

MPAS-Ocean Simulation Quality for Variable-Resolution North American Coastal Meshes

Kristin E. Hoch^{1*}, Mark R. Petersen¹, Steven R. Brus², Darren Engwirda^{4,5},
Andrew F. Roberts², Kevin L. Rosa^{1,3}, Phillip J. Wolfram²

¹ Computational Physics and Methods (CCS-2), Los Alamos National Laboratory, Los Alamos, NM, USA

² Fluid Dynamics and Solid Mechanics (T-3), Los Alamos National Laboratory, Los Alamos, NM, USA

³ Graduate School of Oceanography, University of Rhode Island, Narragansett, Rhode Island, USA

⁴ NASA Goddard Institute for Space Studies, New York, NY, USA

⁵ Center for Climate Systems Research, Columbia University, New York, NY, USA

Key Points:

- Regionally-refined MPAS-Ocean simulations are comparable to global high resolution simulations for numerous metrics.
- Variable-resolution unstructured Voronoi meshes created using JIGSAW are evaluated for quality.
- Simulation quality remains high for steep resolution transitions and intentionally degraded meshes.

*

Corresponding author: Kristin E. Hoch, kehoch@lanl.gov

Abstract

Climate model components utilizing unstructured meshes enable variable-resolution, regionally enhanced simulations within global domains. Here we investigate the relationship between mesh quality and simulation statistics using the JIGSAW unstructured meshing library and the Model for Prediction Across Scales-Ocean (MPAS-Ocean). In the base configuration, the refined region employs 8 km cells that extend 400 km from the coast of North America. This coastal refined region is embedded within a low-resolution global domain, with cell size varying latitudinally between 30 and 60 km. The resolution transition region between the refined region and background mesh is 600 km wide. Three sensitivity tests are conducted: 1) the quality of meshes are intentionally degraded so that horizontal cells are progressively more distorted; 2) the transition region from high to low resolution is steepened; and 3) resolution of the coastal refinement region is varied from 30 km to 8 km. Overall, the ocean simulations are shown to be robust to mesh resolution and quality alterations. Meshes that are substantially degraded still produce realistic currents, with Southern Ocean transports within 0.4% and Gulf Stream transports within 12% of high-quality mesh results. The narrowest transition case of 100 km did not produce any spurious effects. Refined regions with high resolution produce eddy kinetic energy and sea surface height variability that are similar to the high-resolution reference simulation. These results provide heuristics for the design criteria of variable-resolution climate model domains.

Plain Language Summary

Computer simulations used to study the ocean use grids that cover the ocean's surface, and computations are conducted in each grid cell. The smaller these cells are, the more detailed the simulation is, but simulations with more cells are more expensive to run. We experiment with adding small cells in the region of interest, in this case the North American coast, and larger cells in the rest of the ocean. We conducted three series of tests: 1) We wanted to know how much adding these small cells improved the simulation. We changed the size of the coastal cells from 30 km wide (less detailed) to 8 km wide (more detailed). Smaller cells improved the results along the North American coast. 2) We cannot go straight from the small to large cells, and must have intermediate-size cells in between. We experiment with different numbers of these intermediate transition cells. The more intermediate cells we added, the better the results were. 3) We wanted to know whether the cells have to be a regular shape in order to get good results. We experimented with irregular cell shapes. The irregular cells produced results that were very similar to the regular cells.

1 Introduction

Climate models based on unstructured horizontal meshes have matured in recent years. Unstructured global simulations of historical periods compare well when validated against observations and against other future climate projections [Golaz *et al.*, 2019; Petersen *et al.*, 2019; Scholz *et al.*, 2019]. Unstructured meshes offer great freedom in placing resolution in the areas of interest for regionally-refined simulations and also suggest the possibility of improving global simulation quality with targeted areas of high resolution. However, modelers now have a dizzying array of choices to make in designing their meshes, compared to the limited variations of stretched aspect ratio in latitude-longitude-type quadrilateral grids. Furthermore, the role of regional refinement strategies on simulation quality is currently largely unknown.

There is a pressing need for constraints on mesh design and model configuration criteria that are informed by how local resolution affects simulation quality. However, time constraints and available computational resources generally allow only a limited number of configurations to be rigorously tested. In this study, we explore the role of mesh design and quality on various ocean simulations metrics using the Model for Prediction Across Scales

66 (MPAS) [Ringler *et al.*, 2013] with the goal of providing guidance on the design of meshes
67 for variable resolution climate models.

68 The generation of high quality unstructured meshes for General Circulation Models
69 (GCMs) is a challenging problem, and a new generation of mesh creation tools have been de-
70 veloped to satisfy the needs of high-resolution unstructured-mesh models. This paper docu-
71 ments the use of JIGSAW [Engwirda, 2017] to produce optimized spherical Voronoi/Delaunay
72 meshes for use with MPAS. MPAS-Ocean and MPAS-sea ice are components of the Depart-
73 ment of Energy’s Energy Exascale Earth System Model (E3SM)¹ [Golaz *et al.*, 2019; Pe-
74 tersen *et al.*, 2019; Scholz *et al.*, 2019].

75 An ensemble of horizontal meshes was investigated using the Coastal United States
76 ‘Plus’ (CUSP) configuration, which is designed to enhance the resolution of coastal regions
77 of North and Central America plus Hawaii. Three case studies were performed: one where
78 global mesh quality is intentionally degraded; a second where the resolution transition width
79 is varied in the CUSP mesh; and a third where the coastal-refined region is tested at a num-
80 ber of resolutions. In each case, a family of meshes was generated and the results of a ten-
81 year simulation were analyzed, allowing for the convergence of model metrics to be assessed
82 with respect to perturbations in the underlying grid and model configuration. Using this data,
83 modelers can assess which mesh characteristics are most important for the needs of their ap-
84 plication and inform their choices for the design of future configurations.

85 We aim to highlight the impact of various mesh characteristics on simulation quality
86 and to document how different choices in mesh design feed back onto the simulated state.
87 We focus on the geometric ‘quality’ of a mesh, its rate of transition from regions of low
88 to high resolution, and the placement of high resolution near energetic boundary currents
89 and areas of interest. The configurations used in this paper enhance resolution of the North
90 American coastal region, but the aim is to provide general guidelines that may be applied to
91 the design of any variable-resolution mesh.

92 This paper is structured as follows. Section 2 reviews the state of variable resolution
93 meshes on ocean modeling. Section 3 introduces MPAS-Ocean, JIGSAW, and the details of
94 the meshes created for this work. Section 4 presents the analysis of global simulations for the
95 three sensitivity studies. Based on this evidence, the paper concludes with recommendations
96 for mesh generation criteria in Section 5.

97 **2 Background**

98 There now exists a growing selection of unstructured-mesh models that are used for
99 various global and regionally-focused forecasts and analyses. This includes MPAS [Ringler
100 *et al.*, 2013], FESOM [Androsov *et al.*, 2019], ICON [Korn, 2017], FVCOM [Chen *et al.*,
101 2003], SLIM [White *et al.*, 2008], and Fluidity [Davies *et al.*, 2011]. Models differ in the
102 arrangement of variables on the underlying computational grid and in the numerical tech-
103 niques employed, with both unstructured triangle- and polygon-based finite-volume and
104 finite-element type discretization schemes adopted in various frameworks. As such, differ-
105 ent approaches to the construction and optimization of the models’ underlying unstructured
106 meshes have been explored, including techniques based on Centroidal Voronoi Tessellation
107 (CVTs) [Jacobsen *et al.*, 2013; Yang *et al.*, 2018], optimization via optimal transport [Weller
108 *et al.*, 2016; McRae *et al.*, 2018], as well as triangulation-based refinement schemes [Lam-
109 brechts *et al.*, 2008; Remacle and Lambrechts, 2018]. In the context of MPAS-Ocean, the
110 numerical scheme requires that the mesh define a highly regular, orthogonal tessellation,
111 constraining grid generation choice to algorithms that can generate optimized Voronoi-type
112 meshes [Jacobsen *et al.*, 2013].

¹ <https://e3sm.org>

113 Variable resolution is advantageous in situations where highlighting a region may help
114 to correct a bias or resolve a dynamic condition. In many cases, the resolved region will also
115 be the focus of the investigation, but resolution can also be placed to correct a bias that is
116 impacting a global simulation. They can serve as a replacement for nested grids, with the ad-
117 vantage that variable resolution can be applied in more complex configurations and is more
118 integrated with the global simulation [Hagos *et al.*, 2013; Biastoch *et al.*, 2018]; furthermore,
119 unstructured meshes avoid coupling challenges inherent to nesting. Meshes in which the res-
120 olution varies as a function of latitude have been used to compensate for the changing Rossby
121 radius with latitude. This approach is used in the standard high-resolution MPAS mesh [Pe-
122 tersen *et al.*, 2019]. Variable-resolution meshes are designed to improve the dynamics of a
123 particular region or process, and also to provide good global dynamics. Variable resolution
124 meshes may refine particular regions, for example, the Arctic Ocean [Wang *et al.*, 2018] or
125 a coastal region [Androsov *et al.*, 2018]. They can also place resolution based on a particu-
126 lar parameter. For example, FESOM uses meshes that refine to the local Rossby radius [Sein
127 *et al.*, 2017], a more sophisticated approach than refining based on latitude alone. FESOM
128 also uses meshes that refine based on sea surface height (SSH) variability, which is useful for
129 capturing boundary currents [Biastoch *et al.*, 2018].

130 Because of the computational cost and complexity of global simulations, the majority
131 of variable resolution tests have been performed on idealized or simplified domains. For ex-
132 ample, in order to eliminate the effects of continental geography, many tests have used aqua-
133 planet configurations [Abiodun *et al.*, 2008; Rauscher and Ringler, 2014; Lorant and Royer,
134 2001; Rauscher *et al.*, 2012; Hagos *et al.*, 2013; Zhao *et al.*, 2016]. Others have used two-
135 dimensional domains [Düben and Korn, 2014]. These simplified domains can demonstrate
136 the effects of mesh resolution independent of other variables. Additionally, atmospheric
137 variable-resolution simulations can inform choices in ocean domains [Abiodun *et al.*, 2008;
138 Düben and Korn, 2014; Park *et al.*, 2014; Zarzycki *et al.*, 2015; Rauscher and Ringler, 2014;
139 Zhao *et al.*, 2016]. However, mesh-resolution and design consequences on more-realistic
140 simulations are still largely unknown, even though use of variable resolution in realistic sim-
141 ulations is becoming more widespread.

142 While mesh design is still a developing field, the literature points to several important
143 considerations. In the past, parameter values for sub-grid scale physics were typically tuned
144 for each resolution. Now, for variable-resolution meshes, parameterization schemes must
145 work well across the span of grid-cell sizes. Another consideration is that variable resolu-
146 tion results compared against uniform high-resolution simulations may not necessarily be
147 comparable near mesh transition regions. For example, a current flowing from a non-eddy
148 permitting to an eddy permitting region may not immediately develop eddies. Instead, eddies
149 will develop downstream of the beginning of the high resolution region once perturbations
150 have time to evolve [Danilov and Wang, 2015]. A similar result was found in atmospheric
151 variable resolution aquaplanet simulations, in which precipitation error was decreased in the
152 eastern (downstream) section of the high resolution region, but not in the western (upstream)
153 section [Hagos *et al.*, 2013].

154 A high resolution region will also have effects on the rest of the domain. Most obvi-
155 ously, a high resolution region will have an effect immediately downstream, as the increased
156 variability of the high resolution region is carried into the low resolution region [Danilov and
157 Wang, 2015]. Changes to dynamics within the high resolution region can propagate to other
158 global processes [Lorant and Royer, 2001; Hagos *et al.*, 2013; Sein *et al.*, 2017; Sakaguchi
159 *et al.*, 2016]. Conversely, the impact of the global domain on the high resolution region is
160 also important. A high resolution region can decrease local error, but will have a limited im-
161 pact on processes that are due to causes outside the high resolution region [Zarzycki *et al.*,
162 2015].

3 Methods

3.1 The Model for Prediction Across Scales-Ocean (MPAS-O)

The Model for Prediction Across Scales (MPAS) is an open source framework that provides common functionality for climate model components on unstructured meshes. This includes a mesh specification, decomposition of variables across processors, parallel input and output specified in a run-time streams file, timers, and error handling. Finite volume operators were developed for Voronoi tessellations in *Ringler et al.* [2010] for the shallow water equations using mimetic methods to guarantee that mass, velocity and potential vorticity evolve in a consistent and compatible manner.

MPAS-Ocean solves prognostic equations for momentum, thickness (volume), and tracers using these operators [*Ringler et al.*, 2013] and can be run using both regular and unstructured meshes on Cartesian and spherical domains. The time stepping is split-explicit, where the 2D barotropic equations are sub-cycled within 3D baroclinic time steps. MPAS-Ocean uses an Arbitrary Lagrangian-Eulerian (ALE) method for the vertical coordinate [*Petersen et al.*, 2015; *Reckinger et al.*, 2015]. It is typically run with 60 to 100 layers using a z-star vertical coordinate, where layer thickness varies in proportion with sea surface height, varying from 2 m thick at the surface to 150 m thick at a depth of 5000 m.

