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## 1 Unraveling the Conformational Landscape of Ligand Binding to Glucose/Galactose-Binding Protein by Paramagnetic NMR and MD Simulations

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**S** [Supporting Information](#page-7-0)

 ABSTRACT: Protein dynamics related to function can nowadays be structurally well characterized (i.e., instances obtained by high resolution structures), but they are still ill-defined energetically, and the energy landscapes are only accessible computationally. This is the case for glucose−galactose binding protein (GGBP), where the crystal structures of the apo and holo states provide structural information for the domain rearrangement upon ligand binding, while the time scale and the energetic determinants for such concerted dynamics have been so far elusive. Here, we use GGBP as a paradigm to define a functional conformational landscape, both structurally and energetically, by using an innovative combination of paramagnetic NMR experiments and MD simulations. Anisotropic NMR parameters induced by self-alignment of paramagnetic metal ions was used to characterize the ensemble of conformations adopted by the protein in solution while the rate of interconversion between conformations was elucidated by long molecular



33 dynamics simulation on two states of GGBP, the closed-liganded (holo cl) and open-unloaded (apo op) states. Our results demonstrate that, in its apo state, the protein coexists between open-like (68%) and closed-like (32%) conformations, with an

exchange rate around 25 ns. Despite such conformational heterogeneity, the presence of the ligand is the ultimate driving force to

36 unbalance the equilibrium toward the holo cl form, in a mechanism largely governed by a conformational selection mechanism.

**P**rotein function arises from the delicate interplay among structure, molecular recognition features, and dynamics, but unraveling such contributions is often elusive. The periplasmic binding protein family (PBPs) represents a paradigm for describing functional conformational changes in 42 flexible proteins.<sup>[1](#page-7-0)</sup> In Gram-negative bacteria, PBPs selectively recognize and actively transport various nutrients across the inner membrane. The family is composed of about 100 members, classified according to the recognized ligand: amino acids, carbohydrates, oxyanions, and vitamins.<sup>[2](#page-7-0)</sup> Almost all of them share a common structural fold consisting of two globular Rossman fold domains connected by three short linkers, thus

suggesting interdomain flexibility.  $3-5$  $3-5$  This hypothesis is further 49 supported by the different interdomain orientations found in 50 the X-ray structures. For instance, one of the paradigmatic <sup>51</sup> glycan-binding protein family, the glucose/galactose binding <sup>52</sup> proteins (GGBP) from different organisms have generated a <sup>53</sup> plethora of crystal structures that trap the biomolecule at <sup>54</sup> distinct conformational instances: some unloaded structures are <sup>55</sup> canonically open  $(apo$  op), while others are closed  $(apo$  cl) 56

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57 and resemble the holo-ligand-bound state  $(holo \; cl).^{6-10}$  In the past few years, segmental interdomain reorientations in periplasmic binding proteins have been extensively investigated by solution NMR spectroscopy.[11](#page-7-0) Clore and co-workers have demonstrated that a conformational selection process under-62 goes an open-to-closed transition in MBP, $^{12}$  $^{12}$  $^{12}$  while Tjandra and co-workers have shown that an induced fit mechanism well describes the open−closed transition of another PBP, 65 glutamine-binding protein.<sup>[13](#page-7-0)</sup> In a comparative NMR study of GGBP and the structurally homologous ribose binding protein, it was shown that the (apparent) ligand affinity can be modulated by redesigning the flexible hinge region, thus 69 emphasizing the functional role of interdomain dynamics.<sup>[14](#page-7-0)</sup> However, the time scale and amplitude of these motions are experimentally ill-defined for all the investigated cases. Molecular dynamics simulations have also been widely used  $73$  to characterize the conformational landscape of PBPs. $3$  For instance, advanced sampling techniques have been used to 75 study the allosteric equilibrium of the ribose-binding protein, $1$  while accelerated MD simulations provided a detailed picture of the transition between the open and partially closed states in MBP.[16](#page-7-0) Moreover, these proteins have been the target of 79 intense studies in protein engineering, $17$  and the computational redesign of PBPs to build up nanobiosensors have raised great 81 expectations.<sup>[18](#page-7-0),[19](#page-7-0)</sup> For instance, Daunert and co-workers have 82 proposed GGBP as a possible biosensor of glucose in blood.<sup>[20](#page-8-0)</sup> However, ironically, the main limitation of the method is the high affinity for the substrate (nM range). Then, despite the extensive use of MD simulations in the study of PBPs, integrative approaches of protein design with experimental data are still largely unedited.

Here, we propose an integrated approach by combining experimental NMR data with molecular dynamics simulations to quantitatively characterize interdomain dynamics in GGBP. First, pseudocontact shifts (PCSs) and residual dipolar couplings (RDCs) allowed for disentangling the population distribution of conformers in the open−closed transition that GGBP undergoes. Next, the time scale for the open−closed transition is defined by detailed molecular dynamics simu- lations. Finally, the energy barrier in the protein landscape has been estimated using nonequilibrium molecular dynamics calculations. Our results demonstrate that, in its apo state, the protein coexists between the open (68%) and closed (32%) conformations. The time scale for closed−open interconversion is around 25 ns. The presence of the ligand is the driving force for closing, largely through a conformational selection mechanism.

