

# A classification of CEMP stars based on neutron density that reveals the important role of the i process and the need for better nuclear physics data

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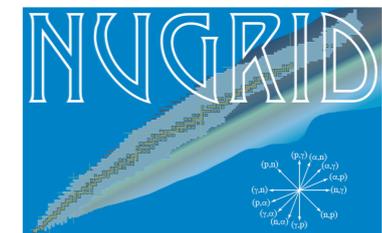
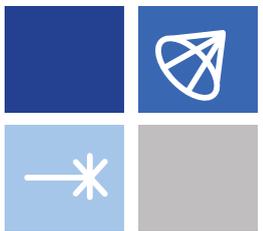
Pavel Denissenkov, Ondrea Clarkson, David Stephens

Hydro collaboration: Paul Woodward, Huaqing Mao (Minnesota), Robert Andrassy (HITS)

NuGrid/JINA collaboration: Marco Pignatari, Benoit Côté, Hendrik Schatz, Gorgios Perdikakis

NuGrid@NPA IX: Jacqueline den Hartog, Sam Lloyd, Umberto Battino, Raphael Hirschi, Deniz Kurtulgil, Alison Laird, Rene Reifarh, Richard Stancliffe, Adria Casanova, Athanasios Psaltis, Adria Casanovas, César Domingo Pardo

**JINA-CEE**  
NSF Physics Frontier Center

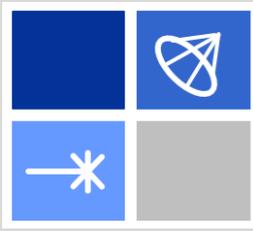


University of Victoria



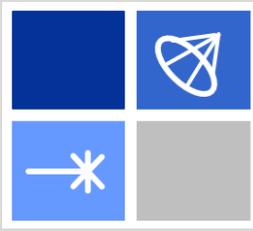
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NPA IX, Schloß Waldhausen, 16 Sep 2019



# Outline

- C-enhanced metal poor stars: Classification with constant neutron-density models (nuclear physics only!) → the need for the i process
- Rapidly accreting white dwarfs as a likely site for the i process
- Nuclear physics impact studies



# C-enhanced metal-poor stars

## Classification of CEMP stars

**TABLE 2** Definition of subclasses of metal-poor stars

### Neutron-capture-rich stars

r-I	$0.3 \leq [\text{Eu}/\text{Fe}] \leq +1.0$ and $[\text{Ba}/\text{Eu}] < 0$
r-II	$[\text{Eu}/\text{Fe}] > +1.0$ and $[\text{Ba}/\text{Eu}] < 0$
s	$[\text{Ba}/\text{Fe}] > +1.0$ and $[\text{Ba}/\text{Eu}] > +0.5$
r/s	$0.0 < [\text{Ba}/\text{Eu}] < +0.5$

### Carbon-enhanced metal-poor stars

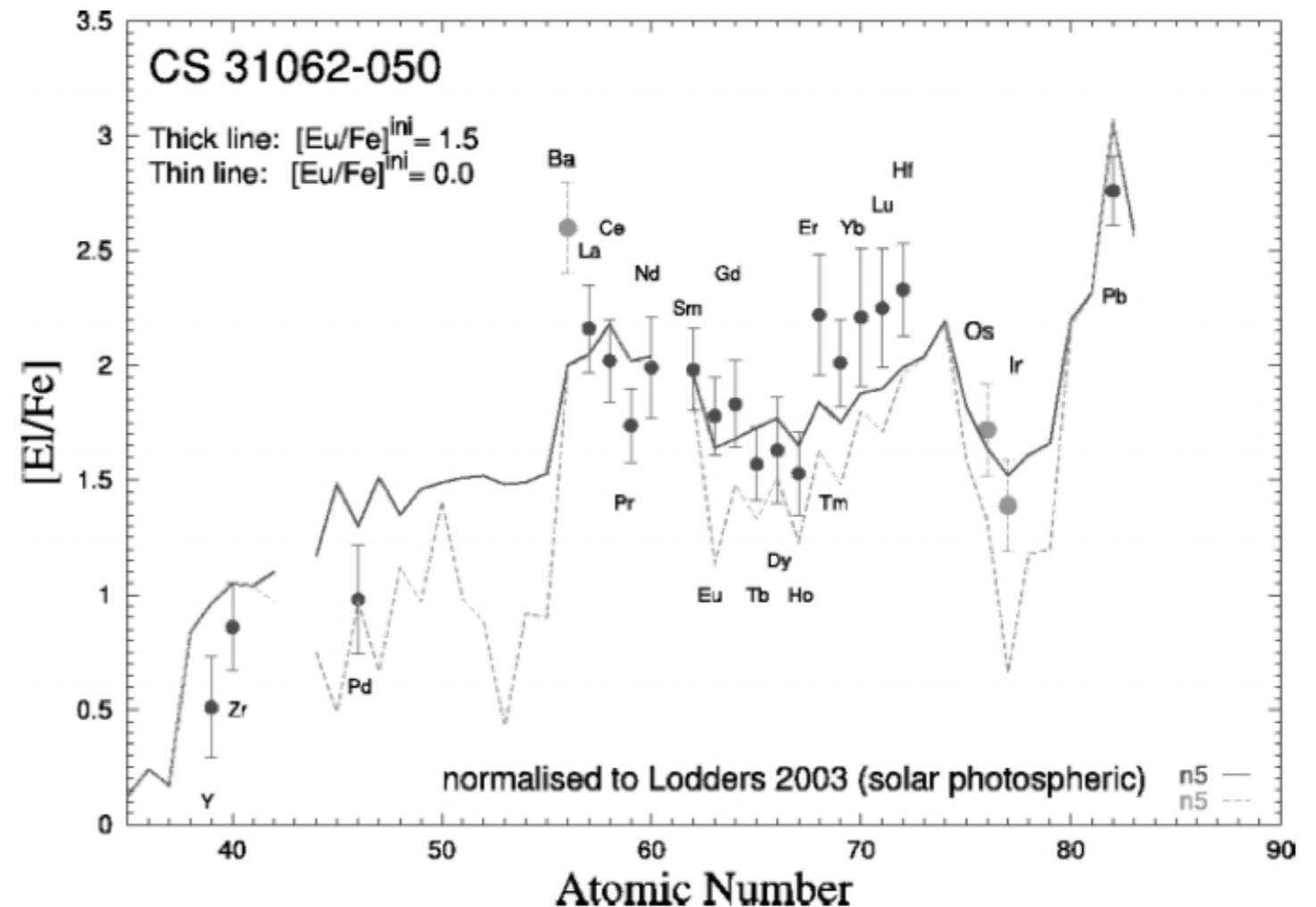
CEMP	$[\text{C}/\text{Fe}] > +1.0$
CEMP-r	$[\text{C}/\text{Fe}] > +1.0$ and $[\text{Eu}/\text{Fe}] > +1.0$
CEMP-s	$[\text{C}/\text{Fe}] > +1.0$ , $[\text{Ba}/\text{Fe}] > +1.0$ , and $[\text{Ba}/\text{Eu}] > +0.5$
CEMP-r/s	$[\text{C}/\text{Fe}] > +1.0$ and $0.0 < [\text{Ba}/\text{Eu}] < +0.5$
CEMP-no	$[\text{C}/\text{Fe}] > +1.0$ and $[\text{Ba}/\text{Fe}] < 0$

- Classification-scheme based on Ba and Eu (which can be observed in many stars)

*Beers & Christlieb 05 (ARAA)*

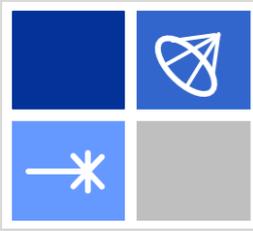
- CEMP-r/s stars have both Ba and Eu
- previously best model assumes superposition of “known” r- and s-process sources

Is there a more plausible explanation?



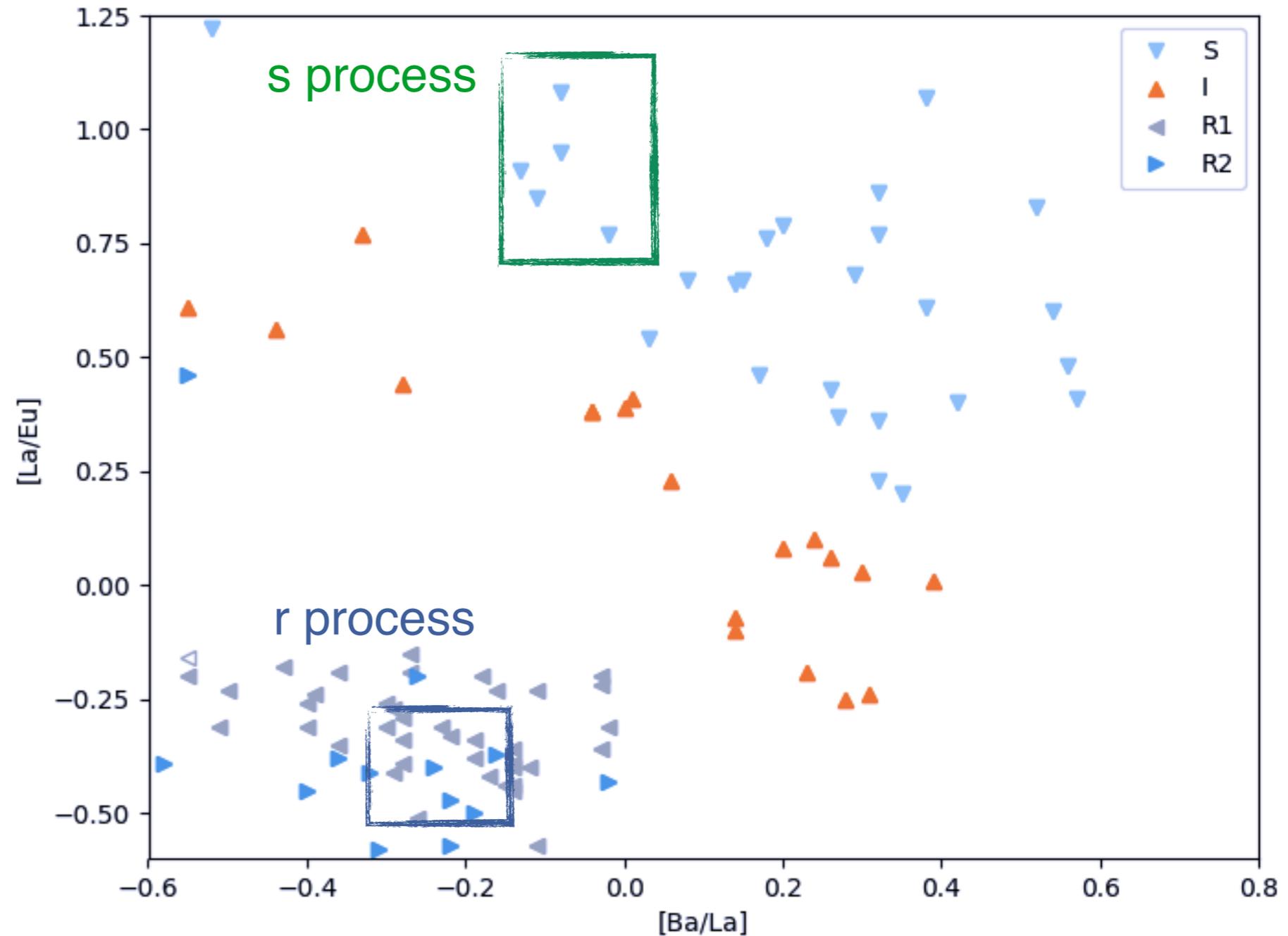
*Bisterzo+ 12; observations Aoki+ 02 & Johnson & Bolte 04*