The vertical mixing scheme is the K-Profile Parameterization (KPP) [*Van Roedel et al.*, 2018]), calculated in the CVMix library² and applied implicitly. The horizontal eddy mixing scheme is Gent-McWilliams thickness advection [*Gent and McWilliams*, 1990], applied to variable-resolution meshes with a coefficient of $600 \text{ m}^2 \text{ s}^{-1}$ at gridcells larger than 30 km, and tapering linearly to zero between 30 and 20 km. Viscosity (del-2) and hyperviscosity (del-4) are applied to the momentum equation with coefficients that depend on the grid cell size as

$$\nu_2 = 1000[\text{m}^2 \text{ s}^{-1}] \frac{\Delta x}{30[\text{km}]} \quad (1)$$

$$\nu_4 = 1.2e11[\text{m}^4 \text{ s}^{-2}] \left(\frac{\Delta x}{30[\text{km}]} \right)^3, \quad (2)$$

respectively, where Δx is the horizontal gridcell width [*Ringler et al.*, 2013; *Petersen et al.*, 2015]. The tracer advection scheme is Flux Corrected Transport [*Skamarock and Gassmann*, 2011], and no horizontal diffusion is explicitly applied to the tracers.

For this study MPAS-Ocean was run in stand-alone mode with idealized, constant atmospheric forcing, where wind forcing is averaged over a 65-year CORE cycle [*Griffies et al.*, 2009]. The simulation is spun up for one year from an initial climatology of Polar Science Center Hydrographic Climatology, version 3 (PHC3.0, *Steele et al.* [2001]). Surface salinity and temperature restoring to yearly-averaged PHC3.0 is conducted with a piston velocity of $1.585e-5 \text{ m s}^{-1}$ to represent surface fluxes. Sea-ice is not included in these simulations.

This idealized set-up was chosen to evaluate the effects of a large number of mesh variations in short, standardized simulations using the MPAS-Ocean stand-alone configuration. Simulations with more realistic atmospheric forcing (six-hourly CORE winds and surface fluxes) and active sea ice have been run within E3SM using the coastal-refined mesh (CUSP8) are currently underway and will be presented in a future publication.

3.2 JIGSAW mesh generation

JIGSAW is an unstructured meshing library designed to generate high quality grids for computational simulation, with a focus on constructing optimized Voronoi-type grids for

² <https://github.com/CVMix/CVMix-src>, <https://doi.org/10.5281/zenodo.1000800>

205 unstructured-mesh GCM's. JIGSAW is a hybrid algorithm that combines both Delaunay-
206 refinement and Voronoi optimization type approaches to enable the rapid generation of very
207 high quality, high resolution Voronoi/Delaunay meshes on the sphere. A key advantage of
208 this combined strategy is efficiency and guaranteed mesh quality. Previous mesh generation
209 methods used in MPAS [Jacobsen *et al.*, 2013] used an iterative Lloyd's method, and were
210 extremely slow.

211 With JIGSAW, highly optimized, large-scale variable resolution Voronoi-type meshes
212 can be generated in the order of minutes, allowing model users to easily create and explore
213 a range of alternative configurations, investigate mesh quality and resolution dependence,
214 and tailor the overall mesh and model configuration to their simulation needs. This capability
215 was exploited in the present study to design and assess a range of coastal-enhanced MPAS-
216 Ocean configurations and to explore various model/mesh feedbacks.

217 Meshes can be generated in local two-dimensional domains and over general spheroidal
218 surfaces. Mesh resolution can be adapted to follow complex user-defined metrics, including
219 topographic contours, solution profiles and/or coastal features. This flexibility enables the
220 construction of complex, variable resolution model configurations, offering enhanced simula-
221 tion fidelity in regions of interest or importance.

222 Given a particular geometry definition and resolution specification, JIGSAW proceeds
223 to assemble the unstructured mesh incrementally—first creating a conforming Delaunay
224 triangulation of the domain using a 'frontal' Delaunay-refinement strategy [Engwirda and
225 Ivers, 2016], before optimizing the resulting Voronoi/Delaunay tessellation using Optimal
226 Delaunay Tessellation (ODT) type techniques [Chen and Holst, 2011; Engwirda, 2017].
227 The final mesh is guaranteed to consist of high quality triangular and polygonal cells that
228 form a locally orthogonal unstructured C-grid staggering. The final meshes are heavily op-
229 timized, typically satisfying the stringent mesh quality requirements imposed by the TRiSK
230 discretization scheme [Ringler *et al.*, 2010] used in MPAS-Ocean.

231 For TRiSK-based schemes, a complex array of geometrical and topological constraints
232 must be satisfied [Engwirda, 2018], requiring tessellations be orthogonal, centroidal, well-
233 centered and smoothly varying. These criteria require that the vertices of the triangular and
234 polygonal grid cells lie close to the centroids of their enclosing control-volumes, that the
235 staggered Voronoi and Delaunay edges intersect near their midpoints, that the Delaunay tri-
236 angles contain their own circumcenters, and that the cell angles and edge-lengths be 'nicely'
237 distributed with respect to the desired mesh resolution constraints. Satisfying such criteria
238 is nontrivial, and failure to do so has been shown to impact on the asymptotic accuracy and
239 stability of the underlying numerical scheme [Peixoto, 2016] in idealized cases.

240 The expected accuracy of the TRiSK formulation is thus a function of both the geome-
241 try and topology of the mesh, and can be quantified by considering the nature of the discrete
242 gradient, divergence, curl and interpolation operators used to discretize the continuous PDE's
243 [Ringler *et al.*, 2010; Engwirda, 2018]. Based on theoretical analysis, it is expected that the
244 accuracy of TRiSK is maximized (achieving quasi 2nd-order scaling) only for 'perfect' tes-
245 sellations consisting of regular hexagons and equilateral triangles. For general unstructured
246 meshes incorporating irregular and/or deformed polygonal and triangular cells, numerical
247 accuracy is expected to degrade—leading to quasi 1st-order behavior in many practical con-
248 figurations [Peixoto, 2016]. The goal of mesh optimization is to construct a tessellation that
249 serves to minimize these numerical errors, thus maximizing the quality of the resulting simu-
250 lation.

251 A key question in the current study is to assess what impact mesh quality has on practi-
252 cal MPAS-Ocean simulations and to define an associated set of 'best practice' guidelines for
253 mesh generation. To this end, an 'ensemble' of meshes was considered in the current work—
254 exploring the impact of different mesh quality perturbations and variable-resolution designs
255 on the characteristics of spun-up ocean simulations.

256

3.3 Meshes and simulations

257

258

259

260

All the meshes used are based on two base configurations, a global low resolution mesh and a mesh with refinement along the coast of North America. The global low resolution mesh, EC60to30, varies from 30 km resolution at the equator and poles to 60 km resolution at the mid-latitudes and uses 100 vertical layers.

263

264

265

266

267

268

269

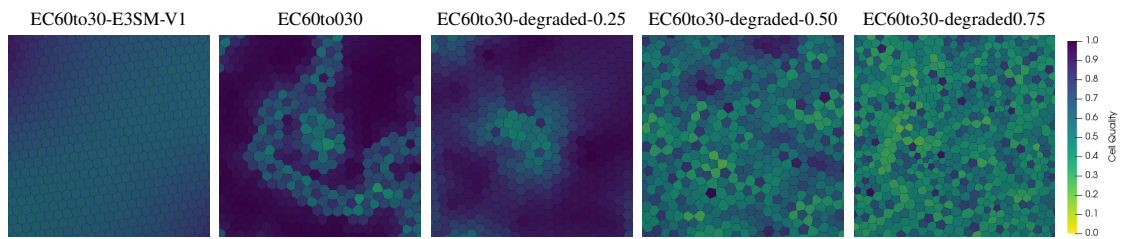
270

271

272

273

The base EC60to30 mesh created using JIGSAW was compared against the EC60to30-E3SM-V1 mesh created using a parallel Lloyd's algorithm [Jacobsen *et al.*, 2013], which was used in previously published E3SM simulations [Petersen *et al.*, 2019; Golaz *et al.*, 2019]. Images of the two EC60to30 meshes can be seen in the first two panels of Figures 1 and 2, which show two different metrics for measuring cell quality. Figure 2 shows the percent change between the size of neighboring cells and Figure 1 shows close up images of the mesh and the ratio of the smallest to largest sides of the cells. These metrics show the different strategies used by each of the mesh creation methods. In order to cover the sphere, the mesh must deviate from regular hexagons. E3SM-V1 spreads these imperfections between large numbers of cells, resulting in smooth regions of lower quality cells. JIGSAW concentrates the imperfections into "seams" of low quality cells separating regions of very high quality cells.

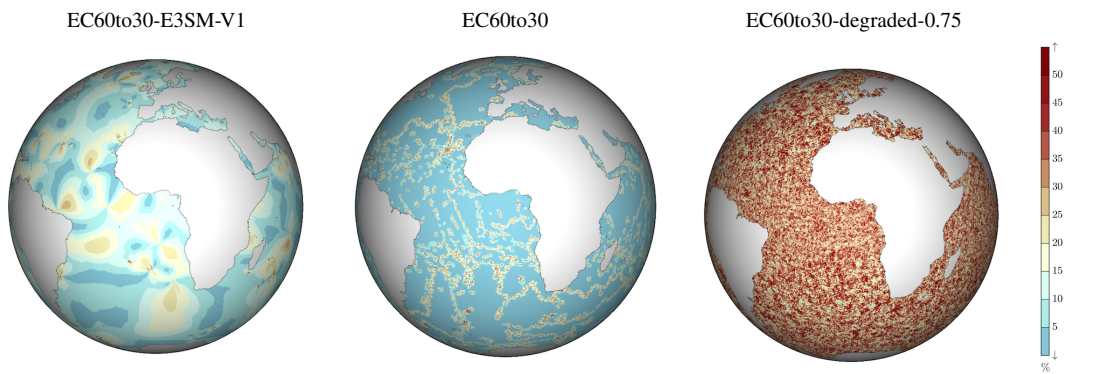


261

262

Figure 1. Cell quality of the degraded meshes. A small region of the mesh is shown. Cell quality is the ratio of the smallest to largest sides of a cell, 1.0 being a perfect polygon.

274



275

Figure 2. Percent change in grid cell area between neighboring cells.

276

277

278

279

The second base mesh is the North American refined mesh, created to investigate processes affecting North American coastal regions at high resolution while avoiding the cost of running a global high resolution model. In addition to the improvements in the dynamics of the Gulf Stream investigated in this study, using the CUSP8 mesh will allow improved simu-

280 lation of a variety of coastal processes around North America. The CUSP8 mesh (Coastal
 281 United States ‘Plus’ with 8 km coastal resolution) has high resolution along the Atlantic
 282 and Pacific coasts from Central America to the Arctic, with additional high resolution in the
 283 Caribbean and around Greenland, Hawaii and the Bering Strait (see Figure 3). The CUSP8
 284 mesh is built on top of a background low resolution EC60to30 mesh. It uses 80 vertical lay-
 285 ers.

286 In the CUSP8 mesh, the transition between the high resolution region and background
 287 mesh begins 400 km off the coast and is 600 km wide according to the following functions,

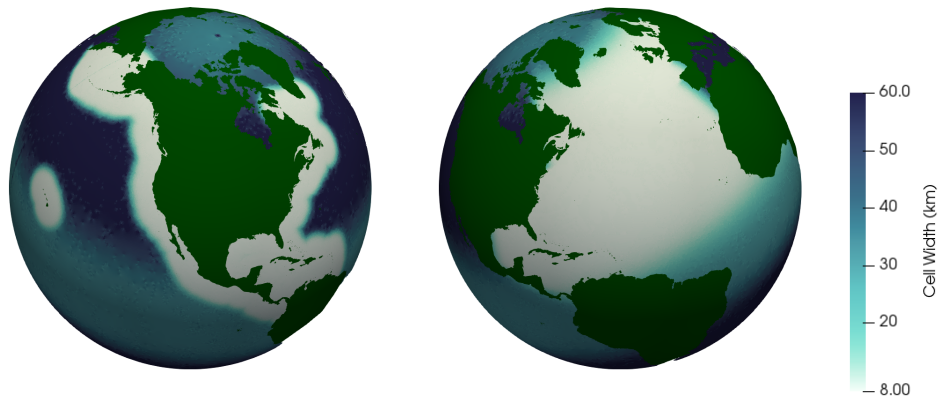
$$W = 0.5 \left(\tanh \frac{D - D_{start} - 0.5D_{width}}{0.2D_{width}} + 1 \right) \quad (3)$$

$$C = C_{coast} (1.0 - W) + C_{back} W \quad (4)$$

288 where W is the weight, D is the distance from the coast, D_{start} is the distance from the coast
 289 where the transition region begins, and D_{width} is the transition width. The final cell width,
 290 C , shown in Figs. 3 and 4, is simply a linear combination of the coastal and background cell
 291 widths, C_{coast} and C_{back} .

292 In addition to these two base meshes, a mesh with 8 km resolution spanning the full
 293 North Atlantic basin (NA8) was created. Like the CUSP8 mesh, it was built on a background
 294 EC60to30 mesh (see Figure 3). A global high resolution simulation was not feasible for this
 295 study, but the NA8 mesh provides high resolution within the region of interest in the North
 296 Atlantic, providing a benchmark for the performance of the CUSP8 mesh.

