# Large Area MOVPE Growth of Topological Insulator Bi<sub>2</sub>Te<sub>3</sub> Epitaxial Layers on i-Si (111)

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#### Abstract

Bi<sub>2</sub>Te<sub>3</sub> is attracting a renewed interest, due to its topological insulator properties; however, even using advanced physical and chemical deposition techniques, the growth of high-crystal quality layers on substrates allowing its technological employment, such as Si, is very challenging, due to the structural complexity of Bi<sub>2</sub>Te<sub>3</sub>. In this work, we present the optimized large area growth of topological insulator Bi<sub>2</sub>Te<sub>3</sub> epitaxial layers on un-buffered i-Si (111) substrates via Metal-Organic Vapor Phase Epitaxy (MOVPE), which is of crucial importance for future integration into CMOS compatible spintronic devices. We found that the key to maximizing the layer quality requires a balanced control of the reactor pressure (P), growth temperature (T), and growth time (t). Within a proper parameter window, the grown Bi<sub>2</sub>Te<sub>3</sub> thin layers are crystalline, stoichiometric, and highly uniform, also at the local scale. They exhibit a rhombohedral crystalline structure, and they are [001] out-of-plane oriented on the i-Si (111) substrate. Low temperature magnetoresistance measurements revealed clear Weak Antilocalization (WAL) effects, demonstrating that the optimized MOVPE - grown Bi<sub>2</sub>Te<sub>3</sub> is a topological insulator, hence opening further possibilities for its technology-transfer to innovative devices.

KEYWORDS: Bi<sub>2</sub>Te<sub>3</sub>, MOCVD, MOVPE, epitaxial layers, topological insulators

### Introduction

Bismuth telluride (Bi<sub>2</sub>Te<sub>3</sub>) thin layers are being intensively investigated, mainly due to their exceptional performances as thermoelectric materials and perspective applications as topological insulators (TIs) [1–5]. Bi<sub>2</sub>Te<sub>3</sub> is considered to be the best thermoelectric material at ~100 °C [6,7] and it has been confirmed to have TI properties by angle-resolved photoemission spectroscopy [3,8], magneto-transport measurements [9], and scanning tunnelling microscopy analysis [10]. TIs are insulating materials in the bulk, while presenting metallic conducting surface states, which are topologically protected with spin polarization locked to crystal momentum [5]. The unique electronic and optical properties associated with these surface states make them strong candidates to study innovative quantum phenomena and to explore practical applications for quantum computing, spintronics and nano technological devices [11,12]. Since the most interesting properties of Bi<sub>2</sub>Te<sub>3</sub> is to have homogeneous, wafer-scale, large-domain size and defect-free layer growth, with a high surface-to-bulk ratio [8,13,14].

Various deposition techniques, such as sputtering [15,16], evaporation [17,18], pulsed laser deposition [19,20], physical vapor transport [21–23], electrochemical deposition [24], molecular beam epitaxy (MBE) [8,25–30] have been employed to grow continuous Bi<sub>2</sub>Te<sub>3</sub> films on different substrates. The high cost of the complex equipment with operation in ultrahigh vacuum conditions and the low yield of the MBE process can limit its applications in the industry [14]. Compared to other methods, Metal-Organic Chemical Vapour Deposition (MOCVD) (Metal-Organic Vapour Phase Epitaxy, MOVPE in the case of epitaxial deposition) allows for the preparation of highly crystalline, uniform thin layers, with good control of the composition; by using the optimized conditions, it can be potentially suitable for large-scale industrial production. Indeed, MOCVD has been already employed to grow continuous Bi<sub>2</sub>Te<sub>3</sub> layers [1,31–35] but most of the reported studies are limited to substrates like Pyrex [31], SiO<sub>2</sub> [32], Sapphire [33], GaAs [34,35]. However, the quality of this class of layers has not been significantly improved, which could not be compatible with device integration. Most importantly, the direct growth of high quality Bi<sub>2</sub>Te<sub>3</sub> layers on silicon substrates would be a great advantage for their application in microelectronic devices.

Motivated by its potential technological impact, in this work we present a method for growing controlled epitaxial  $Bi_2Te_3$  layers on i-Si (111) by MOVPE on large area substrates. To our knowledge, there are no reports in the literature on a successful growth of this type. Beside the role of the [111] orientation of the substrate, we report how the fine-tuning of the precursor's

partial pressure ratio, growth pressure, temperature, and time are necessary to achieve the required crystalline quality. Our work revealed step by step growth of high crystal quality  $Bi_2Te_3$  layers in the typical rhombohedral crystalline structure having *c* - lattice out-of-plane to the i-Si (111) substrate. To study the electrical conduction properties of the grown layers, the grown  $Bi_2Te_3$  layers were also analyzed by Magnetoresistance (MR) measurements. A clear Weak Antilocalization (WAL) effect was found, thus proving the topological insulating character of the  $Bi_2Te_3$  layers.

