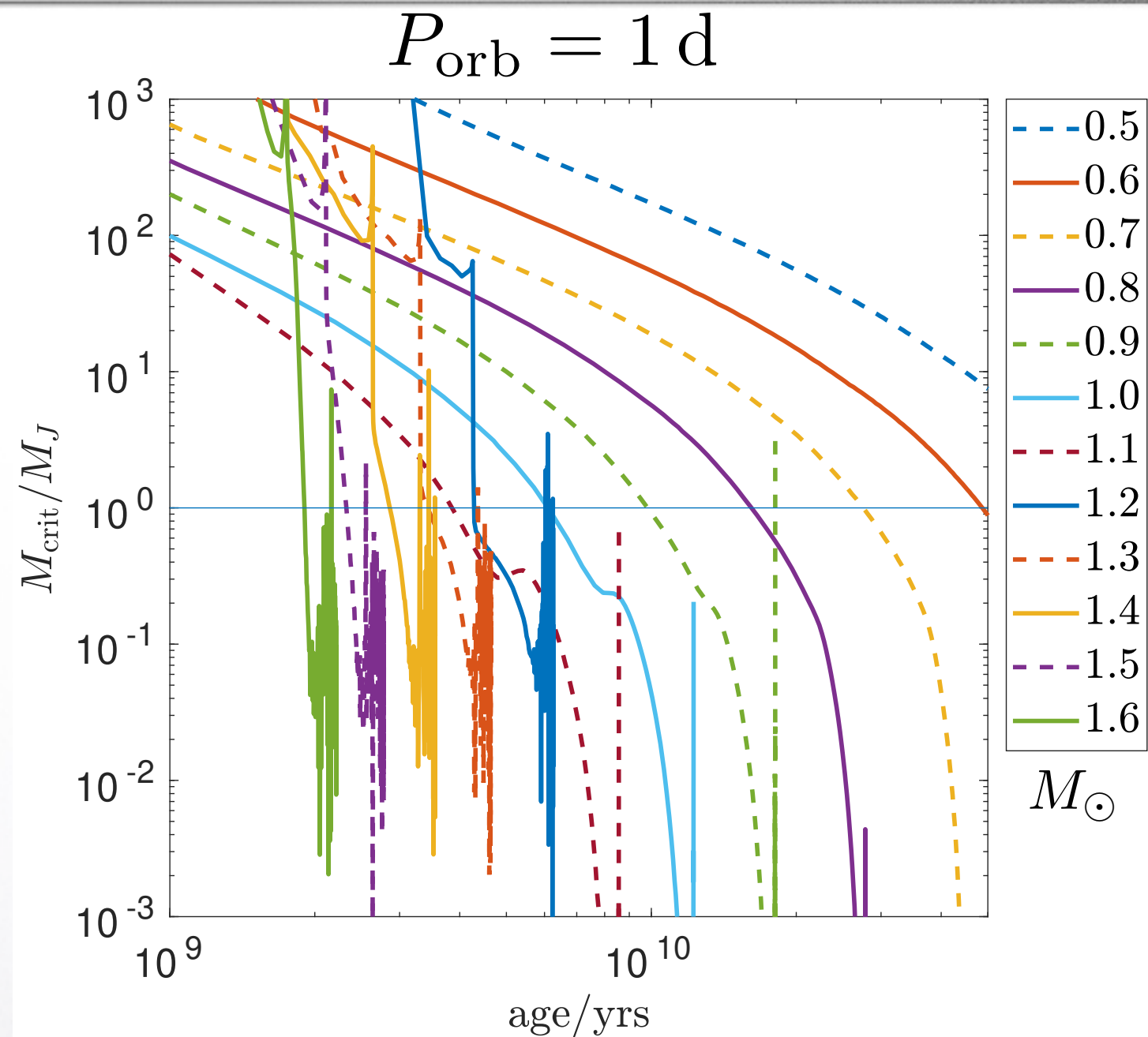


Barker & Ogilvie 2010

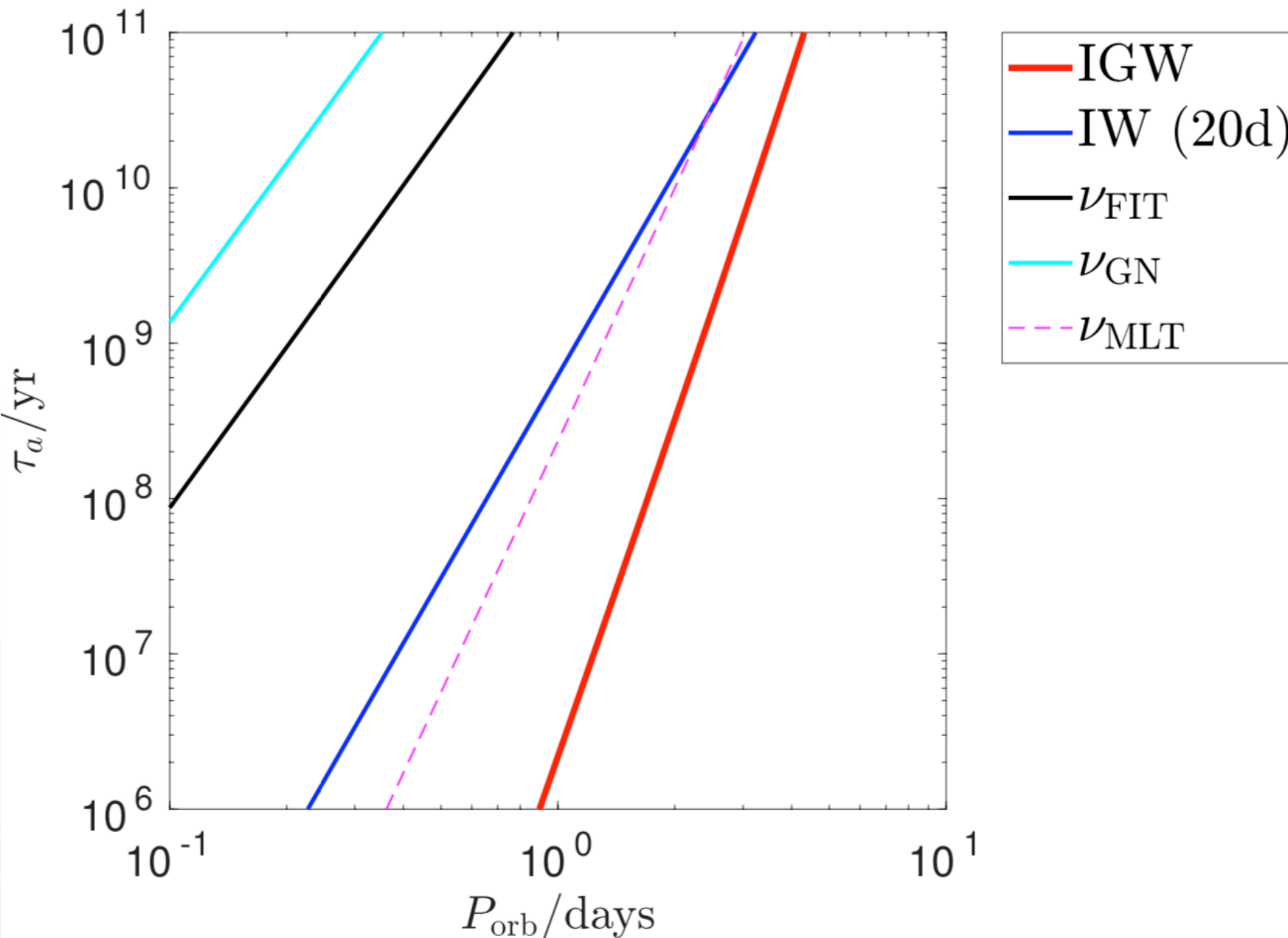


### 3. Dynamical tides: internal gravity waves in RZs

- Wave breaking predicted if planetary mass exceeds  $M_{\text{crit}}$ , which depends strongly on mass and age
- For the current Sun + 1 day orbit, we predict
$$M_{\text{crit}} \approx 3.3M_J \left( \frac{g_{\odot}}{g} \right)^{\frac{1}{2}} \left( \frac{C_{\odot}}{C} \right)^{\frac{5}{2}}$$
- Wave breaking predicted for lower mass planets as the star evolves
- Many short-period hot Jupiters can be destroyed by this mechanism near the end of the MS.  
Consistent with observational results from Hamer & Schlaufman (2019, 2020) and Mustill et al. 2021



# Summary: orbital decay timescales for HJs



(c)  $M/M_{\odot} = 1$ , age/yr =  $4.70 \times 10^9$

- Hot Jupiter orbital decay slowly rotating stars is primarily due to internal gravity waves in RZs — if these waves are efficiently damped (e.g. by wave breaking)
- Inertial waves in CZs are important around rapidly rotating young stars
- Equilibrium tides are probably unimportant on the PMS and MS, *though is probably the dominant mechanism of tidal dissipation in giant stars*