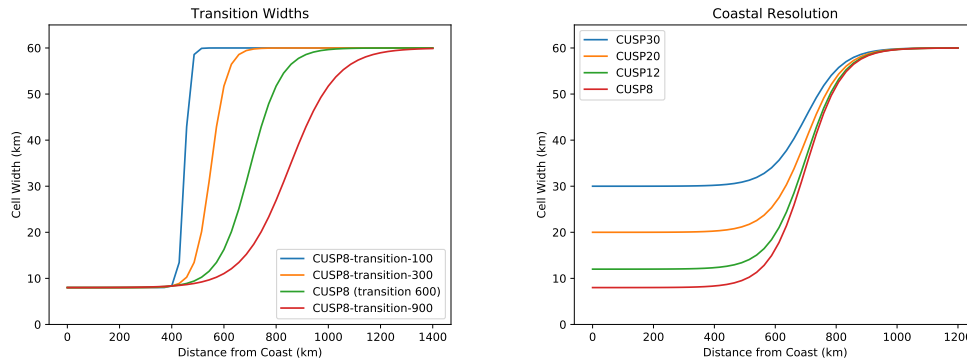
297 In order to ensure that all the meshes could be compared, EC60to30 simulations were
 298 run in each vertical configuration: 60, 80, and 100 layers. All three EC60to30 meshes per-
 299 formed similarly in terms of kinetic energy (KE), sea surface height (SSH), eddy kinetic
 300 energy (EKE) and sea surface height root mean squared (SSH RMS) (see Figure 16 in Ap-
 301 pendix).



302 **Figure 3.** The Coastal United States ‘Plus’ mesh (CUSP8) on the left and the North Atlantic refined mesh
 303 (NA8) on the right. The white areas show the 8 km high resolution regions. The blues show the background
 304 EC60to30 low resolution mesh, with 30 km resolution at the tropics and poles (light blue) and 60 km resolu-
 305 tion in between (dark blue).

306 Three studies were performed to investigate mesh features and their effects on simula-
 307 tion quality.

308 The first study uses the EC60to30 mesh to examine the effect of poor mesh quality on
 309 simulations. Meshes were intentionally degraded, producing poor quality cells. Variable



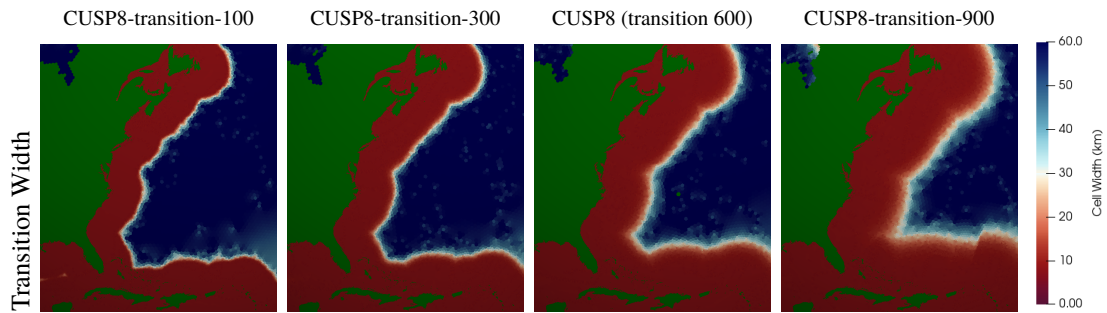
329 **Figure 4.** Plots of the transition function (Equation 4) for the transition width study (left) and the conver-
 330 gence study(right). The background resolution plotted is 60 km, however, the background resolution varies
 331 from 30 km to 60 km depending on latitude.

310 resolution meshes by necessity contain distorted cells within the transition regions. This is a
 311 particular concern when designing complex meshes such as the CUSP8 mesh that have large
 312 variations in resolution and relatively narrow transition regions. Because of the difficulty
 313 of decoupling the effects of poor cell quality from the effects of a change in resolution, the
 314 effect of poor cell quality on simulations was investigated using EC60to30 meshes with cell
 315 quality degraded globally.

316 A mesh degradation heuristic was developed to systematically reduce the quality of
 317 meshes, perturbing the position of vertices and updating topology to effectively ‘de-optimize’
 318 the overall structure of a given mesh and degrade the shape of its cells. Care was taken to en-
 319 sure that degraded meshes inherited the large-scale properties of their parent grids, adhering
 320 to variations in resolution and matching cell counts exactly. The kernel of the degradation
 321 operation consisted of randomly perturbing a subset of vertices toward the centroid of their
 322 largest neighboring triangle. By controlling the magnitude of the average relative vertex per-
 323 turbation, the notion of a ‘ β -degraded’ mesh was introduced — a 0.5-degraded mesh would
 324 re-position vertices (on average) halfway between their current position and the neighbor-
 325 ing centroid location. Mesh topology was updated following the re-positioning of vertices
 326 to ensure the orthogonality of the mesh was preserved. Starting from a fully optimized ini-
 327 tial mesh, several iterations of this process were repeated to ensure that degraded grids were
 328 sufficiently randomized.

332 Three degraded meshes were created, EC60to30-degraded-0.25, EC60to30-degraded-
 333 0.50, and EC60to30-degraded-0.75, with larger degradation fractions indicating a more de-
 334 graded mesh. Figure 1 shows the mesh quality of the standard EC60to30 mesh and the de-
 335 graded meshes.

336 The second study investigates the effects of the steepness of the transition function in
 337 the CUSP8 mesh (Equation 3) by varying the transition width from 100 km to 900 km (Fig-
 338 ure 4). A 10 km transition was attempted as well, but failed early in the spin-up process.
 339 This study was designed to investigate how steep the transition function could be without
 340 negatively affecting the simulation quality. In addition to exploring the steepness of the tran-
 341 sition function, this study also investigates the impact of the size of the higher resolution re-
 342 gion. Because the beginning of the transition region was kept fixed, the center of the tran-
 343 sition region and the beginning of the low resolution region were closer to the coast for steeper
 344 transitions, effectively shrinking the higher resolution region (see Figure 5).



345 **Figure 5.** A view of the East Coast showing the different transition widths used. The transition begins at
 346 400 km off the coast for all transition widths. Note that the size of the higher resolution region is expanded
 347 with a wider transition.

348 The third study investigates different coastal resolutions ranging from 8km (CUSP8)
 349 to 30km (CUSP30) in order to explore the improvements in dynamics with increased res-
 350 olution. The computational performance of the meshes was also examined in order to give
 351 a better sense of the trade-off between higher resolution and higher simulation cost. These
 352 meshes were compared against the EC60to30 and NA8 meshes. Ideally, the CUSP8 mesh
 353 would show dynamics comparable to the NA8 mesh within the high resolution region with a
 354 much lower cost than a global high resolution mesh.

355 Table 1 shows the parameter values used for each simulation. These values were cho-
 356 sen based on the highest resolution region of the simulation. The EC60to30-degraded-0.50
 357 and EC60to30-degraded-0.75 meshes had to be run at a smaller timestep than the standard
 358 EC60to30 meshes due to the smaller cell sizes introduced by the degradation process. All
 359 meshes were run with a 7 day spin up except the EC60to30-E3SM-V1, EC60to30-degraded-
 360 0.50, and EC60to30-degraded-0.75 meshes. The EC60to30-E3SM-V1 mesh used a 21 day
 361 spin up process. The EC60to30-degraded-0.50 and EC60to30-degraded-0.75 meshes re-
 362 quired longer spin ups and were spun up to a different point because of the smaller timestep
 363 required.

367 4 Results and Discussion

368 The analysis focuses on the Gulf Stream because it is the most prominent feature within
 369 the high resolution region of the CUSP simulations. The Gulf Stream also crosses out of
 370 the high resolution region, allowing the effect of the transition in resolution to be inves-
 371 tigated. The sea surface height, kinetic energy, sea surface height root mean squared, and
 372 eddy kinetic energy were analyzed for all simulations. Transport through transects along the
 373 Gulf Stream was calculated (see Figure 6 for a map of the Gulf Stream transects). Transport
 374 through Southern Ocean transects were also calculated in order to see if the high resolution
 375 region had an impact on global dynamics (see Table 2 and Figure 7 for the transect results).
 376 SSH RMS and EKE were averaged along the Gulf Stream region (see Figure 6). These re-
 377 sults are not expected to closely match observations, both because of the idealized forcing
 378 used and because of the differences between the sampling techniques used to calculate obser-
 379 vational estimates and those used in our calculations. Global analysis was also run looking
 380 at global temperature, salinity, SSH, and EKE. However, because of the extremely similar
 381 results for all simulations, this paper focuses only on analysis of the areas within and around
 382 the high resolution region. Preliminary results from simulations with realistic climatological
 383 forcing are also used to give an indication of how CUSP meshes perform in realistic climate
 384 simulations. Further results will follow in subsequent papers.

study	mesh name	refined resolution km	number of cells thousands	vertical layers	transition width km	degradation factor	time step min:sec	barotropic step min:sec
reference meshes	EC60to30	none	236	100	none	none	30:00	1:00
	CUSP8	8	649	80	600	none	7:30	00:15
	NA8	8	842	80	600	none	7:30	00:15
	EC60to30-E3SM-V1	none	235	100	none	none	20:00	1:00
degraded meshes	EC60to30 (not degraded)	none	236	100	none	<i>none</i>	30:00	1:00
	EC60to30-degraded-0.25	none	237	100	none	<i>0.25</i>	30:00	1:00
	EC60to30-degraded-0.50	none	248	100	none	<i>0.50</i>	20:00	0:40
	EC60to30-degraded-0.75	none	338	100	none	<i>0.75</i>	2:00	0:06
transition width	CUSP8-transition-900	8	700	80	<i>900</i>	none	7:30	00:15
	CUSP8 (transition 600)	8	649	80	<i>600</i>	none	7:30	00:15
	CUSP8-transition-300	8	603	80	<i>300</i>	none	7:30	00:15
	CUSP8-transition-100	8	574	80	<i>100</i>	none	7:30	00:15
coastal resolution	CUSP8	8	649	80	600	none	7:30	00:15
	CUSP12	<i>12</i>	414	80	600	none	12:00	00:24
	CUSP20	<i>20</i>	295	80	600	none	20:00	00:40
	CUSP30	<i>30</i>	256	80	600	none	30:00	1:00

364 **Table 1.** Simulation parameters. The reference simulations, EC60to30 and CUSP8, are bold. The varied
365 parameter for each study is in italics. Timestep values were chosen based on the smallest resolution present in
366 the mesh.

385 The comparison of the JIGSAW EC60to30 mesh and the EC60to30-E3SM-V1 mesh
386 showed that they performed very similarly, confirming that the meshes created using JIG-
387 SAW produce comparable results to those used in previous MPAS studies (see Figure 15 in
388 the Appendix).

405 4.1 Study 1: Degraded meshes

406 Though the degradation factor for the degraded mesh study and the transition widths
407 for the transition width study were chosen independently, the degraded meshes were found
408 to be a good proxy for the transition regions (see Figure 9). The cell quality in the transition
409 region for the 100 km transition width is comparable to the cell quality in the 0.75 degraded
410 mesh, and the cell quality in the transition region for the 900 km transition is comparable
411 to the cell quality in the 0.25 degraded mesh. Thus, the results of the degraded mesh study
412 should also be considered when interpreting the results within the transition regions of the
413 CUSP meshes.

420 The SSH and KE analysis of the degraded meshes can be found in Figure 11. The de-
421 graded meshes were compared against the normal EC60to30 mesh. Overall, mesh degrada-
422 tion did not significantly affect the quality of the simulations. The pattern and magnitude of
423 sea surface height and kinetic energy for the degraded meshes and EC60to30 mesh are nearly
424 visually identical. The more degraded meshes have higher average sea surface height vari-
425 ability and eddy kinetic energy (see Figure 8). Transport through all the transects measured
426 showed no significant variation between the degraded meshes (see Figure 7). Overall, the de-
427 graded meshes seem to have had no significant impact on the simulations beyond the need
428 for smaller timesteps, which does not impact the results.