### **Experimental**

The employed substrates were i-Si (111) and i- Si (100) oriented wafers. Before deposition, the substrates were treated with HF (5% in deionized H<sub>2</sub>O) for 3 min, thoroughly rinsed with deionized H<sub>2</sub>O, and N<sub>2</sub>-dried. Then, they were rapidly loaded into the MOVPE chamber. The Bi<sub>2</sub>Te<sub>3</sub> layers were grown with an AIXTRON AIX 200/4 MOCVD tool, equipped with an IR-heated 4" rotating graphite susceptor. Electronic grade Trimethylbismuth (BiMe<sub>3</sub>) and Di-isopropyltellurium (DiPTe) were employed as Bi and Te precursors, respectively. The depositions were carried out in the temperature (T) range of 300 - 400 °C for a duration (t) of 30-240 min, at a pressure (P) of 25-150 mbar pressure, with a total flow of 2000-4500 *l* min<sup>-1</sup> and setting the precursors vapor pressures at 4.20 and 7.80 x 10<sup>-3</sup> mbar for BiMe<sub>3</sub> and DiPTe, respectively.

X-ray Diffraction (XRD) patterns were acquired with an HRXRD IS2000 tool, equipped with a Cu K $\alpha$  radiation source, a four-circle goniometer, and a curved 120° position-sensitive detector (Inel CPS-120).

Atomic Force Microscopy (AFM) images were obtained by a Bruker Dimension Edge instrument in non-contact mode, using a sharp silicon AFM probe (TESPA, Bruker) with a typical radius of curvature in the 8-12 nm range. A polynomial background correction was applied to the raw data. Surface roughness was reported as Root Mean Square roughness (RMS roughness) and expressed in nm.

Raman characterization was performed in a z-backscattering geometry by using a Renishaw In-via Spectrometer, equipped with the 1.96 eV (wavelength 633 nm) line of a continuous He-Ne laser (class 3B). The laser beam was focused onto the sample by a  $50 \times (0.75$  numerical aperture) Leica objective. The beam power at the sample was maintained below 1 mW to avoid laser-induced heating on the surface and sample damage.

Scanning Electron Microscopy (SEM) images were acquired with a ZEISS Supra 40 field emission scanning electron microscope, at an acceleration voltage of 15 kV. Cross-section images were collected at a tilting angle of 10°.

Transmission Electron Microscopy (TEM) analyses were performed using a JEOL 2200FS microscope working at 200 kV, equipped with an Energy Dispersive X-ray Spectrometer (EDX), two high-angle annular dark-field (HAADF) detectors, and an in-column energy (Omega) filter. The cross-sectional view samples were prepared using standard mechanochemical procedures and finished by ion beam thinning.

Magnetoresistance (MR) measurements were carried out at 5.5 K using a four-point probe in the van der Pauw configuration. Magnetic fields up to 0.8 T were applied both perpendicularly and parallelly to the sample's plane (and current). The measurements were conducted at a constant applied current of 100  $\mu$ A, by measuring the longitudinal voltage changes during the magnetic field scan. In particular, the measurements of the sheet resistance (Rs) as a function of the magnetic field (B) were recorded on (1x1) cm<sup>2</sup> substrates, cleaved from the grown samples without any patterning or capping, to avoid any damage to the layers. The values obtained for the sheet resistances is then reversed, referred to its value at zero field and expressed in unit of e<sup>2</sup>/h. The signal obtained from this operation is defined as  $\Delta\sigma$  and it was used to visualize and compare directly the analysed data.

#### **Results & Discussions**

A systematic preliminary growth study was performed as a function of the main deposition parameters (temperature, pressure and time), after identifying the proper super-saturation conditions leading to material deposition (see Supporting Information Figure S1, S2, S3). The effect of an i-Si (100) oriented substrate was also investigated, as described in the Supporting Information (see Figure S4). The optimized growth conditions of Bi<sub>2</sub>Te<sub>3</sub> layers on i-Si (111) were found to be a temperature of 350 °C, a pressure of 75 mbar, and for 180 min growth time. The sample morphology is reported in Figure 1 (a, b), where the plan and cross-sectional SEM micrographs are shown. From the cross-section SEM images, it can be seen that the Bi<sub>2</sub>Te<sub>3</sub> layers have a uniform thickness of about 90 nm and continuous, flat surfaces.

