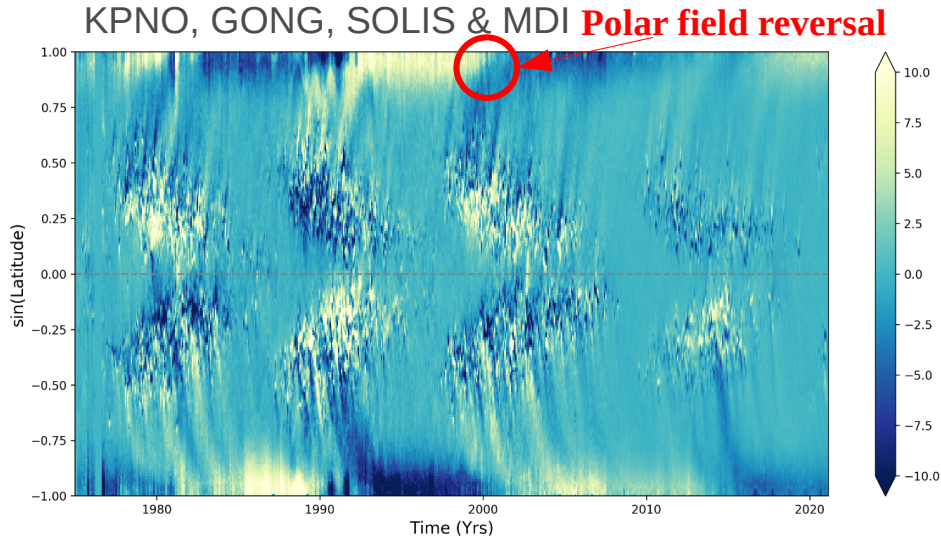


Strugarek et al. 2017, Science

# How does solar-type stars magnetism evolve?

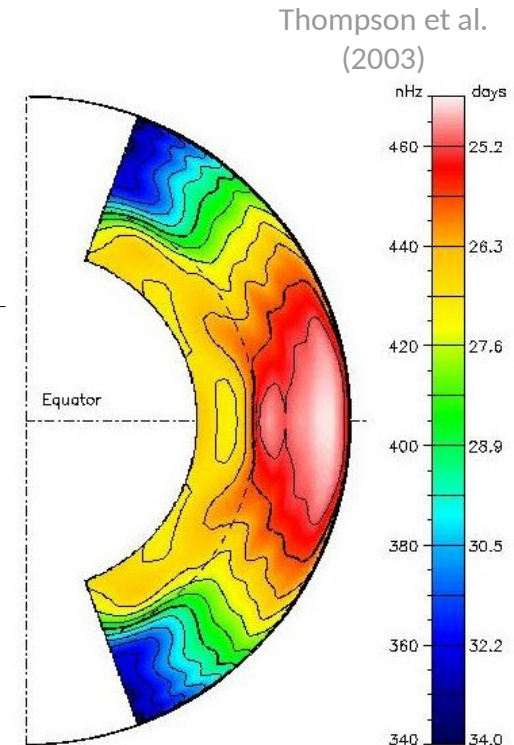
# Our star: the Sun

## Current Observations :



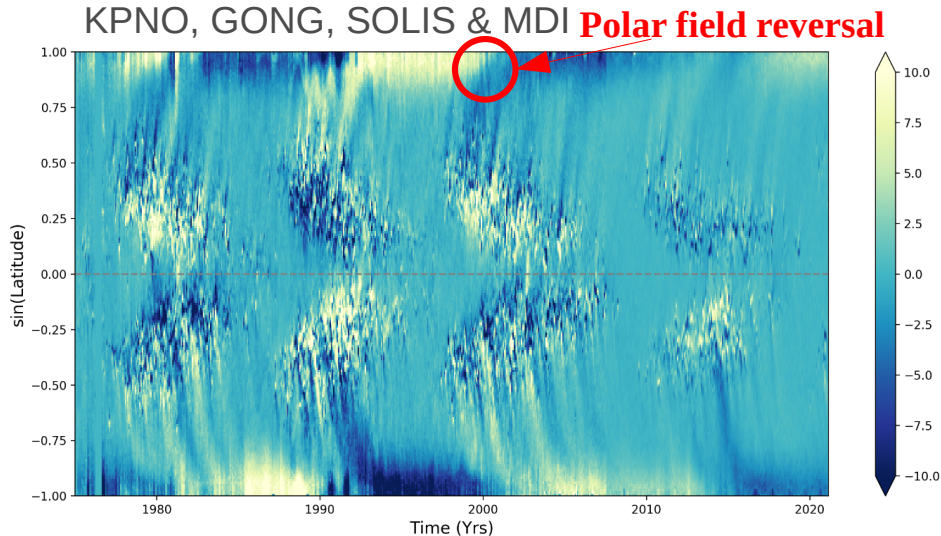
- 11-years activity cycle
- Migration of structures toward equator
- Opposite hemisphere polarities
- Alternating dipole/quadrupole
- 22-years magnetic cycle

A **solid-body rotation** of the **radiative core**, surrounded by a **differentially rotating convective envelope**



# Our star: the Sun

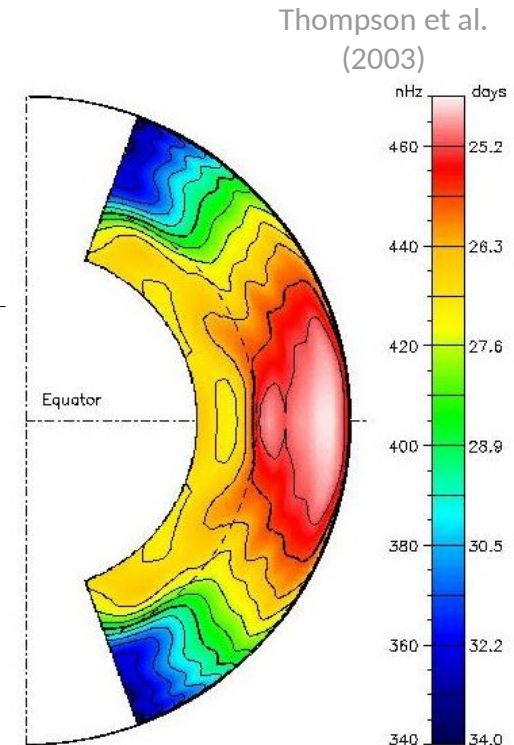
## Current Observations :

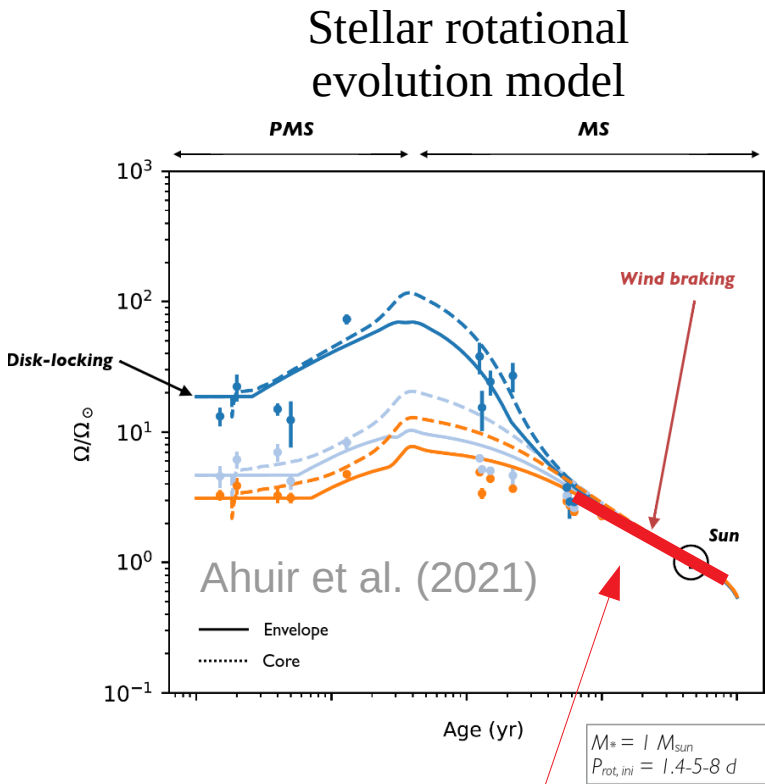


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- Alternating dipole/quadrupole
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A **solid-body rotation** of the **radiative core**, surrounded by a **differentially rotating convective envelope**

**Differential rotation** plays a key role for **dynamo mechanism**

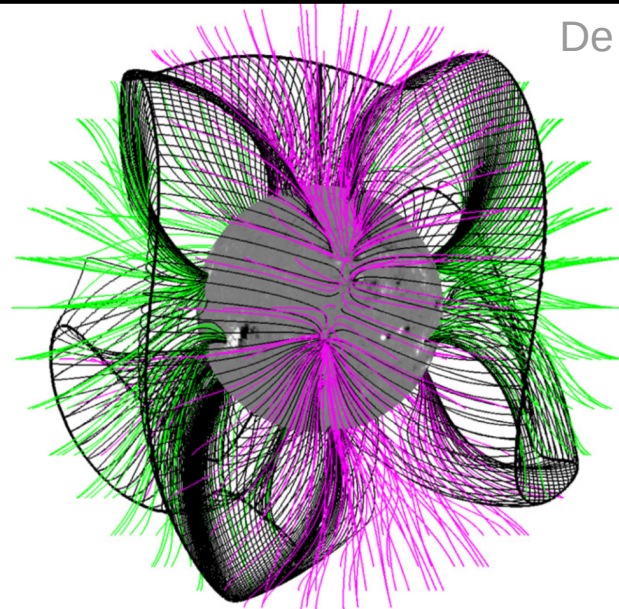




**Skumanich law:**  $\Omega(t) \propto t^{-1/2}$

Dynamo

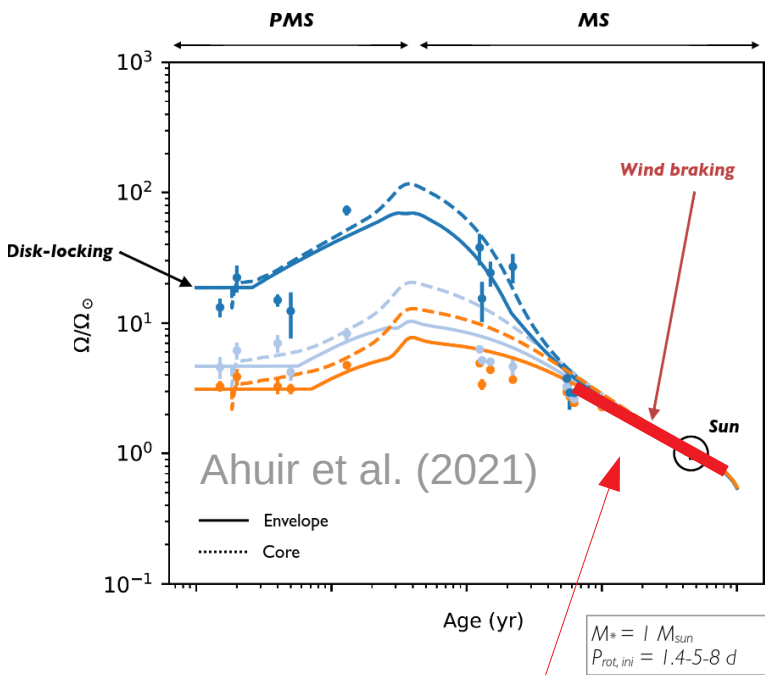
Wind



De Rosa et al. (2012)