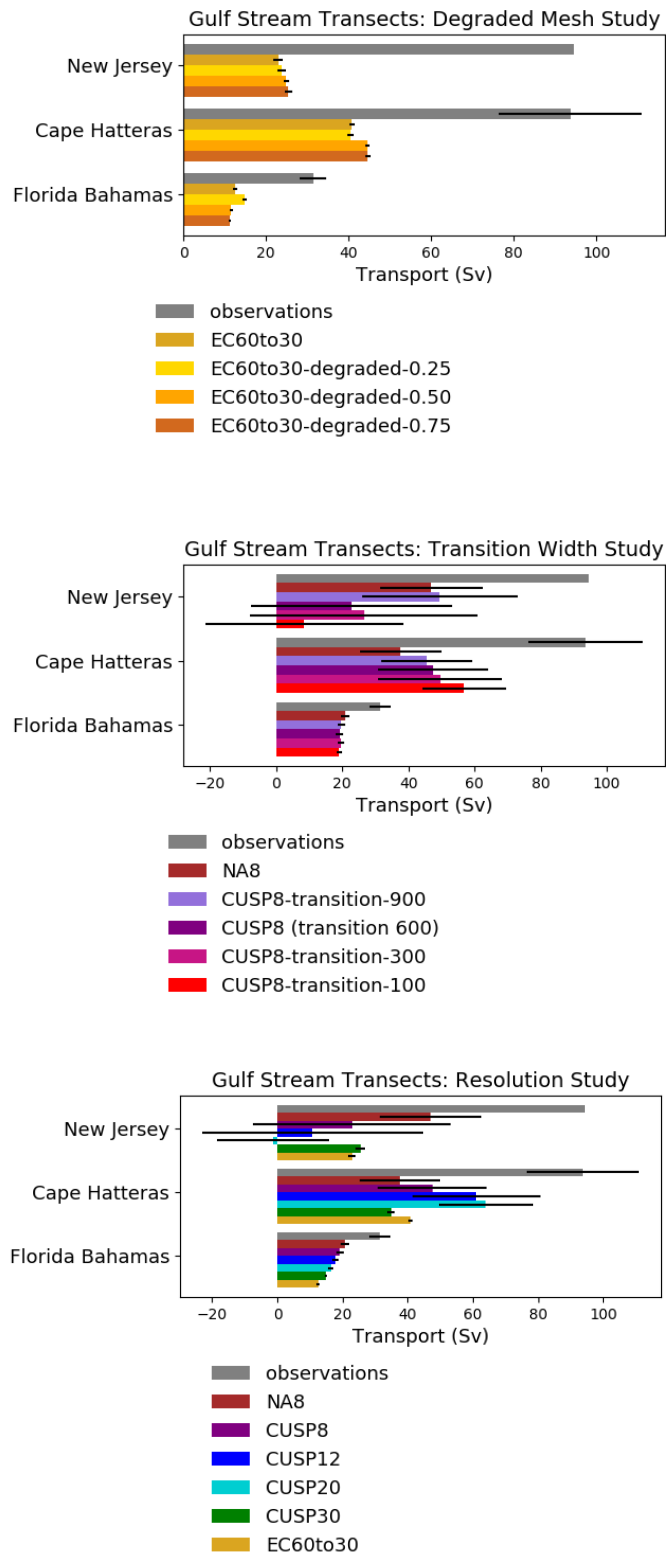


389 **Figure 6.** Shaded region indicates the area in which SSH RMS and EKE averages were computed. Yellow
 390 sections show the locations of transects along the Gulf Stream.

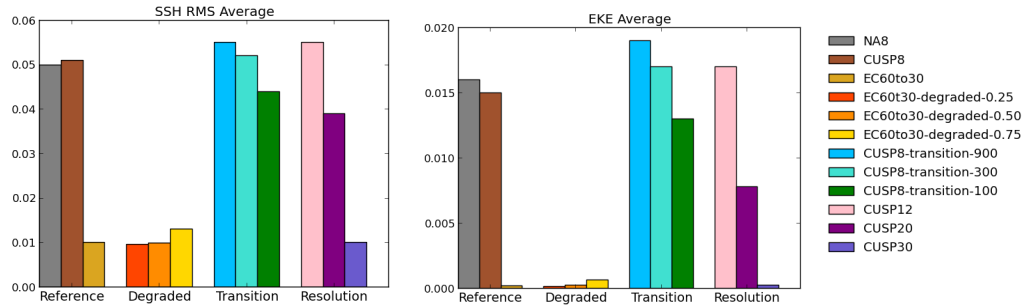
429 4.2 Study 2: Transition width

430 The analysis of the transition width study can be found in Figure 12. As the transition
 431 width increases, the dynamics of the simulations improve. The simulations with wider tran-
 432 sition regions show greater SSH RMS and EKE (see Figure 8). This is to be expected, both
 433 because the transition is less steep, leading to higher quality cells in the transition region, and
 434 because the higher resolution area is effectively larger with a greater transition width (note
 435 the locations of the center of the transition region in Figure 12). The three widest transitions
 436 (900 km, 600 km and 300 km) have closer average values. The 100 km transition, where the
 437 Gulf Stream is meandering into the low resolution region, shows a more significant decline
 438 in average SSH RMS and EKE. Even within 400 km of the coast, where the resolution is 8
 439 km in all the simulations, the dynamics were improved by a wider transition region. With a
 440 narrower transition, meanders and eddies from the Gulf Stream cross into regions of lower
 441 resolution. It appears that these features are then smoothed out and do not have time to re-
 442 cover even when returning to the high resolution region. This result is consistent with that
 443 found by *Danilov and Wang* [2015], in which eddies did not develop at the beginning of the
 444 eddy-permitting region but instead developed only once perturbations had developed further
 445 downstream. This is clearly seen in Figure 12. It also appears possible that the transition re-
 446 gion is affecting the path of the Gulf Stream, "trapping" it within the high resolution
 447 region. However, there is not a wide enough spread of transition widths in this study to say anything
 448 definitive about this effect. The transport through the Gulf Stream transects increases with
 449 wider transition widths, with the exception of the Cape Hatteras transect, which shows the
 450 opposite pattern.

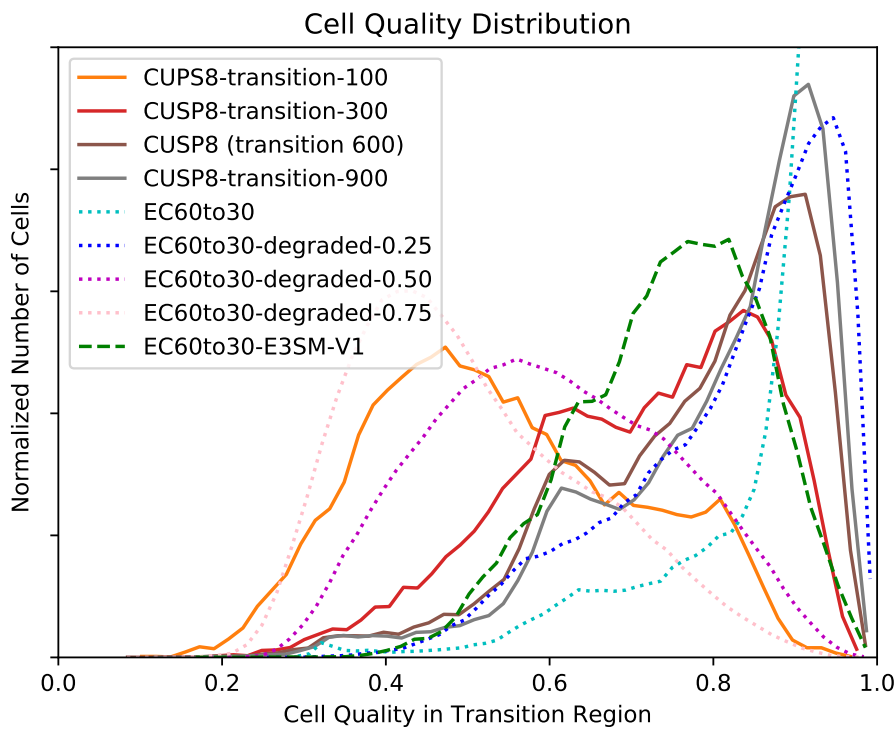
451 In addition to the results examined here, the results of the degraded mesh study should
 452 be considered as a proxy for the transition regions. Although the cell quality within the CUSP8-
 453 transition-100 transition region is comparable to that of the EC60to30-degraded-0.75 mesh,
 454 the CUSP8-transition-100 mesh did not require the smaller timesteps that the EC60to30-
 455 degraded-0.50 and EC60to30-degraded-0.75 meshes did. The results of the degraded mesh
 456 study indicate that mesh quality does not have a large impact on simulation results. The
 457 variation between the CUSP meshes is probably due primarily to other effects, such as the
 458 smaller region of higher resolution, rather than the cell quality within the transition region.



397 **Figure 7.** Transport through transects along the gulf stream. Table 2 shows the data and Figure 6 shows the
 398 locations of the transects.



399 **Figure 8.** Plot of the average surface SSH RMS and EKE over the region shown in Figure 6 for years 2-10.
 400 The CUSP meshes have significantly higher average SSH RMS and EKE than the EC60to30 meshes. The
 401 variability increased as the mesh degradation increased. As the transition width was narrowed, the variability
 402 decreased, though this effect was small between CUSP8-transition-900 and CUSP8-transition-300. As the
 403 resolution decreased, the variability decreased, reaching the same values as the EC60to30 mesh for CUSP30,
 404 as would be expected.



414 **Figure 9.** Plot of cell quality (the ratio of the largest to smallest sides of a cell) in the transition region
 415 and, for comparison, the global cell quality for the global low resolution mesh and the degraded meshes.
 416 The degraded meshes can serve as a proxy for the impact of cell quality in the transition region. Notice
 417 that the cell quality in the transition region of the CUSP8-transition-900 mesh is comparable to that of the
 418 EC60to30-degraded-0.25 mesh and that of the CUSP8-transition-100 mesh is comparable to the EC60to30-
 419 degraded-0.75 mesh.

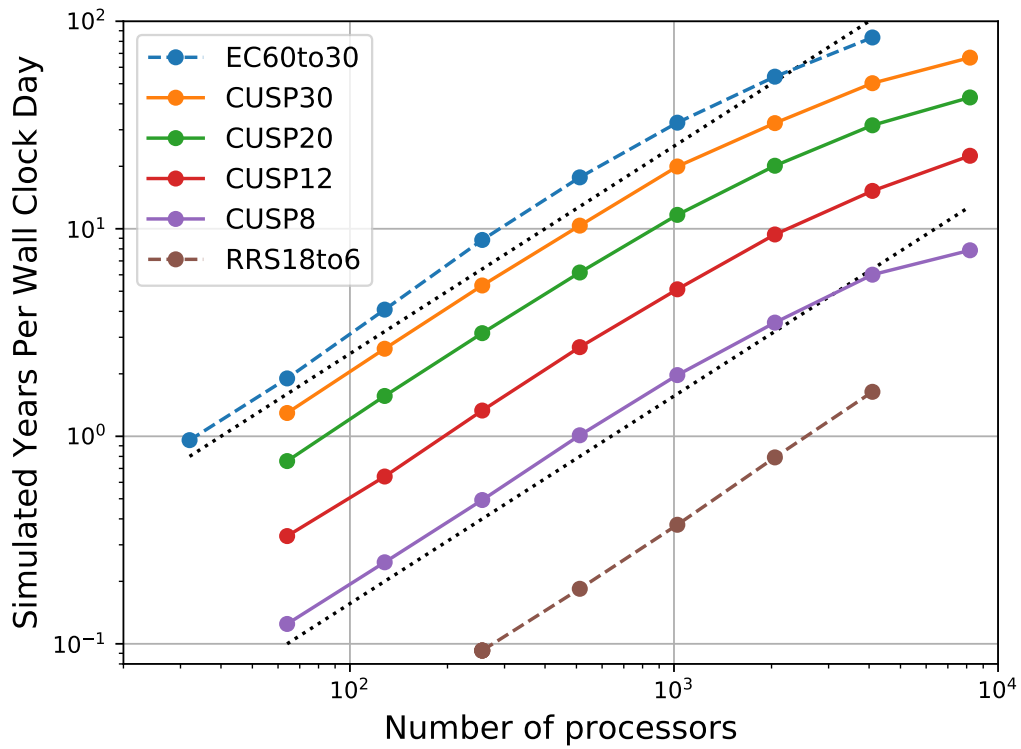
	Florida-Cuba	Florida-Bahamas	Cape Hatteras	New Jersey	Drake Passage	Tasmania-Ant	Africa-Ant
Observation	31.0 ± 1.5	31.5 ± 1.5	87.8 ± 17.3	94.5	173.0 ± 10	157 ± 10	150.0 ± 30
CUSP8	16.43 ± 1.21	19.17 ± 1.13	47.52 ± 16.67	22.87 ± 30.44	174.42 ± 2.11	190.51 ± 2.87	174.49 ± 2.04
NA8	17.46 ± 1.28	20.83 ± 1.35	37.66 ± 12.33	46.99 ± 15.45	174.29 ± 1.60	188.58 ± 2.37	173.74 ± 1.50
CUSP8-transition-100	16.45 ± 0.93	19.04 ± 0.85	56.86 ± 12.56	8.45 ± 29.96	174.12 ± 1.56	189.67 ± 2.24	174.79 ± 1.44
CUSP8-transition-300	16.95 ± 1.19	19.50 ± 0.99	49.63 ± 18.68	26.51 ± 34.42	172.50 ± 2.27	187.14 ± 3.10	172.30 ± 2.17
CUSP8-transition-900	17.20 ± 1.12	19.66 ± 1.11	45.48 ± 13.72	49.51 ± 23.45	176.22 ± 1.31	191.36 ± 1.98	176.41 ± 1.23
CUSP12	15.66 ± 0.99	17.89 ± 0.99	61.07 ± 19.56	10.70 ± 33.87	171.01 ± 2.14	186.81 ± 2.93	171.25 ± 2.06
CUSP20	14.76 ± 0.72	16.37 ± 0.74	64.01 ± 14.37	-1.37 ± 17.20	170. ± 1.97	186.66 ± 2.91	170.82 ± 1.99
CUSP30	13.62 ± 0.41	14.89 ± 0.47	34.91 ± 1.19	25.47 ± 1.50	173. ± 1.88	187.82 ± 2.64	171.87 ± 1.80
EC60to30-degraded-0.25	12.31 ± 0.50	14.84 ± 0.52	40.49 ± 0.70	23.89 ± 1.00	175.20 ± 1.95	190. ± 2.69	175.58 ± 1.90
EC60to30-degraded-0.50	10.75 ± 0.39	11.56 ± 0.39	44.56 ± 0.61	24.92 ± 0.69	173.97 ± 1.68	189.82 ± 2.45	173.73 ± 1.58
EC60to30-degraded-0.75	10.91 ± 0.37	11.25 ± 0.36	44.58 ± 0.65	25.43 ± 0.83	172.18 ± 1.41	187.95 ± 2.16	172.99 ± 1.33
EC60to30	10.13 ± 0.40	12.55 ± 0.53	40.84 ± 0.73	22.93 ± 1.14	172.81 ± 2.37	188.21 ± 3.10	173.52 ± 2.36
EC60to30-E3SM-V1	10.51 ± 0.51	10.57 ± 0.45	43.07 ± 0.68	23.22 ± 0.53	173.18 ± 1.79	188.85 ± 2.55	172.30 ± 1.83