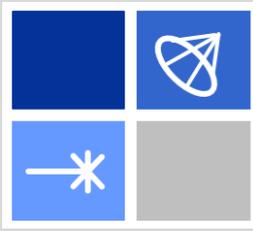


# Observed abundances in CEMP stars

- The CEMP stars have (sometimes) very large enhancements of heavy elements, such as La, Ba, Eu
- C-enhanced metal poor (CEMP) stars carry the nucleosynthesis signature of the **rise of the elements in the early universe**
- CEMP-s, -r, -r/s stars: [Ba/La] and [La/Eu] elemental ratios very different for known **s process** and **r process** nucleosynthesis

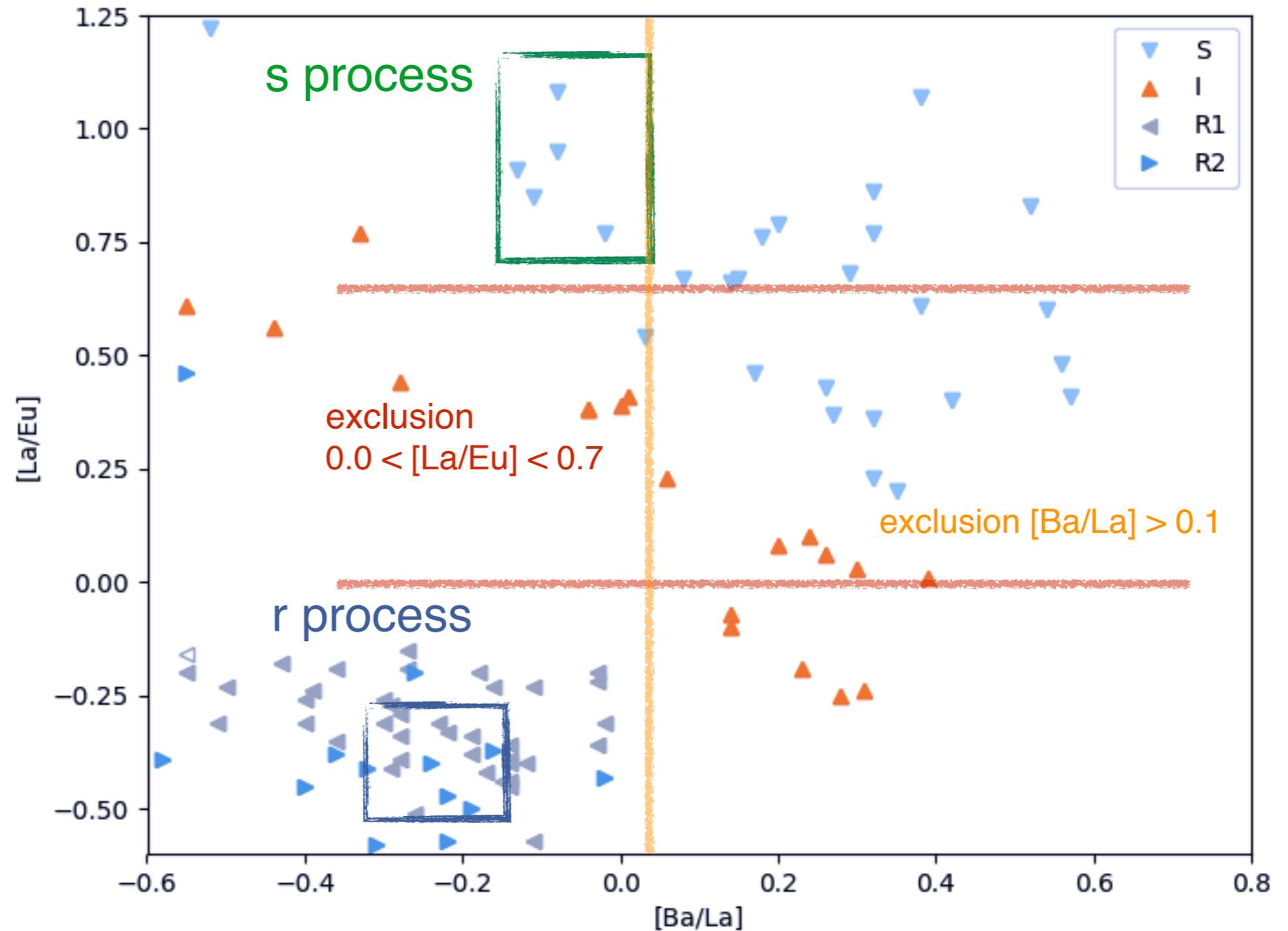


Observations JINAbase (Abohalima & Frebel, 2018, <http://jinabase.pythonanywhere.com>)



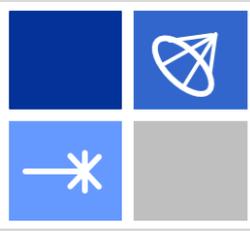
# Observed abundances in CEMP stars

- Neither s nor r can make  $0.0 < [\text{La}/\text{Eu}] < 0.7$ .
- With a superposition of just s and r it is impossible to make  $[\text{Ba}/\text{La}] > 0.1$ .
- But many CEMP stars have element ratios in exclusion zones!!!!



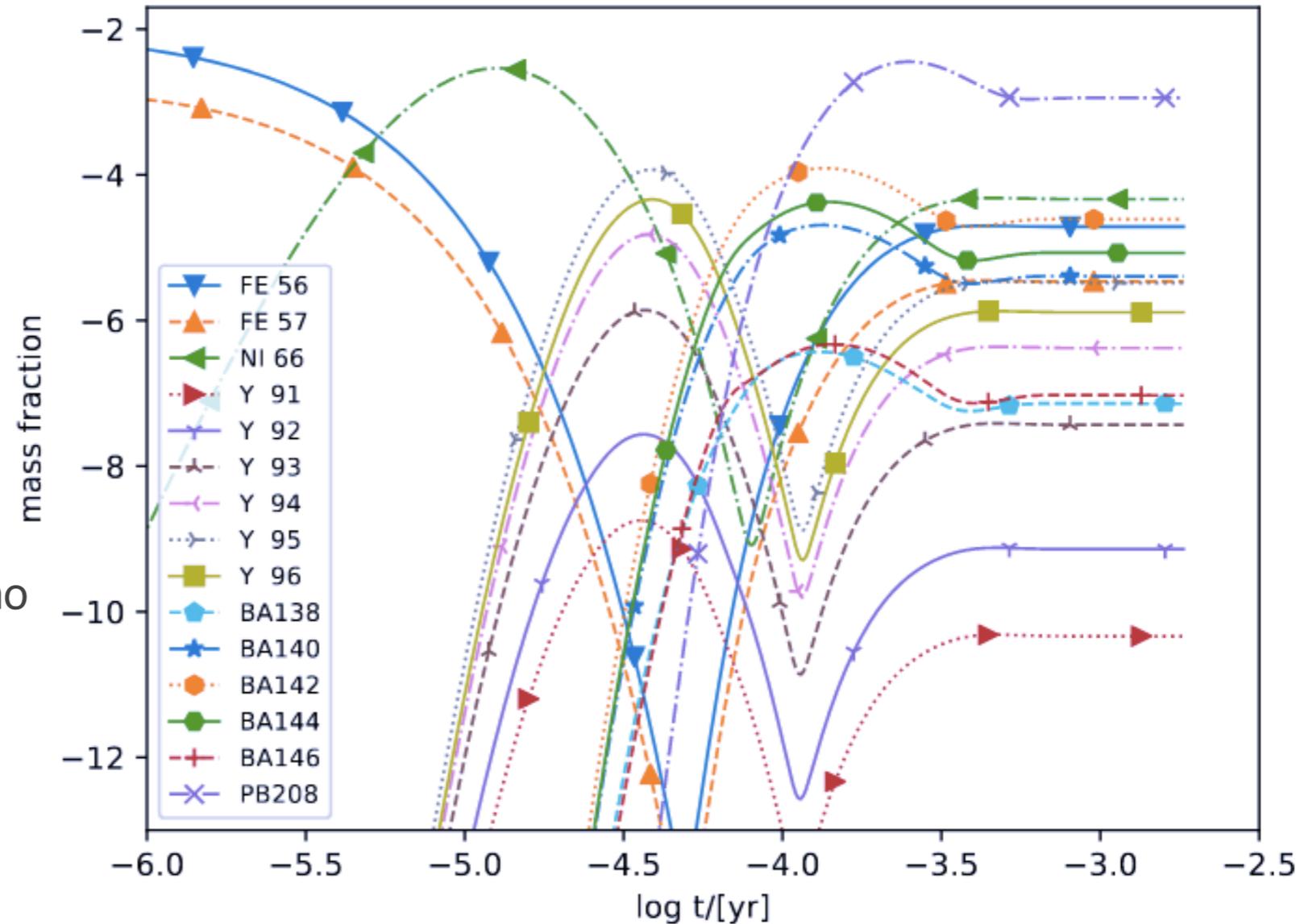
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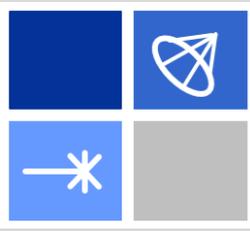