Magnetic activity

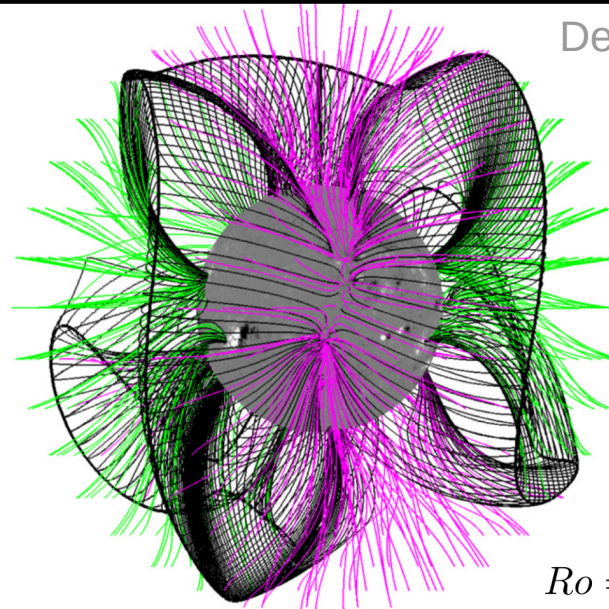
## Stellar rotational evolution model



**Skumanich law:**  $\Omega(t) \propto t^{-1/2}$

Dynamo

Wind



De Rosa et al. (2012)

Magnetic activity

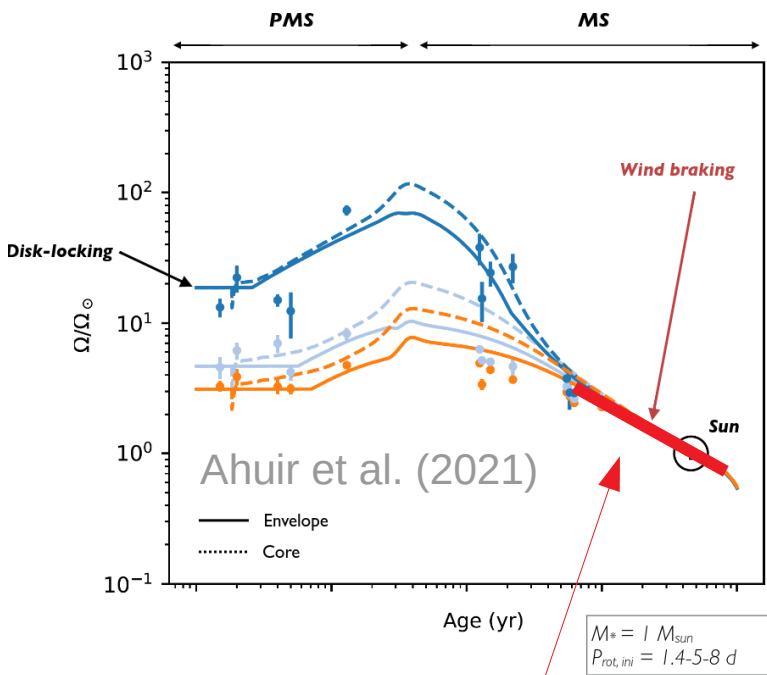
$$B_{ZDI} \propto Ro^{-1.4}$$

See et al. (2019)

Vidotto et al. (2021)

$$Ro = \frac{v}{2\Omega L} \sim \frac{\text{Advection}}{\text{Coriolis}}$$

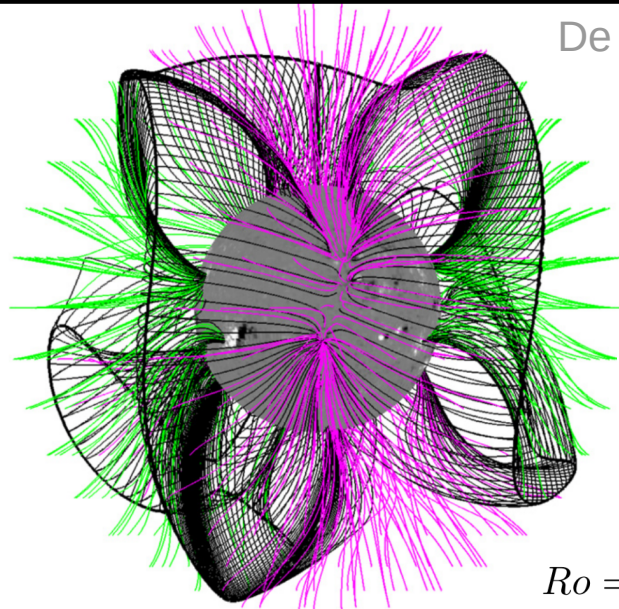
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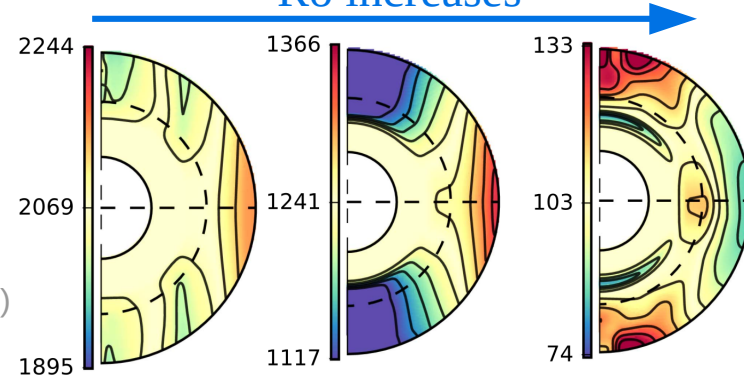
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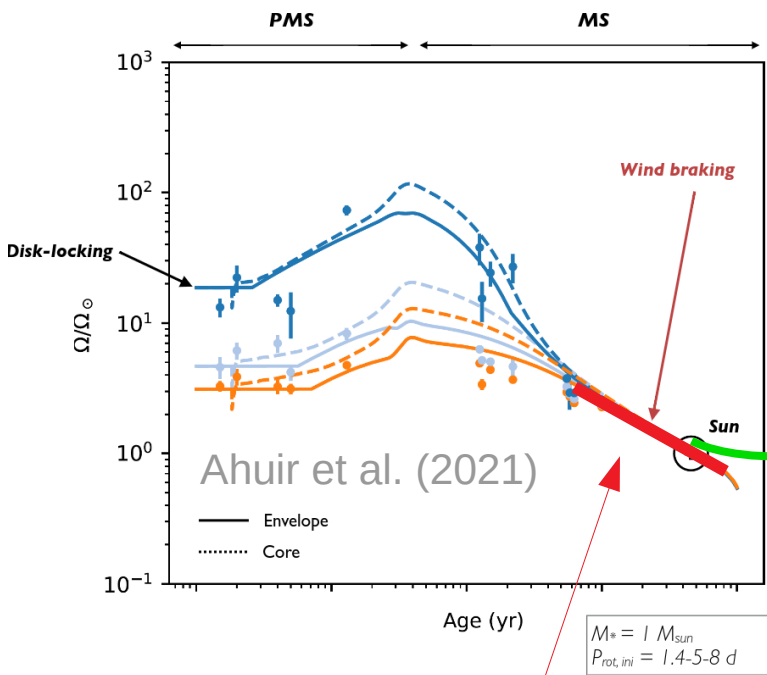
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Ro increases



Brun et al. (2017)

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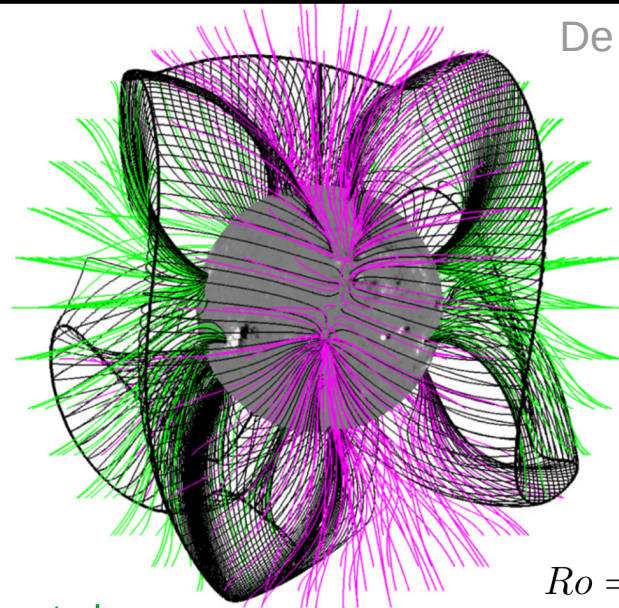


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Dynamo

Wind

van Saders et al. 2016



De Rosa et al. (2012)

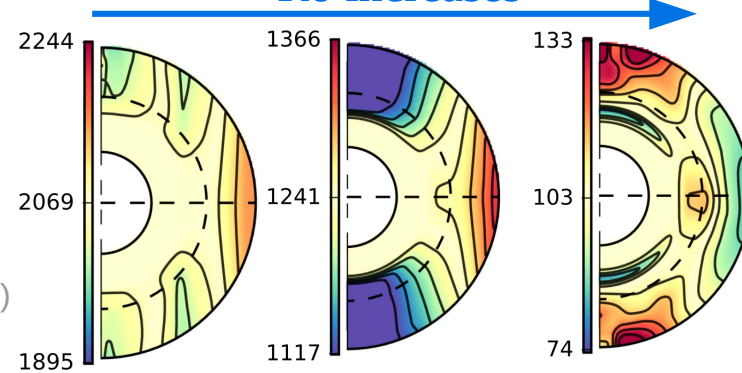
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Brun et al. (2017)

**Dynamo effect:**

The ability of a conductive fluid (plasma) to **amplify and maintain** a magnetic field against its **ohmic dissipation**.