### 104 **RESULTS AND DISCUSSION**

<sup>105</sup> Self-Alignment with a Paramagnetic Tag Reveals <sup>106</sup> Conformational Heterogeneity in apoGGBP. GGBP <sup>107</sup> consists of two globular domains, the C-terminal domain <sup>108</sup> (residues 112−254 and 297−306) and the smaller N-terminal <sup>109</sup> domain (residues 3−108 and 258−291), linked by a three-<sup>110</sup> strand hinge (residues 109−111, 255−258, and 292−296) 111 (Figure 1, panel A). The  $\beta$  anomer of glucose binds to GGBP 112 through an extensive network of hydrogen bonds and CH- $\pi$ 113 interactions with high affinity and specificity ( $K_D = 290$  nM at 114 37 °C and pH 7.0) into the cleft near the hinge region, invoking <sup>115</sup> a large conformational change from the open unbound 11[6](#page-7-0) (apo\_op) to closed bound (holo\_cl) state.<sup>6</sup> This segmental 117 interdomain reorientation is well described by a rotation of 40° 118 in the twist  $(\phi)$  angle accompanied by a 23° rotation in the



Figure 1. (A) Crystal structures of apo opGGBP (left, 2FW0) and holo clGGBP (right, 2FVY). The bound  $\beta$  anomer of D-glucopyranose and residues Asp14, Asn91, His152, Asp154, Arg158, Asn211, Asp 236, and Asn256 forming stabilizing H-bonds with the ligand are drawn as stick models (ligand in violet). Inset: schematic representation of globular domains (cylinders) and the hinge region (lines). The difference in N-terminal domain position highlights the difference in closure angle, according to X-ray structures. (B) Domain reorientation of GGBP. Left, side view illustrating hinge domain rearrangement between apo\_op (blue) and holo cl (orange) GGBP. Right, front view illustrating twist motion. Inset: the angle between the segments connecting the center of mass of the hinge region and those of the C-terminal domain and N-terminal domain is defined as hinge angle, while the center of mass of the N-terminal domain, the base of the N-terminal domain, the C-terminal domain and the base of the Cterminal domain define three segments. The dihedral angle formed by these three segments is defined as twist angle.

hinge ( $\theta$ ) angle<sup>[14](#page-7-0),[15](#page-7-0)</sup> (Figure 1, panel B). Since closed unloaded 119 structures are available for some PBPs, an open question is <sup>120</sup> whether apoGGBP can make excursions to the holo cl 121 conformation in the absence of ligand. Structural data for <sup>122</sup> GGBP in solution can be integrated by anisotropic NMR <sup>123</sup> parameters that are induced by self-alignment of paramagnetic <sup>124</sup> molecules. Such self-alignment has been achieved by binding <sup>125</sup> paramagnetic metal ions to a small molecule chemical metal <sup>126</sup> chelating tag, covalently attached to the biomolecule.<sup>[21](#page-8-0)</sup> The  $127$ synthesis of the novel tag is described in the [Supporting](http://pubs.acs.org/doi/suppl/10.1021/acschembio.6b00148/suppl_file/cb6b00148_si_001.pdf) <sup>128</sup> [Information.](http://pubs.acs.org/doi/suppl/10.1021/acschembio.6b00148/suppl_file/cb6b00148_si_001.pdf) The alignment tensors for the tagged and <sup>129</sup> nontagged domains have been determined by the combined <sup>130</sup> use of residual dipolar couplings (RDCs) and pseudocontact <sup>131</sup> shifts (PCSs). Actually, owing to the large magnetic moment of <sup>132</sup> the unpaired electrons on the paramagnetic lanthanide ion, the <sup>133</sup> paramagnetic effects are detectable up to large distances (>40 <sup>134</sup> Å). $^{22}$  $^{22}$  $^{22}$  In a system comprising two or more species in rapid 135 exchange, the observed effect is a population-weighted average <sup>136</sup> of the component conformers. As a result, PCSs and RDCs <sup>137</sup> provide a unique way to describe complex mixtures of <sup>138</sup> translational and rotational interdomain motions, simply <sup>139</sup>

 aligning the tagged domain by the paramagnetic tag and 141 determining the induced alignment on the other moieties.<sup>[23](#page-8-0)</sup> To that end, the paramagnetic probe has been conjugated through nucleophilic substitution to an engineered cysteine residue, M182C, located in the C-terminal domain at the periphery of the interface between the two domains. Several factors were considered for the insertion of the tag molecule: (i) the minimal biorelevant mutation Cys instead of Met, (ii) an adequate distance between the tag molecule and the target binding site, so the spin-label does not perturb the backbone structure nor the ligand-binding site, and (iii) surface-accessible amino acids that experience minimum variation in chemical environment upon sugar binding [\(Figure S1](http://pubs.acs.org/doi/suppl/10.1021/acschembio.6b00148/suppl_file/cb6b00148_si_001.pdf)). The NMR signals for the residues within the shell around the para- magnetic center (C182 and A181) were broadened beyond detection due to paramagnetic relaxation enhancement (PRE). Nevertheless, the chemical nature of the tag molecule (the 157 metal is located at a distance  $>16$  Å away from the protein f2 158 backbone) allowed collecting 135 measurable PCSs (Figure 2)



Figure 2. Pseudocontact shifts (PCSs) obtained as the difference in the chemical shift of the protons signals in diamagnetic (lanthanum) and paramagnetic (dysprosium) conditions, for the observable  $H<sup>N</sup>$ , N nuclei. Resonances of residues within the shell around the paramagnetic center were broadened beyond detection due to paramagnetic relaxation enhancement (PRE).

 and 20 RDCs. Clearly, two regions of the protein orient differently with respect to the paramagnetic metal ion. One region undergoes negative chemical shift perturbation, close in space to the negative lobe of the magnetic susceptibility tensor. The other region experiences opposite changes due to its orientation toward the positive lobe of the paramagnetic metal isosurface.