# Constant neutron density one-zone nucleosynthesis network calculations

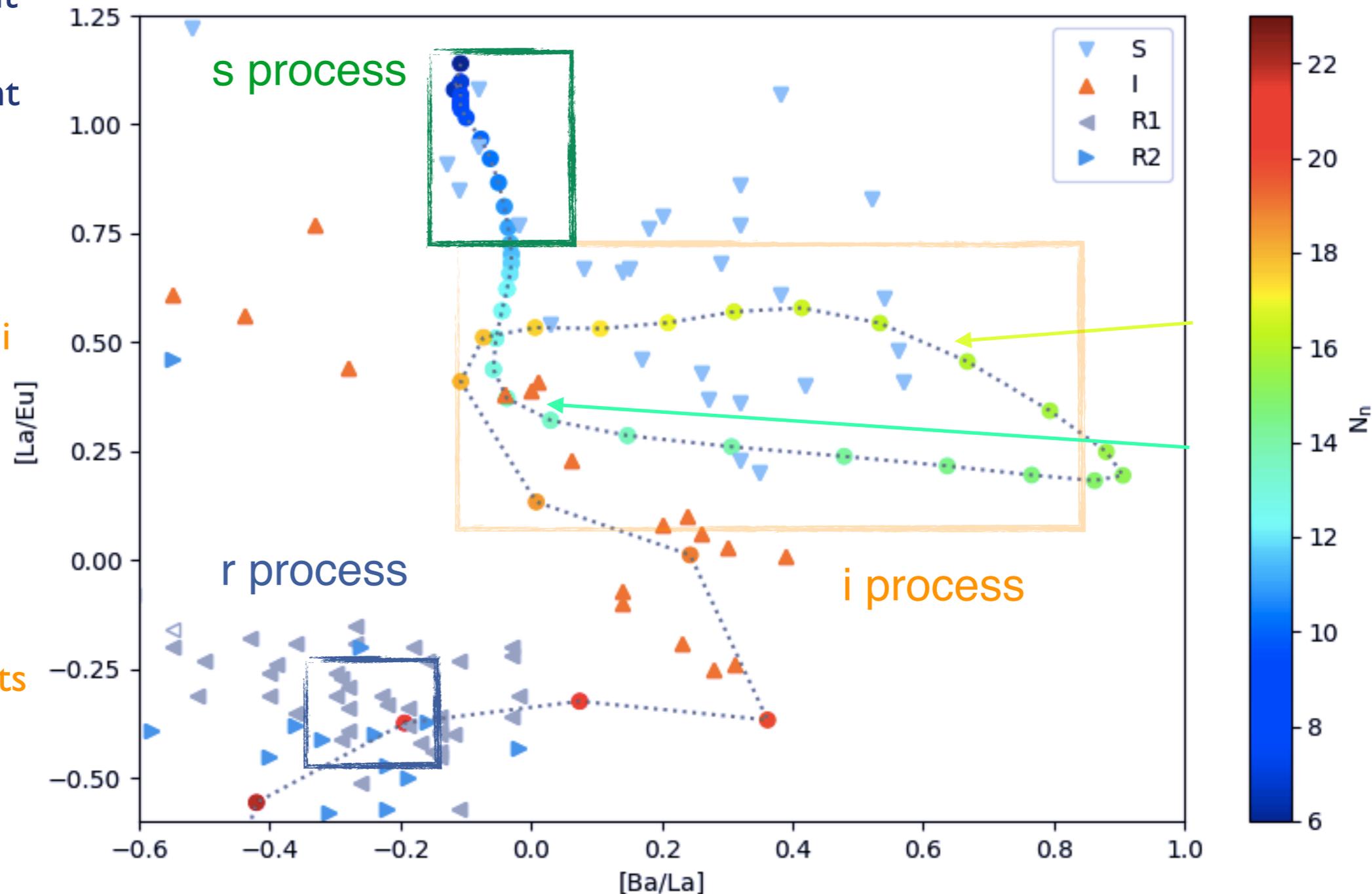
- Neutron density  $N_n$  is input, feed network endpoint  $^{209}\text{Bi}$  back to  $^{56}\text{Fe}$  seed
- This is the method used in the past for the classical s process
- Once equilibrium is reached we decay, and the result is a single elemental abundance distribution for the given value of  $N_n$
- Equilibrium models are independent of neutron exposure  $\rightarrow$  no information on ratio of elements in different peaks, e.g. no Ba/Sr
- Instead astrophysics-independent (pure nuclear physics) predictions (only!) for nearby elemental ratios, such as Ba/La
- Potentially a power-full new method for nuclear physics-based classification

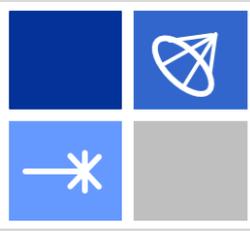




# The i-process covers many observations that can't be explained by s and r process alone → intermediate process

- Characterize different neutron capture regimes with constant neutron density equilibrium models. Depend *only* on nuclear physics
- The intermediate or i process at neutron densities higher than for s and lower than for r allows us to explain many of the CEMP stars with n-capture enhancements





# [Ba/Eu] & CEMP star classification & s-process origin

**TABLE 2** Definition of subclasses of metal-poor stars

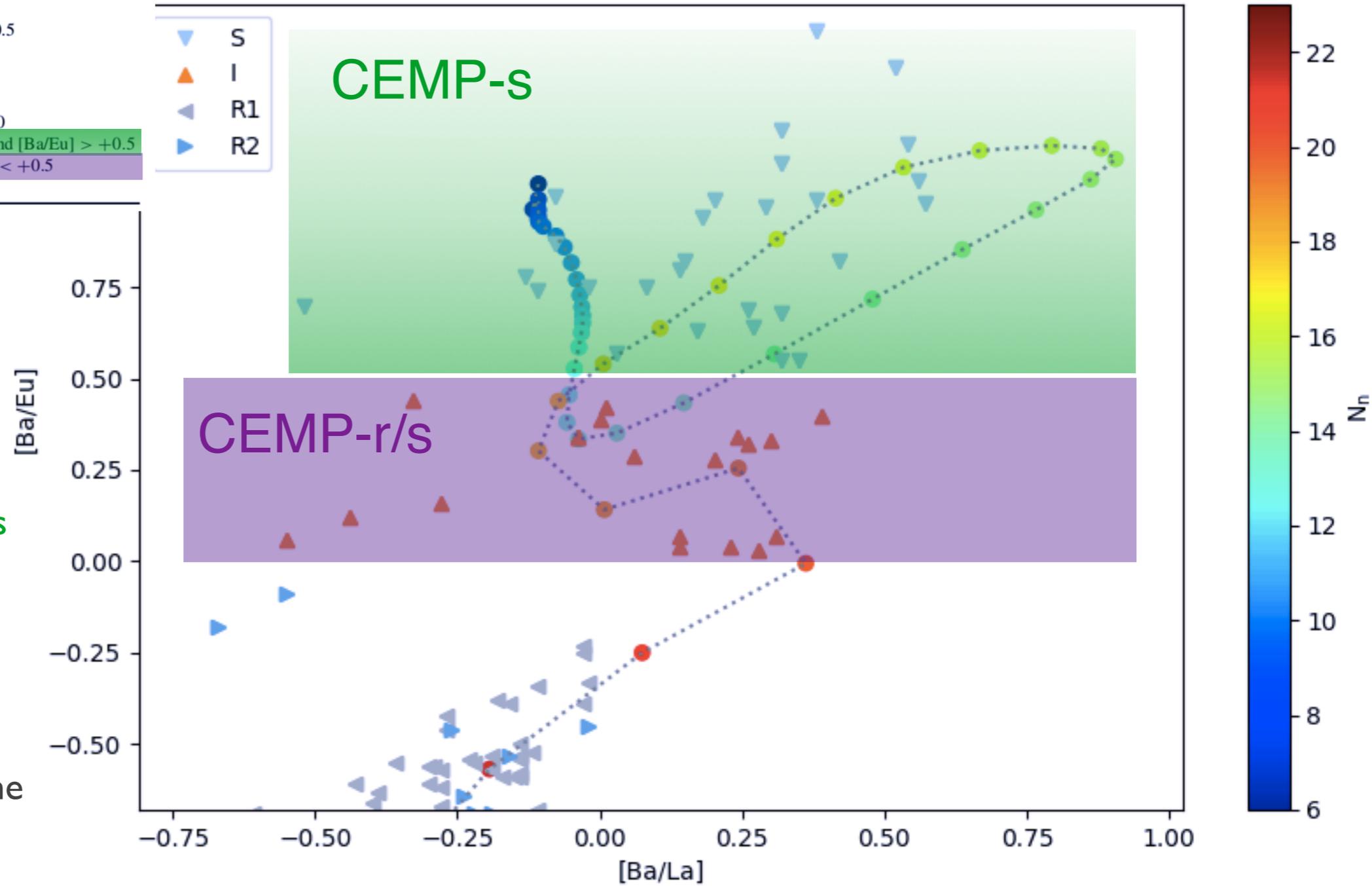
Neutron-capture-rich stars

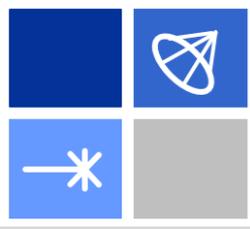
r-I	$0.3 \leq [\text{Eu}/\text{Fe}] \leq +1.0$ and $[\text{Ba}/\text{Eu}] < 0$
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s	$[\text{Ba}/\text{Fe}] > +1.0$ and $[\text{Ba}/\text{Eu}] > +0.5$
r/s	$0.0 < [\text{Ba}/\text{Eu}] < +0.5$

Carbon-enhanced metal-poor stars

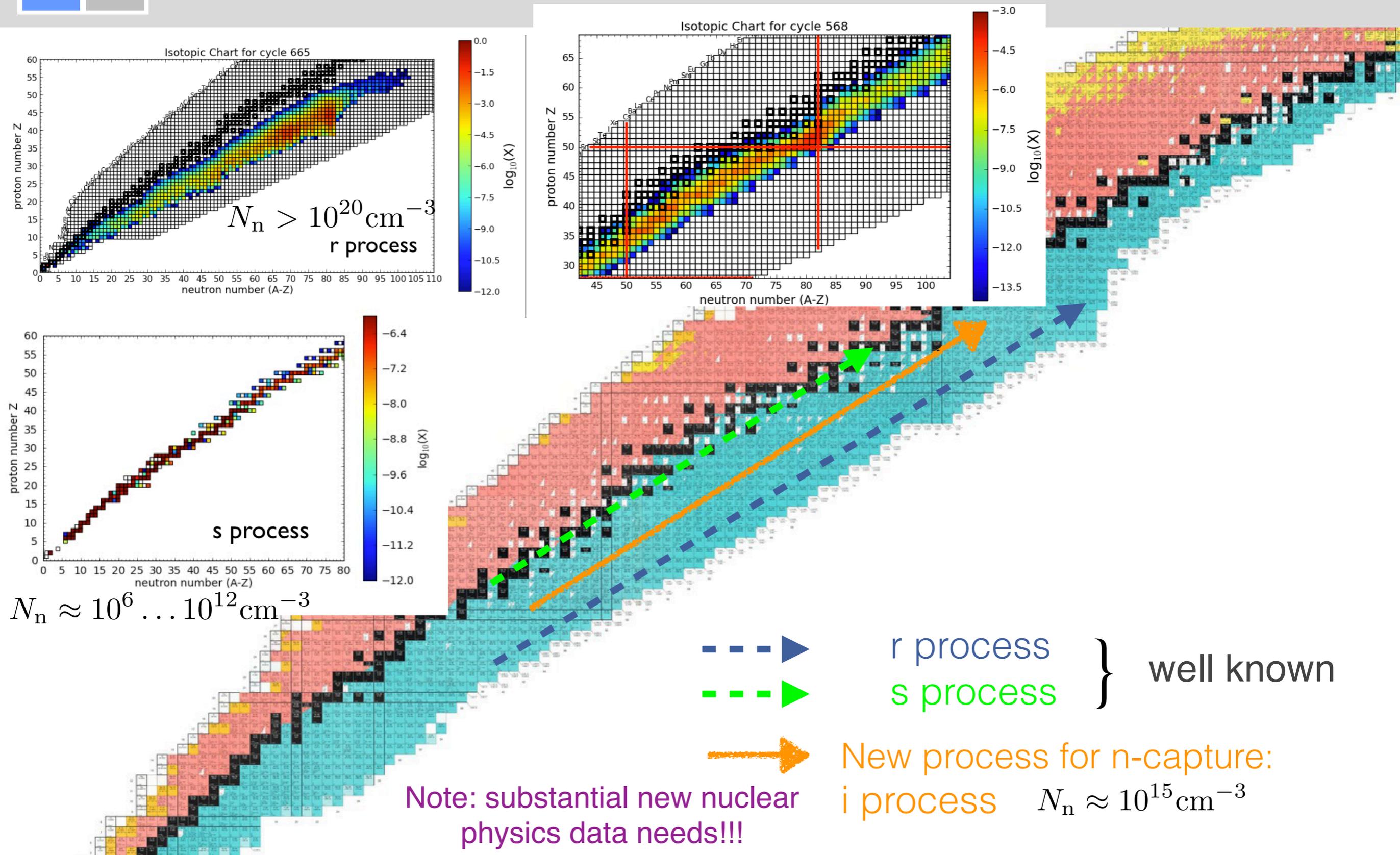
CEMP	$[\text{C}/\text{Fe}] > +1.0$
CEMP-r	$[\text{C}/\text{Fe}] > +1.0$ and $[\text{Eu}/\text{Fe}] > +1.0$
CEMP-s	$[\text{C}/\text{Fe}] > +1.0$ , $[\text{Ba}/\text{Fe}] > +1.0$ , and $[\text{Ba}/\text{Eu}] > +0.5$
CEMP-r/s	$[\text{C}/\text{Fe}] > +1.0$ and $0.0 < [\text{Ba}/\text{Eu}] < +0.5$
CEMP-no	$[\text{C}/\text{Fe}] > +1.0$ and $[\text{Ba}/\text{Fe}] < 0$