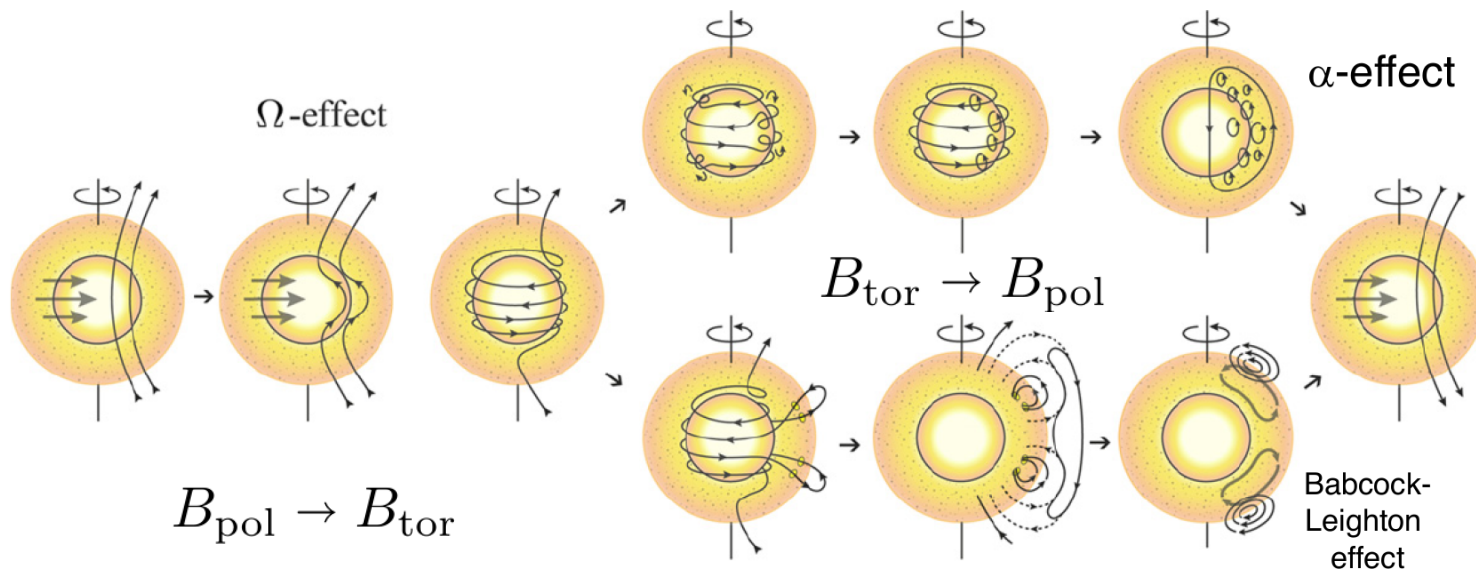
$$\frac{\partial \mathbf{B}}{\partial t} = \underbrace{\nabla \times (\mathbf{v} \times \mathbf{B})}_{\text{Induction}} - \underbrace{\nabla \times (\eta \nabla \times \mathbf{B})}_{\text{Dissipation}}$$



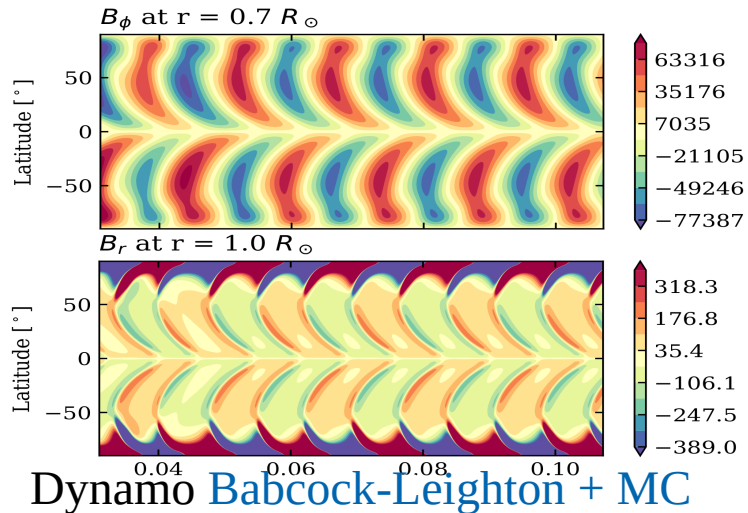
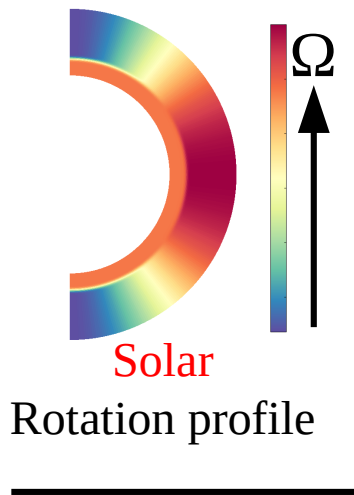
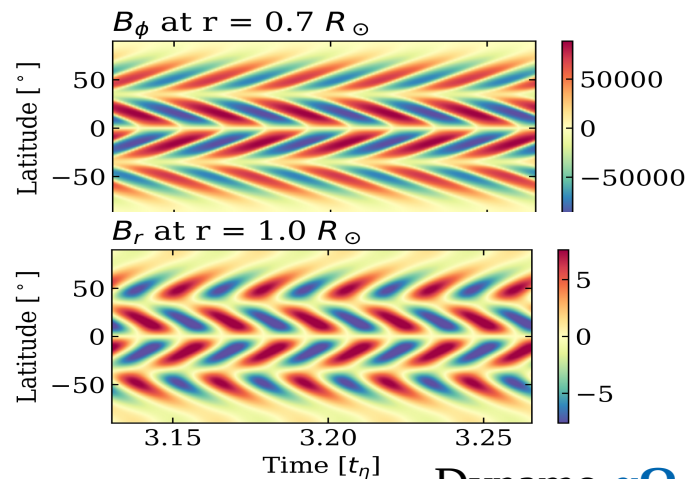
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# 2D Kinematic Approach: Solar reference cases

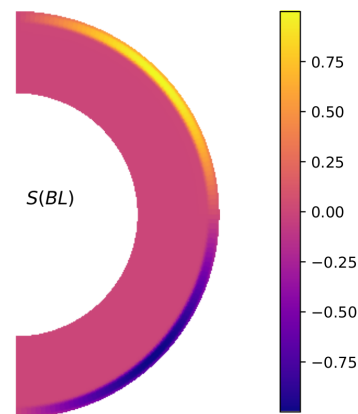
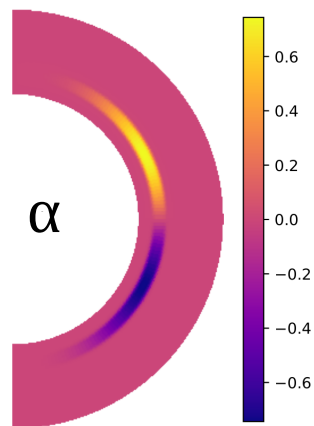


$$s = r \sin \theta \alpha \cdot e_\phi \times \nabla \Omega$$

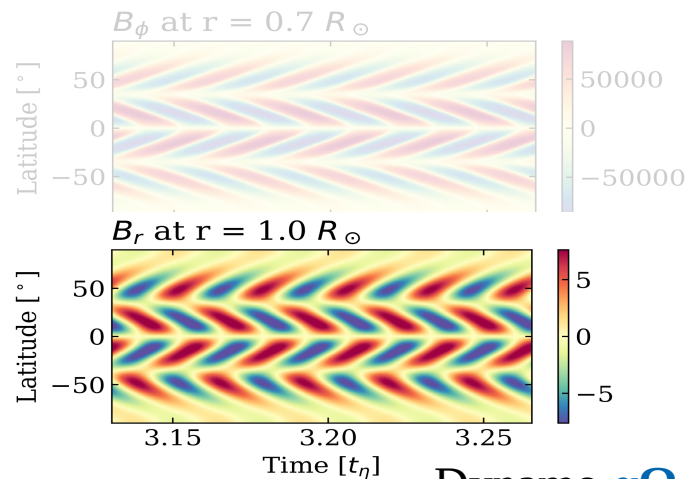
Parker-Yoshimura

Dynamo  $\alpha\Omega$

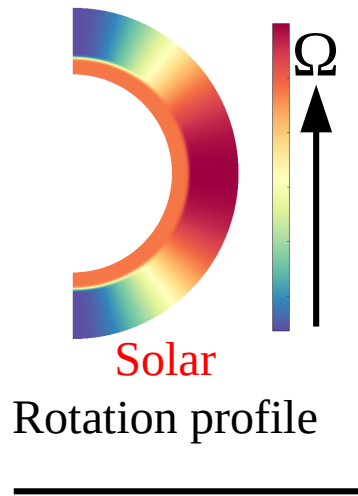
Dynamo Babcock-Leighton + MC



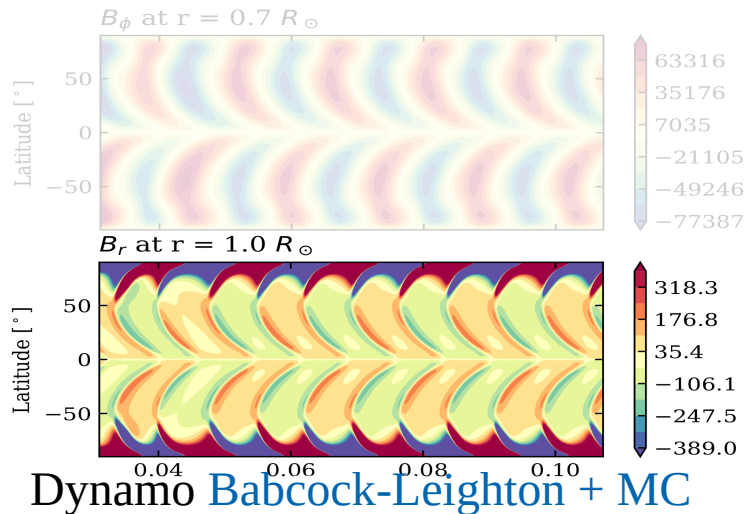
# 2D Kinematic Approach: Solar reference cases



Dynamo  $\alpha\Omega$



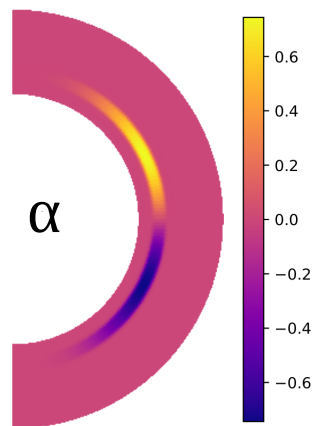
Solar  
Rotation profile



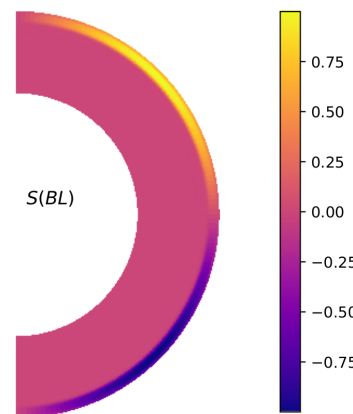
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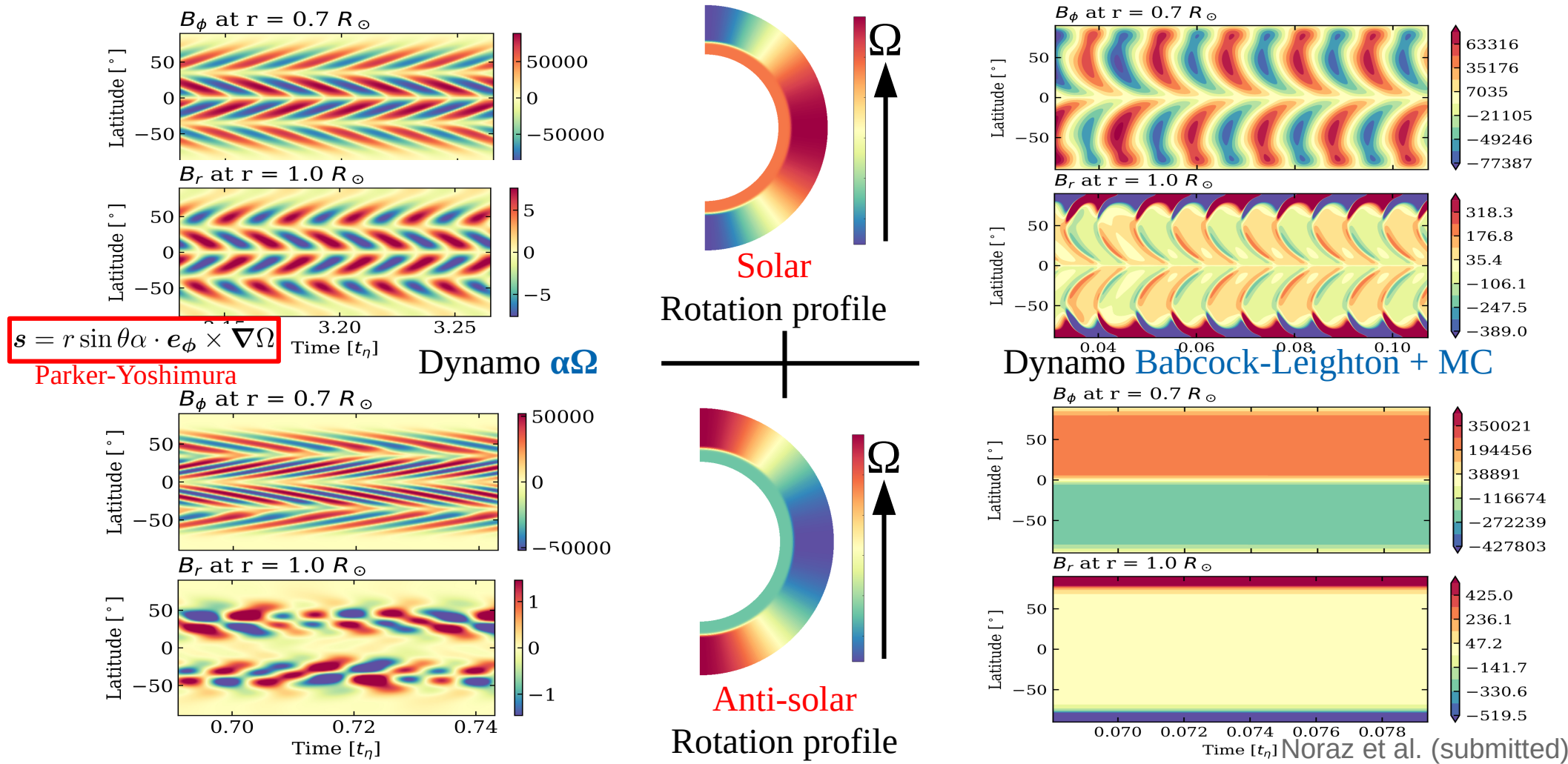


