

# PROBLEME N°10

# PARTICLE DETECTOR FOR DUMMIES

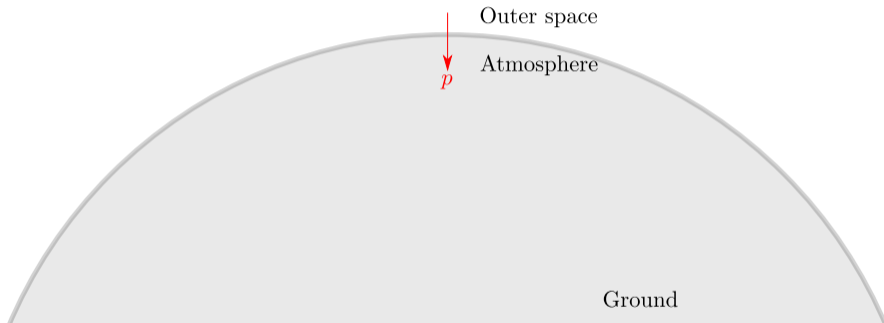
Team Ecole polytechnique

# The problem

*Build a simple device that can **detect cosmic ray particles**.  
Characterize the **particle identification** capabilities of your device.  
Try to test your device in **different conditions** and also try to obtain  
the **energy spectrum of the cosmic ray particles**.*

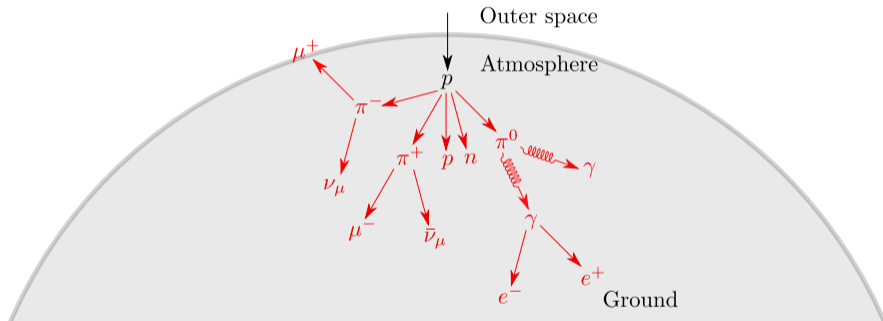
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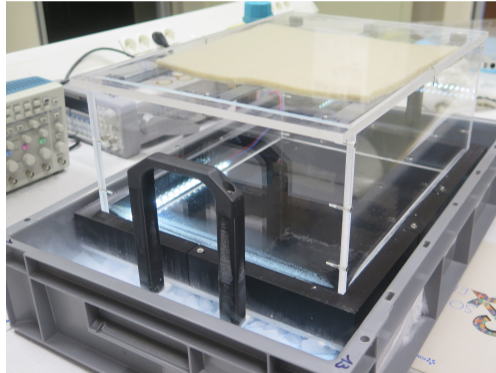
**Secondary cosmic rays : 75 % muons**

Muons of 4 GeV have a Lorentz coefficient  $\gamma \simeq 40$ , allowing them to reach ground before decaying thanks to special relativity.

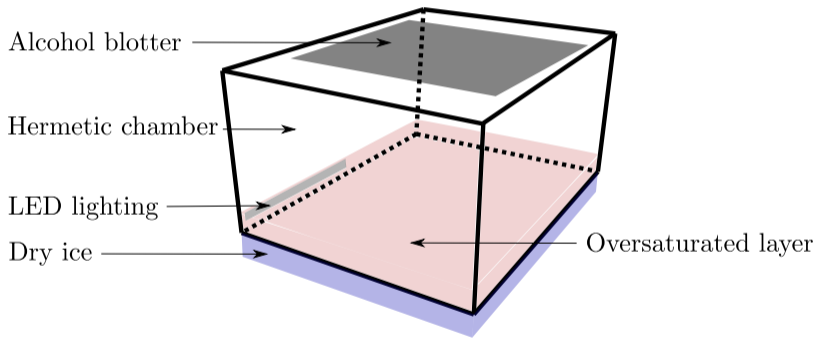
# What are cosmic rays ?

