

# Supernova Neutrino Light Curves from Cooling of High/Low-Mass Proto-Neutron Stars

Ken'ichiro Nakazato\* (Kyushu Univ.)

for the nuLC Collaboration †

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We present the model spectra of neutrinos emitted from proto-neutron star (PNS) cooling. So as to obtain the time evolution of neutrino spectra, quasi-static evolutionary calculations of PNS cooling are performed. The masses of the PNS models are chosen to reflect recent observations of high-mass and low-mass pulsars in binary systems. For each, we adopt two initial conditions with high- and low-entropy profiles. The details of these models are described in our paper [1]. These data are open for use in any scientific research, provided that our paper is referenced in your publication.

The names of data files are as below:

- `spectobHmaxS.data`:  $M_b = 2.35M_\odot$  ( $M_{\text{NS,g}} = 2.05M_\odot$ ), High initial entropy
- `spectobHminS.data`:  $M_b = 1.29M_\odot$  ( $M_{\text{NS,g}} = 1.20M_\odot$ ), High initial entropy
- `spectobLmaxS.data`:  $M_b = 2.35M_\odot$  ( $M_{\text{NS,g}} = 2.05M_\odot$ ), Low initial entropy
- `spectobLminS.data`:  $M_b = 1.29M_\odot$  ( $M_{\text{NS,g}} = 1.20M_\odot$ ), Low initial entropy

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\*e-mail: nakazato@artsci.kyushu-u.ac.jp

†nuLC stands for *neutrino-light curve*, which is neutrino analogue of the light curves of supernovae. The nuLC collaboration began in 2017 aiming to build a methodology to derive supernova and neutron star physics out of supernova neutrinos. Through the combined efforts of theoretical and experimental researchers, the collaboration is constructing a real data analysis framework for supernova neutrinos that will be observed in the near future.

where  $M_b$  is the baryon mass of PNS and  $M_{\text{NS,g}}$  is the gravitational mass of the remnant neutron star. The format of these data is the same with that of Supernova Neutrino Database [2], except that the number of energy bin is increased from 20 to 25. The data are arranged as follows:

$$\begin{array}{cccccccc}
t_0 & & & & & & & \\
E_0 & E_1 & \frac{\Delta N_{1,\nu_e}(t_0)}{\Delta E_1} & \frac{\Delta N_{1,\bar{\nu}_e}(t_0)}{\Delta E_1} & \frac{\Delta N_{1,\nu_x}(t_0)}{\Delta E_1} & \frac{\Delta L_{1,\nu_e}(t_0)}{\Delta E_1} & \frac{\Delta L_{1,\bar{\nu}_e}(t_0)}{\Delta E_1} & \frac{\Delta L_{1,\nu_x}(t_0)}{\Delta E_1} \\
E_1 & E_2 & \frac{\Delta N_{2,\nu_e}(t_0)}{\Delta E_2} & \frac{\Delta N_{2,\bar{\nu}_e}(t_0)}{\Delta E_2} & \frac{\Delta N_{2,\nu_x}(t_0)}{\Delta E_2} & \frac{\Delta L_{2,\nu_e}(t_0)}{\Delta E_2} & \frac{\Delta L_{2,\bar{\nu}_e}(t_0)}{\Delta E_2} & \frac{\Delta L_{2,\nu_x}(t_0)}{\Delta E_2} \\
\dots & & & & & & & \\
E_{24} & E_{25} & \frac{\Delta N_{25,\nu_e}(t_0)}{\Delta E_{25}} & \frac{\Delta N_{25,\bar{\nu}_e}(t_0)}{\Delta E_{25}} & \frac{\Delta N_{25,\nu_x}(t_0)}{\Delta E_{25}} & \frac{\Delta L_{25,\nu_e}(t_0)}{\Delta E_{25}} & \frac{\Delta L_{25,\bar{\nu}_e}(t_0)}{\Delta E_{25}} & \frac{\Delta L_{25,\nu_x}(t_0)}{\Delta E_{25}} \\
t_1 & & & & & & & \\
E_0 & E_1 & \frac{\Delta N_{1,\nu_e}(t_1)}{\Delta E_1} & \frac{\Delta N_{1,\bar{\nu}_e}(t_1)}{\Delta E_1} & \frac{\Delta N_{1,\nu_x}(t_1)}{\Delta E_1} & \frac{\Delta L_{1,\nu_e}(t_1)}{\Delta E_1} & \frac{\Delta L_{1,\bar{\nu}_e}(t_1)}{\Delta E_1} & \frac{\Delta L_{1,\nu_x}(t_1)}{\Delta E_1} \\
\dots & & & & & & & 
\end{array}$$

where  $t_n$  [s] is a time measured from the onset of the computation and  $E_k$  [MeV] is a neutrino energy. Note that,  $E_k$  is defined on the interface between  $k$ -th and  $(k+1)$ -th energy bins while  $E_0 = 0$  MeV. For  $k$ -th energy bin,  $\frac{\Delta N_{k,\nu_i}(t_n)}{\Delta E_k}$  [/s/MeV] and  $\frac{\Delta L_{k,\nu_i}(t_n)}{\Delta E_k}$  [erg/s/MeV] are differential neutrino number flux and differential neutrino luminosity, respectively, where  $\nu_x = (\nu_\mu + \bar{\nu}_\mu + \nu_\tau + \bar{\nu}_\tau)/4$ . Thus, the luminosity of  $\bar{\nu}_e$  [erg/s] is given by

$$L_{\bar{\nu}_e}(t_n) = \sum_{k=1}^{25} (E_k - E_{k-1}) \times \frac{\Delta L_{k,\bar{\nu}_e}(t_n)}{\Delta E_k}, \quad (1)$$

and the number luminosity of  $\bar{\nu}_e$  [/s] is given by

$$N_{\bar{\nu}_e}(t_n) = \sum_{k=1}^{25} (E_k - E_{k-1}) \times \frac{\Delta N_{k,\bar{\nu}_e}(t_n)}{\Delta E_k}. \quad (2)$$

Therefore the mean energy of emitted  $\bar{\nu}_e$  at the time  $t_n$  is given by

$$\langle E_{\bar{\nu}_e}(t_n) \rangle = \frac{L_{\bar{\nu}_e}(t_n)}{N_{\bar{\nu}_e}(t_n)} \times \frac{\text{MeV}}{1.6022 \times 10^{-6} \text{ erg}}. \quad (3)$$

If you find some strange problem, please contact us. We would appreciate it very much if you could give us comments or suggestions on our data. The correspondence address is

- Ken'ichiro Nakazato  
Faculty of Arts and Science, Kyushu University  
744 Motooka, Nishi-ku, Fukuoka 819-0395, Japan  
E-mail: nakazato@artsci.kyushu-u.ac.jp

## References

- [1] Y. Suwa, K. Sumiyoshi, K. Nakazato, Y. Takahira, Y. Koshio, M. Mori, and R. A. Wendell, *Astrophys. J.* **881** (2019) 139, arXiv:1904.09996 [astro-ph.HE]
- [2] K. Nakazato, K. Sumiyoshi, H. Suzuki, T. Totani, H. Umeda, and S. Yamada, *Astrophys. J. Supp.* **205** (2013) 2, arXiv:1210.6841 [astro-ph.HE]  
<http://asphwww.ph.noda.tus.ac.jp/snn/>