Novel Signatures of Dark Matter in Laser-Interferometric Gravitational-Wave Detectors

Yevgeny Stadnik

Kavli Fellow

Kavli IPMU, University of Tokyo, Japan



"Gravitational Wave Probes of Fundamental Physics", Amsterdam, November 2019

Motivation

Strong astrophysical evidence for existence of **dark matter** (~5 times more dark matter than ordinary matter).



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• Low-mass spin-0 particles form a coherently oscillating classical field $\varphi(t) = \varphi_0 \cos(m_{\varphi}c^2 t/\hbar)$, with energy density $<\rho_{\varphi}> \approx m_{\varphi}^2 \varphi_0^2/2 \ (\rho_{\text{DM,local}} \approx 0.4 \text{ GeV/cm}^3)$



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 - Wave-like signatures [cf. particle-like signatures of WIMP DM]

[Stadnik, Flambaum, *PRL* **114**, 161301 (2015); *PRL* **115**, 201301 (2015)], [Hees, Minazzoli, Savalle, Stadnik, Wolf, *PRD* **98**, 064051 (2018)]

$$\mathcal{L}_{\gamma} = \frac{\phi}{\Lambda_{\gamma}} \frac{F_{\mu\nu} F^{\mu\nu}}{4} \implies \frac{\delta\alpha}{\alpha} \approx \frac{\phi_0 \cos(m_{\phi} t)}{\Lambda_{\gamma}}$$

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Solid material



$$L \sim Na_{\rm B} = N/(m_e \alpha)$$

Laser Interferometry Searches for Oscillating Variations in Fundamental Constants due to Dark Matter

[Grote, Stadnik, arXiv:1906.06193]



Michelson interferometer (GEO 600)

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• Geometric asymmetry from beam-splitter: $\delta(L_x - L_y) \sim \delta(nI)$

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- Geometric asymmetry from beam-splitter: $\delta(L_x L_y) \sim \delta(nI)$
- Both broadband and resonant narrowband searches possible: $f_{DM} \approx f_{vibr,BS} \sim v_{sound}/I$, $Q \sim 10^6$ enhancement

Michelson vs Fabry-Perot-Michelson Interferometers

[Grote, Stadnik, arXiv:1906.06193]

Michelson interferometer (GEO 600, Fermilab holometer) Fabry-Perot-Michelson interferometer (LIGO, VIRGO, KAGRA)



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Linear Interaction of Scalar Dark Matter with the Electron



Linear Interaction of Scalar Dark Matter with the Electron



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Summary

- Existing laser-interferometric gravitational-wave detectors are sensitive probes of scalar dark-matter fields oscillating at audio-band frequencies
- Changing arm mirror thicknesses by ~10% can greatly boost the sensitivity of Fabry-Perot-Michelson interferometers (LIGO, VIRGO, KAGRA) to dark matter
- (Small-scale) Interferometry experiments can be adapted to perform resonant narrowband searches
- Existing interferometers also sensitive to the passage of macroscopic dark-matter objects through detectors

Back-Up Slides

Temporal Coherence

• Low-mass spin-0 particles form a coherently oscillating classical field $\varphi(t) = \varphi_0 \cos(m_{\varphi}c^2t/\hbar)$, with energy density $<\rho_{\varphi}> \approx m_{\varphi}^2 \varphi_0^2/2 \ (\rho_{\text{DM,local}} \approx 0.4 \text{ GeV/cm}^3)$

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Probability distribution function of φ_0



[Stadnik, Flambaum, PRL 114, 161301 (2015); PRL 115, 201301 (2015)], [Hees, Minazzoli, Savalle, Stadnik, Wolf, PRD 98, 064051 (2018)]