	Fermions			Bosons
Quarks	$u$ up	$c$ charm	$t$ top	$g$ gluon
	$d$ down	$s$ strange	$b$ bottom	$\gamma$ photon
Leptons	$e$ electron	$\mu$ <b>muon</b>	$\tau$ tau	$Z$ Z boson
	$\nu_e$ electron neutrino	$\nu_\mu$ muon neutrino	$\nu_\tau$ tau neutrino	$W$ W boson
	<b>1<sup>st</sup> generation</b>	<b>2<sup>nd</sup> generation</b>	<b>3<sup>rd</sup> generation</b>	$H$ Higgs boson

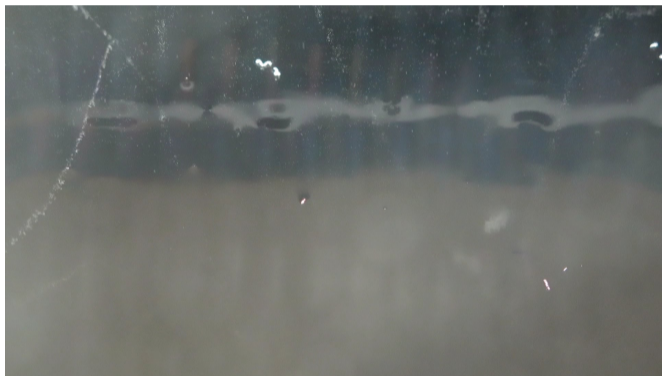
# First device : a cloud chamber



# How does it works



# Observations





# Particle identification capabilities

Observable particles :  $e^-$ ,  $e^+$ ,  $\mu^-$ ,  $\mu^+$ ,  $p^+$ ,  $He^{2+}$ ,  $K$ ,  $\Lambda$ ,  $\Xi$ , ...

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For some charged particles with energies between 1 and 100 MeV, the energy loss is given by Bethe formula :

$$\frac{dE}{dx} = -\frac{4\pi}{m_e c^2} \frac{nq^2}{\beta^2} \left( \frac{1}{4\pi\epsilon_0} \right)^2 \left[ \ln \left( \frac{2m_e c^2 \beta^2}{I(1-\beta^2)} - \beta^2 \right) \right]$$

With :  $m_e$  the mass of electrons,  $c$  the speed of light,  $n$  the electron density of the material,  $q$  the charge of the particle,  $\epsilon_0$  the vacuum permittivity,  $\beta$  the boost of the particle and  $I$  the mean extraction potential of the material.

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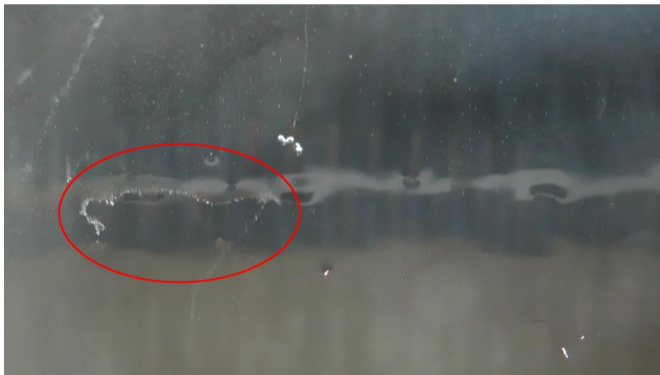
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# Alpha particle and delta ray



# Low energy electron



# Muon or high energy electron



# Gamma ray?

Pair production :  $\gamma + n \longrightarrow e^+ + e^- + n$





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We could get the **momentum**  $p$  of the particle, then from :

$$E = \sqrt{(pc)^2 + (mc^2)^2}$$

With  $E$  the energy of the particle,  $c$  the speed of light and  $m$  the mass of the particle

We could get it's **energy**  $E$ .