 To take into account the flexibility of the linker, two main conformations with staggered dihedral angles around the disulfide bond have been generated (90° and −90°). Only the −90° conformer fits well the experimental data, confirming the impossible population of the 90° conformer due to steric 171 clashes (see [Supporting Information](http://pubs.acs.org/doi/suppl/10.1021/acschembio.6b00148/suppl_file/cb6b00148_si_001.pdf)). PCSs of  $H<sup>N</sup>$ , N nuclei, 172 and RDCs of  $H^{\tilde{N}}$ -N pairs for paramagnetic lanthanide  $(Dy^{3+})$ 173 were determined from an <sup>I</sup>H,<sup>15</sup>N HSQC. A La<sup>3+</sup>-loaded complex was used as diamagnetic reference, as this metal has an ionic radius similar to that of the dysprosium ion. The 176 protein shows excellent signal dispersion in both <sup>1</sup>H,<sup>15</sup>N-HSQC 177 spectra  $(Dy^{3+}$  and La<sup>3+</sup>), ([Figure 3\)](#page-3-0). Chemical shifts between unloaded and diamagnetic metal ion loaded molecules are virtually identical, and nearly complete assignments could be 180 obtained based on previously published data.<sup>[14](#page-7-0)</sup> The lack of chemical shift perturbation in the GGBP backbone amide signals after metalation of the sample indicates that all the

lanthanide ions are bound to the tag and not directly to the <sup>183</sup> protein, consistent with the high affinity of the tag for <sup>184</sup> lanthanides (in the 10−18 M range).<sup>[24](#page-8-0),[25](#page-8-0)</sup> Representative 185 structures for the apo\_op and apo\_cl conformations (with  $5^\circ$  186 stepwise changes in the closure and twisting angles) were <sup>187</sup> extracted from a molecular dynamic simulation (vide infra) and <sup>188</sup> used for the alignment tensor estimation, using RDCs and <sup>189</sup> PCSs as experimental restraints. The structures providing the <sup>190</sup> lowest Q-factor for the tagged domain were selected and <sup>191</sup> subsequently used for the prediction of the NMR parameters in <sup>192</sup> the tagged-free domain. Two structures showed the lowest <sup>193</sup> quality factors (Q-factor) when fitting the experimental data for <sup>194</sup> the tagged domain: one representative for the *apo* op form ( $\theta$  195 = 145° and  $\phi$  = 64°) and another for the *apo* cl conformation 196  $(\theta = 137^{\circ}$  and  $\phi = 20^{\circ})$ . When analyzed independently, PCSs 197 and RDCs (the latter to a lower extent) provided excellent fits <sup>198</sup> for each tagged domain for both structures, as reflected in the <sup>199</sup> range of the Q-factor values: 0.084−0.099 (PCSs) and 0.494− <sup>200</sup> 0.420 (RDCs). 201

These results demonstrate that the Tag molecule and the C- <sup>202</sup> terminal domain behave as a rigid body ([Figure 4,](#page-4-0) panel A and 203 f4 intra Q-factor in panels C−E). Interestingly, Q-factors <sup>204</sup> drastically increase when both domains are included in the <sup>205</sup> analysis ([Figure. 4,](#page-4-0) panel B and overall Q in panels C,D). When <sup>206</sup> fitting the PCS data set, the quality factor referred to the open <sup>207</sup> structure rises up to 0.231, while when referred to a closed-like <sup>208</sup> structure, it reports a value of 0.281. These results demonstrate <sup>209</sup> that a single structure is not able to explain the experimental <sup>210</sup> data set, likely because the N-terminal domain fluctuates with <sup>211</sup> respect to the C-terminal. A model contemplating an average <sup>212</sup> ensemble of differently populated states was tested, and a <sup>213</sup> combination of the two above-mentioned conformations (68% <sup>214</sup> for the apo\_op and 32% for the holo\_cl) provides a very good 215 correlation with the experimental data ([Figure 4,](#page-4-0) panel E). <sup>216</sup> Thus, the RDC and PCS experimental data reported here on <sup>217</sup> apoGGBP fully agree with a model where the apo state 218 undergoes a rapid equilibrium between a major and a minor <sup>219</sup> species, the latter one occupying a region of the conformational <sup>220</sup> landscape similar to the ligand bound form [\(Figure 5,](#page-4-0) panel C). 221 f5 Our results also demonstrate the existence of the postulated <sup>222</sup> dynamic equilibrium between open and partially closed apo <sup>223</sup> states and gives credit to the hypothesis that large-scale domain <sup>224</sup> rearrangements are already present in many two-domain <sup>225</sup> periplasmic proteins. <sup>226</sup>