- According to present classification CEMP-r/s stars are indeed incompatible with s process.
- But many CEMP-s stars are also incompatible with an s-process signature because they have  $[\text{Ba}/\text{La}] > 0.1$ .
- $\rightarrow$   $[\text{Ba}/\text{Eu}]$  alone insufficient to determine s-process nucleosynthesis origin

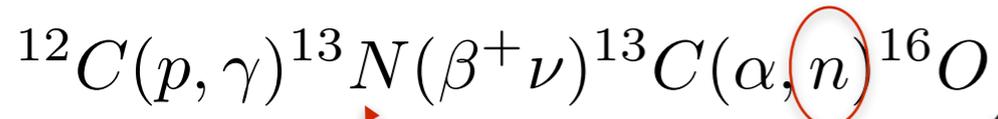




# Neutron-capture: slow, rapid and intermediate!



The convective He-burning shell contains ~40%  $^{12}\text{C}$

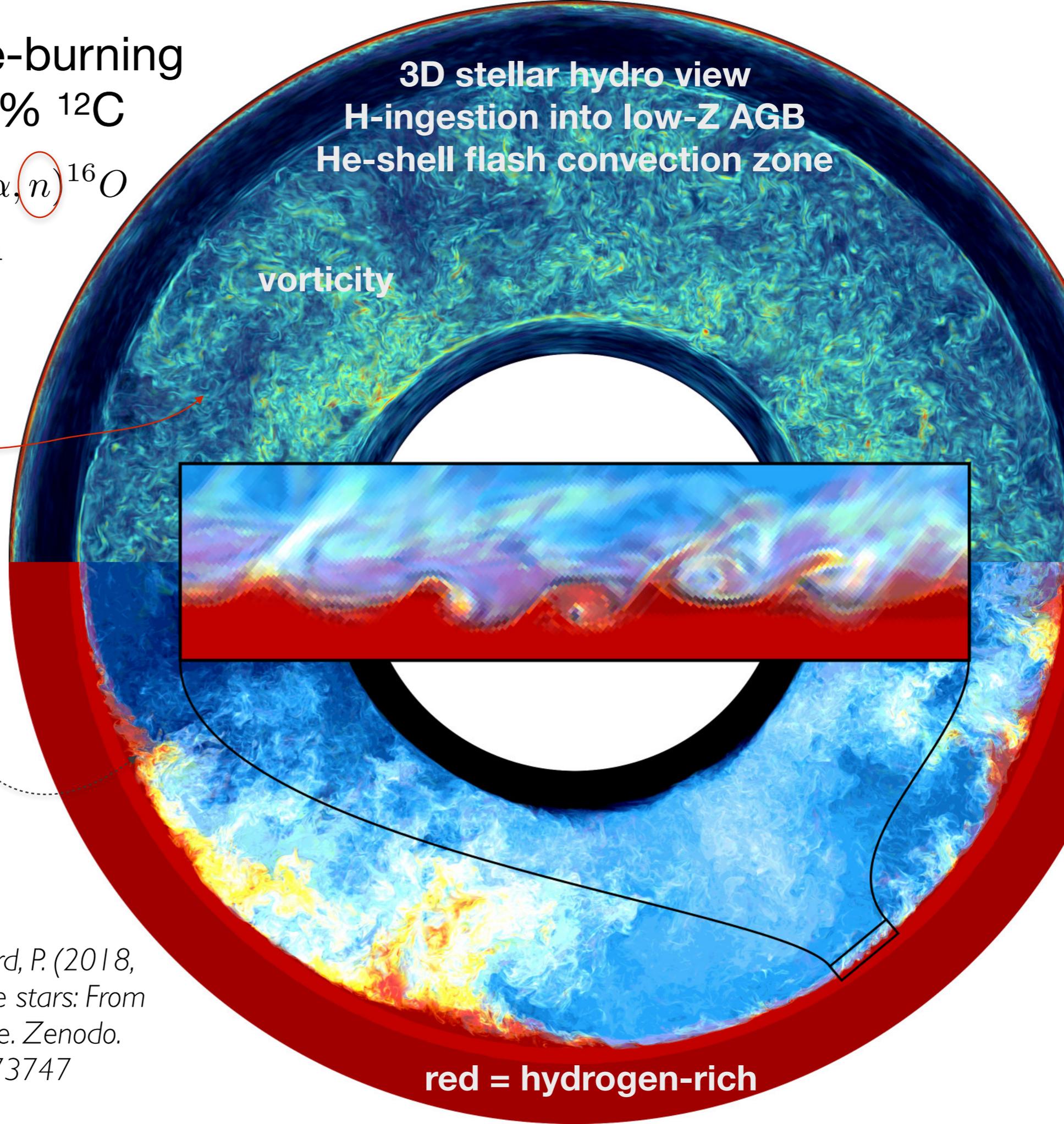


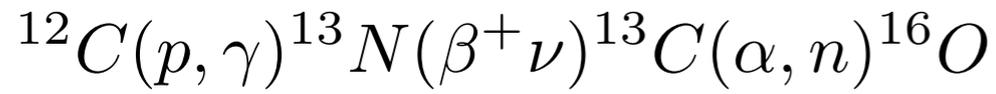
$\tau_{\frac{1}{2}} = 9.6\text{m}$

$\tau_{\text{conv}} \sim 15\text{m}$

entrainment/  
ingestion of  
H in the He-  
shell  
convection

Andrassy, R., Herwig, F., & Woodward, P. (2018, October). 3D convection in massive stars: From the main sequence to core collapse. Zenodo. <http://doi.org/10.5281/zenodo.1473747>





$$\tau_{\text{conv}} \sim 15\text{m} \longleftrightarrow \tau_{\frac{1}{2}} = 9.6\text{m}$$

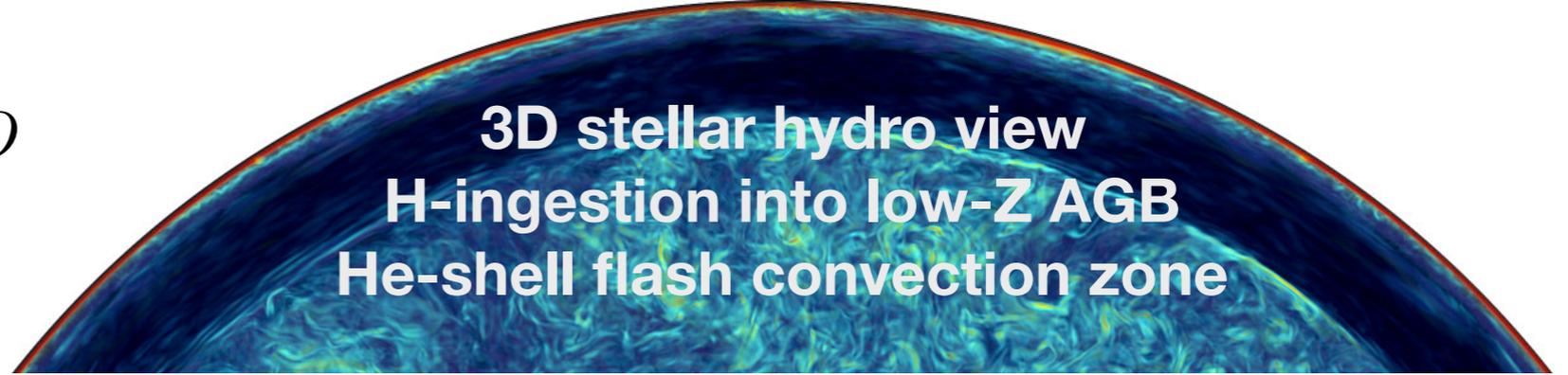
Nuclear and hydrodynamic timescales are the same order

→ convective-reactive nucleosynthesis.