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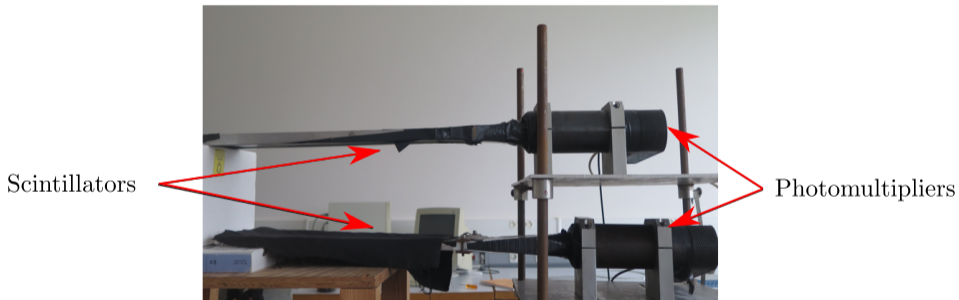
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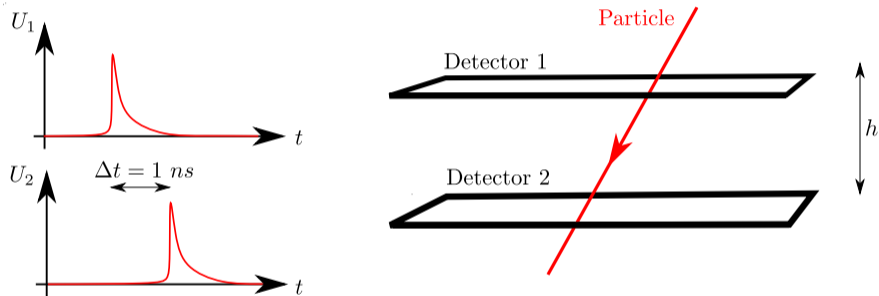
- To deflect a 1 MeV muon with a radius of 10 cm, we would need a **uniform field of 4 T** on the surface of the chamber
- We need to **assume the charge and mass** of the particle to get it's energy

## Second device : scintillation detector



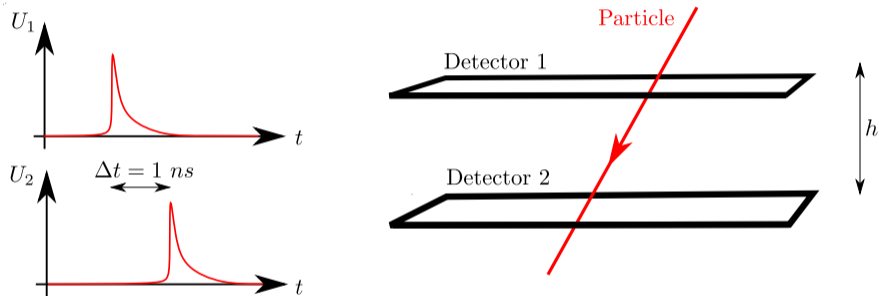
[1.]C.Lagoute; BUP, Réalisation d'un détecteur de muons : une approche de physique du XXème siècle au lycée, 2009

# Coincidence detection



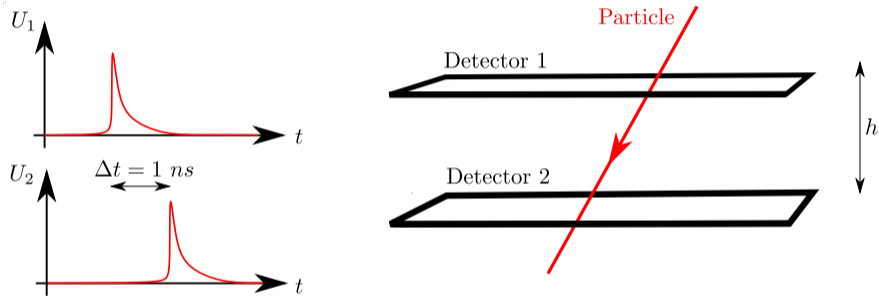


# Coincidence detection



We detect particles going downward with at least 2 % the speed of light

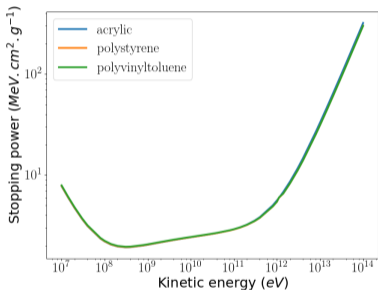
# Coincidence detection



We detect **particles going downward with at least 2 % the speed of light** and we can get the **energy spectrum** of those particles