$\alpha$

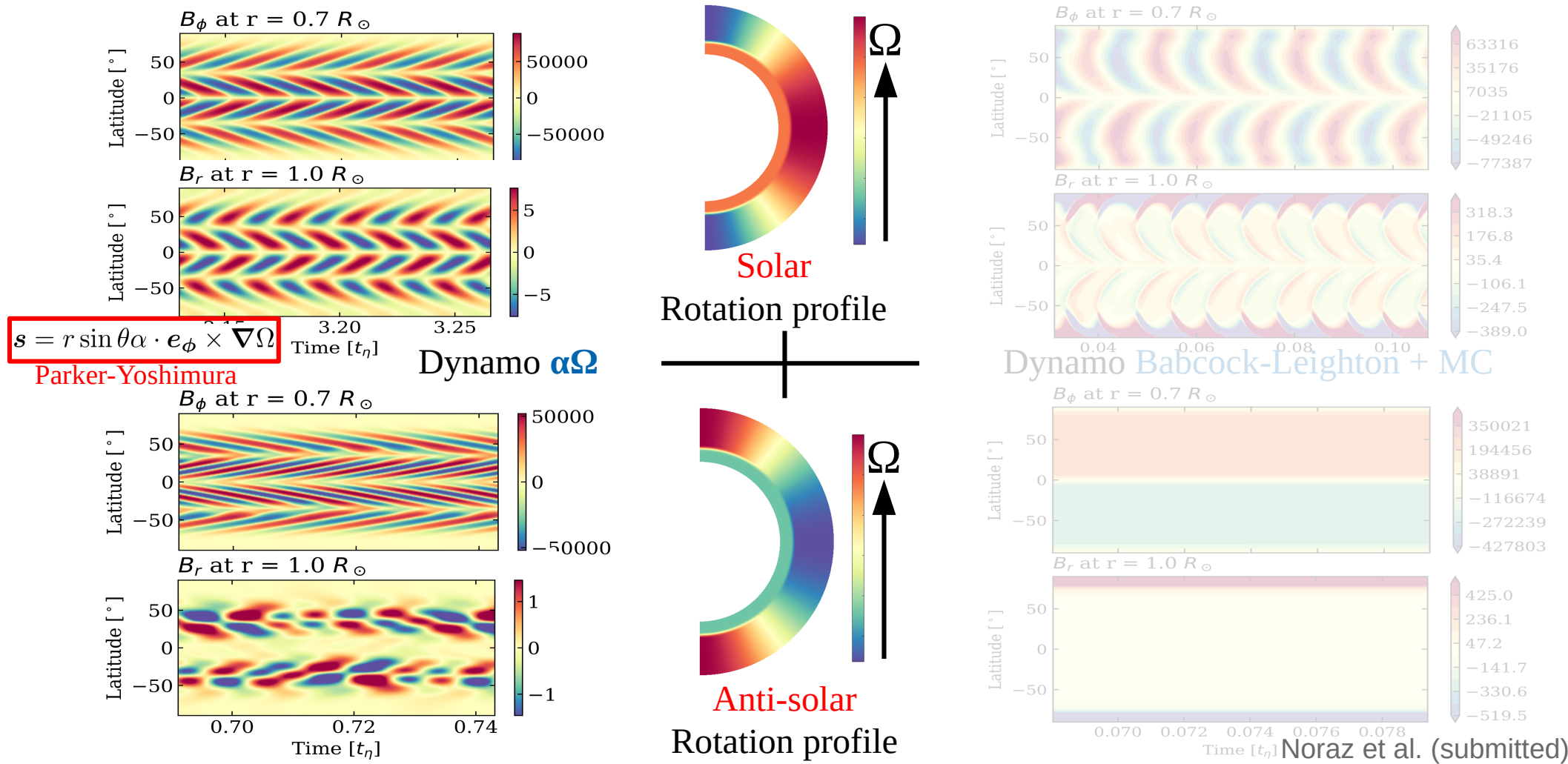


$S(BL)$

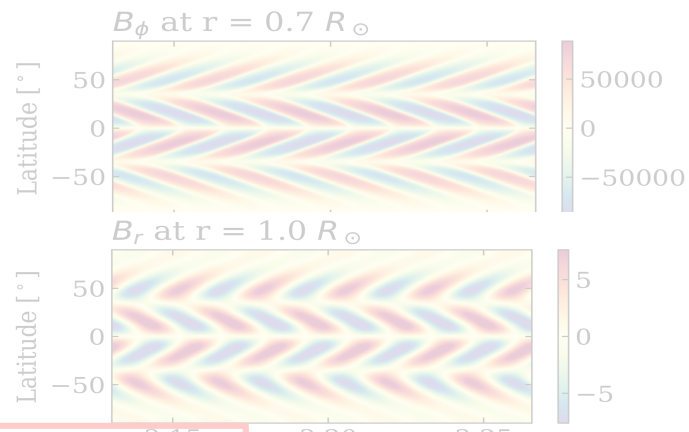
# 2D Kinematic Approach: Application to the anti-solar DR



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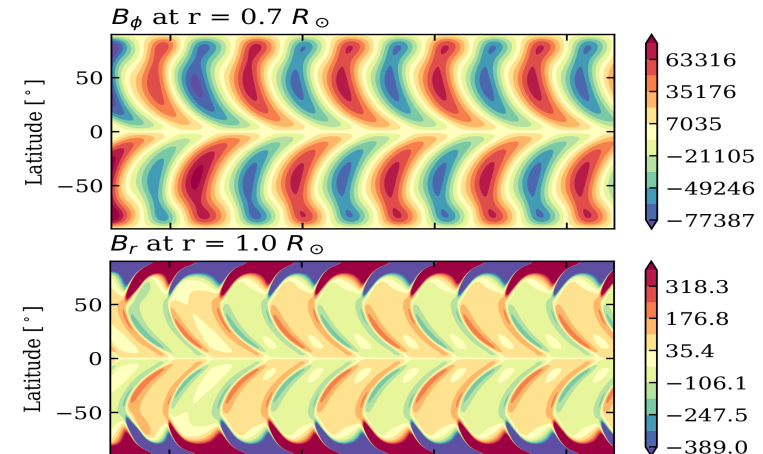
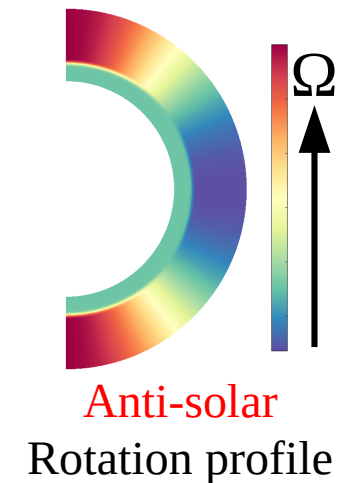
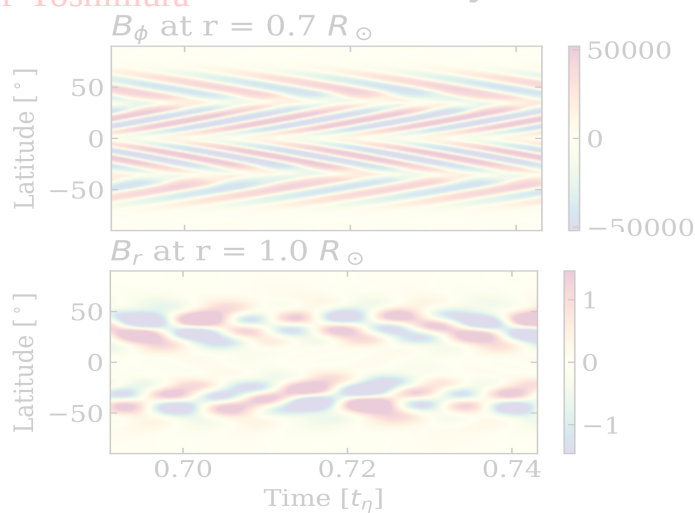
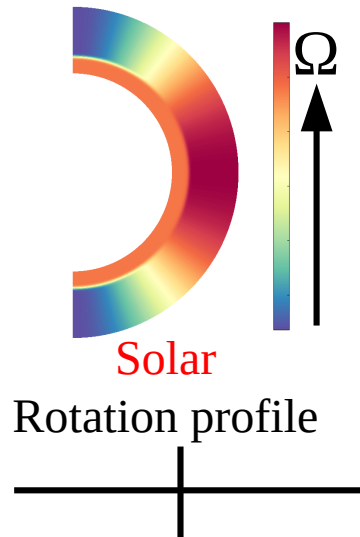
# 2D Kinematic Approach: Application to the anti-solar DR



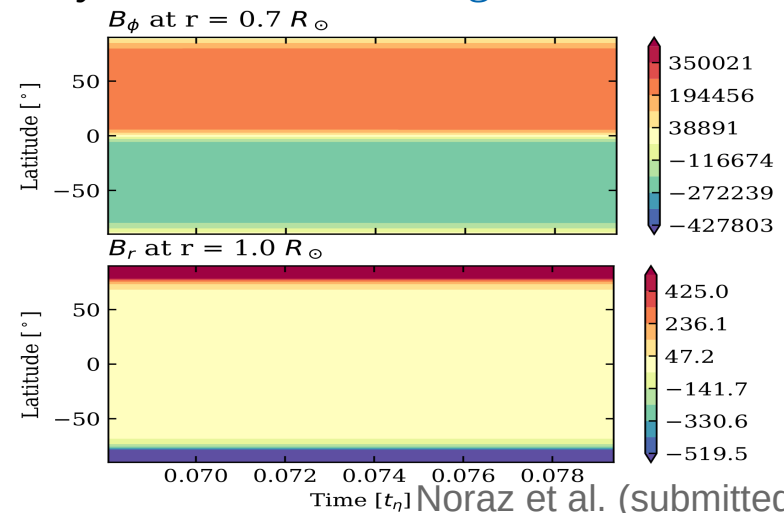
$$s = r \sin \theta \alpha \cdot e_\phi \times \nabla \Omega$$