Unraveling the Conformational Landscape of GGBP <sup>227</sup> by MD Simulations. Because of the dynamic nature of 228 apoGGBP, MD simulations have been employed to investigate <sup>229</sup> its conformational landscape. In a first set of calculations, the <sup>230</sup> experimental RDCs were included as restraints in a conjoined <sup>231</sup> rigid body-torsion angle simmulated annealing followed by an <sup>232</sup> MD simulation of 200 ps. (see [Methods](#page-6-0)). An intermediate <sup>233</sup> conformation, partially closed, complies well with the <sup>234</sup> experimental restraints and can be interpreted in terms of a <sup>235</sup> combination of the open (68%) and closed (32%) states. <sup>236</sup> Specifically, the values for the interdomain hinge/twist angle <sup>237</sup> that satisfies the NMR restraints lies around  $135 \pm 5^{\circ}/40 \pm 10^{\circ}$  238 (see [Figure 5\)](#page-4-0). The values are similar but statistically different <sup>239</sup> from the consensus conformation obtained from fitting <sup>240</sup> experimental diamagnetic RDCs in weakly aligned media ( $\theta$  241 127°;  $\phi$  32°),<sup>[14](#page-7-0)</sup> thus suggesting the existence of interdomain 242 dynamics. <sup>243</sup>

To further characterize such motions, in a second calulation <sup>244</sup> experimentally determined NMR constraints (i.e., PCSs and <sup>245</sup>

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Figure 3. (A) <sup>1</sup>H,<sup>15</sup>N HSQC for GGBP loaded with diamagnetic lanthanum (orange) and paramagnetic dysprosium (blue). Spectra were recorded at 310 K and pH 7.0 in a 600 MHz NMR spectrometer. (B) Selected spectral region of the <sup>1</sup>H,<sup>15</sup>N HSQC spectra.

 RDCs) were not introduced to minimize biases in the conformational landscape. Two MD simulations, starting from X-ray structures of apo\_op (2FW0) and holo\_cl GGBP bound to glucose (2FVY), were performed with the ff10\_Amber and GLYCAM\_06h force fields in explicit water at 310 K to generate atomic coordinate data sets that describe the protein ensemble. Because of the expected higher dynamic complexity, the MD simulation for apo\_op was run for 300 ns, while holo\_cl was run for a time of 100 ns. It is important to emphasize that both free-MD simulations are long enough to explore a wide conformational space and to define the periodic dynamic behavior of GGBP (apo and holo) in solution.

 After excluding the preparatory steps from the trajectory, MD runs have been analyzed in terms of the interdomain hinge and twist angles with a defined hinge position ([Figure 5](#page-4-0), panels A,B). For apo\_op GGBP, a range of twist and hinge angles are accessible, oscillating between a partially closed conformation 263 (θ 135°; φ 40°) and a widely open conformation (θ 170°; φ  $264$  $264$   $160^\circ$ ), consistent with previous studies.<sup>6</sup> Conformational interconversion in apo\_op GGBP involves concerted changes in both angles: the twist angle fluctuates between 30° and 160°, with equally low energy conformers, and it is always accompanied by a hinge angle oscillation between 140° and 160°. The time scale for such collective dynamics, between crest and wave of the periodic motion, is around 25 ns. As 271 expected, MD simulation for holo cl GGBP structure is characterized by minor excursions within the conformational 273 space. The hinge and the twist angles oscillate between  $\pm 10^{\circ}$ and  $\pm 20^\circ$ , respectively, around the starting values.

 The overall QRDC value estimated for the ensemble 276 structure between the *apo*  $op(X-ray)$  (label 4 in [Figure 6\)](#page-5-0) and *holo cl*, (label 8 in [Figure 6\)](#page-5-0) markedly improves with respect to the single conformation ones (see [Figure 4](#page-4-0), panel C−E). The results demonstrate that the conformational behavior of GGBP in its unbound state is compatible only with a dynamic phenomenon that involves partially closed conformations suggesting that the ligand recognition event

cannot be described by pure "induced-fit" or pure "conforma- <sup>283</sup> tional selection" models.

The free energy of the conformational landscape explored by <sup>285</sup> the MD simulations has been estimated from the populations <sup>286</sup> by dividing the conformational space into regular intervals of <sup>287</sup>  $(2^{\circ} \times 4^{\circ})$  hinge/twist angles. The population for each interval 288 has been converted into the energy difference with respect to <sup>289</sup> the highest populated one according to the following <sup>290</sup> expression:

$$
\Delta G_{\text{PA-PB}} = -RT \ln([P_{\text{A}}]/[P_{\text{B}}])
$$

The free energy landsacpe for apo\_opGGBP [\(Figure 5,](#page-4-0) panel 292 C) shows a wide global minimum corresponding to an <sup>293</sup> ensemble of highly variable structures in terms of twist/hinge <sup>294</sup> angles. Remarkably, a variety of conformations expanding the <sup>295</sup> open-to-close conformational coordinate (i.e., superopened and <sup>296</sup> closed-like conformations) have a free energy excess of only 1− <sup>297</sup> 2 kcal·mol<sup>-1</sup> as compared to the global minimum, indicating  $298$ that these structures are also accessible. As inferred from [Figure](#page-4-0) <sup>299</sup> [5](#page-4-0), panels A,B, these conformations are periodically revisited, <sup>300</sup> demonstrating that they are not transiently populated high <sup>301</sup> energy states. However, the free energy landscape for <sup>302</sup> holo\_clGGBP fluctuates around a global minimum, structurally 303 corresponding to the starting conformation. Large scale <sup>304</sup> conformational changes are multiple trajectory processes. <sup>305</sup> Nevertheless, the two independent molecular dynamic <sup>306</sup> simulations seem to energetically coalesce, defining the <sup>307</sup> conformational coordinate for the most likely open-to-close <sup>308</sup> trajectory.