The surface morphology of the optimized  $Bi_2Te_3$  thin layers was characterized by AFM. A high-resolution topographic image (500 nm x 500 nm) of the layer and a representative height line profile are reported in Figure 1 (c) and (d), respectively. The topography shows a terrace-step surface morphology uniformly distributed over the whole sample area. Terraces extend up

to the micron range in the lateral dimension (see Supporting Information Figure S5), while the line profile shows step height of the order of 1 nm, compatible with the thickness of a single quintuple layer (QL) of the Bi<sub>2</sub>Te<sub>3</sub> crystal structure. Indeed, the binary Bi<sub>2</sub>Te<sub>3</sub> compound presents a rhombohedral layered structure with space group R3m (SG #166) and lattice parameters a = 0.4395 nm and c = 3.044 nm (ICSD 74348). Its unit cell stacks three building blocks along the [001] direction, separated by van der Waals gaps, and each one rotated by 120° with respect to the one below. Each block (QL), is composed by five atomic planes in the sequence -Te-Bi-Te-Bi-Te-, as reported in Supporting Information Figure S6. There is a strong affinity to complete each QL before a new QL nucleates on the surface, which results in a surface structure that features QL steps [36,37]. Additionally, the RMS roughness extracted from the AFM analysis (~ 0.5 nm) revealed that the optimized growth leads to an flat surface morphology, comparable with those reported for Bi<sub>2</sub>Te<sub>3</sub> layers grown by MBE [26,27,38–40].

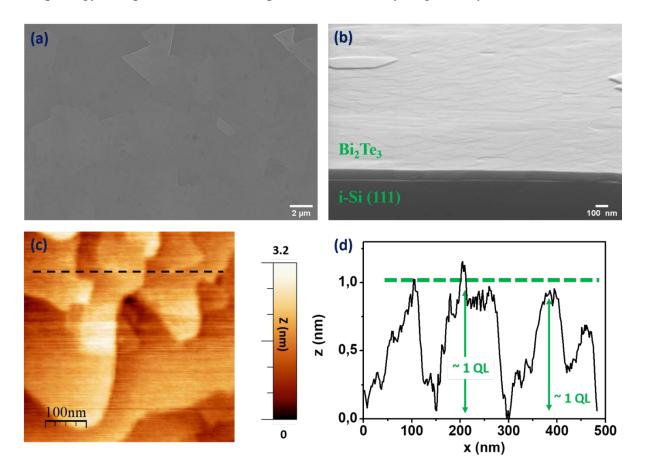


Figure 1. (a) Plan view SEM image, (b) Tilted cross-sectional SEM image, (c) AFM image, and (d) its corresponding height line profile for an optimized  $Bi_2Te_3$  layer, grown on i-Si (111) at T = 350 °C, P = 75 mbar, t = 180 min.

The grown materials were also studied by micro-Raman spectroscopy, as shown in Figure 2(a). The Raman spectrum exhibits two main vibrational peaks located at ~ 102 and ~133 cm<sup>-1</sup>. This finding is compelling evidence of the high quality of the grown Bi<sub>2</sub>Te<sub>3</sub> since the observed peaks are well-consistent with the previously reported in-plane  $E_g^2$  and out-of-plane  $A_{1g}^2$  phonon modes of the Bi<sub>2</sub>Te<sub>3</sub> single crystal, respectively [41].

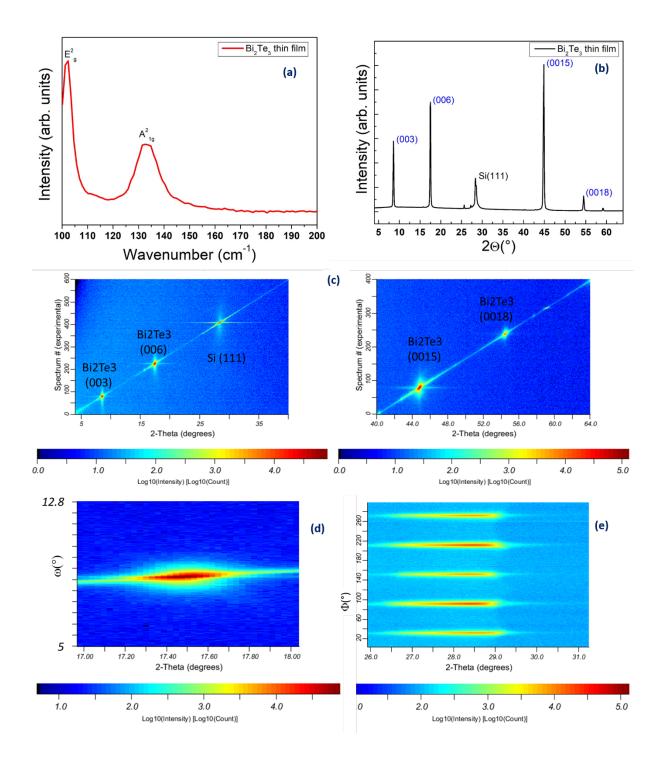


Figure 2. (a) Raman spectrum of Bi<sub>2</sub>Te<sub>3</sub> layer, (b) Bragg-Brentano pattern, (c) and (d) Bragg-Brentano maps, (e) phi-scan around the [015] direction of the [001] out-of-plane oriented crystallites.