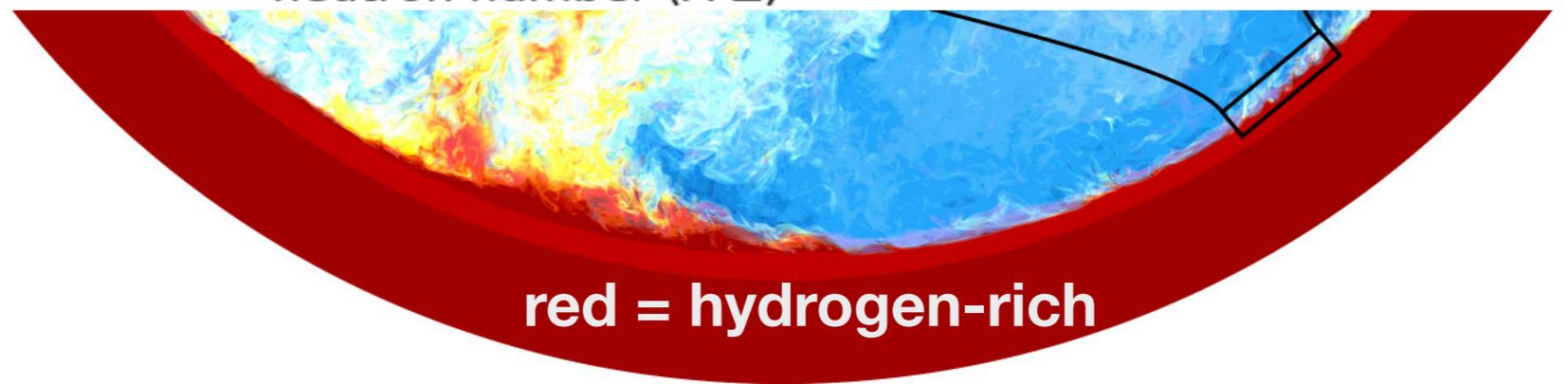
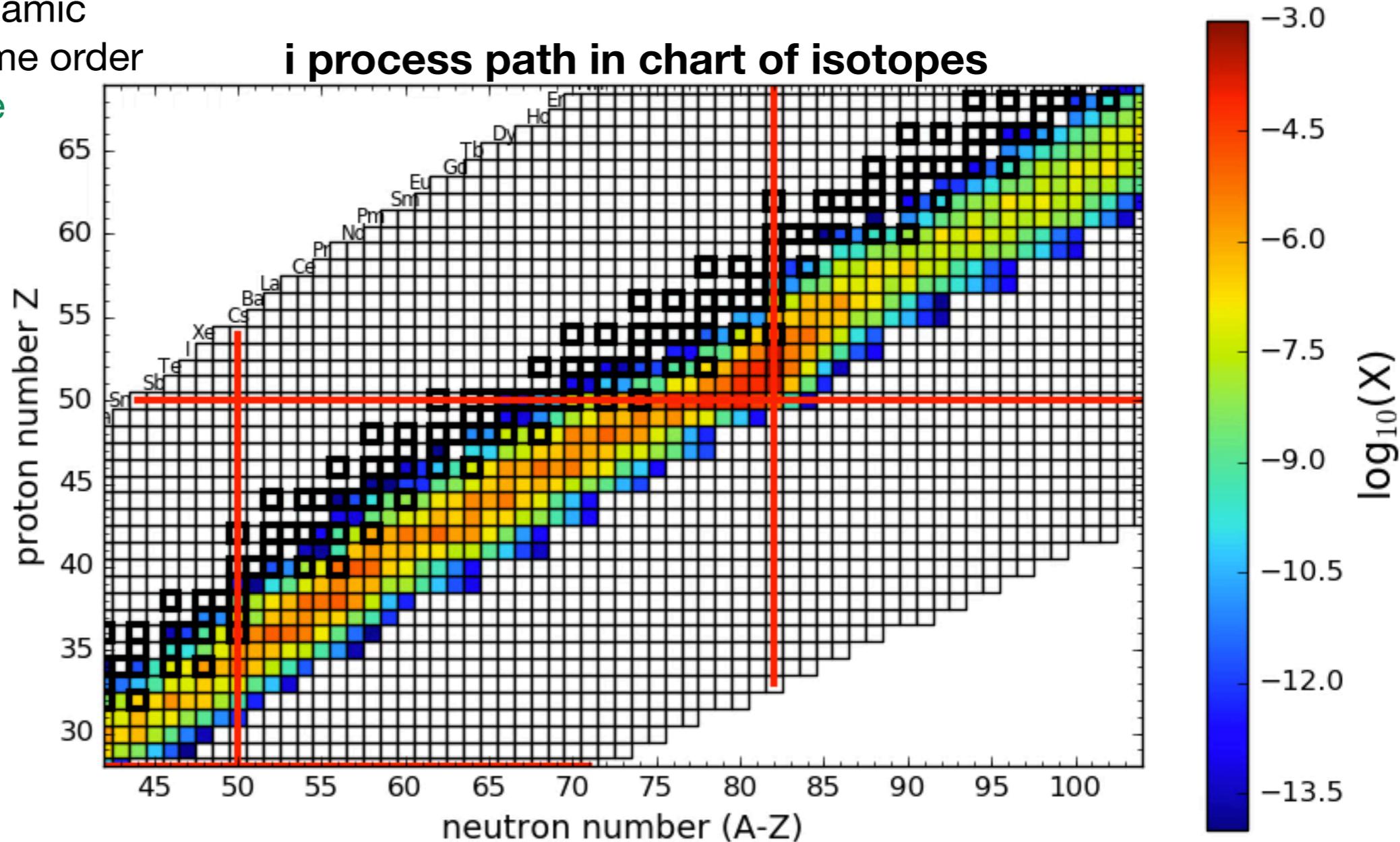
Neutrons released with intermediate neutron density

→ i process element production.

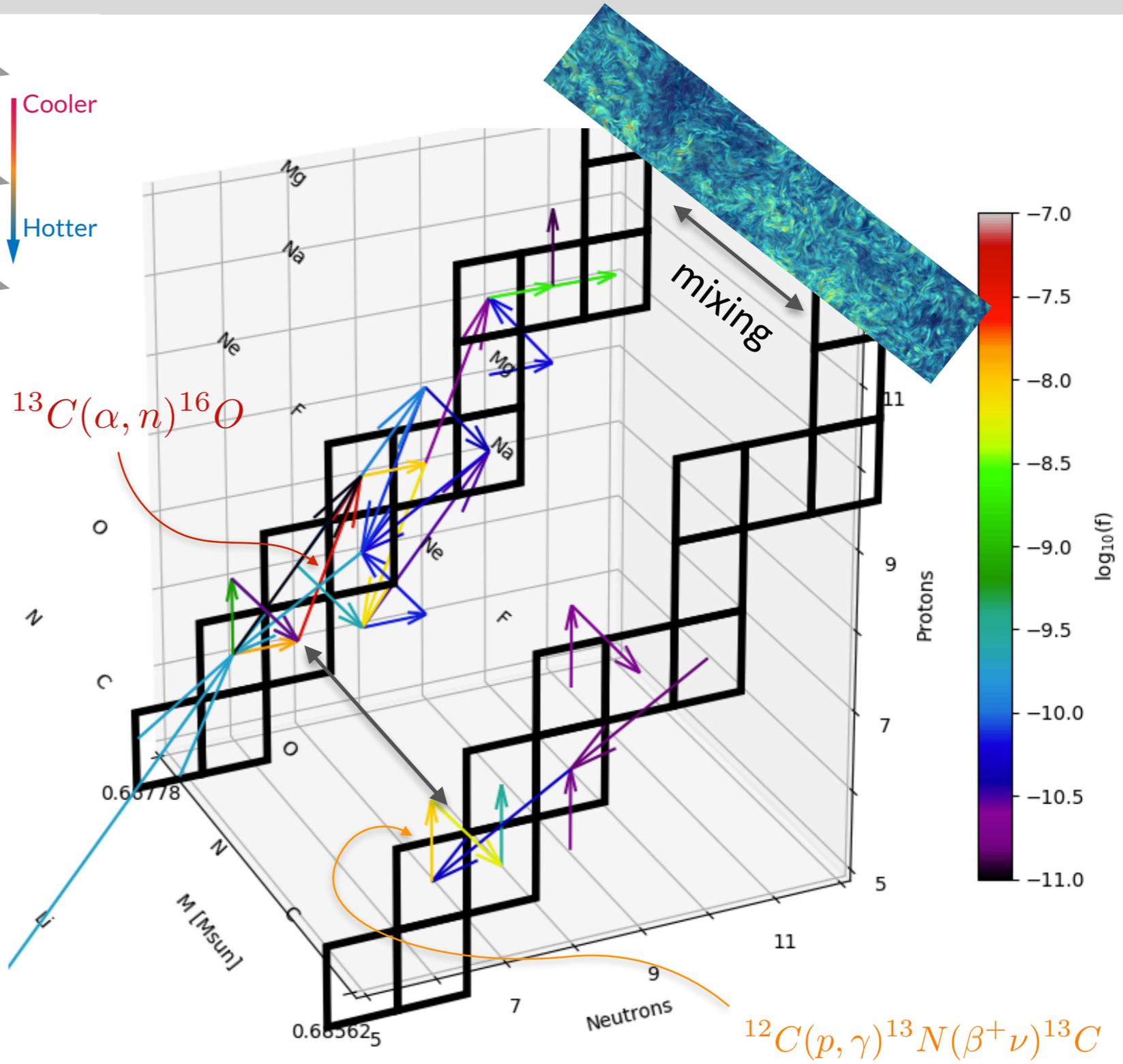
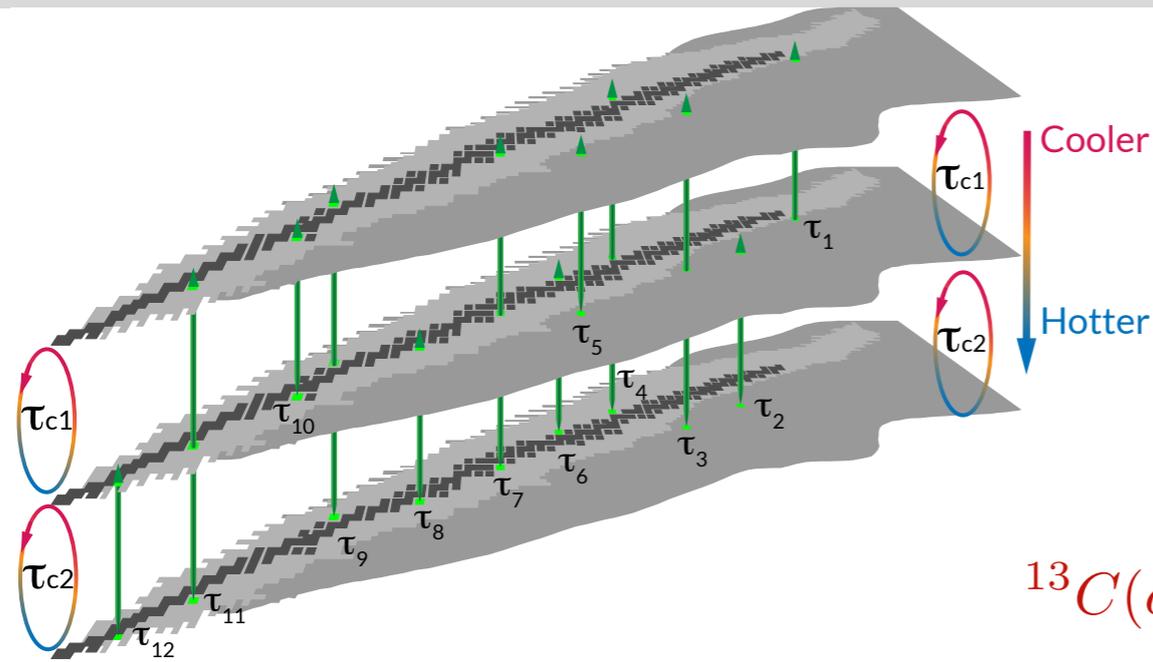
Need multi-physics, multi-method approach. How to combine detailed 3D hydro with detailed n-rich nucleosynthesis involving thousands of species?