# Theory of energy deposition in matter

Mean energy deposition depends on total energy [4] :



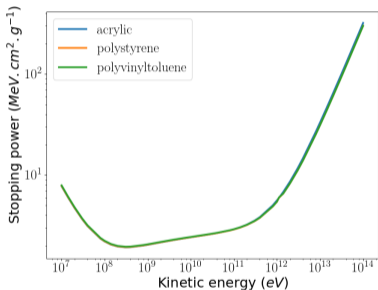
Muon stopping power

[4.]Groom, Mokhov, Striganov; A.D.N.D.T **Muon stopping power and range tables** 10 MeV – 100 TeV, 2001

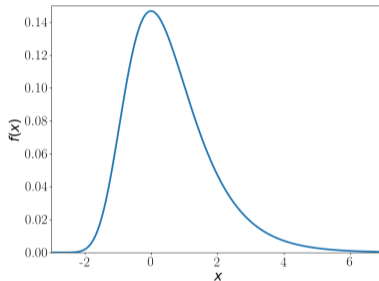
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For a given stopping power, the true deposited energy is statistic [3] :



Muon stopping power



Landau distribution

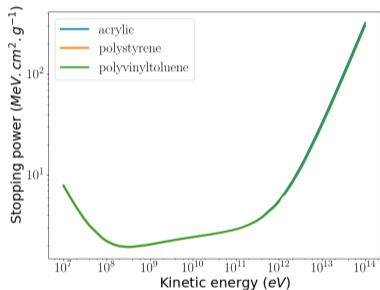
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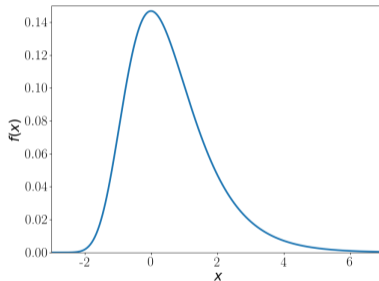
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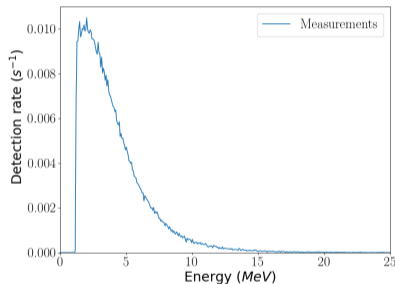


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# Energy spectrum

With three days of measurement and 150964 detections

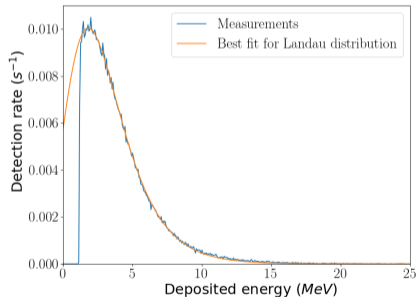


Our energy spectrum

The cut for low energies is due to a threshold reducing electronic noise.

# Interpretation of our spectrum

The spectrum we measure is only the Landau distribution :

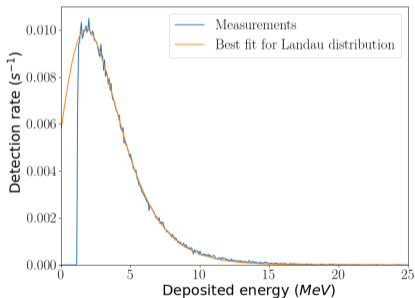


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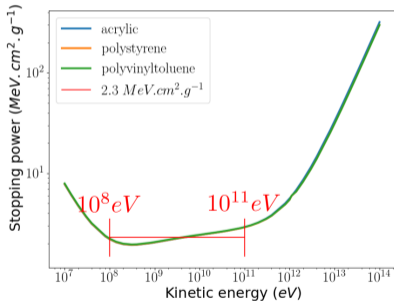
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The stopping power is the same for all the particles :