In order to well define the energy profile associated with the <sup>310</sup> open-to-close transition in GGBP, SMD have performed. This <sup>311</sup> computational approach has been extensively used to calculate <sup>312</sup> the free energy associated with unfolding/refolding pathways of <sup>313</sup> macromolecules,  $26.27$  ligand–receptor binding events, <sup>[28](#page-8-0)</sup> or 314  $DNA$  starching.<sup>[29](#page-8-0)</sup> SMD employs a pulling force to cause a 315 structural change so that different conformations may be <sup>316</sup> sampled along a given pathway. To gain access to time scales <sup>317</sup>

<span id="page-4-0"></span>

#### **Experimental restraints**

Figure 4. Correlation between experimental and back-calculated restraints relative to apo\_opGGBP (blue), holo\_clGGBP (orange), and an average of 68% and 32% of both structures, respectively (violet). (A) PCSs of H<sup>N</sup>, N nuclei belonging to only the protein tagged domain. (B) PCSs of H<sup>N</sup>, N nuclei of the entire protein. (C) RDCs of H<sup>N</sup>-N pairs of different selected amino acids in apo\_opGGBP. Label 4 in Figure 5. (D) RDCs of H<sup>N</sup>-N pairs of different selected amino acids in holo clGGBP. Label 8 in Figure 5. (E) RDCs of H<sup>N</sup>-N pairs of different selected amino acids in apo\_clGGBP. This latter structure is an ensamble average structure of 68% apo\_op and 32% holo\_cl GGBP. The values for hinge and twist angles are 137° and 40°, respectively, and its position in the conformational space is specified in Figure 5.



Figure 5. Analysis of interdomain hinge (A) and twist (B) angles along the MD simulations for apo\_opGGBP (300 ns) and holo\_clGGBP (100 ns). The derived 200 ns of collective motion for holo clGGBP are marked with a tenuous line. (C) Free-energy landscapes of apo\_opGGBP (blue) and holo\_clGGBP (orange) as a function of hinge  $(\theta)$  and twist  $(\phi)$  angles. Selected snapshot along the trajectory, red dashed line, are labeled with numbers in circles. From apo\_op MD, 1-3 and 5,6; 4 is the X-ray structure (2FW0); 4 and 8 have been used to derive the ensamble average population from experimental NMR restraints. From holo cl MD, 8,9; from steered molecular dynamic simulations (SMD), 7 is the high energy structure corresponding to the transition state.

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Figure 6. Potential of mean force as a function of the hinge angle  $\theta$  (deg) in protein/ligand binding. The calculated energy profile relative to the holo cl structure (orange) indicates that the free energy minimum for closed conformation is lower than that of the open state by roughly 6 kcal· mol<sup>-1</sup>. Superimposed is the energy profile for the apo\_op structure (blue). The result indicates that the energy of the open structure increases with the closing of the hinge angle. On top, structural change associated with interdomain closure and ligand binging. The structures correspond to snapshots along the trajectory.

 that would be otherwise computationally too demanding, SMD simulations accelerate (or force) the dynamic process, thus providing the required energy associated with the induced change: while the force is executed and the motion occurs along a given coordinate, the potential energy of the system is calculated, and the potential of mean force (PFM) is related to the free energy profile of the process.[30](#page-8-0) Even if we cannot exclude the existence of other transition coordinates for the open-to-close event, the method ensures that the work done on the system is a function of the activation energy associated with the process. Two SMD simulations have been carried out. In the first one, the value of the critical hinge angle was fixed to 124°, which corresponds to the X-ray structure determined for 331 the holo cl conformation, and a pulling force was applied to 332 reach the holo\_op conformation. In the second one, the starting geometry corresponds to the X-ray structure of the unbound 334 open GGBP (apo op), with a hinge angle of  $145^\circ$ , in order to 335 achieve the closed unbound (apo cl) structure. Figure 6, orange line plots the resistance applied by the system in 337 opening holo clGGBP as a function of the hinge angle. According to the SMD simulations, the transition state corresponds to a hinge angle of 132°, and its height is only  $340 \approx 3.5$  kcal·mol<sup>-1</sup> higher than the free energy of the open structure.