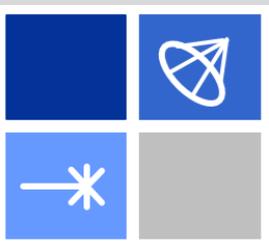
i process path in chart of isotopes



# Convective-reactive i-process nucleosynthesis

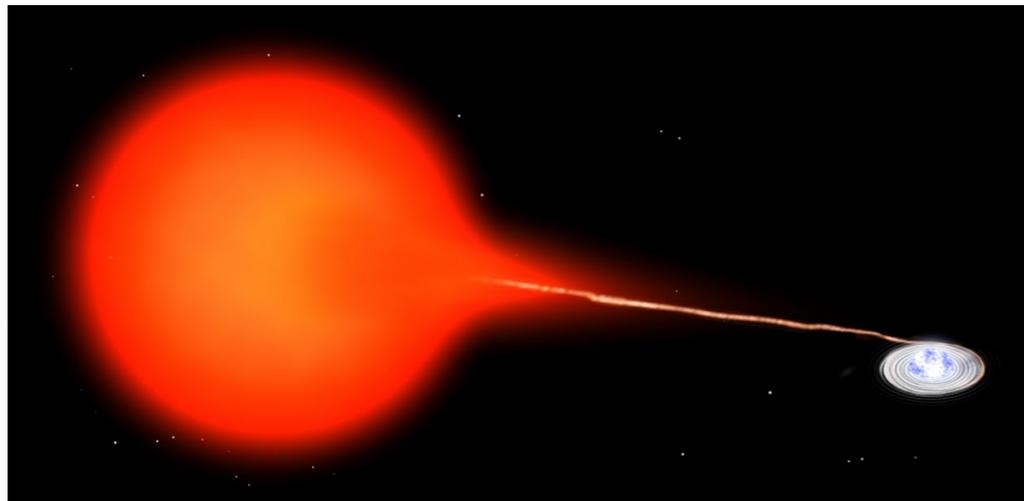


Rapid production of neutrons requires convective mixing connection of two different T regimes for  $^{12}\text{C}(p, \gamma)$  and  $^{13}\text{C}(\alpha, n)$  to operate on same time scale



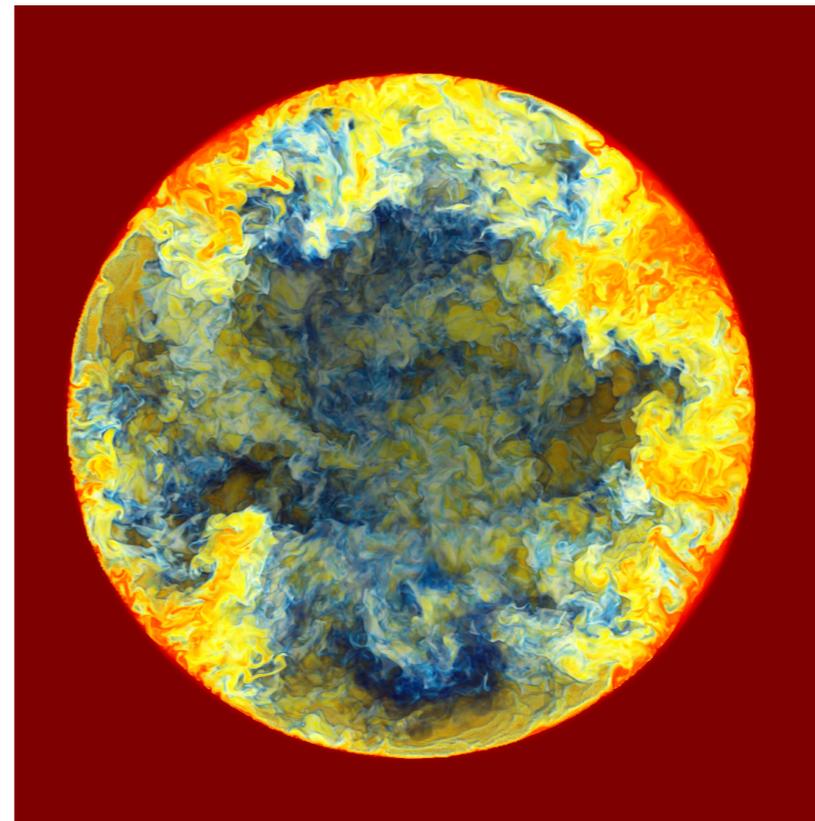
## Where does it happen?

# One promising option for CEMP-i stars: Rapidly Accreting White Dwarfs



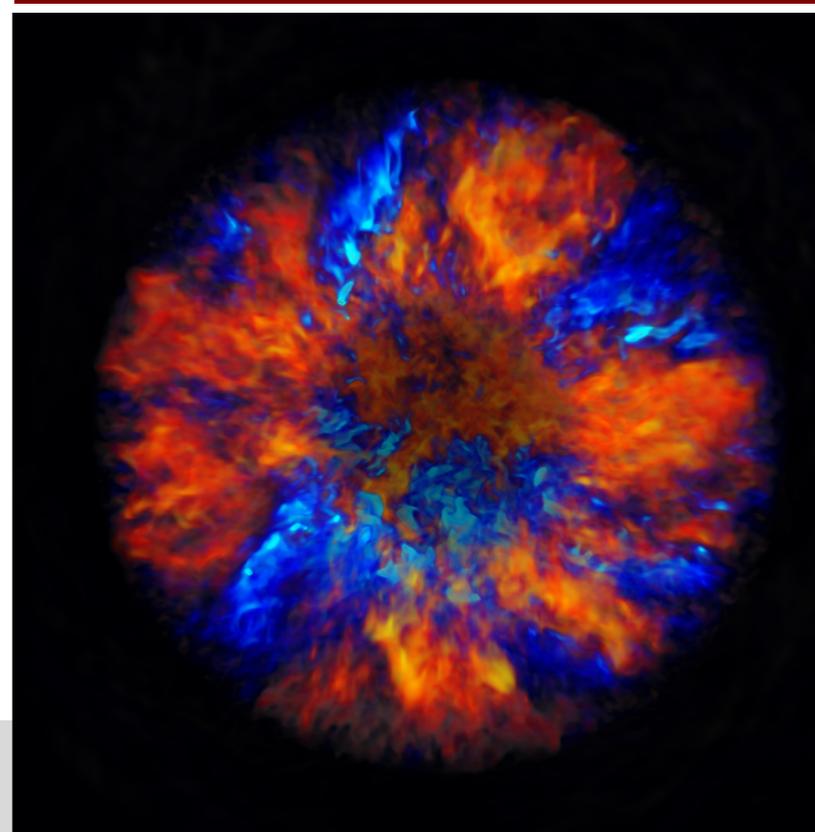
- Artist impression of accreting white dwarf, like novae!
- But unlike nova here accretion rates are high and allow stable H burning!
- However, these accreting WDs then experience **He-shell flashes!** (Cassisi+ 98)
- In these convective He-shell flashes: H-entrainment, **convective reactive i process!**