391 **Table 2.** The average transport in Sverdrups through transects for years 2-10, followed by standard deviation
392 for simulations and error for observations. See Figure 7 for plots of the data and Figure 6 for a map of the
393 Gulf Stream transects. Observational references: Florida-Cuba: *Johns et al.* [2002], Florida-Bahamas: *Johns*
394 *et al.* [2002], Cape Hatteras: *Halkin and Rossby* [1985], New Jersey: *Rossby et al.* [2014], Drake Passage:
395 *Donohue et al.* [2016], Tasmania-Ant: *Ganachaud and Wunsch* [2000], Africa-Ant: *Ganachaud and Wunsch*
396 [2000]

459 4.3 Study 3: Coastal resolution

460 The analysis of the coastal resolution study can be found in Figure 13. The meshes
461 with higher coastal resolution showed significantly improved dynamics, particularly in eddy
462 kinetic energy and sea surface height variability, which were almost non-existent in CUSP30
463 (see Figure 13). The Gulf Stream within the high resolution region in CUSP8 is similar to
464 that of NA8. However, as noted in the transition width study, features that cross into the
465 lower resolution transition region and then back into the high resolution region, such as me-
466 anders and eddies, are less well resolved in the CUSP8 simulation.

467 Figure 14 shows the path of the Gulf Stream in the coastal resolution study. The NA8
468 simulation shows very little variability in the path of the Gulf Stream, while the CUSP8,
469 CUSP12 and CUSP20 simulations show much more. The CUSP30 simulation also does
470 not show much variability in the Gulf Stream path, but this is expected as the resolution
471 is too low to be eddy permitting. The variation in the Gulf Stream path is also apparent in
472 the transport through the transects along the Gulf Stream. In the southernmost transects
473 (Florida-Cuba and Florida-Bahamas) where the flow is geographically constrained, the trans-
474 port increases with increased resolution. The Cape Hatteras and New Jersey transects do not
475 show this pattern. Figure 14 and Table 2 show that in the CUSP8, CUSP12, an CUSP20 sim-
476 ulations, there is significant variability in the path of the flow in the region of the New Jer-
477 sey transect. Periods of very low or negative transport are probably due to North-South flow
478 through the transect as the Gulf Stream separates from the coast. This can be seen in some of
479 the monthly Gulf Stream paths seen in Figure 14.



480 **Figure 10.** Performance for resolution study, showing simulated years per wall clock day (SYPD). Black
 481 dotted lines show perfect scaling. The SYPD values for 1024 processors are: CUSP8: 2.0, CUSP12: 5.1,
 482 CUSP20: 11.7, CUSP30: 20.0, EC60to30: 32.5, RRS18to6: 0.38. CUSP8 is 16 times slower than EC60to30,
 483 but 5.3 times faster than global high resolution with cell sizes ranging from 18 to 6 km (RRS18to6). All
 484 simulations use 80 layers, except the EC60to30, which is 60 layers.

485 This high variability is not due only to high resolution, as the NA8 simulation, which
486 has the same coastal resolution, shows very little variability in the path of the Gulf Stream.
487 The very low variability in the NA8 simulation is probably due largely to the idealized forc-
488 ing used, as this effect did not show up in global high resolution simulations with realistic
489 climatological forcing. Initial results from global high resolution simulations with clima-
490 tological forcing show that the Gulf stream had a realistic path and variability (see Figure
491 17 in the Appendix). It appears that the lack of variation in the forcing or in the mesh itself
492 prevents the NA8 mesh from developing meandering features. Variable forcing appears to
493 resolve this problem. The variability in cell size and quality in the CUSP meshes may allow
494 these features to develop. For all the transects, the variability increased significantly between
495 CUSP30 and the higher resolution meshes, as would be expected when transitioning to an
496 eddy permitting resolution.

497 The CUSP simulations are also a significant improvement on global high resolution
498 in terms of cost (see Figure 10). CUSP8 offers an order of magnitude improvement in speed
499 when compared to global high resolution simulations with cell sizes ranging from 18 km to
500 6 km. EC60to30, while an order of magnitude faster than CUSP8, lacks the improvements
501 in coastal dynamics that motivated the creation of the CUSP8 mesh. Performance tests were
502 run on Grizzly at Los Alamos National Laboratory. Grizzly is an Institutional Computing
503 (IC) cluster, running on the TOSS operating system (Tri-Lab Operating System Stack) and
504 using the Intel OmniPath interconnect. Each processor is a 2.1GHz Broadwell with 45MB
505 cache, with 36 processors per node.

506 5 Conclusion

507 Overall, this mesh resolution case study indicates that simulations are robust to changes
508 in the mesh. Changes to mesh quality were found to have little impact on the simulation
509 quality and statistics. Problems with the stability of the simulations at large timesteps oc-
510 curred in spinning-up the two most degraded configurations, but with a modified timestep,
511 these simulations were found to perform similarly to the undegraded cases. Such behavior
512 is consistent with the expected reduction in CFL limits associated with heavily degraded
513 meshes that incorporate small grid cells. Despite previous theoretical analysis suggesting a
514 strong link between mesh quality and numerical discretization error [Peixoto, 2016], it was
515 found that simulation quality was not obviously diminished with increasing levels of mesh
516 degradation. In this sense, it appears the TRiSK formulation used in MPAS-Ocean may out-
517 perform its theoretical bounds in many practical cases. Changes to the transition width were
518 also found to have relatively little impact on the quality of the simulations.

519 It is likely that much of the variation in the transition width study was due to the change
520 in the size of the higher resolution portion of the transition region rather than the transition
521 itself. The difference between the steady Gulf Stream path in the NA8 simulation and the
522 variable paths in the CUSP simulations shows that the transition region has some impact on
523 variability. It is not clear if this is due to mesh quality or to the effect of changing resolution.
524 In this case, the CUSP meshes had more realistic variability, but it is not clear that this added
525 variability would be desirable in a simulation with realistic forcing. This study also demon-
526 strated that higher coastal resolution improved the dynamics of the Gulf Stream at a much
527 lower cost than a high resolution global model.

528 When designing a mesh, the effect of processes outside of the resolution region is es-
529 sential. The transition width study showed that processes within the high resolution region
530 cannot be properly resolved if they interact with processes in the low resolution region. For
531 example, in the CUSP8-transition-100 simulation, meanders and eddies crossing into the low
532 resolution region had a strong impact on the dynamics present along the Gulf Stream within
533 the high resolution region. More broadly, it is important to evaluate the dependence of the
534 coastal dynamics on basin scale or global dynamics. A coastal high resolution model may
535 be of limited use if the ultimate drivers of the coastal dynamics are not modeled accurately.

536 For instance, flooding during a hurricane requires that off-shore storm surges are modeled at
537 appropriate resolution in order to predict accurate coastal surges.

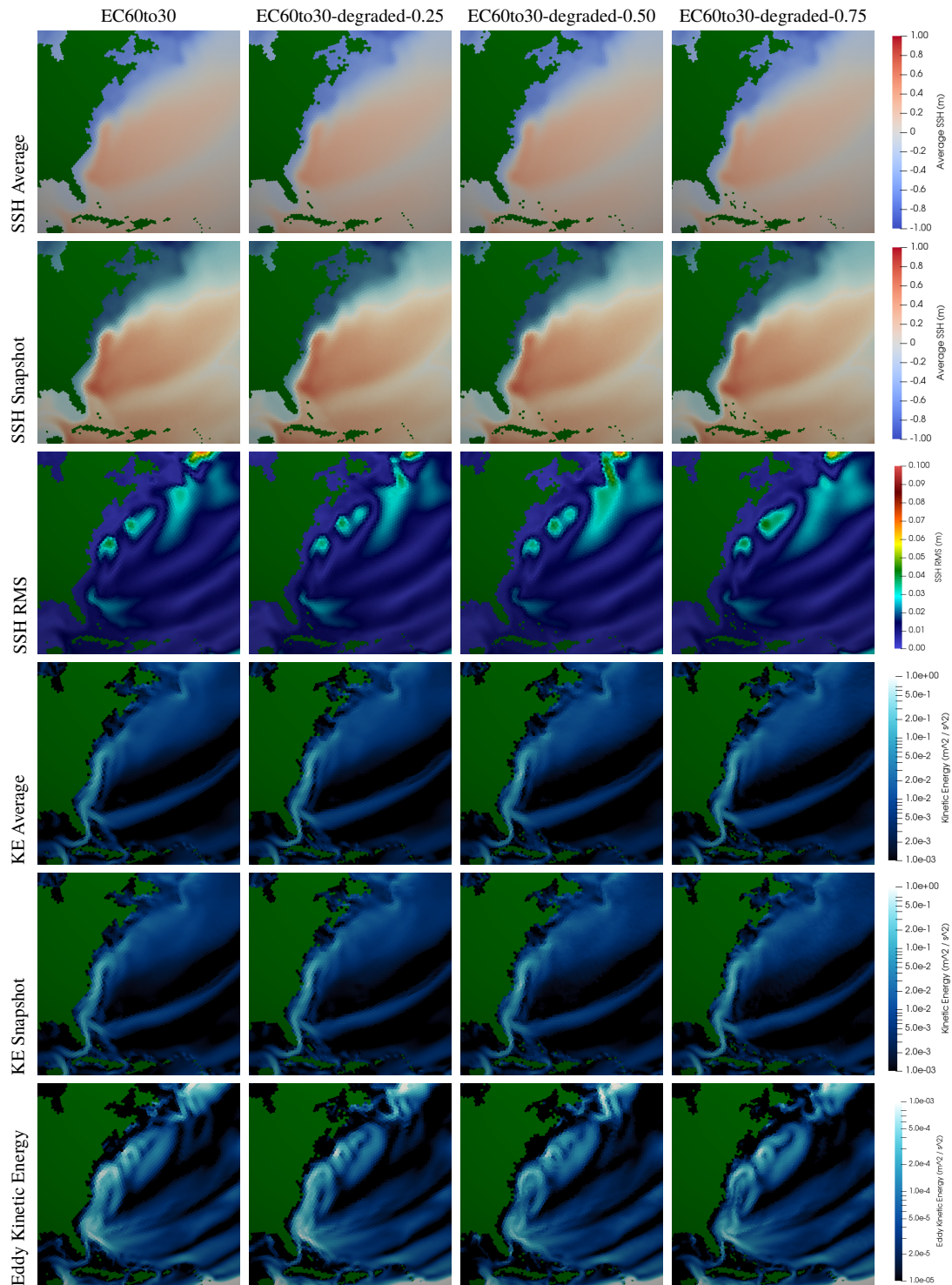
538 Physical dynamics considerations appear to be much more important than mesh met-
539 rics considerations in these stand-alone ocean simulations. Future studies will look in more
540 detail at CUSP8 simulations with realistic atmospheric forcing and in coupled configura-
541 tions, which may have more stringent mesh quality requirements due to cross-component
542 feedbacks. Similar variable resolution meshes are in development for investigating other re-
543 gions of interest, including the Arctic and Southern Oceans. Our results suggest robust capa-
544 bilities inherent in the MPAS-Ocean discretization and mesh generation approaches. These
545 provide the capability to create a diverse range of variable resolution configurations, which
546 will allow modelers to accurately resolve additional physical processes at lower computa-
547 tional costs.