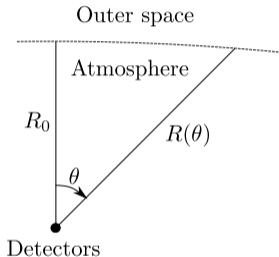
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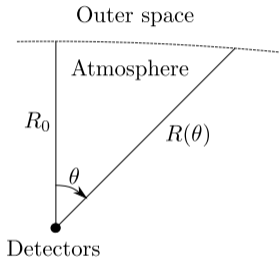
Muon stopping power [4]



**Model** : grazing incidence rays take more time to reach the surface and have more chance to decay.



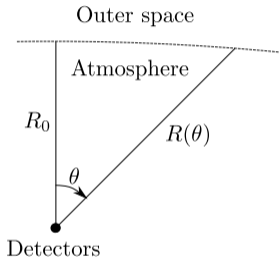
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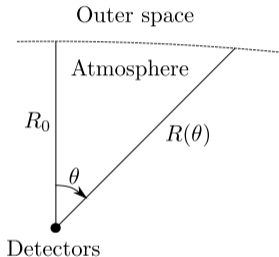
- $R_0$  : thickness of atmosphere
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We have :

- $R(\theta) = \frac{R_0}{\cos(\theta)}$
- $P(t_{decay} > t) = P_0 \exp(-\frac{t}{\gamma\tau_0})$

With  $\gamma$  the Lorentz coefficient and  $\tau_0$  the proper lifetime of muons

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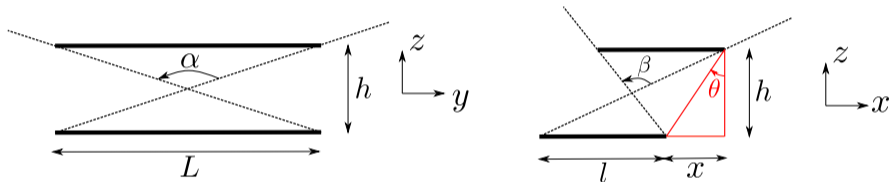
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We get the flux of muons  $\phi_{th}(\theta)$  as :

$$\phi_{th}(\theta) = \phi_0 \exp\left(\frac{-\lambda}{\cos(\theta)}\right), \text{ where } \lambda = \frac{m_\mu c R_0}{E_0 \tau_0} \text{ and } \phi_0 \text{ is a parameter}$$

# Measuring the flux anisotropy

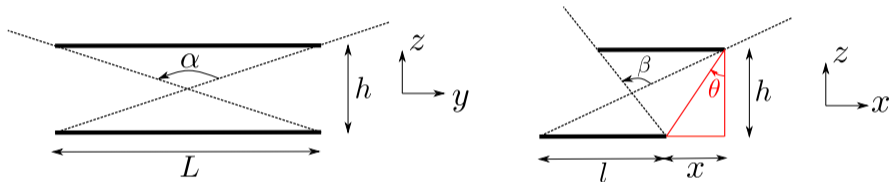
We simply slide the one detector relative to the other :



Geometry of the scintillators

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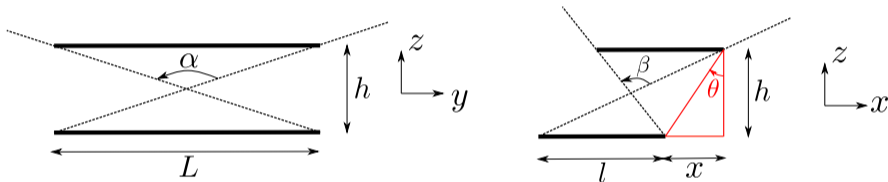


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Hypothesis : for a given angle  $\theta$ , muons only come from **one direction**.