*Denissenkov+ 17, ApJ Letters*

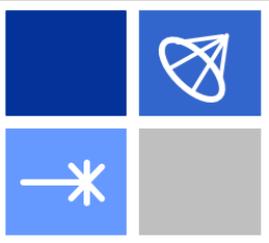


Concentration of entrained material

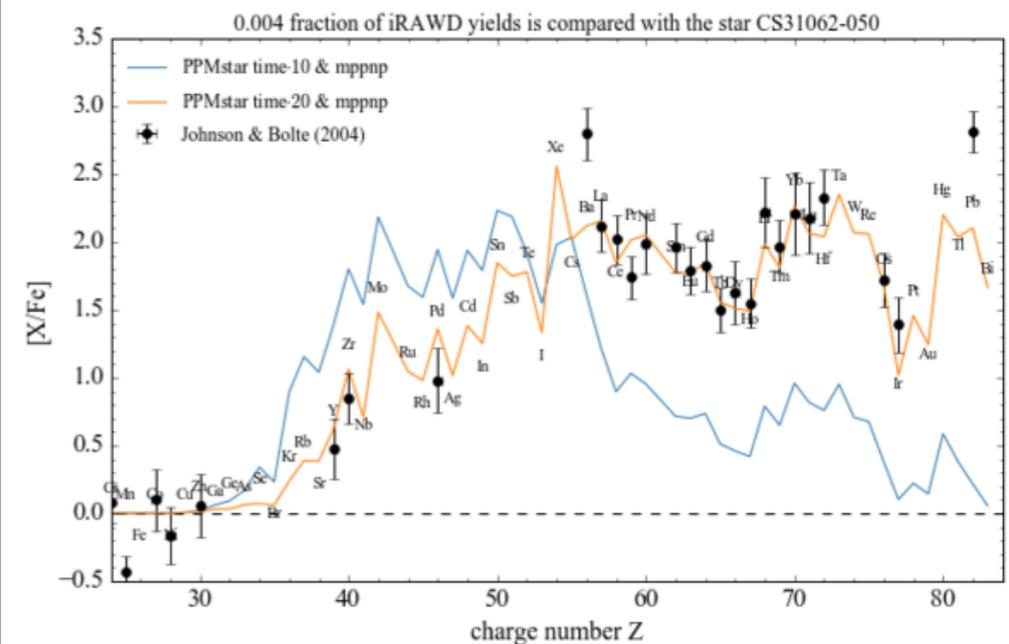
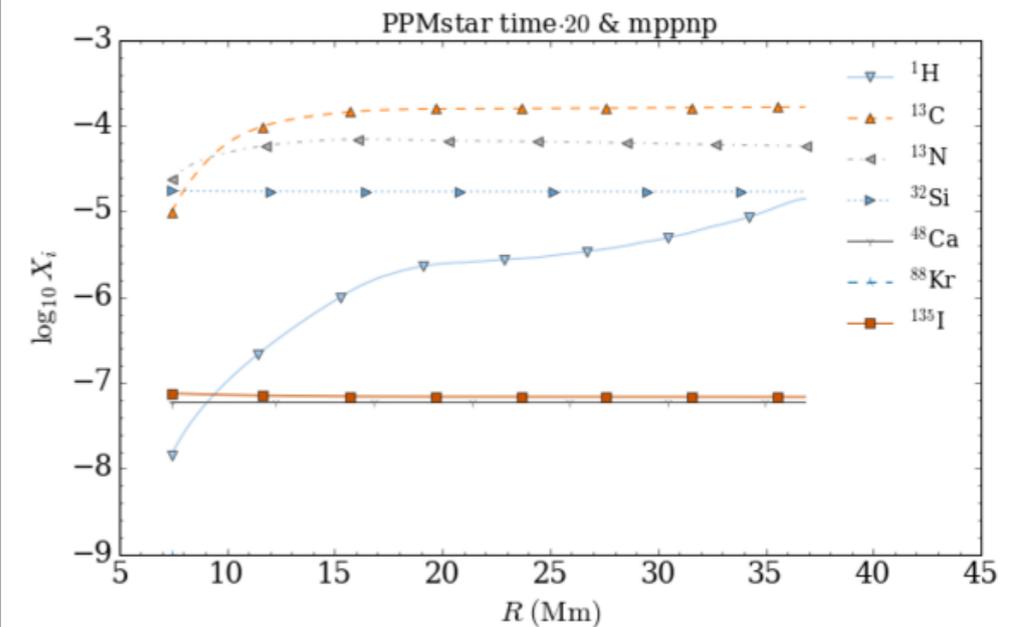
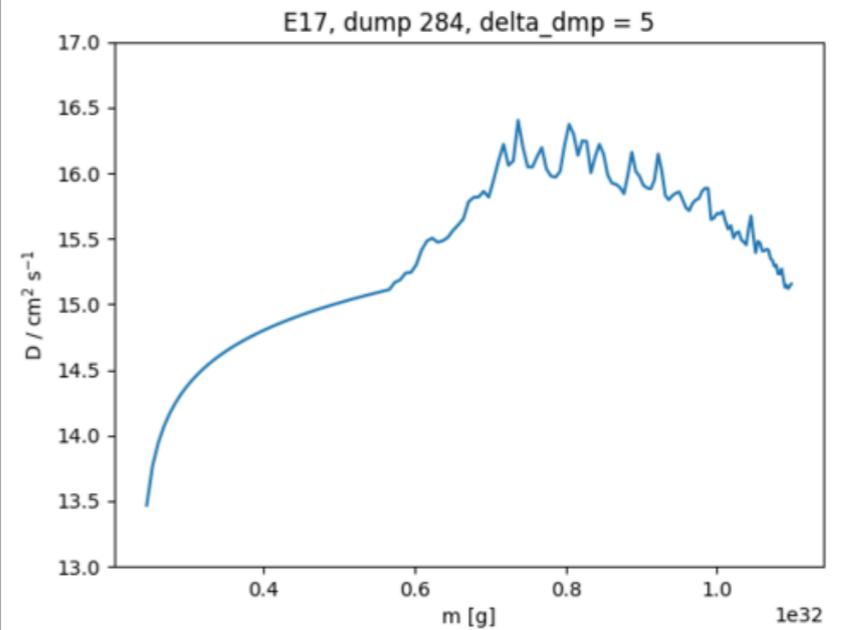
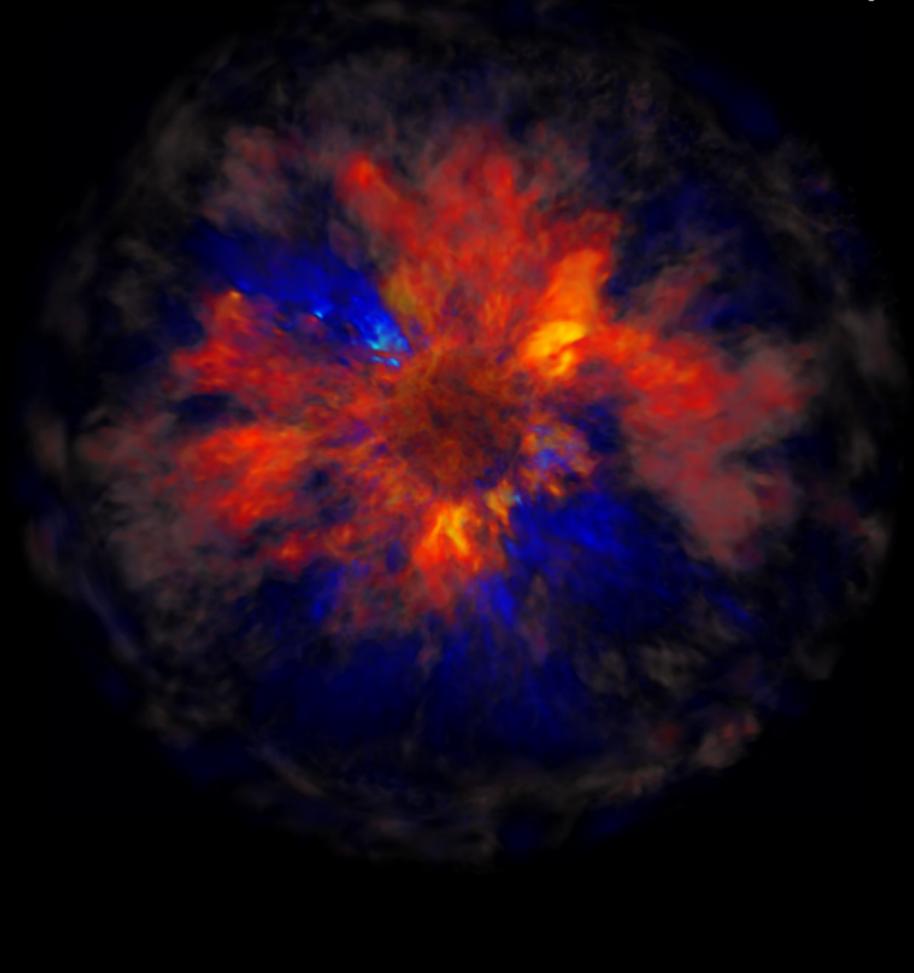
3D hydrodynamic simulations of H ingestion into He-shell flash convection on rapidly accreting white dwarfs



Radial velocity component



Radial velocity



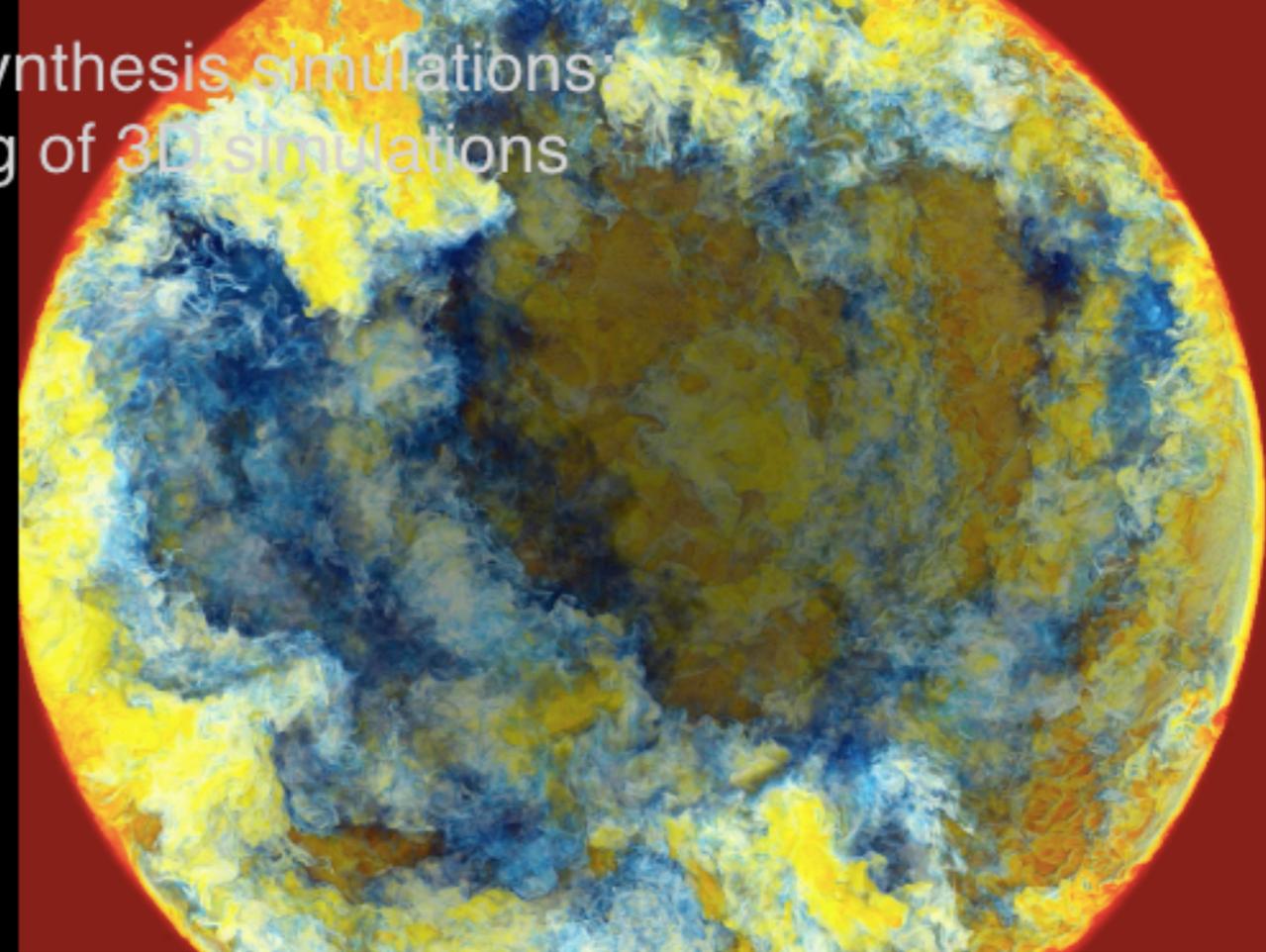
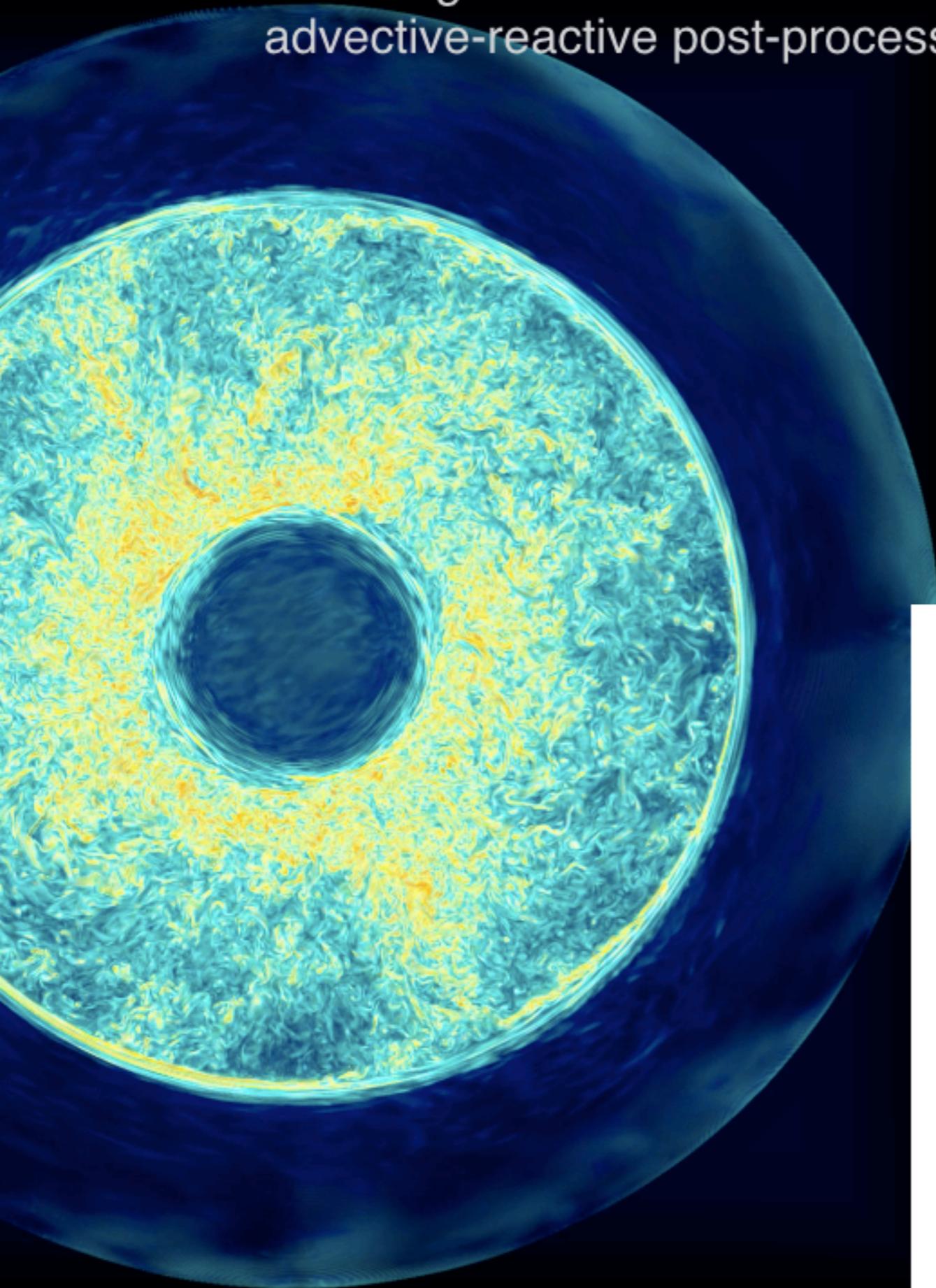
## 3D ID hydro-nucleosynthesis simulations of iRAWD