Parker-Yoshimura

Dynamo  $\alpha\Omega$



Dynamo Babcock-Leighton + MC



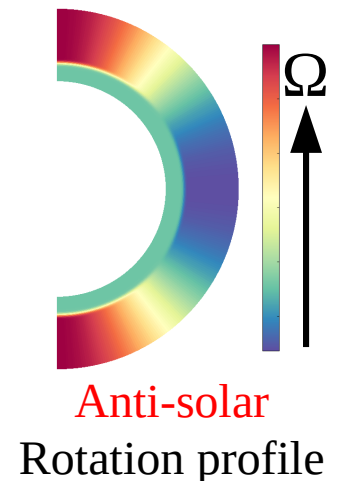
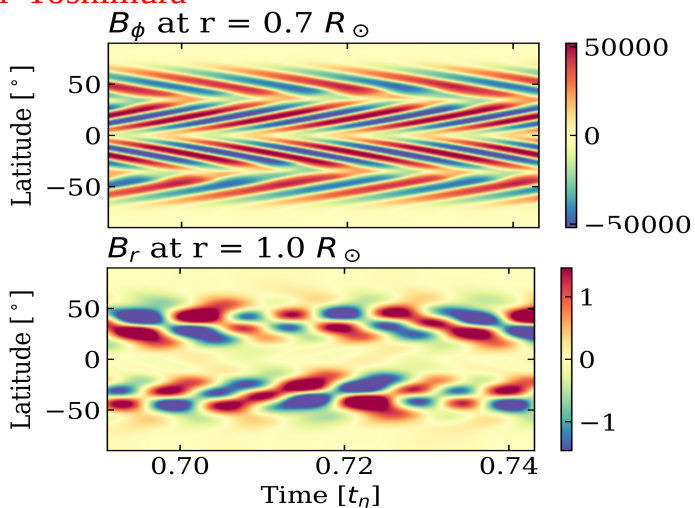
The dynamo still occurs,  
but how is the magnetic cycle lost?



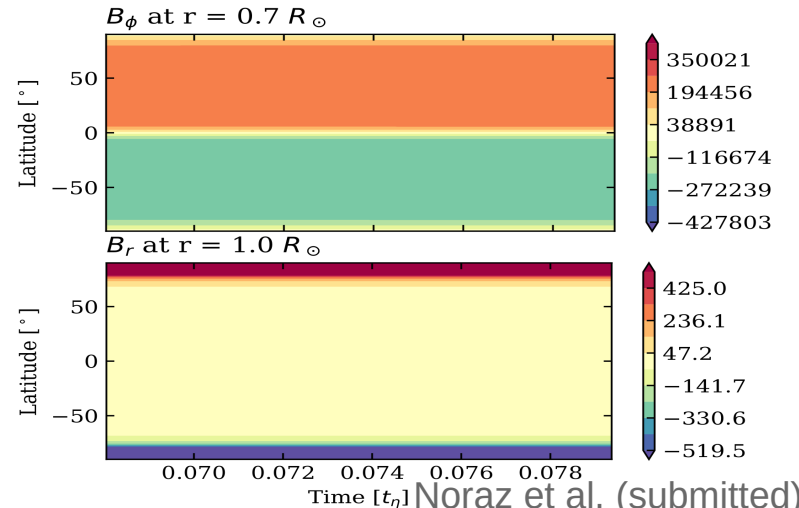
$$s = r \sin \theta \alpha \cdot e_\phi \times \nabla \Omega$$

Parker-Yoshimura

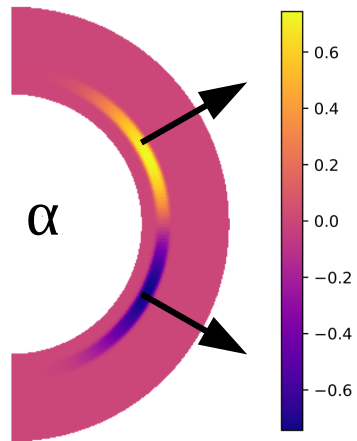
Dynamo  $\alpha\Omega$



Dynamo Babcock-Leighton + MC



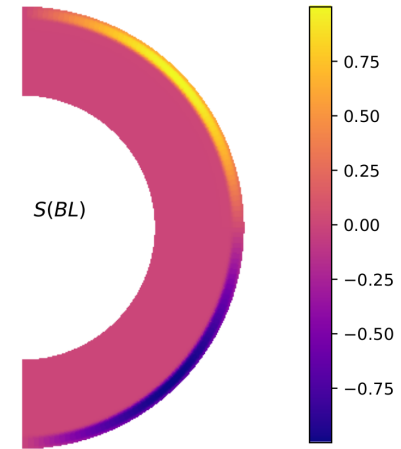
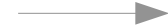
Noraz et al. (submitted)



$\alpha$ -source



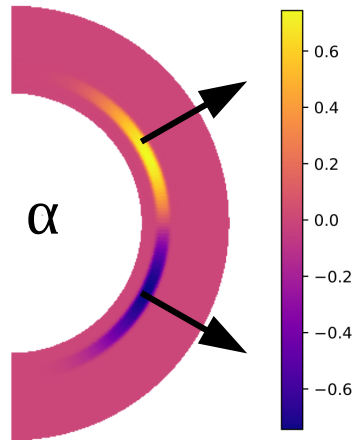
$\alpha$ -effect location:  
From the tachocline to the surface



Babcock-Leighton  
source



For **anti-solar DR** the dynamo becomes **stationnary** once  **$\alpha$  is leaving the tachocline** (ie. the radial shear)



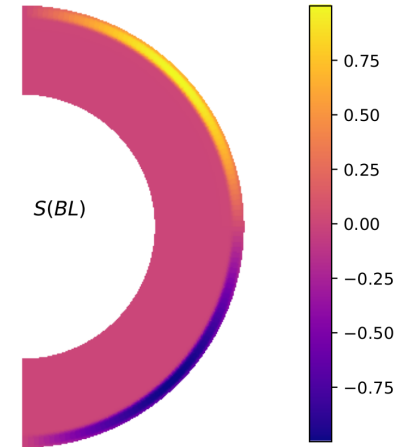
$\alpha$ -source



$r_1 - r_2$	$\odot$	anti- $\odot$
0.9 - 1	7.3	×
0.8 - 1	25.1	×
0.7 - 1	32.4	×
0.8 - 0.9	15.5	×
0.7 - 0.9	26.8	×
0.7 - 0.8	64.9 (9.7)	×
0.68 - 0.74	15.8	9.0
<b>0.67 - 0.73</b>	<b>31.8</b>	<b>17.7</b>