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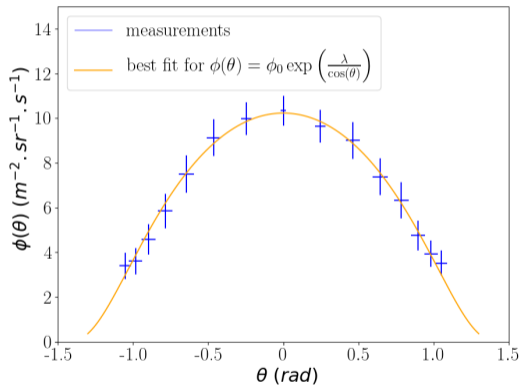


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We then get from the number of detection to the flux with **time and geometric normalization**.

# Result of the measurement



Muon flux anisotropy



# Another measurement of muon's energy

We had :

$$\phi_{th}(\theta) = \phi_0 \exp\left(\frac{-\lambda}{\cos(\theta)}\right), \text{ where : } \lambda = \frac{m_\mu c R_0}{E_0 \tau_0}$$

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- $m_\mu = 2.10^{-28}$  kg the mass of muons
- $c = 3.10^8$  m.s<sup>-1</sup> the speed of light
- $R_0 = 10 - 100$  km atmosphere's thickness
- $\tau_0 = 2.2$   $\mu$ s the lifetime of muons

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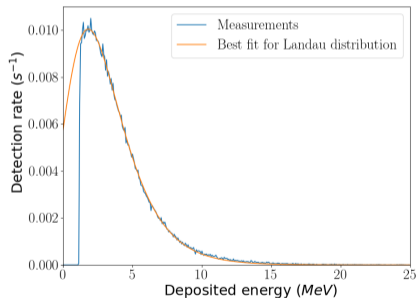
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We can deduce that  $E_0$  is between 1 GeV and 10 GeV, litterature giving a mean ground energy of 4 GeV [5].

[5.]G.Remmen, E.McCreary; Journal of Undergraduate Research in Physics , **Measurement of the speed and energy distribution of cosmic ray muons**, 2012

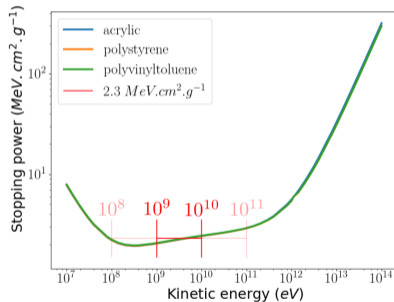
# It is coherent with our spectrum

The spectrum we measure is only the Landau distribution :



Our energy spectrum

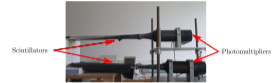
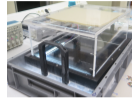
The stopping power is the same for all the particles :



Muon stopping power [4]

# Conclusion

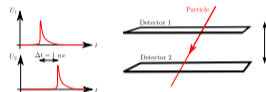
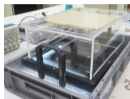
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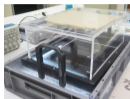
Characterize the particle **identification capabilities** of your device





# Conclusion

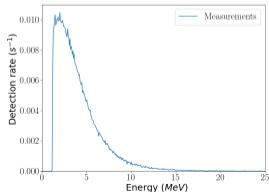
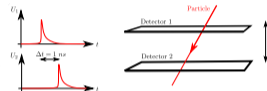
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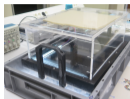
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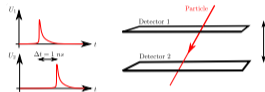
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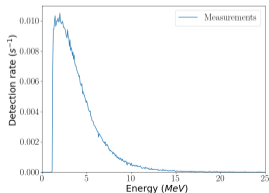
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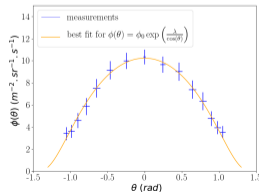
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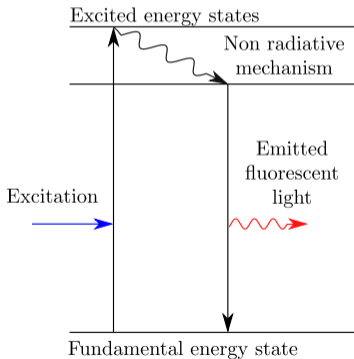