The Bi<sub>2</sub>Te<sub>3</sub> layers were measured by XRD in Bragg Brentano configuration. In Figure 2(b) and 2(c), the XRD pattern exhibited intense reflections at  $2\theta = 8.46^{\circ}$ ,  $17.3^{\circ}$ ,  $44.8^{\circ}$ , and  $54.4^{\circ}$  corresponding to the 003, 006, and 0015 and 0018 reflections, attributed to the rhombohedral crystalline structure of Bi<sub>2</sub>Te<sub>3</sub> in the R3m space group, (ICSD 74348). The XRD measurements confirmed that the obtained Bi<sub>2</sub>Te<sub>3</sub> layers are relaxed, single crystalline, with the *c* lattice parameter equal to  $30.4\pm0.1$  Å and [001] out-of-plane oriented on the i- Si (111) substrate.

The XRD patterns collected in the Bragg-Brentano geometry (Figure 2(c)) revealed the scattering of the reflections' intensity across  $\omega$ . The full width at half maximum (FWHM) of the  $\omega$  scan, collected for the [006] reflection (see Figure 2d and Supporting Information, Figure S7) showed a very small value of 0.26°, indicating a mosaicity as low as the one that has been obtained by the MOCVD of Bi<sub>2</sub>Te<sub>3</sub> on GaAs [42], confirming the excellent quality of the layer. The in-plane orientation of the [001] out of plane-oriented crystallites was investigated by complete  $\Phi$  scans around the [015] asymmetric reflection of the Bi<sub>2</sub>Te<sub>3</sub> structure. The  $\Phi$  scan is shown in Figure 2(e). Only six spots were observed, confirming the in-plane alignment of the crystallites, with a six-fold geometry.

High-Resolution TEM (HRTEM) analysis was performed to investigate the material properties at the sub-nanometric scale. A representative cross-sectional HRTEM image of the Bi<sub>2</sub>Te<sub>3</sub> thin layer is reported in Figure 3(a). The overview image reveals the high morphological and crystalline quality of the layer with the typical van der Waals gaps separating the QLs along the [001] direction [27]. The Fast Fourier Transform of the image, taken on the Bi<sub>2</sub>Te<sub>3</sub> region only, shown as inset, also highlights the single-crystal nature of the layer and the high degree of order along the c axis. Figure 3(b) reports the high magnification of the interface in the boxed region in Figure 3(a), while Figure 3(c) shows the superposition of the FFTs of Bi<sub>2</sub>Te<sub>3</sub> and Si across the interface. The relative orientations of the FFT patterns demonstrate the epitaxial relationship Bi<sub>2</sub>Te<sub>3</sub> [001]||Si[111] between the layer and the substrate. The out-of-plane and inplane lattice constants of the Bi<sub>2</sub>Te<sub>3</sub> layer, calculated from the FFT, were found to be c = 30.5Å and a = 4.4 Å, respectively, in accordance with the bulk lattice parameters of Bi<sub>2</sub>Te<sub>3</sub> (ICSD 74348) and the values calculated by the XRD analysis. The same value of the unit cell along the [001] axis was found by measuring the distance between the van der Waals gaps in highresolution HAADF (Supporting Information, Figure S8). The QL spacing indicated for  $Bi_2Te_3$  is 1.00 nm. The large overview image reveals the high crystalline quality of the thin layer.