548

5.1 Plots

549

Degraded Mesh Study

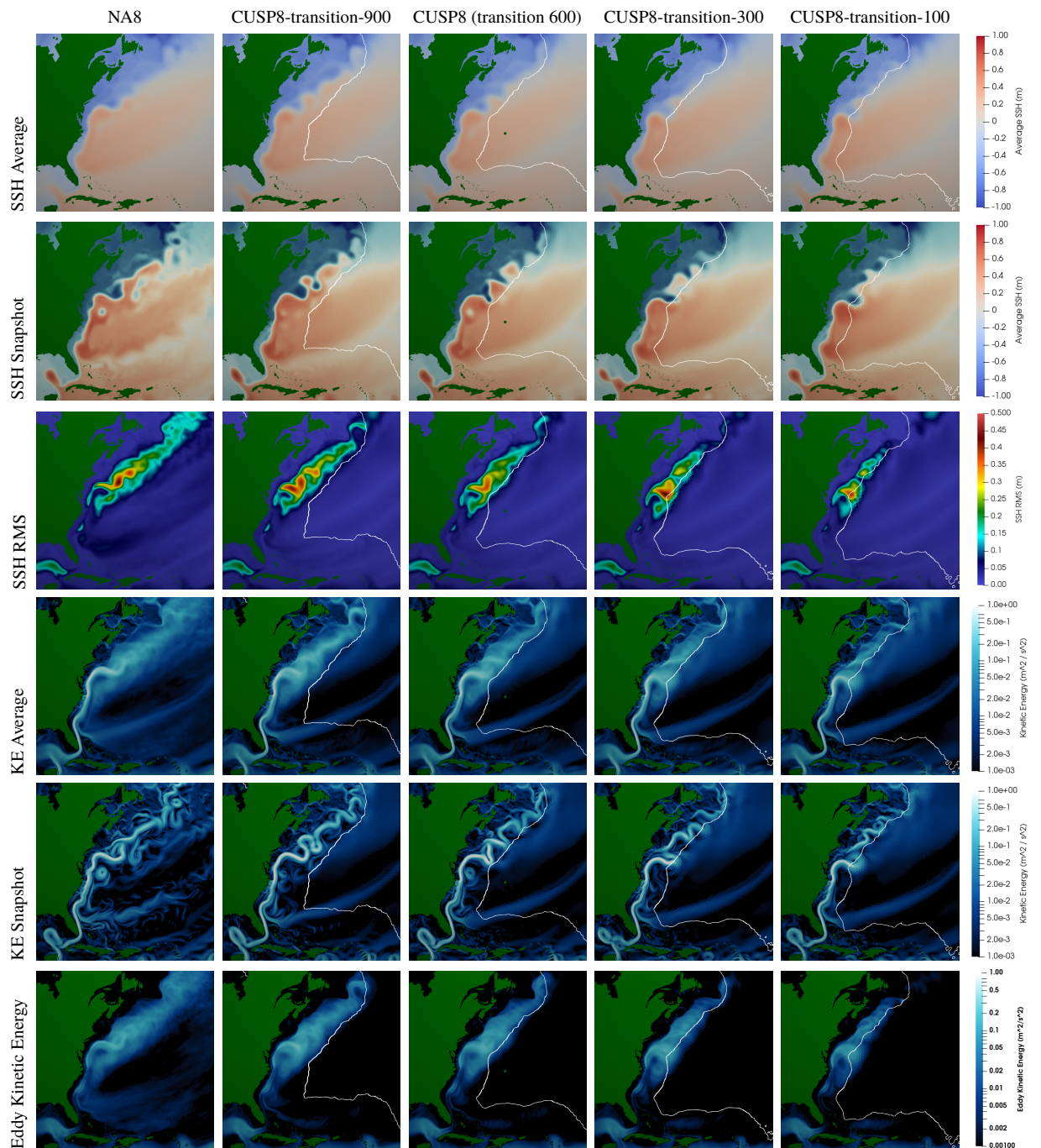


550

Figure 11. Degraded Mesh Study: Averages are taken from years 2-10 of the simulation, snapshots from 0002-06-01. The degraded meshes have a minimal impact on simulation quality.

551

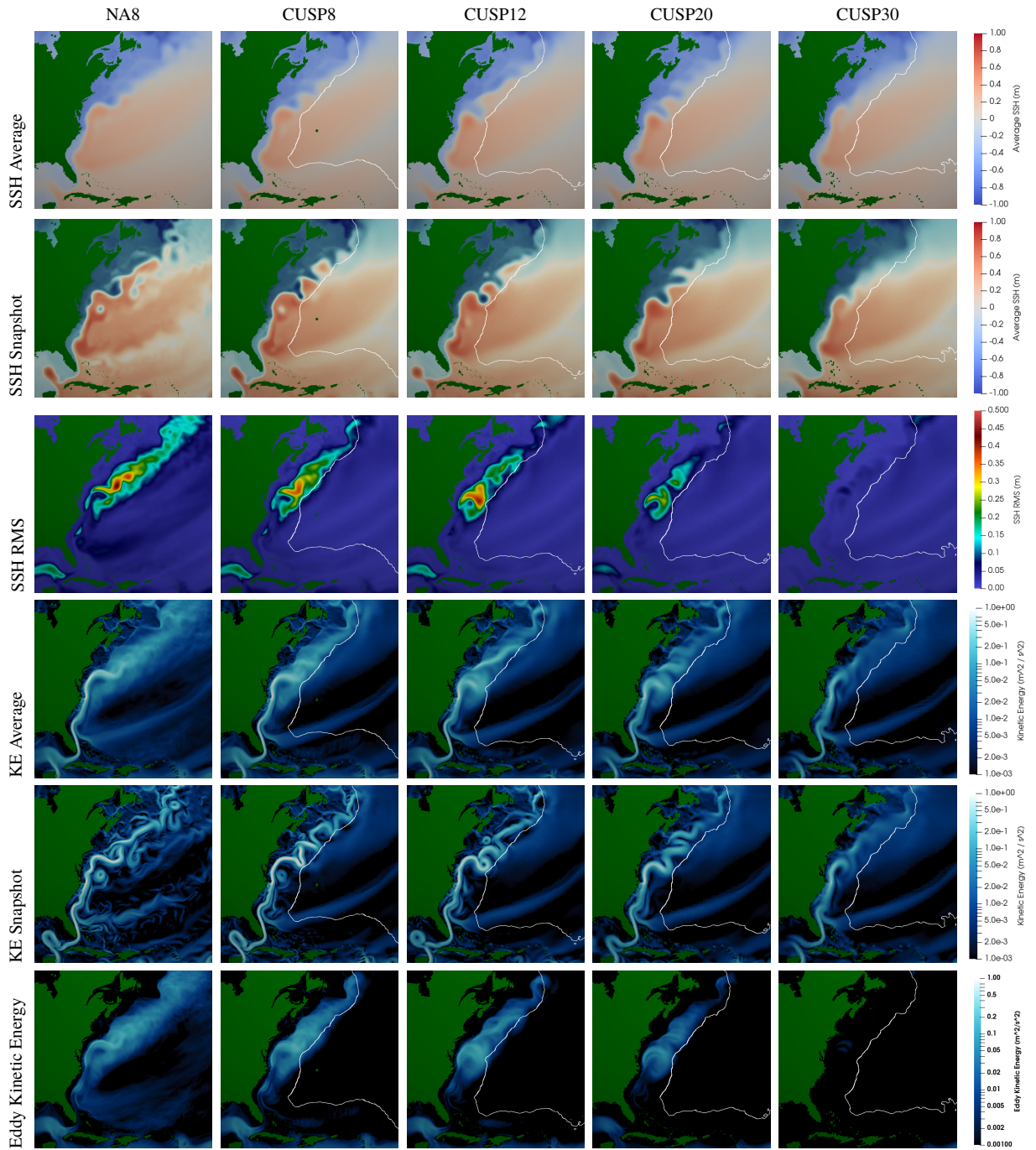
Transition Width Study



553 **Figure 12. Transition Width Study:** A wider transition improves simulation quality and increases vari-
 554 ability within the coastal region. This appears to be less a function of the transition itself and rather a function
 555 of the size of the higher resolution region (see position of the center of the transition region). The white line
 556 shows the center of the transition region. See Figure 11 for details.

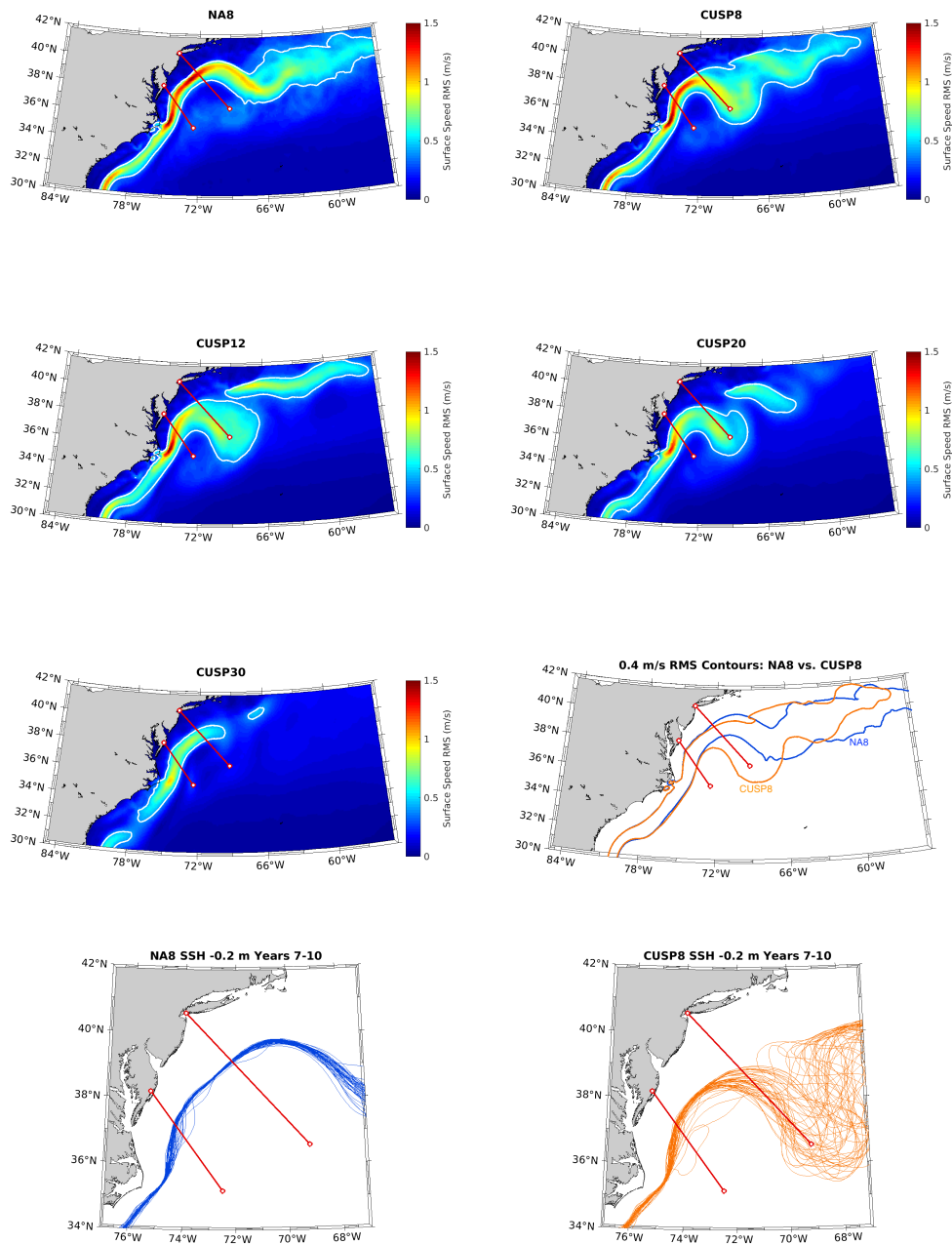
557

Coastal Resolution Study



558 **Figure 13. Resolution Study:** Same as Figure 12. A higher resolution improves the dynamics of the Gulf
 559 Stream significantly, with CSUP8 approaching the dynamics of NA8 along the coast. Variability increases
 560 with increasing resolution.

Coastal Resolution Study



562 **Figure 14.** Root mean square (RMS) of surface speed for resolution convergence runs. Surface speed is
 563 taken from the first 5-day snapshot of each month for years 2 through 10. White contour on (a-e) indicates
 564 0.4 m/s surface speed RMS contour. These contours are seen in (f) for NA8 (blue) and CUSP8 (orange). The
 565 northern and southern red lines on each plot indicate the New Jersey and Cape Hatteras transects, respectively.
 566 Snapshots of the Gulf Stream path are shown for NA8 (g) and CUSP8 (h). These pathlines follow the -0.2 m
 567 SSH contour. Paths are taken from the first 5-day snapshot of each month for years 7 through 10. Notice how
 568 the path of CUSP8 (h) frequently loops back through the New Jersey transect. This is the cause of the low
 569 transport through this transect for CUSP8 relative to NA8 and observations.

Acknowledgments

This research was supported as part of the Energy Exascale Earth System Model (E3SM) project, funded by the U.S. Department of Energy, Office of Science, Office of Biological and Environmental Research as well as the U.S. Department of Energy Advanced Research Projects Agency - Energy (ARPA-E) Macroalgae Research Inspiring Novel Energy Resources (MARINER) program (Funding Opportunity No. DE-FOA-0001726, MARINER Award 17/CJ000/09/01, Pacific Northwest National Laboratory, prime recipient). E3SM simulations are conducted at: Argonne Leadership Computing Facility (contract DE-AC02-06CH11357); National Energy Research Scientific Computing Center (DE-AC05-00OR22725); Oak Ridge Leadership Computing Facility (DE-AC05-00OR22725); Argonne Nat. Lab. high-performance computing cluster, provided by BER Earth System Modeling; and Los Alamos Nat. Lab. Institutional Computing, US DOE NNSA (DE-AC52-06NA25396). E3SM data is freely available through the Earth System Grid Federation (ESGF) distributed archives. See details at <https://e3sm.org/data>. JIGSAW can be obtained freely online at <https://github.com/dengwirda/jigsaw-geo-matlab>.