 We can conclude that the energetic barrier for open-to- closed interconversion can be easily crossed via thermal fluctuations or with the help of water molecules that attack interdomain hydrogen bonds. Not surprisingly, the closed structure is significantly stabilized by intermolecular inter-347 actions with the ligand (hydrogen bond and  $CH/\pi$  interactions) that contribute with 6 kcal·mol<sup>-1</sup> to the binding energy of the complex. The thick hydrogen bond network between the polar face of the sugar and the polar amino acids on both domains of GGBP and the two aromatic residues Trp181 and Phe16 on the C-terminal and N-terminal domain, respectively, stacked against the less polar faces of the sugar provide the necessary enthalpy contribution to stabilize the 355 holo clGGBP structure. The second SMD simulation starting 356 in the apo opGGBP is also shown in Figure 6, blue line. Here, a quasi-linear energy dependence is observed in the transition

from the open to the closed states of GGBP, with an associated <sup>358</sup> Gibbs free energy that is inversely proportional to the closing <sup>359</sup> angle. The stability of the protein decreases at a rate of about <sup>360</sup> 300 cal·mol<sup>-1</sup>·deg<sup>-1</sup>. Considering a range of 3 kcal·mol<sup>-1</sup> 361 accessible by thermal fluctuations at 310 K, a domain closure <sup>362</sup> of  $10^{\circ}$  to  $20^{\circ}$  is expected. Then, the *apo* cl conformation of the 363 protein, which is almost 6 kcal·mol<sup>-1</sup> higher in energy than the  $364$ apo op, would be inaccessible by thermal fluctuations, 365 consistent with previous results reported for MBP. $^{11}$  $^{11}$  $^{11}$  The 366 intersection between the two independent SMD define the <sup>367</sup> transition state, where only the incoming ligand provides the <sup>368</sup> necessary energy stabilization to shift the equilibrium in favor of <sup>369</sup> the most stable ligand−protein complex. <sup>370</sup>

Synergy between Structure and Dynamics Is Essential <sup>371</sup> for Protein Function. Segmental dynamics in proteins is 372 often functional and is the basis of protein allosterism. Yet, they <sup>373</sup> are often very loosely characterized due to the lack of <sup>374</sup> experimental tools available. Here, we demonstrate that a <sup>375</sup> combination of NMR spectroscopy and MD simulations <sup>376</sup> successfully unravels the energy landscape for the functional <sup>377</sup> conformational coordinate of GGBP. <sup>378</sup>

A single structural model unsatisfactory describes the vast <sup>379</sup> conformational space covered by apo opGGBP in solution. 380 According to unrestricted MD simulations, the conformational <sup>381</sup> dynamics is well characterized by large amplitude motions <sup>382</sup> between hemiclosed and open structures in the nanosecond <sup>383</sup> time scale. Actually, the first reported X-ray structure of <sup>384</sup> apo  $GGBP$  is partially closed.<sup>[6](#page-7-0)</sup> The authors noticed the 385 presence of the citrate ion in the binding cleft and hypothesized <sup>386</sup> that they may be thermodynamically relevant for crystal growth. <sup>387</sup> Our results agree, demonstrating that the sodium ion <sup>388</sup> coordinates at the hinge position, thus stabilizing the <sup>389</sup> hemiclosed structures in apo GGBP. Importantly, these 390 conformational interconversions are experimentally validated, <sup>391</sup> while the combined use of PCSs and RDCs induced by self- <sup>392</sup> alignment of paramagnetic metal provides a fine method to <sup>393</sup> unveil the existence of functional closed conformations that are <sup>394</sup> consistent with a 32% of the total population of apo GGBP. 395

Albeit the limited range of closure angles for the deposited X- <sup>396</sup> ray structure of *apo*  $opGGBP$  spans between  $147^{\circ}$  and  $149^{\circ}$ , 397

<span id="page-6-0"></span> our data suggest that excursions toward more open states are lightly permitted (150° and 170°, "superopen" structures). Such conformations have been also predicted and observed for other bidomain proteins like uroporphyrinogen III synthase,  $31$  and it is probably dependent on the hinge region structure and composition since it reflects the maximum torque force that this region allows. However, closed-like states are also represented, between 132° and 147°, supporting the existence of closed-like structures for apoGGBP as observed via X-ray 407 crystallography.<sup>[32](#page-8-0)</sup>

408 The interdomain excursions found for apo GGBP are largely quenched in holo\_GGBP, but interestingly, they are not completely abrogated ([Figure 5\)](#page-4-0). Actually, when bound to 411 glucose, holo clGGBP still holds some degree of flexibility, with interdomain librations of up to 10−15° in the closure and hinge angles. According to MD simulations, such motions also fall in the nanosecond time scale, and they are consistent with the increased dynamics previously observed in the order parameter 416 of holo  $GGBP^{14}$  $GGBP^{14}$  $GGBP^{14}$  that likely interfered on the spectral density function motional parameters.