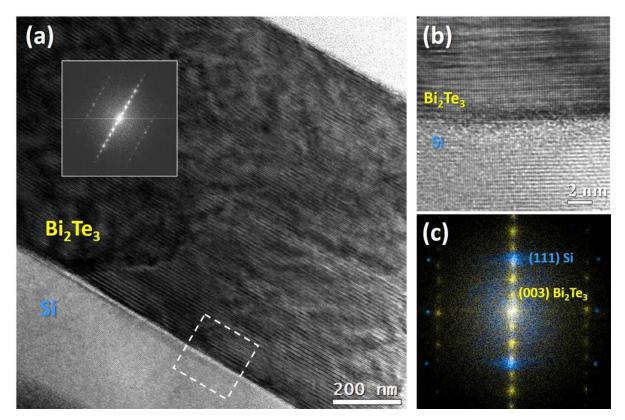


Figure 3. (a) HRTEM image of a Bi<sub>2</sub>Te<sub>3</sub> layer on Si substrate with the corresponding FFT in the inset;
(b) magnified HRTEM image of the interface in the boxed region in (a); (c) superposition of the FFT patterns of [-1-20] zone axis of Bi<sub>2</sub>Te<sub>3</sub> (yellow) and [11-2] zone axis of Si (blue).

The layers were found to be stoichiometric also at the local scale by performing TEM-EDX measurements (see Supporting Information Figure S9 and Table S1). In Figure S8, a typical image of a sample cross-section, taken in the HAADF (Z contrast) mode, is shown. The picture in Figure S9(a) shows a very good thickness uniformity, observed all over the investigated sample. Figure S9(b) and S9 (c), represent the elemental mapping of Te and Bi, obtained with their X-ray L lines, respectively. The two elements appeared to be homogeneously distributed and the stoichiometry preserved everywhere in the sample.

In the literature, both theoretical and experimental results have demonstrated that stoichiometric Bi<sub>2</sub>Te<sub>3</sub> is also a TI with time-reversal-symmetry protected surface states that are present in its bulk insulating gap [3,43,44]. Figure 4 shows the comparison between the magneto conductance ( $\Delta\sigma$ ) obtained when the field is applied perpendicularly to the plane of the sample ( $\theta = 90^{\circ}$ ) and parallelly to it ( $\theta = 0^{\circ}$ ). The  $\theta = 90^{\circ}$  condition shows a cusp-like shape close to 0 T that overlaps with the ordinary Lorentzian B<sup>2</sup> term. The comparison between

the  $\theta = 90^{\circ}$  and  $0^{\circ}$  conditions points to an interpretation in terms of a 2D-type of conduction mechanism emerging at low temperature [45,46]. In fact, as clearly shown in the inset of Figure 4, by following the subtraction of the Lorentzian B<sup>2</sup> term from the MC at  $\theta = 90^{\circ}$ , the typical cusp due to WAL [47] is now clear-cut. This shape of the  $\Delta\sigma$  is accepted as a marker of a 2Dtype of conduction promoted by the contribution of topological surface states to the transport mechanism [48]. Upon comparing with the published literature [46], we observed that the amplitude of the  $\Delta\sigma$  presented for perpendicular field is similar to ours results at the same applied field, expressed in unit of  $e^2/h$ .

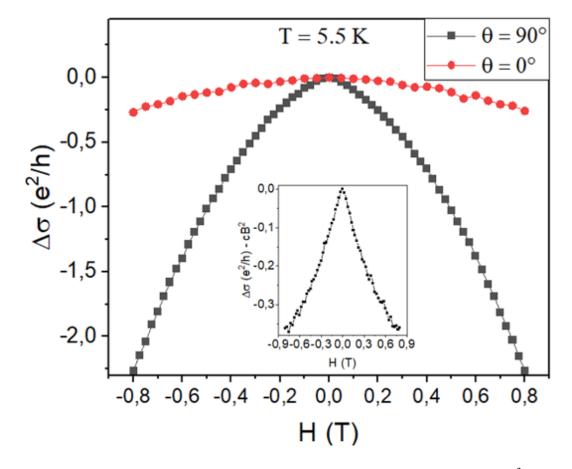


Figure 4: Magneto conductance referred to its zero field value and expressed in unit of e<sup>2</sup>/h. The data are collected at 5.5 K on a square macroscopic sample of Bi<sub>2</sub>Te<sub>3</sub>.

## Conclusion

We obtained a large area epitaxy of  $Bi_2Te_3$  thin layers directly on i-Si (111) substrates by an MOVPE process and reported the control and optimization of growth conditions, followed by the assessment of the high crystal quality, flat morphology, composition uniformity, as well as the topological insulator character of the electrical transport. The developed MOVPE process provides a route to enable the large-area  $Bi_2Te_3$  technology-transfer and to exploit its topological conduction behaviour in future spintronic devices.

#### Acknowledgement

The authors acknowledge Dr. Enzo Rotunno of CNR-Nano, Modena (Italy) for the STEM image simulation.

#### Funding

This project has received funding from the European Union's Horizon 2020 project SKYTOP "Skyrmion - Topological Insulator and Weyl Semimetal Technology" under grant agreement No. 824123.

Notes The authors declare no conflict of interest.

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