References

- Abiodun, B. J., J. M. Prusa, and W. J. Gutowski (2008), Implementation of a non-hydrostatic, adaptive-grid dynamics core in CAM3. Part I: comparison of dynamics cores in aquaplanet simulations, *Climate Dynamics*, *31*(7), 795–810, doi:10.1007/s00382-008-0381-y.
- Androsov, A., V. Fofonova, I. Kuznetsov, S. Danilov, N. Rakowsky, S. Harig, and K. H. Wiltshire (2018), FESOM-C: coastal dynamics on hybrid unstructured meshes, *Geoscientific Model Development Discussions*, pp. 1–32, doi:https://doi.org/10.5194/gmd-2018-112.
- Androsov, A., V. Fofonova, I. Kuznetsov, S. Danilov, N. Rakowsky, S. Harig, H. Brix, and K. H. Wiltshire (2019), FESOM-C v.2: coastal dynamics on hybrid unstructured meshes, *Geoscientific Model Development*, *12*(3), 1009–1028, doi:10.5194/gmd-12-1009-2019.
- Biaostoch, A., D. Sein, J. V. Durgadoo, Q. Wang, and S. Danilov (2018), Simulating the Agulhas system in global ocean models — nesting vs. multi-resolution unstructured meshes, *Ocean Modelling*, *121*, 117–131, doi:10.1016/j.ocemod.2017.12.002.
- Chen, C., H. Liu, and R. C. Beardsley (2003), An unstructured grid, finite-volume, three-dimensional, primitive equations ocean model: Application to coastal ocean and estuaries, *Journal of Atmospheric and Oceanic Technology*, *20*(1), 159–186, doi:10.1175/1520-0426(2003)020<0159:AUGFVT>2.0.CO;2.
- Chen, L., and M. Holst (2011), Efficient mesh optimization schemes based on optimal delaunay triangulations, *Computer Methods in Applied Mechanics and Engineering*, *200*(9), 967–984, doi:https://doi.org/10.1016/j.cma.2010.11.007.
- Danilov, S., and Q. Wang (2015), Resolving eddies by local mesh refinement, *Ocean Modelling*, *93*, 75–83, doi:10.1016/j.ocemod.2015.07.006.
- Davies, D. R., C. R. Wilson, and S. C. Kramer (2011), Fluidity: A fully unstructured anisotropic adaptive mesh computational modeling framework for geodynamics, *Geochemistry, Geophysics, Geosystems*, *12*(6), doi:10.1029/2011GC003551.
- Donohue, K. A., K. L. Tracey, D. R. Watts, M. P. Chidichimo, and T. K. Chereskin (2016), Mean Antarctic Circumpolar Current transport measured in Drake Passage, *Geophysical Research Letters*, *43*(22), 11,760–11,767, doi:10.1002/2016GL070319.
- Düben, P. D., and P. Korn (2014), Atmosphere and ocean modeling on grids of variable resolution—a 2d case study, *Monthly Weather Review*, *142*(5), 1997–2017, doi:10.1175/MWR-D-13-00217.1.
- Engwirda, D. (2017), JIGSAW-GEO (1.0): locally orthogonal staggered unstructured grid generation for general circulation modelling on the sphere, *Geoscientific Model Development*, *10*(6), 2117–2140, doi:https://doi.org/10.5194/gmd-10-2117-2017.
- Engwirda, D. (2018), Generalised primal-dual grids for unstructured co-volume schemes, *Journal of Computational Physics*, *375*, 155–176, doi:https://doi.org/10.1016/j.jcp.2018.07.025.

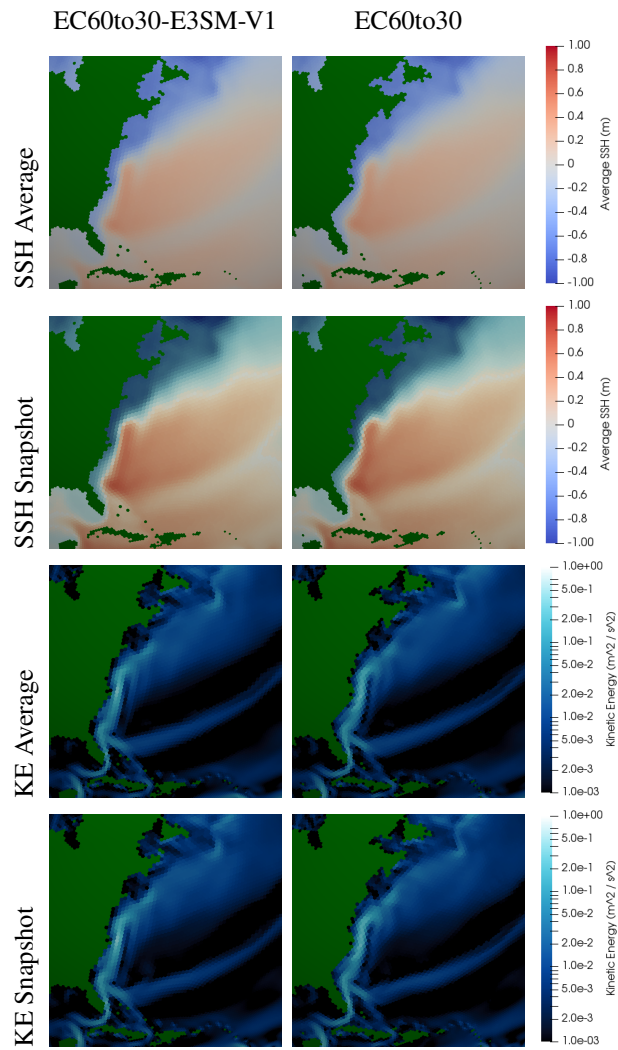
- 622 Engwirda, D., and D. Ivers (2016), Off-centre steiner points for delaunay-refinement on
 623 curved surfaces, *Computer-Aided Design*, *72*, 157 – 171, doi:https://doi.org/10.1016/j.cad.
 624 2015.10.007, 23rd International Meshing Roundtable Special Issue: Advances in Mesh
 625 Generation.
- 626 Ganachaud, A., and C. Wunsch (2000), Improved estimates of global ocean circulation, heat
 627 transport and mixing from hydrographic data, *Nature*, *408*(6811), 453–457, doi:10.1038/
 628 35044048.
- 629 Gent, P. R., and J. C. McWilliams (1990), Isopycnal mixing in ocean circulation models,
 630 *Journal of Physical Oceanography*, *20*(1), 150–155, doi:10.1175/1520-0485(1990)
 631 020<0150:IMIOCM>2.0.CO;2.
- 632 Golaz, J.-C., P. M. Caldwell, L. P. Van Roekel, M. R. Petersen, Q. Tang, J. D. Wolfe,
 633 G. Abeshu, V. Anantharaj, X. S. Asay-Davis, D. C. Bader, S. A. Baldwin, G. Bisht, P. A.
 634 Bogenschutz, M. Branstetter, M. A. Brunke, S. R. Brus, S. M. Burrows, P. J. Cameron-
 635 Smith, A. S. Donahue, M. Deakin, R. C. Easter, K. J. Evans, Y. Feng, M. Flanner, J. G.
 636 Foucar, J. G. Fyke, B. M. Griffin, C. Hannay, B. E. Harrop, E. C. Hunke, R. L. Jacob,
 637 D. W. Jacobsen, N. Jeffery, P. W. Jones, N. D. Keen, S. A. Klein, V. E. Larson, L. R. Le-
 638 ung, H.-Y. Li, W. Lin, W. H. Lipscomb, P.-L. Ma, S. Mahajan, M. E. Maltrud, A. Mamet-
 639 janov, J. L. McClean, R. B. McCoy, R. B. Neale, S. F. Price, Y. Qian, P. J. Rasch, J. J.
 640 Reeves Eyre, W. J. Riley, T. D. Ringler, A. F. Roberts, E. L. Roesler, A. G. Salinger,
 641 Z. Shaheen, X. Shi, B. Singh, J. Tang, M. A. Taylor, P. E. Thornton, A. K. Turner,
 642 M. Veneziani, H. Wan, H. Wang, S. Wang, D. N. Williams, P. J. Wolfram, P. H. Worley,
 643 S. Xie, Y. Yang, J.-H. Yoon, M. D. Zelinka, C. S. Zender, X. Zeng, C. Zhang, K. Zhang,
 644 Y. Zhang, X. Zheng, T. Zhou, and Q. Zhu (2019), The DOE E3SM coupled model version
 645 1: Overview and evaluation at standard resolution, *Journal of Advances in Modeling Earth*
 646 *Systems*, *11*(ja), doi:10.1029/2018MS001603.
- 647 Griffies, S. M., A. Biastoch, C. Böning, F. Bryan, G. Danabasoglu, E. P. Chassignet, M. H.
 648 England, R. Gerdes, H. Haak, R. W. Hallberg, W. Hazeleger, J. Jungclaus, W. G. Large,
 649 G. Madec, A. Pirani, B. L. Samuels, M. Scheinert, A. S. Gupta, C. A. Severijns, H. L.
 650 Simmons, A. M. Treguier, M. Winton, S. Yeager, and J. Yin (2009), Coordinated Ocean-
 651 ice Reference Experiments (COREs), *Ocean Modelling*, *26*(1-2), 1–46, doi:10.1016/j.
 652 ocemod.2008.08.007.
- 653 Hagos, S., R. Leung, S. A. Rauscher, and T. Ringler (2013), Error characteristics of two grid
 654 refinement approaches in aquaplanet simulations: MPAS-A and WRF, *Monthly Weather*
 655 *Review*, *141*(9), 3022–3036, doi:10.1175/MWR-D-12-00338.1.
- 656 Halkin, D., and T. Rossby (1985), The Structure and Transport of the Gulf Stream at 73°W,
 657 *Journal of Physical Oceanography*, *15*(11), 1439–1452, doi:10.1175/1520-0485(1985)
 658 015<1439:TSATOT>2.0.CO;2.
- 659 Jacobsen, D. W., M. Gunzburger, T. Ringler, J. Burkardt, and J. Peterson (2013), Parallel
 660 algorithms for planar and spherical Delaunay construction with an application to cen-
 661 troidal Voronoi tessellations, *Geoscientific Model Development*, *6*(4), 1353–1365, doi:
 662 10.5194/gmd-6-1353-2013.
- 663 Johns, W. E., T. L. Townsend, D. M. Fratantoni, and W. D. Wilson (2002), On the Atlantic
 664 inflow to the Caribbean Sea, *Deep Sea Research Part I: Oceanographic Research Papers*,
 665 *49*(2), 211–243, doi:10.1016/S0967-0637(01)00041-3.
- 666 Korn, P. (2017), Formulation of an unstructured grid model for global ocean dynamics, *Jour-
 667 nal of Computational Physics*, *339*, 525 – 552, doi:https://doi.org/10.1016/j.jcp.2017.03.
 668 009.
- 669 Lambrechts, J., R. Comblen, V. Legat, C. Geuzaine, and J.-F. Remacle (2008), Multi-
 670 scale mesh generation on the sphere, *Ocean Dynamics*, *58*(5), 461–473, doi:10.1007/
 671 s10236-008-0148-3.
- 672 Lorant, V., and J.-F. Royer (2001), Sensitivity of equatorial convection to horizontal resolu-
 673 tion in aquaplanet simulations with a variable-resolution GCM, *Monthly Weather Review*,
 674 *129*(11), 2730–2745, doi:10.1175/1520-0493(2001)129<2730:SOECTH>2.0.CO;2.