1

 $\Gamma \mu \nu$

$$\mathcal{L}_{\gamma} = \frac{\phi}{\Lambda_{\gamma}} \frac{F_{\mu\nu} F^{\mu\nu}}{4} \implies \frac{\delta\alpha}{\alpha} \approx \frac{\phi_{0} \cos(m_{\phi}t)}{\Lambda_{\gamma}}$$

$$\mathcal{L}_{f} = -\frac{\phi}{\Lambda_{f}} m_{f} \bar{f}f \implies \frac{\delta m_{f}}{m_{f}} \approx \frac{\phi_{0} \cos(m_{\phi}t)}{\Lambda_{f}}$$

$$\phi = \phi_{0} \cos(m_{\phi}t - \underline{p}_{\phi} \cdot \underline{x}) \implies F \propto \underline{p}_{\phi} \sin(m_{\phi}t)$$

$$\mathcal{L}_{\gamma}' = \frac{\phi^{2}}{(\Lambda_{\gamma}')^{2}} \frac{F_{\mu\nu} F^{\mu\nu}}{4}$$

$$\mathcal{L}_{f}' = -\frac{\phi^{2}}{(\Lambda_{f}')^{2}} m_{f} \bar{f}f$$

$$= \sum \frac{\delta\alpha}{\alpha} \propto \frac{\delta m_{f}}{m_{f}} \propto \delta\rho_{\phi}$$

$$F \propto \nabla\rho_{\phi}$$

[Stadnik, Flambaum, *PRL* **114**, 161301 (2015); *PRL* **115**, 201301 (2015)], [Hees, Minazzoli, Savalle, Stadnik, Wolf, *PRD* **98**, 064051 (2018)]

Consider <u>quadratic couplings</u> of an oscillating classical scalar field, $\varphi(t) = \varphi_0 \cos(m_{\varphi}t)$, with SM fields.

$$\mathcal{L}_{f} = -\frac{\phi^{2}}{(\Lambda_{f}')^{2}} m_{f} \bar{f} f \quad \text{c.f.} \quad \mathcal{L}_{f}^{\text{SM}} = -m_{f} \bar{f} f \quad => \quad m_{f} \to m_{f} \left[1 + \frac{\phi^{2}}{(\Lambda_{f}')^{2}} \right]$$
$$= > \frac{\delta m_{f}}{m_{f}} = \frac{\phi_{0}^{2}}{(\Lambda_{f}')^{2}} \cos^{2}(m_{\phi}t) = \left[\frac{\phi_{0}^{2}}{2(\Lambda_{f}')^{2}} + \frac{\phi_{0}^{2}}{2(\Lambda_{f}')^{2}} \cos(2m_{\phi}t) \right]$$
$$\rho_{\phi} = \frac{m_{\phi}^{2}\phi_{0}^{2}}{2} \implies \phi_{0}^{2} \propto \rho_{\phi}$$

[Stadnik, Flambaum, *PRL* **114**, 161301 (2015); *PRL* **115**, 201301 (2015)], [Hees, Minazzoli, Savalle, Stadnik, Wolf, *PRD* **98**, 064051 (2018)]

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Fifth Forces: Linear vs Quadratic Couplings [Hees, Minazzoli, Savalle, Stadnik, Wolf, PRD 98, 064051 (2018)] Consider the effect of a massive body (e.g., Earth) on the scalar DM field Linear couplings ($\varphi \bar{X} X$) Quadratic couplings ($\varphi^2 \bar{X} X$) $\phi = \phi_0 \cos(m_\phi t) - A \frac{e^{-m_\phi r}}{r} \qquad \phi = \phi_0 \cos(m_\phi t) \left(1 - \frac{B}{r}\right)$

Gradients + screening/amplification



Gradients + screening/amplification



"Fifth-force" experiments: torsion pendula, atom interferometry

Gradients + screening/amplification

Constraints on Linear Interaction of Scalar Dark Matter with the Electron



Quartic Self-Interaction of Scalar



Constraints on Linear Interaction of Scalar Dark Matter with the Higgs Boson

Rb/Cs constraints:

[Stadnik, Flambaum, PRA 94, 022111 (2016)]

2 – 3 orders of magnitude improvement!