 The functional role of these dynamics naturally emerges when the energetics are introduced, and the integrative analysis of MD simulations and NMR analysis comply largely with a mainly conformational selection mechanism for which the presence of the ligand is indispensable for the open-to-closed transition. MD simulations reveal a gap in the conformational 424 landscape between apo opGGBP and holo clGGBP (125− 140° hinge and 30−40° twist, [Figure 5,](#page-4-0) panel C). This high energetic point corresponds to the transition state with an activation barrier ≈3.5 kcal·mol<sup>−</sup><sup>1</sup> as estimated by SMD simulations [\(Figure 6](#page-5-0)). Thus, ligand binding provides the energy to overcome such a barrier and shift the conformational 430 ensemble toward holo clGGBP in a second step that agrees well with an induced fit mechanism. The open-to-closed transition for GGBP is described in a short representative video 433 in [Video 1](http://pubs.acs.org/doi/suppl/10.1021/acschembio.6b00148/suppl_file/cb6b00148_si_002.mp4). Remarkably, apo\_opGGBP is predicted to unfold at low values of the closing angle due to the increase in nonpolar solvent accessible area located on the hinge region on the opposite site with respect to the ligand binding pocket. This negative term, if not balanced by the enthalpic surplus of the binding event, will result in protein unfolding. This mechanism is equivalent to the experimentally found one for maltose binding protein, where the analysis of hinge mutants with different closure angles predicted protein unfolding at low 442 closure angles, within a similar free energy range. $<sup>11</sup>$  $<sup>11</sup>$  $<sup>11</sup>$ </sup>

 In conclusion, the amplitude and time scale of GGBP interdomain dynamics have been unveiled by NMR spectros- copy and detailed MD simulations. The population contribu- tion of biofunctional relevant conformers has been determined by PCSs and RDCs induced by a paramagnetic metal ion. Subsequently, the energetic barrier in open-to-closed transition has been defined by nonequilibrium MD simulations. We 450 conclude that the population of apo cl-like conformations is 451 essential to activate the transition to the *holo* cl form, according to a conformational selection mechanism coupled to a final rearrangement that obeys an induced-fit kinetics process.

 From a general perspective, the protocol exemplified herein can be extended to the study of a variety of molecular recognition processes in which significant molecular rearrange- ments take place, thus expanding the limits of the application of NMR methods to explore binding events.

# ■ METHODS 459

Protein Expression and Purification. Uniformly labeled  $(^{15}N)$  460 GGBP protein was overexpressed in minimum M9 media (1.5 L, 461 purchased from CIL) containing  ${}^{15}NH<sub>4</sub>Cl$  as the only source for 462 nitrogen. GGBP was purified via ion exchange chromatography using <sup>463</sup> fast flow Q-Sepharose followed by ion exclusion chromatography <sup>464</sup> (Superdex 75, GE Healthcare) in 20 mM Tris buffer at pH 7.1 and 150 <sup>465</sup> mM NaCl. To eliminate associated sugars, the protein was dialyzed 466 several times against 3.5 M guanidinium chloride under the same 467 buffer conditions. Final sample conditions: 0.5 mM in 20 mM Tris, <sup>468</sup> 150 mM NaCl, 10 mM CaCl<sub>2</sub>, and 95%/5% H<sub>2</sub>O/D<sub>2</sub>O, pH 7.0. 469

Sample Preparation for Paramagnetic studies. Protein 470 samples for NMR studies were prepared at a final protein <sup>471</sup> concentration of 0.5 mM in 20 mM Tris (pH 7.0), 150 mM NaCl, 472 10  $\mu$ M CaCl<sub>2</sub>, and 1 mM NaN<sub>3</sub> with 10% D<sub>2</sub>O. In order to conjugate 473 the paramagnetic tag to the C-terminal domain cysteine, the protein 474 was titrated with 10 mM solution of lanthanide (Ln) chelating tag 475 previously loaded with the lanthanide  $(Ln = La^{3+}, Dy^{3+})$ . The titrations 476 were performed by monitoring the changes in the chemical shift in a 477  ${}^{1}H, {}^{15}N$  HSQC spectrum. The nucleophilic substitution reaction is 478 instantaneous, and the excess of chelating tag molecule was removed 479 by filtration. 480

NMR Spectroscopy. All of the NMR experiments were carried out 481 on a 600 MHz AVANCE-III Bruker spectrometer. Spectra were 482 acquired at 310 K. All NMR spectra were processed with the software 483 TopSpin. The Program CARA was used for the analysis of the 2D 484 spectra. Lanthanum was used as diamagnetic reference as it has an 485 ionic radius similar to the paramagnetic dysprosium. PCSs were 486 measured as the difference between the chemical shift of the <sup>487</sup> corresponding nuclei in the paramagnetic and diamagnetic samples. 488 Residual dipolar couplings  ${}^{1}D_{HN}$  were measured as the  ${}^{1}H$ -doublet 489 splitting of the paramagnetic sample minus the equivalent splitting 490 difference in the diamagnetic sample. 491

PCS and RDC analyses were performed using MSpin software. In 492 order to assess interdomain dynamics, we used selected conformations 493 from molecular dynamics simulation, including some that resemble the 494 experimentally obtained X-ray open and closed GGBP conformations. 495 The optimal ensemble of conformations was found after evaluating the 496 tensor from the tag domain using both PCSs and RDCs 497 independently. The PCSs and RDCs for the other domain were 498 back calculated for different protein coordinates, and the ensemble of <sup>499</sup> structures that better fits the experimental values was selected on the <sup>500</sup> basis of the best quality factor. 501