Try to test your device in **different conditions**

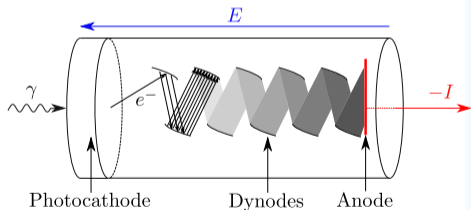


# Scintillator and photomultiplier

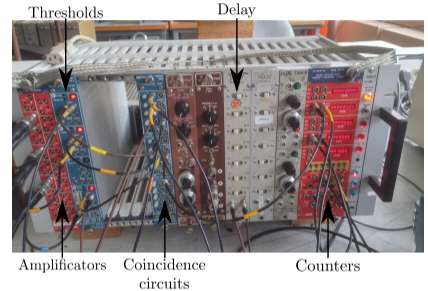
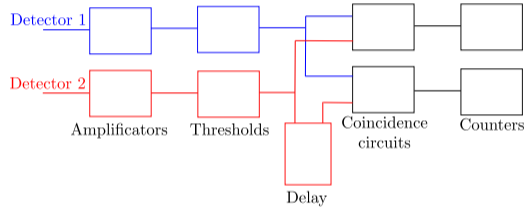
**Scintillator** : fluorescent plastic plate, producing photons from excitations



**Photomultiplier** : association of a photocathode and dynodes, producing a measurable current from single photons.



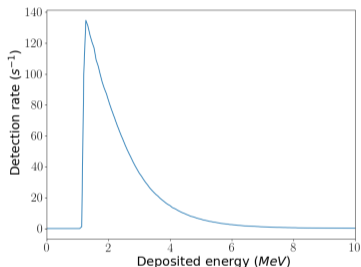
# Coincidence detection



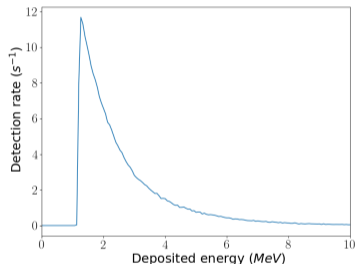
Principe of electronic processing

We detect **particles going downward with at least 2 % the speed of light**

Tentative with radioactive sources :



Cesium 137

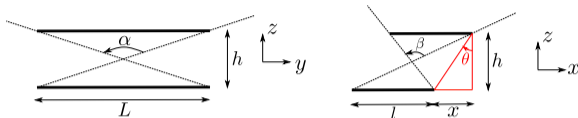


Sodium 22

The expected value of the muon spectrum is the stopping power of the detector [2] :

$$2.3 \text{ MeV.g.cm}^{-1} \times 1 \text{ g.cm}^{-1} \times 1 \text{ cm} = 2.3 \text{ MeV}$$

# Geometric normalization of the flux



With basic geometry :

$$\alpha = 2 \arctan \left( \frac{L}{h} \right), \beta = \arctan \left( \frac{l+x}{h} \right) + \arctan \left( \frac{l-x}{h} \right)$$

Then, the solid angle from the center of the detector is :

$$\Omega(L, l) = 4 \arcsin \left( \sin \left( \frac{\alpha(L)}{2} \right) \sin \left( \frac{\beta(l)}{2} \right) \right)$$

We then have to take the solid angle from any point in the detector, it gives the following integral for  $C(\theta)$  :

$$C(\theta) = 4 \int_{u=0}^L \int_{v=0}^l \Omega(u, v) du dv$$

The normalization follows, considering we are detecting the flux from only one direction at a time :

$$\phi_{\text{measured}}(\theta) = \frac{N_{\text{detections}}(\theta)}{TC(\theta)}$$

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5. G.Remmen, E.McCreary ; Journal of Undergraduate Research in Physics , **Measurement of the speed and energy distribution of cosmic ray muons**, 2012
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