CZ

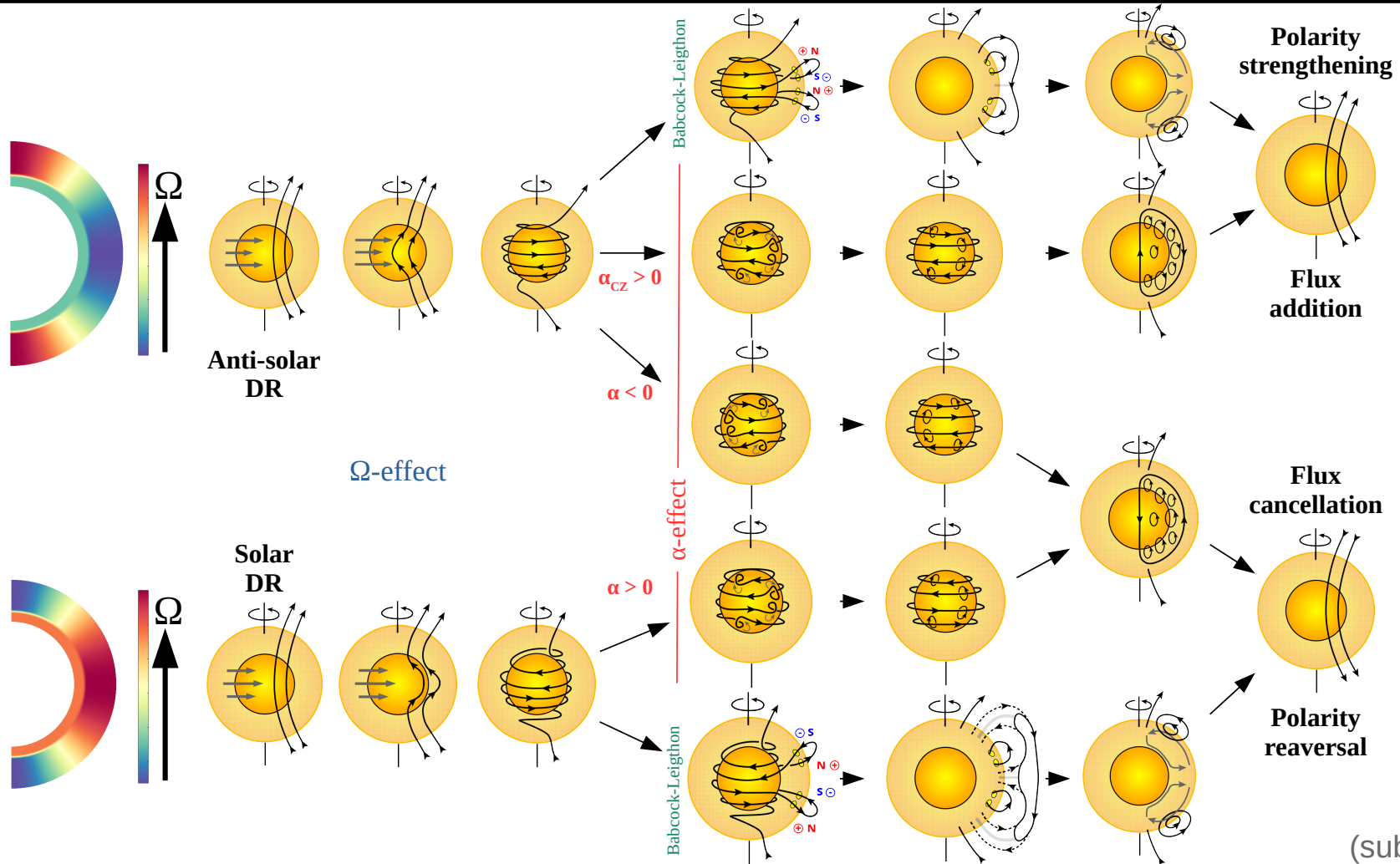
tachocline



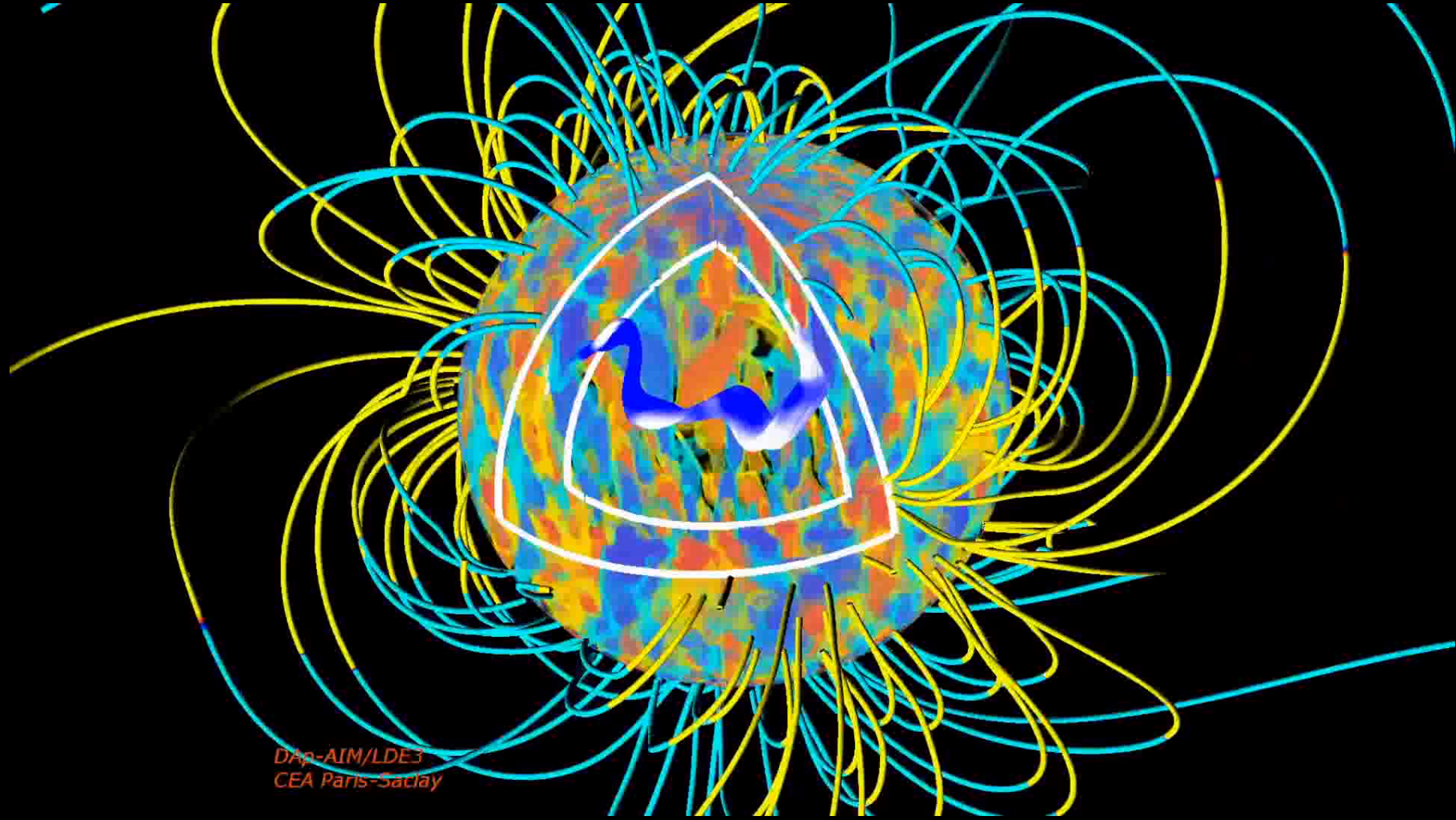
Babcock-Leighton source

**Table 3.** The  $\alpha$ -effect is spread through the tachocline and CZ for solar and anti-solar DR regimes. Magnetic cycle periods are expressed in years.

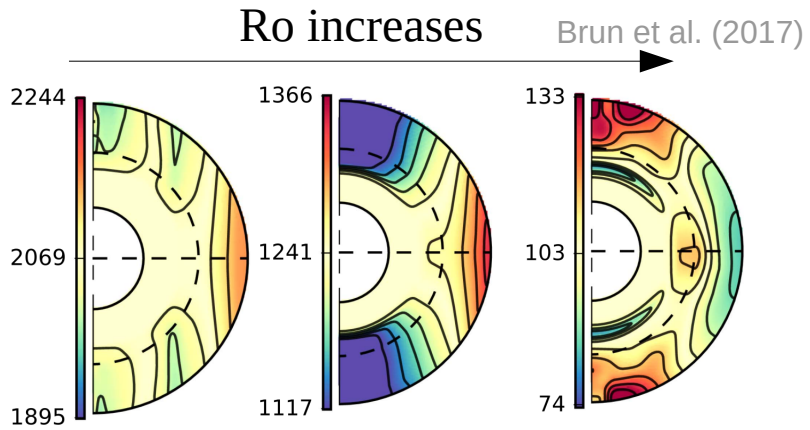
Robustness over other parameters change  
(Meridional circulation, diffusion, shear or process amplitudes...)



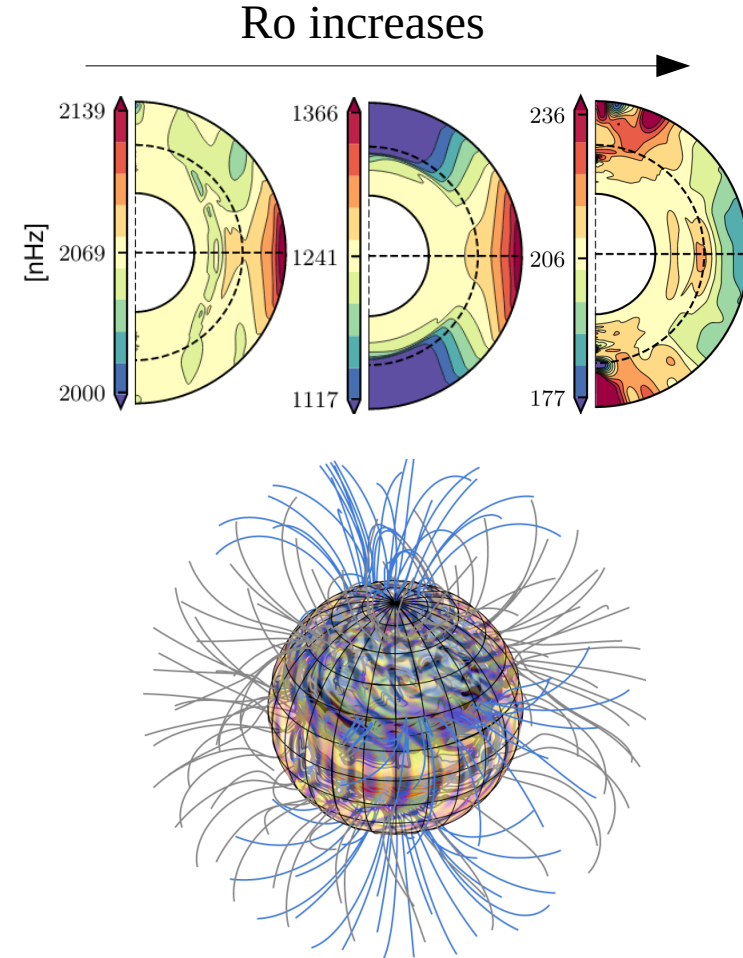
# From 2D Kinematic to Global 3D MHD turbulent dynamo



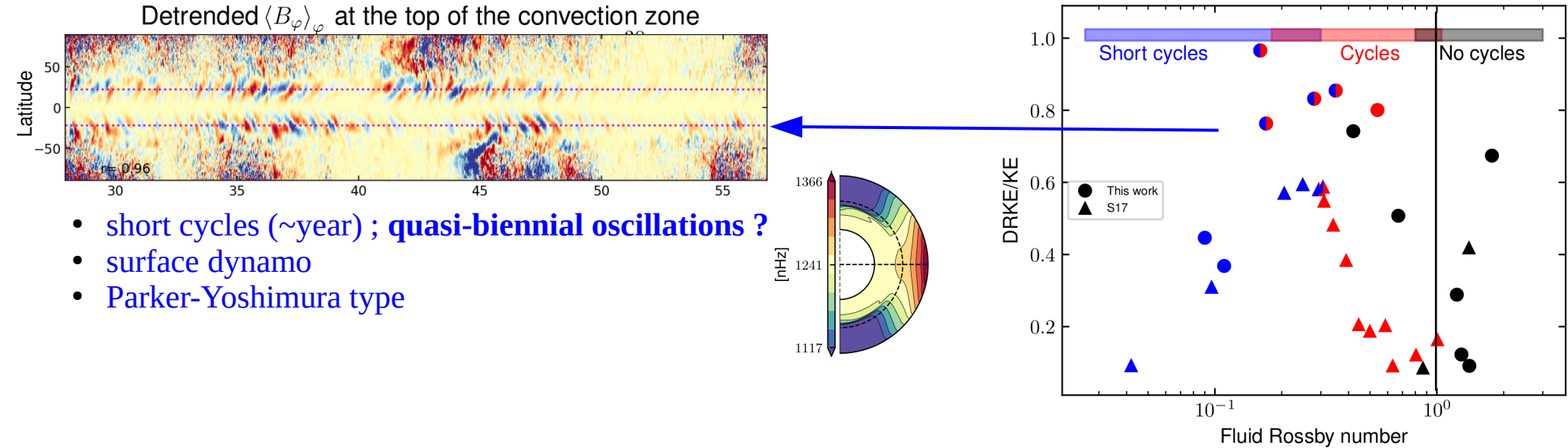
# Differential Rotation along the Rossby number



MHD →  
Brun et al. (2021)

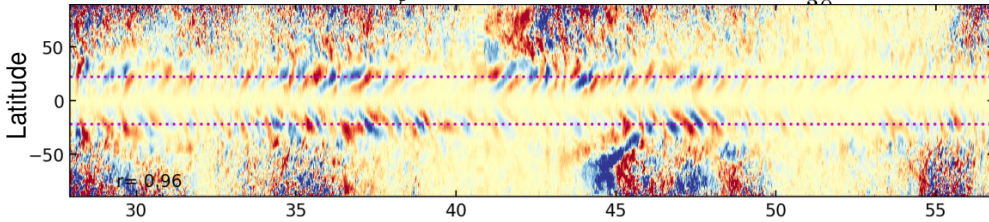


# Role of the rotation on dynamos

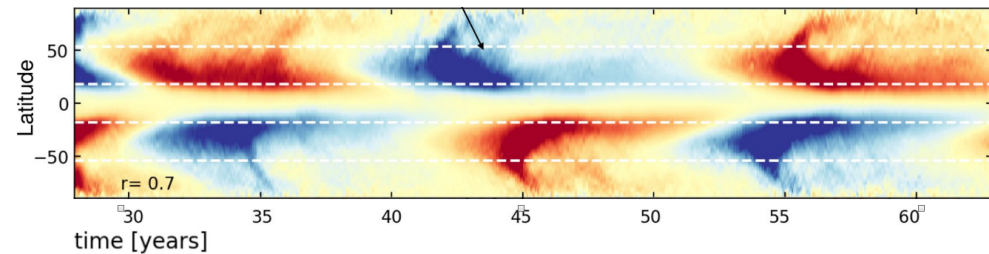


# Role of the rotation on dynamos

Detrended  $\langle B_\varphi \rangle_\varphi$  at the top of the convection zone