- 675 McRae, A., C. Cotter, and C. Budd (2018), Optimal-transport–based mesh adaptivity on
676 the plane and sphere using finite elements, *SIAM Journal on Scientific Computing*, *40*(2),
677 A1121–A1148, doi:10.1137/16M1109515.
- 678 Park, S.-H., J. B. Klemp, and W. C. Skamarock (2014), A comparison of mesh refine-
679 ment in the global MPAS-A and WRF models using an idealized normal-mode baro-
680 clinic wave simulation, *Monthly Weather Review*, *142*(10), 3614–3634, doi:10.1175/
681 MWR-D-14-00004.1.
- 682 Peixoto, P. S. (2016), Accuracy analysis of mimetic finite volume operators on geodesic grids
683 and a consistent alternative, *Journal of Computational Physics*, *310*, 127 – 160, doi:https:
684 //doi.org/10.1016/j.jcp.2015.12.058.
- 685 Petersen, M. R., D. W. Jacobsen, T. D. Ringler, M. W. Hecht, and M. E. Maltrud (2015),
686 Evaluation of the arbitrary Lagrangian–Eulerian vertical coordinate method in the MPAS-
687 Ocean model, *Ocean Modelling*, *86*, 93–113, doi:10.1016/j.ocemod.2014.12.004.
- 688 Petersen, M. R., X. S. Asay-Davis, A. S. Berres, Q. Chen, N. Feige, M. J. Hoffman, D. W. Ja-
689 cobsen, P. W. Jones, M. E. Maltrud, S. F. Price, T. D. Ringler, G. J. Streltetz, A. K. Turner,
690 L. P. Van Roekel, M. Veneziani, J. D. Wolfe, P. J. Wolfram, and J. L. Woodring (2019),
691 An evaluation of the ocean and sea ice climate of E3SM using MPAS and interannual
692 CORE-II forcing, *Journal of Advances in Modeling Earth Systems*, *11*(5), 1438–1458,
693 doi:10.1029/2018MS001373.
- 694 Rauscher, S. A., and T. D. Ringler (2014), Impact of variable-resolution meshes on midlat-
695 itude baroclinic eddies using CAM-MPAS-A, *Monthly Weather Review*, *142*(11), 4256–
696 4268, doi:10.1175/MWR-D-13-00366.1.
- 697 Rauscher, S. A., T. D. Ringler, W. C. Skamarock, and A. A. Mirin (2012), Exploring a global
698 multiresolution modeling approach using aquaplanet simulations, *Journal of Climate*,
699 *26*(8), 2432–2452, doi:10.1175/JCLI-D-12-00154.1.
- 700 Reckinger, S. M., M. R. Petersen, and S. J. Reckinger (2015), A study of overflow simula-
701 tions using MPAS-Ocean: Vertical grids, resolution, and viscosity, *Ocean Modelling*, *96*,
702 291–313, doi:10.1016/j.ocemod.2015.09.006.
- 703 Remacle, J.-F., and J. Lambrechts (2018), Fast and robust mesh generation on the sphere—
704 application to coastal domains, *Computer-Aided Design*, *103*, 14 – 23, doi:https://doi.org/
705 10.1016/j.cad.2018.03.002, 25th International Meshing Roundtable Special Issue: Ad-
706 vances in Mesh Generation.
- 707 Ringler, T., M. Petersen, R. L. Higdon, D. Jacobsen, P. W. Jones, and M. Maltrud (2013),
708 A multi-resolution approach to global ocean modeling, *Ocean Modelling*, *69*, 211–232,
709 doi:10.1016/j.ocemod.2013.04.010.
- 710 Ringler, T. D., J. Thuburn, J. B. Klemp, and W. C. Skamarock (2010), A unified approach
711 to energy conservation and potential vorticity dynamics for arbitrarily-structured C-grids,
712 *Journal of Computational Physics*, *229*(9), 3065–3090, doi:10.1016/j.jcp.2009.12.007.
- 713 Rossby, T., C. N. Flagg, K. Donohue, A. Sanchez-Franks, and J. Lillibridge (2014), On the
714 long-term stability of Gulf Stream transport based on 20 years of direct measurements,
715 *Geophysical Research Letters*, *41*(1), 114–120, doi:10.1002/2013GL058636.
- 716 Sakaguchi, K., J. Lu, L. R. Leung, C. Zhao, Y. Li, and S. Hagos (2016), Sources and path-
717 ways of the upscale effects on the Southern Hemisphere jet in MPAS-CAM4 variable-
718 resolution simulations, *Journal of Advances in Modeling Earth Systems*, *8*(4), 1786–1805,
719 doi:10.1002/2016MS000743.
- 720 Scholz, P., D. Sidorenko, O. Gurses, S. Danilov, N. Koldunov, Q. Wang, D. Sein, M. Smo-
721 lentseva, N. Rakowsky, and T. Jung (2019), Assessment of the finite volume sea ice ocean
722 model (FESOM2.0), part I: Description of selected key model elements and comparison
723 to its predecessor version, *Geoscientific Model Development Discussions*, *2019*, 1–42, doi:
724 10.5194/gmd-2018-329.
- 725 Sein, D. V., N. V. Koldunov, S. Danilov, Q. Wang, D. Sidorenko, I. Fast, T. Rackow, W. Ca-
726 bos, and T. Jung (2017), Ocean modeling on a mesh with resolution following the local
727 rossby radius, *Journal of Advances in Modeling Earth Systems*, doi:http://doi.org/10.1002/
728 2017MS001099.

- 729 Skamarock, W. C., and A. Gassmann (2011), Conservative Transport Schemes for Spherical
730 Geodesic Grids: High-Order Flux Operators for ODE-Based Time Integration, *Monthly*
731 *Weather Review*, *139*(9), 2962–2975, doi:10.1175/MWR-D-10-05056.1.
- 732 Steele, M., R. Morley, and W. Ermold (2001), Phc: A global ocean hydrography with a
733 high-quality arctic ocean, *14*(9), 2079–2087, doi:10.1175/1520-0442(2001)014<2079:
734 PAGOHW>2.0.CO;2.
- 735 Van Roekel, L., A. J. Adcroft, G. Danabasoglu, S. M. Griffies, B. Kauffman, W. Large,
736 M. Levy, B. G. Reichl, T. Ringler, and M. Schmidt (2018), The KPP boundary layer
737 scheme for the ocean: Revisiting its formulation and benchmarking one-dimensional sim-
738 ulations relative to les, *Journal of Advances in Modeling Earth Systems*, *10*(11), 2647–
739 2685, doi:10.1029/2018MS001336.
- 740 Wang, Q., C. Wekerle, S. Danilov, X. Wang, and T. Jung (2018), A 4.5 km resolution Arctic
741 Ocean simulation with the global multi-resolution model FESOM1.4, *Geosci. Model Dev.*,
742 *11*, 1229–1255, doi:http://doi.org/10.5194/gmd-11-1229-2018.
- 743 Weller, H., P. Browne, C. Budd, and M. Cullen (2016), Mesh adaptation on the sphere using
744 optimal transport and the numerical solution of a monge—ampÁire type equation, *Journal*
745 *of Computational Physics*, *308*, 102 – 123, doi:https://doi.org/10.1016/j.jcp.2015.12.018.
- 746 White, L., E. Deleersnijder, and V. Legat (2008), A three-dimensional unstructured mesh
747 finite element shallow-water model, with application to the flows around an island and in
748 a wind-driven, elongated basin, *Ocean Modelling*, *22*(1), 26 – 47, doi:https://doi.org/10.
749 1016/j.ocemod.2008.01.001.
- 750 Yang, H., M. Gunzburger, and L. Ju (2018), Fast spherical centroidal Voronoi mesh gen-
751 eration: A Lloyd-preconditioned LBFGS method in parallel, *Journal of Computational*
752 *Physics*, *367*, 235 – 252, doi:https://doi.org/10.1016/j.jcp.2018.04.034.
- 753 Zarzycki, C. M., C. Jablonowski, D. R. Thatcher, and M. A. Taylor (2015), Effects of local-
754 ized grid refinement on the general circulation and climatology in the Community Atmo-
755 sphere Model, *Journal of Climate*, *28*(7), 2777–2803, doi:10.1175/JCLI-D-14-00599.1.
- 756 Zhao, C., L. R. Leung, S.-H. Park, S. Hagos, J. Lu, K. Sakaguchi, J. Yoon, B. E. Harrop,
757 W. Skamarock, and M. G. Duda (2016), Exploring the impacts of physics and resolution
758 on aqua-planet simulations from a nonhydrostatic global variable-resolution modeling
759 framework, *Journal of Advances in Modeling Earth Systems*, *8*(4), 1751–1768, doi:10.
760 1002/2016MS000727.

761

6 Appendix

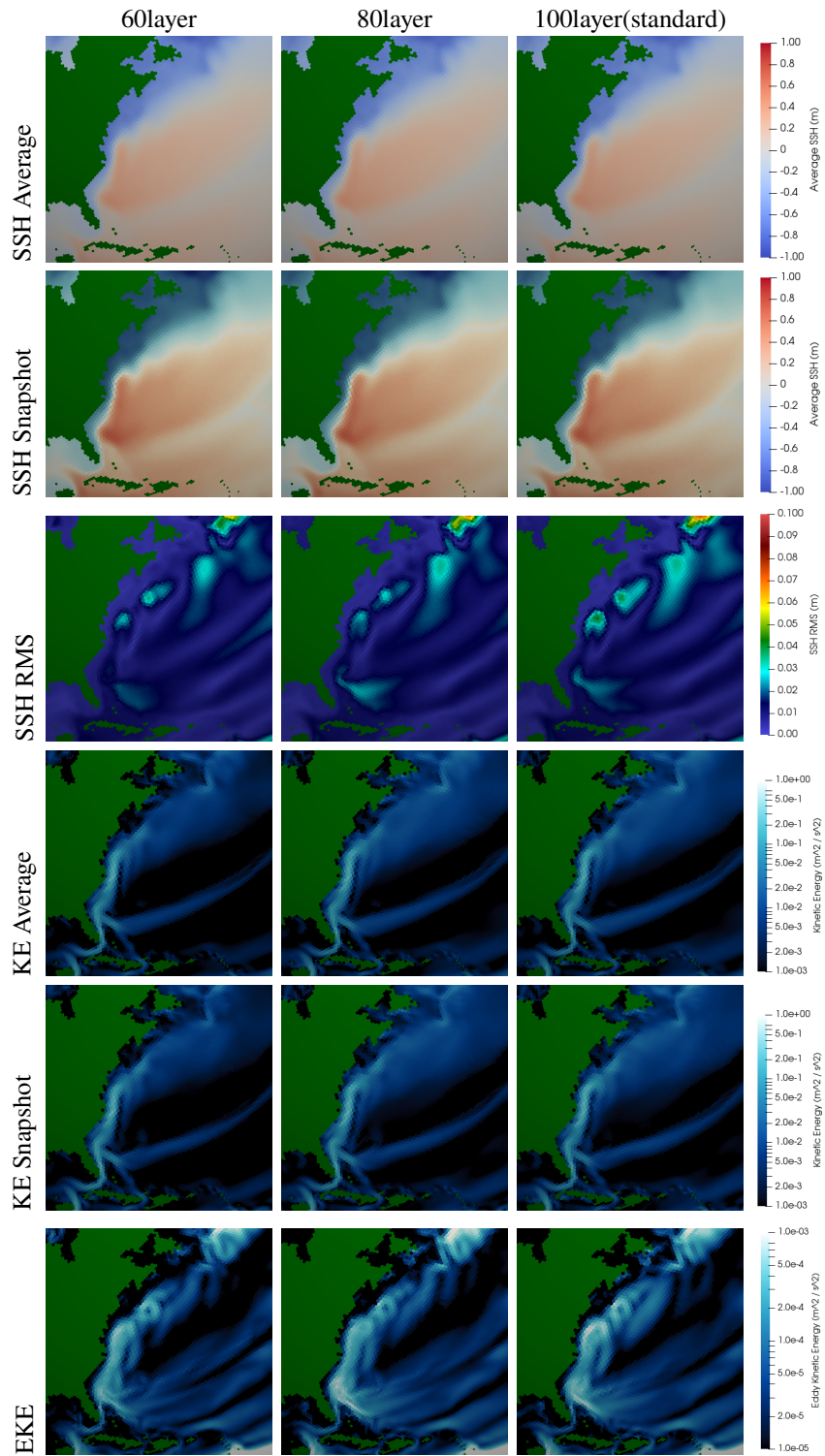


762

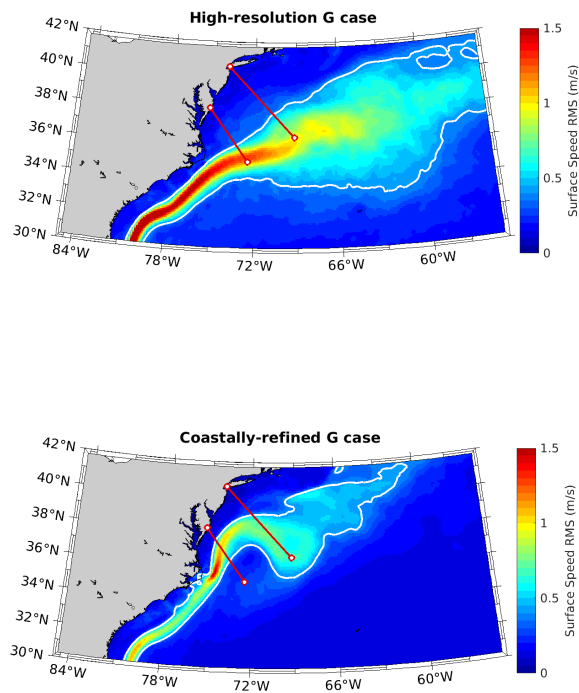
Figure 15. EC60to30-E3SM-V1 vs EC60to30: Averages are taken from years 2-10 of the simulation,

763

snapshots from 0002-06-01.



764 **Figure 16.** EC60to30 Layers: A comparison of EC60to30 meshes with different numbers of vertical
 765 layers. The mesh used in this paper was the 100 layer mesh. The CUSP meshes used an 80 layer mesh.



766 **Figure 17.** Root mean square (RMS) of surface speed for runs forced with CORE realistic atmosphere.
 767 High-resolution 18 - 6 km eddy-permitting run (a) and coastal-refined 8 km run (b). White contour indicates
 768 0.4 m/s surface speed RMS contour. In the North Atlantic, the high-resolution grid in (a) is similar to the
 769 NA8 grid. The Gulf Stream separation, variability and transport is much more realistic in the CORE-forced
 770 high-resolution run (a) than in any of the climatology-forced runs (Figure 14).