# Summary: predictions for orbital decay

Name	$M_p/M_J$	$P_{\text{orb}}/\text{d}$	$M/M_\odot$	$R/R_\odot$	$Z_{\text{init}}$	$T_{\text{eff}}$	$P_{\text{rot}}/\text{d}$	age/Gyr	$Q'_{\text{obs}}$	$Q'_{\text{IGW}}$	$\tau_a/\text{Myr}$	$T_{\text{shift}}/s$
WASP-4b	1.2	1.34	0.93	0.9	0.02	5542	23	5.8	$4.5 - 8.5 \times 10^4$	$3.3 - 3.8 \times 10^5$	14 - 17	13-15
WASP-12b (MS1)	1.47	1.09	1.43	1.68	0.03	6376	38	1.62	$2 \times 10^5$	$4.5 \times 10^6$	11	20.5
WASP-12b (MS2)	1.47	1.09	1.32	1.69	0.025	6072	38	3.1	$2 \times 10^5$	$2.2 \times 10^5$	0.42	522
WASP-12b (SG)	1.47	1.09	1.24	1.69	0.023	6126	38	4.1	$2 \times 10^5$	$2.7 \times 10^5$	0.58	384
WASP-18b	11.4	0.94	1.24	1.29	0.02	6306	6	1.37	$> 1.3 \times 10^6$	$2.6 \times 10^6$	1.2	200
WASP-19b	1.14	0.79	0.94	1.01	0.02	5624	13	9.28	$3.5 - 7.5 \times 10^5$	$0.6 - 0.8 \times 10^5$	0.13 - 0.3	675-1000
WASP-43b	2.03	0.81	0.72	0.67	0.02	4462	6	5.03	$> 0.7 - 3.5 \times 10^5$	$1.3 \times 10^5$	0.98	230
WASP-72b	1.55	2.22	1.39	2.01	0.01	6876?	17	2.3	$> 2.1 \times 10^3$	$> 10^{12}$		
WASP-103b	1.51	0.93	1.21	1.43	0.02	6115	7	3.48	$> 1.1 \times 10^5$	$4 \times 10^5$	0.68	322
WASP-114b	1.77	1.55	1.29	1.42	0.03	6206	12	2.12		$3 \times 10^6$	48	4.7
WASP-121b	1.18	1.27	1.353	1.49	0.025	6429	5.5	1.42		$2.4 \times 10^7$	225	1
WASP-122b	1.284	1.71	1.24	1.50	0.04	5895	23	4.2		$3.5 \times 10^5$	8.4	27
WASP-128b	37.2	2.21	1.16	1.16	0.02	6108	3	1.57		$1.7 \times 10^8$	1400	0.2
NGTS-6b	1.33	0.88	0.787	0.74	0.025	4774		9.01		$> 0.99 \times 10^5$	1.2	182
NGTS-7Ab	62.0	0.676	0.48?	0.645	0.02	3736	sync?	0.0055		$> 0.9 \times 10^5$		
NGTS-10b	2.16	0.77	0.696	0.68	0.02	4428	8.8	10.06		$0.99 \times 10^5$	0.5	440
HAT-P-23b	2.09	1.21	1.13	1.22	0.03	5916	7.5	4.3	$> 4.5 \times 10^5$	$3.5 \times 10^5$	2.5	91
HATS-18b	1.98	0.84	1.04	1.03	0.03	5735	8.3	4.26		$1.1 \times 10^5$	0.33	686
KELT-16b	2.75	0.97	1.21	1.38	0.02	6180	9	2.97	$> 0.7 \times 10^5$	$7 \times 10^5$	0.98	228
TRES-3b	1.91	1.306	0.924	0.845	0.009	5699	27	1.23	$1.1 \times 10^5$	$6.1 \times 10^5$	23.6	9.5
OGLE-TR-56b	1.39	1.21	1.23	1.38	0.02	6235	23	2.53	$> 5 \times 10^5$	$1.8 \times 10^6$	14	16.5
WTS-2b	1.12	1.018	0.82	0.74	0.02	4761	17	0.43		$1.9 \times 10^5$	6.4	35

- For further details please see the paper below or ask me!

# Conclusions

- Tidal dissipation in stars drives orbital decay of hot Jupiters, as well as orbital circularization & spin synchronization in stellar binaries.
- I have presented theoretical calculations of tidal dissipation in stars with masses  $0.1 \leq M/M_{\odot} \leq 1.6$  following their evolution
- Dominant tidal mechanism for planetary orbital decay around slowly rotating stars is internal gravity waves in RZs if waves are efficiently damped (e.g. by wave breaking)
- Dominant tidal mechanism for tidal evolution around rapidly rotating young MS and PMS stars is inertial waves in CZs
- Equilibrium tides are probably unimportant on the PMS and MS, *but are expected to be the dominant mechanism in giant stars*
- For further details and for predictions for planetary orbital decay for current hot Jupiters, see: Barker (2020, MNRAS, 498, 2270) <https://arxiv.org/abs/2008.03262.pdf>
- Further work is required to study:
  1. Direct simulations of turbulent convection & tidal flows in realistic stellar models
  2. Direct simulations of tidal dissipation due to inertial waves in realistic stellar models