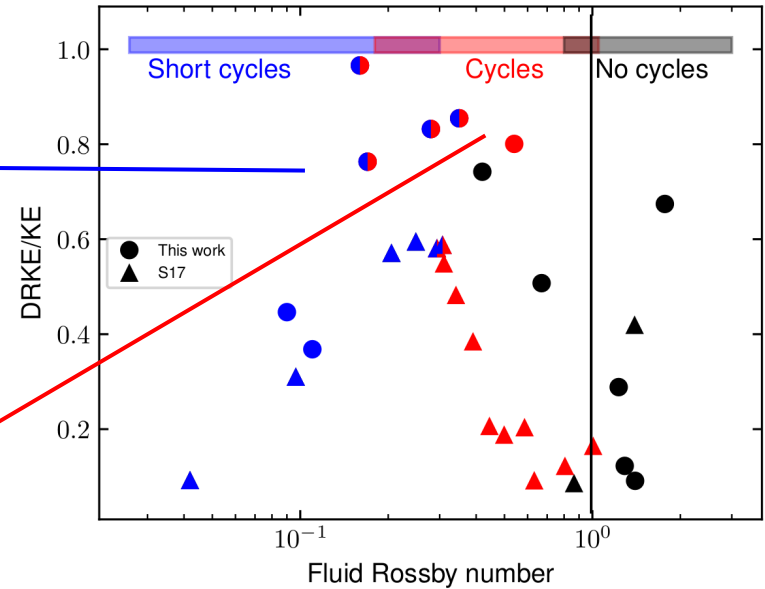
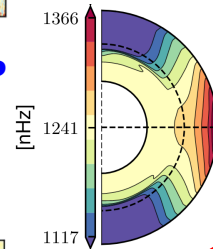


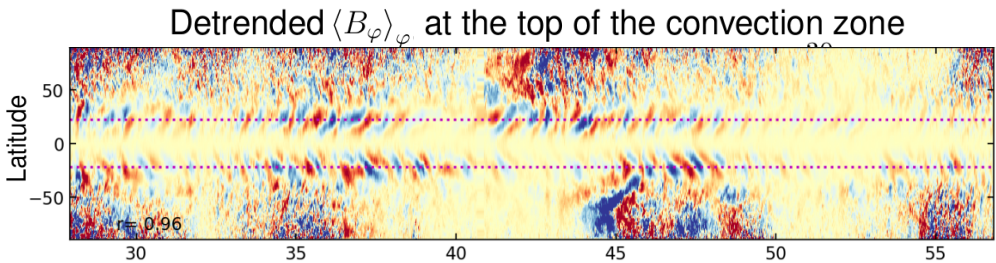
- short cycles ( $\sim$ year) ; **quasi-biennial oscillations ?**
- surface dynamo
- Parker-Yoshimura type



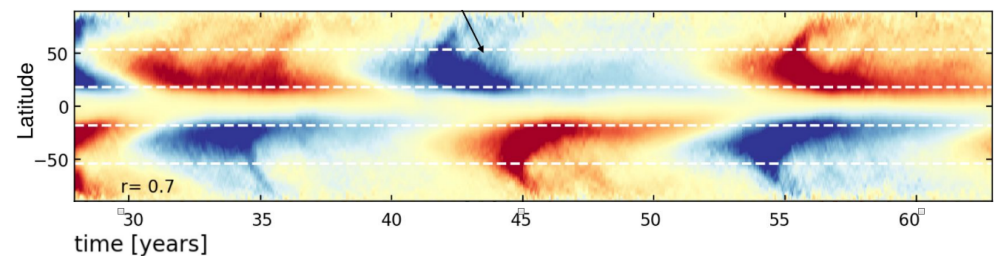
$\langle B_\varphi \rangle_\varphi$  at the base of the convection zone

- long cycles (**decadal solar-like**)
- deeper dynamo
- non-linear retroaction

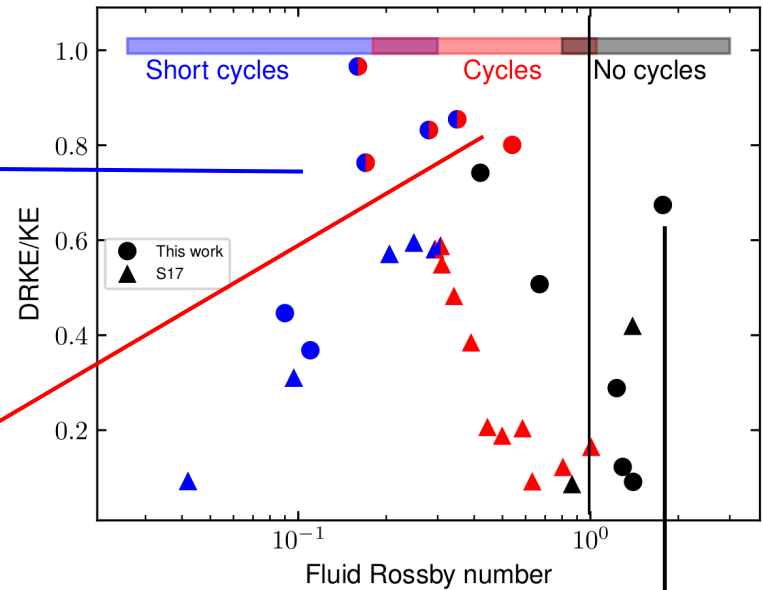
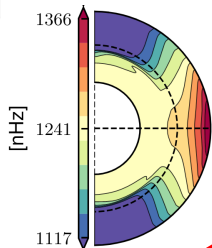




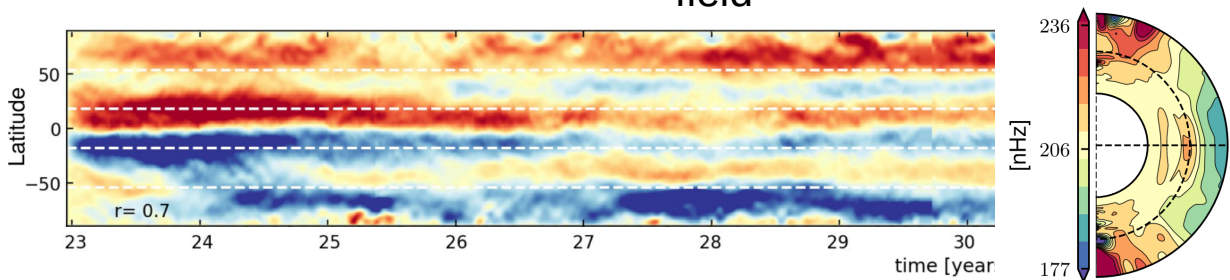
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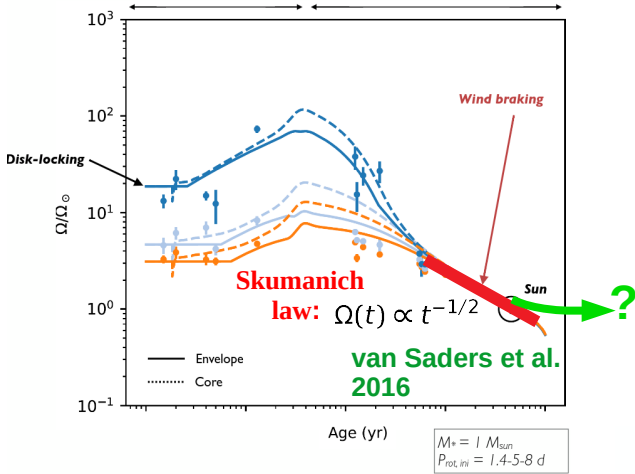


- Stationary dynamo
- Strong deep toroidal field

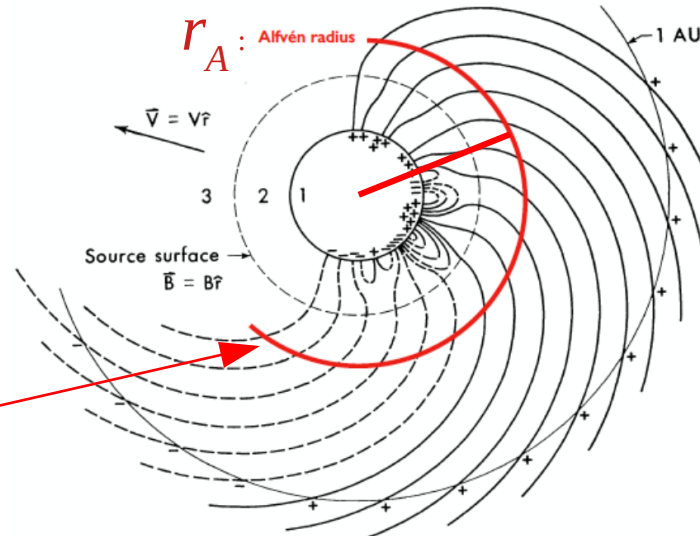


# Weakened Magnetic Wind Breaking ?

Ahuir et al. (2021)



Schatten et al. (1969)



Angular momentum loss

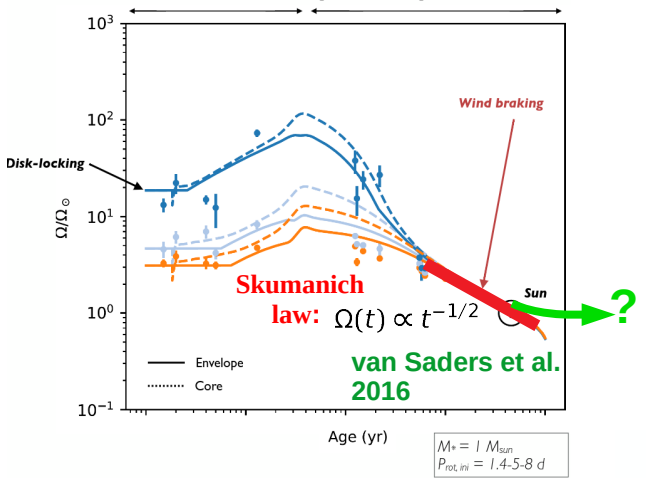
$$\dot{J} \propto \dot{M} \Omega_* \langle r_A \rangle^2$$



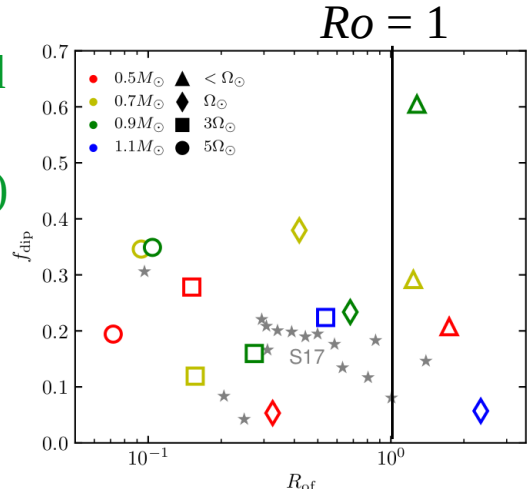


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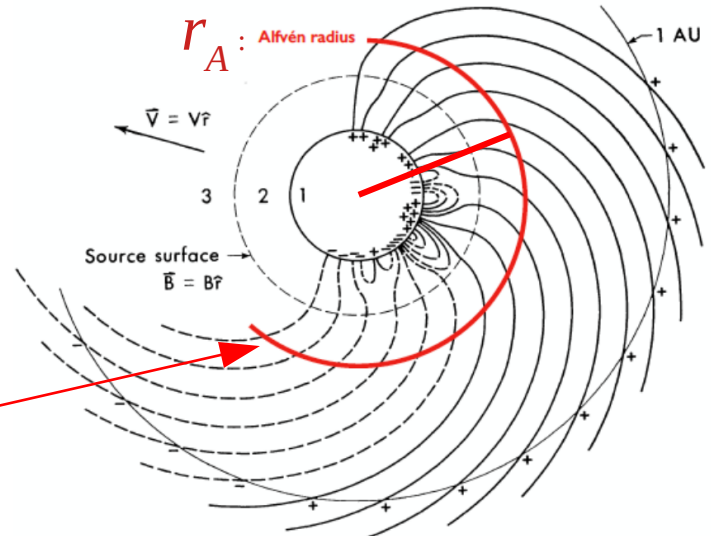
Ahuir et al. (2021)