Molecular Dynamics Simulations. The coordinates from the X- 502 ray structure of apo\_op GGBP (PDB code 2FW0) and holo\_cl GGBP 503 (PDB code 2FVY) were used as starting points to generate 504 intermediate models by molecular dynamics (MD) simulations at 505<br>310 K applying the ff10<sup>[33](#page-8-0)</sup> and GLYCAM\_06h<sup>[34](#page-8-0)</sup> force fields. Missing 506 hydrogens were added to the starting PDB structures using the 507 program LEAP. The N- and C-terminal residues were acetylated and 508 amidated according to the AMBER standard database. The latter 509 structure was solvated in a cubic TIP3P water box, $35$  and eight sodium 510 ions were added to neutralize the system. In order to fill all of the <sup>511</sup> proteins cavities with water molecules, a previous minimization for 512 only solvent and ions was made. Moreover, to reach a reasonable 513 starting structure, the entire system was minimized with a higher 514 number of cycles, using the accurate steepest descent algorithm. The 515 system was subjected to two rapid molecular dynamic simulations of 516 20 and 100 ps, respectively, before starting the real dynamic simulation 517 of 270 ns for apo op and 100 ns for holo clGGBP. During these two 518 preparatory steps, the structure was slowly heated from 0 to 310 K. 519 Fifty-thousand additional steps were performed to switch from 520 constant volume to constant pressure. A relaxation time of 2 ps was 521 used in order to equilibrate the entire system in each step. The final <sup>522</sup> simulations of 270 and 100 ns were performed starting from 523 equilibrated structures. Coordinates and energy values were recorded 524 every 2 ps for a total of 135000 MD models for apo\_op and 50000 for 525 holo clGGBP. For the SMD simulations, the starting structures, with 526 hinge values of 124° and 147° have been extracted by unrestricted MD 527

<span id="page-7-0"></span>528 simulations of holo cl and apo opGGBP, respectively; so the entire system was already equilibrated. The center of mass of the N-terminal domain together with the center of mass of the hinge segment has been fixed, while the center of mass of the C-terminal domain has been 532 pulled with a constant force  $K = 500$  kcal·mol<sup>-1</sup>·Å<sup>2</sup>. The total time for the molecular dynamics simulation was of 10 ns with an angle opening or closing of 2° per ns. Atomic coordinates were saved every 2 ps and the energy information extracted.

 Constrained MD simulations were initiated from the X-ray 537 structures apo op (2FW0) and holo cl (2FVY), shown as orange and blue circles, respectively. Conjoined rigid body-torsion angle 539 simmulated annealing was performed as previously described.<sup>[36,](#page-8-0)14</sup> The hinge region for the GGBP is defined by residues 109−111, 253−256, and 293−296. The starting structure (apo\_op, F2W0) was heated to 600 K for 3 ps with tautp equal to 0.4. Then, the system was cooled to 543 100 K for 297 ps (tautp = 4.0). The final cooling to 0 K was carried 544 out for 100 ps with tautp varying from 1.0 to 0.1. The  $\gamma$  tensor anisotropy parameters extracted from MSpin software were those were derived from the lowest QF structures. MD simulations were performed with ff10\_Amber and GLYCAM\_06h force fields integrated with the experimentally derived NMR restraints. The restrained MD calculations were performed in explicit water solvent and using a simulated annealing approach.

 Analysis of the Trajectories. Root mean-square deviation (simulated) and thermodynamic data were monitored throughout the whole trajectory to confirm that all simulations evolved along a stable plateau. For the analysis of the collective motions, the closure 555 (hinge), twisting, and bending coordinate system was used.<sup>[37](#page-8-0)</sup> Hinge and twist angles were obtained from clusters of one in every 100 models (1350 in total) for apo\_op and one in every 50 models for the 558 closed bound GGBP conformation (holo clGGBP). For all structures, 559 the values of hinge  $\theta$  and twist  $\phi$  angles were calculated using an in- house program in Matlab to adequately represent the conformational landscape in terms of the hinge and twist angles. For the evaluation of the tag-domain data, the structures were aligned with respect to the backbone of residues 112−254. The tag molecule was accommodated for all of the structures, and a rapid minimization on the tag region was made. The structures were visualized and evaluated by using the programs VMD and Discovery Studio.

#### <sup>567</sup> ■ ASSOCIATED CONTENT

#### 568 **6 Supporting Information**

<sup>569</sup> The Supporting Information is available free of charge on the <sup>570</sup> [ACS Publications website](http://pubs.acs.org) at DOI: [10.1021/acschem-](http://pubs.acs.org/doi/abs/10.1021/acschembio.6b00148)<sup>571</sup> [bio.6b00148](http://pubs.acs.org/doi/abs/10.1021/acschembio.6b00148).

- <sup>572</sup> Synthesis and chemical characterization of the tag <sup>573</sup> molecule; description of the MSpin calculations; and <sup>574</sup> detailed tables with the PCSs and RDCs measured with 575 the two metal ions  $(La^{3+}$  and  $Dy^{3+})$  and parameters of
- 576 the calculated  $\Delta \chi$  tensor [\(PDF](http://pubs.acs.org/doi/suppl/10.1021/acschembio.6b00148/suppl_file/cb6b00148_si_001.pdf))
- <sup>577</sup> Movie illustrating the dynamic behavior of GGBP free in <sup>578</sup> solution and in the presence of the ligand [\(MP4\)](http://pubs.acs.org/doi/suppl/10.1021/acschembio.6b00148/suppl_file/cb6b00148_si_002.mp4)

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#### 583 Notes

<sup>584</sup> The authors declare no competing financial interest.

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