No sign of weakened large-scale field (i.e. no weakened  $r_A$ )



Schatten et al. (1969)

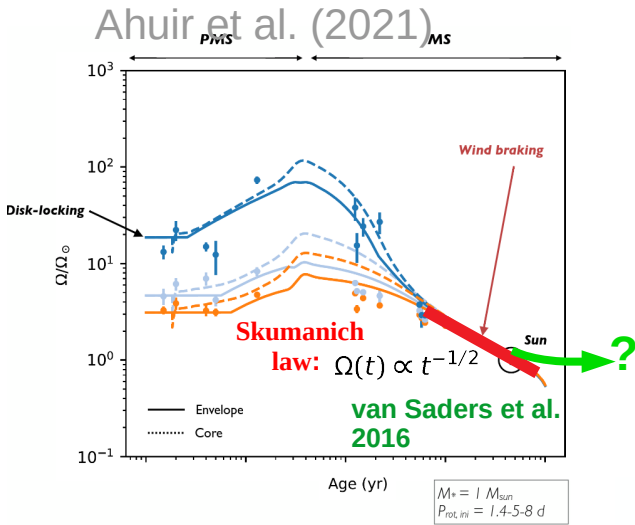


Angular momentum loss

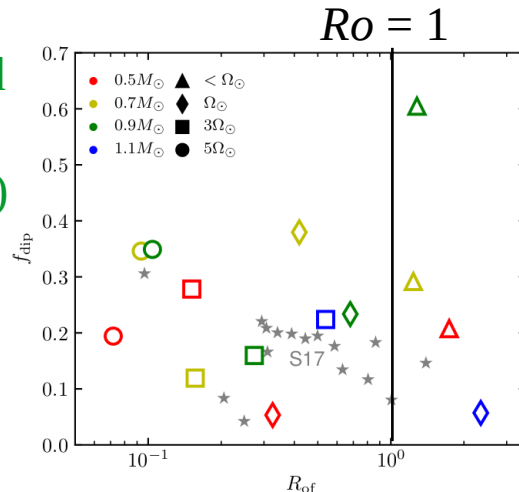
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Brun et al. (2021)

# Weakened Magnetic Wind Breaking ?



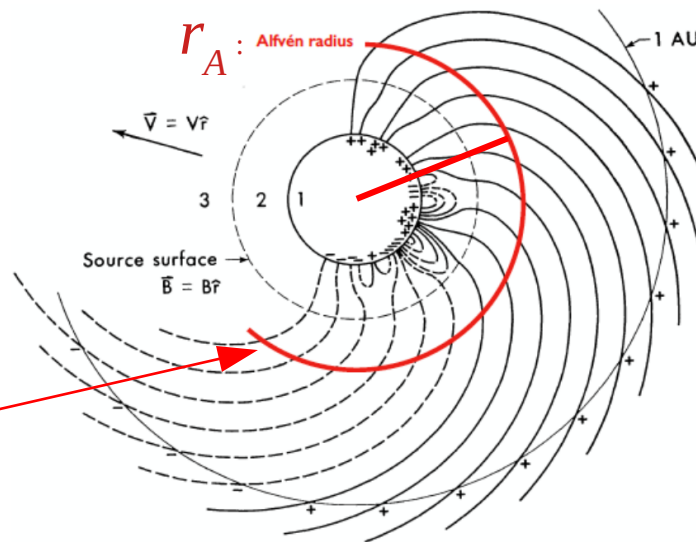
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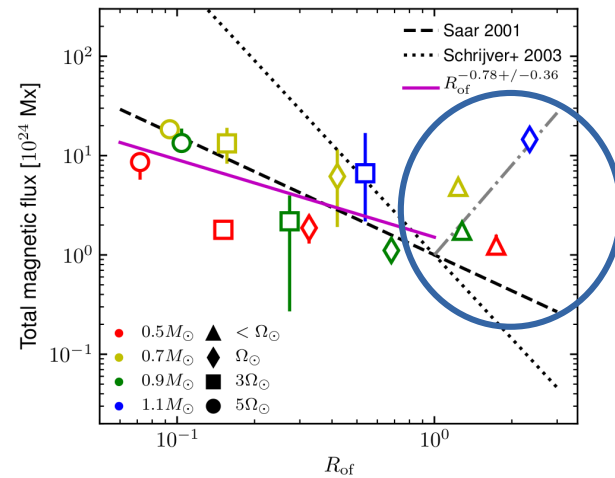
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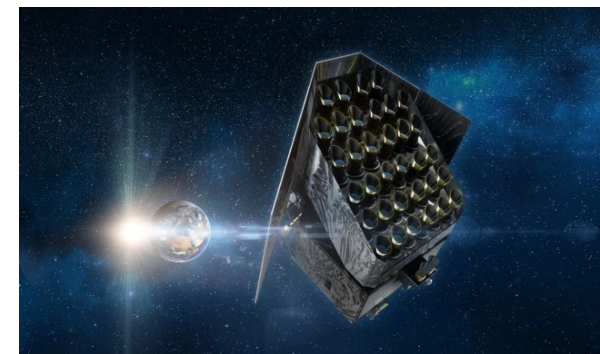
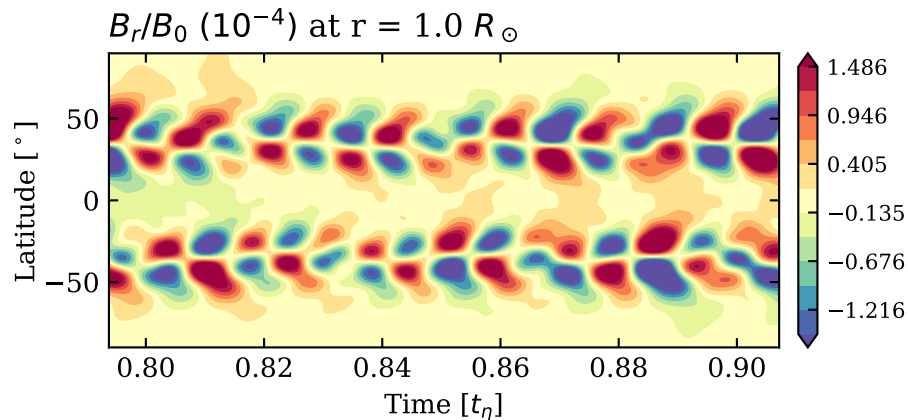
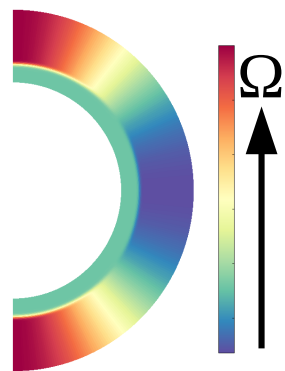
Schatten et al. (1969)



Brun et al. (2021)

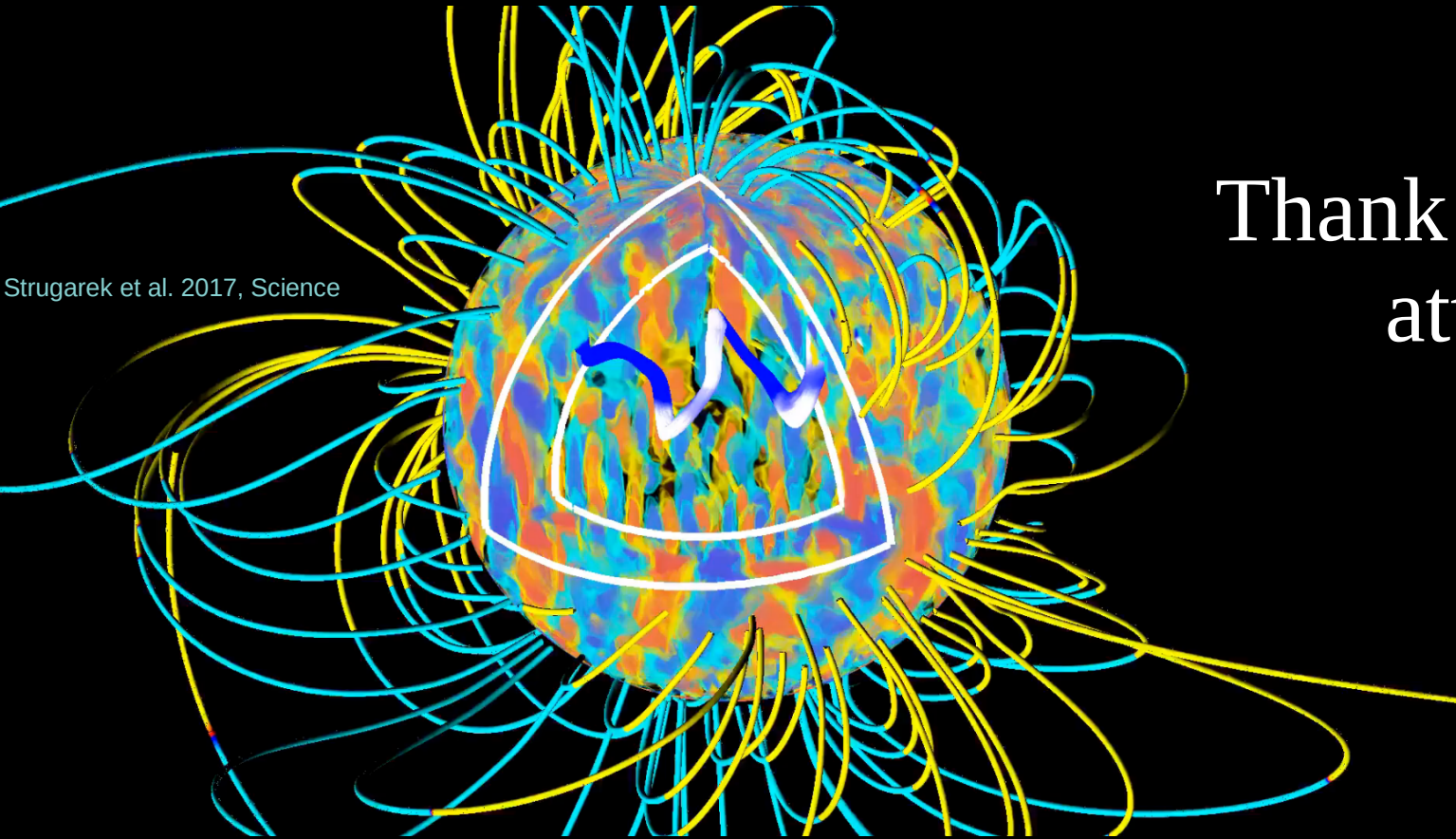


- We present a numerical **multi-D study** with the STELEM and ASH codes to understand the **magnetic field generation** of solar-like stars under **various differential rotation** regimes, and focus on the existence of **magnetic cycles**.
- **Magnetic cycles** can be produced with **anti-solar differential rotation only if the alpha effect is fine tuned** for this purpose in mean-field models.



WP123400  
differential rotation and dynamo

- A **detection of magnetic cycles for such stars** (or lack of thereof) would therefore be a **tremendous constrain** on deciphering **what type of dynamo** is actually acting in the Sun and solar-type stars.



Strugarek et al. 2017, Science

Thank you for your  
attention !