Left panels: 768<sup>3</sup>-grid simulation of H-ingestion in RAWD

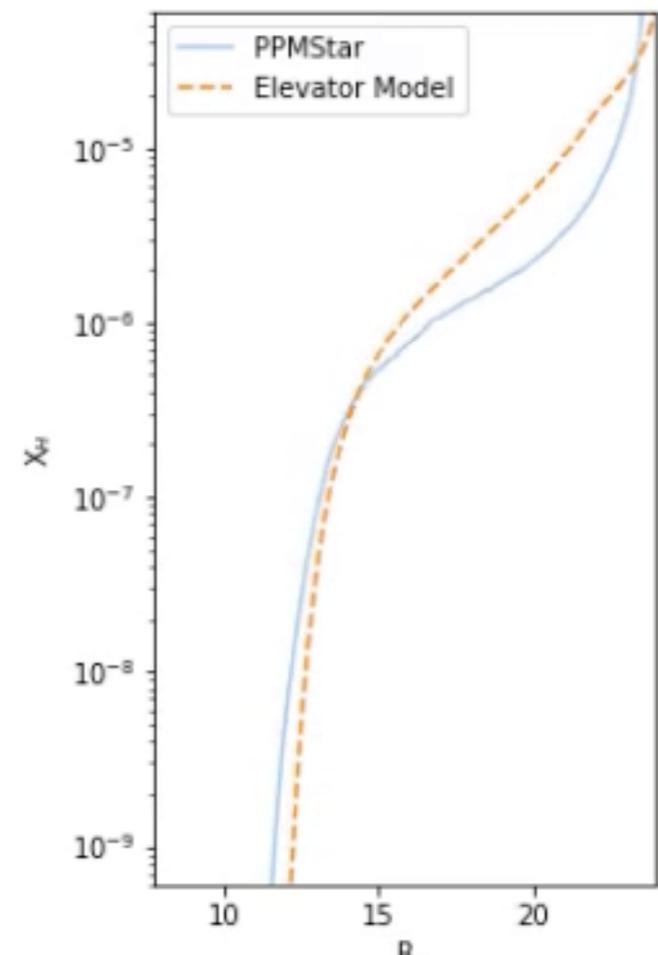
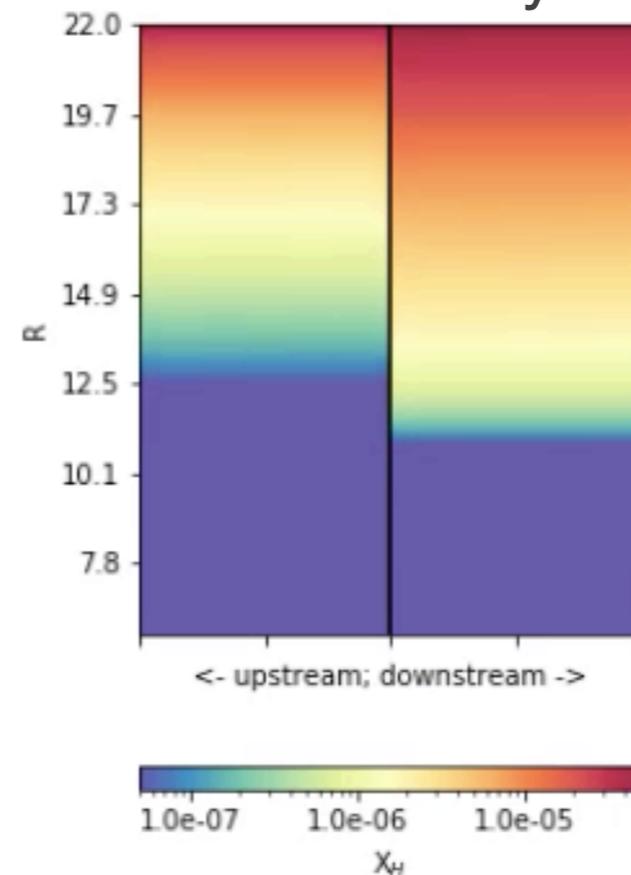
Right (top to bottom):

- ID diffusion coefficient determined from spherical averaged abundance profile from 3D simulations
- abundance profiles of H, <sup>13</sup>C and unstable i-process isotopes from multi-zone ID post-processing of evolving hydro stratification
- abundance distribution compared to CEMP-i star observation
- [time stretched in post-process to compensate for shorter hydro simulation]

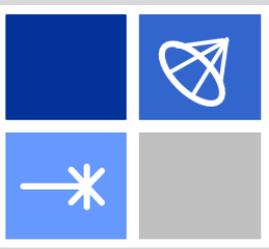
The next generation 3D1D nucleosynthesis simulations:  
advective-reactive post-processing of 3D simulations



simTime=25636.9s  
Advective conveyor belt 1D post-processing

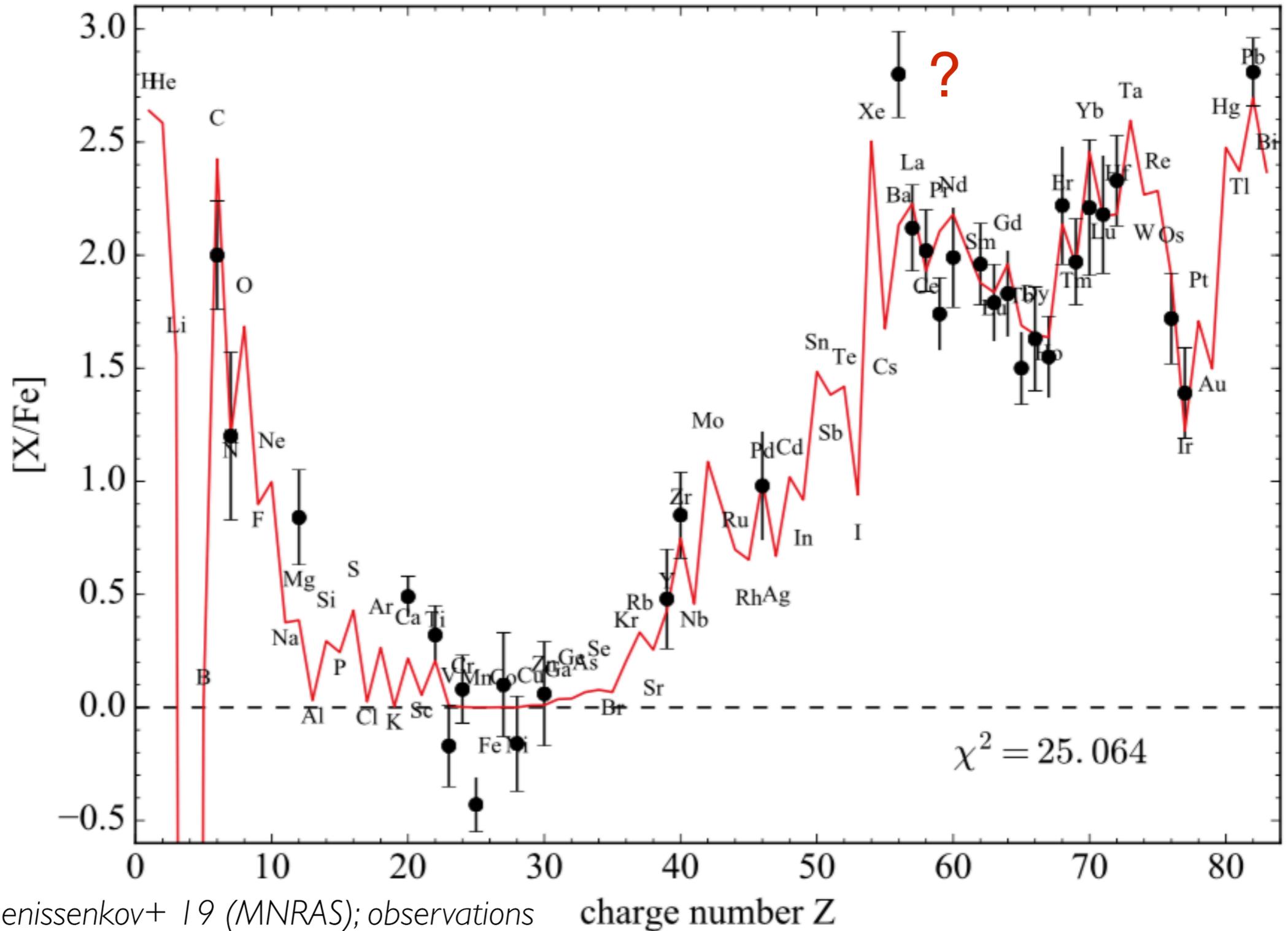


*Stephens+ in prep*



# i process in low-Z RAWD - origin of CEMP-i stars

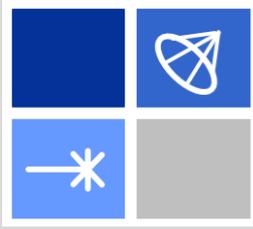
- Rapidly Accreting White Dwarf model G from DI9.
- These i-process simulations are essentially parameter free.
- These new models can for the first time explain these CEMP-r/s — or in fact CEMP-i — star abundances within a realistic astrophysical site.



Denissenkov+ 19 (MNRAS); observations  
Aoki+ 02 & Johnson & Bolte 04

charge number Z

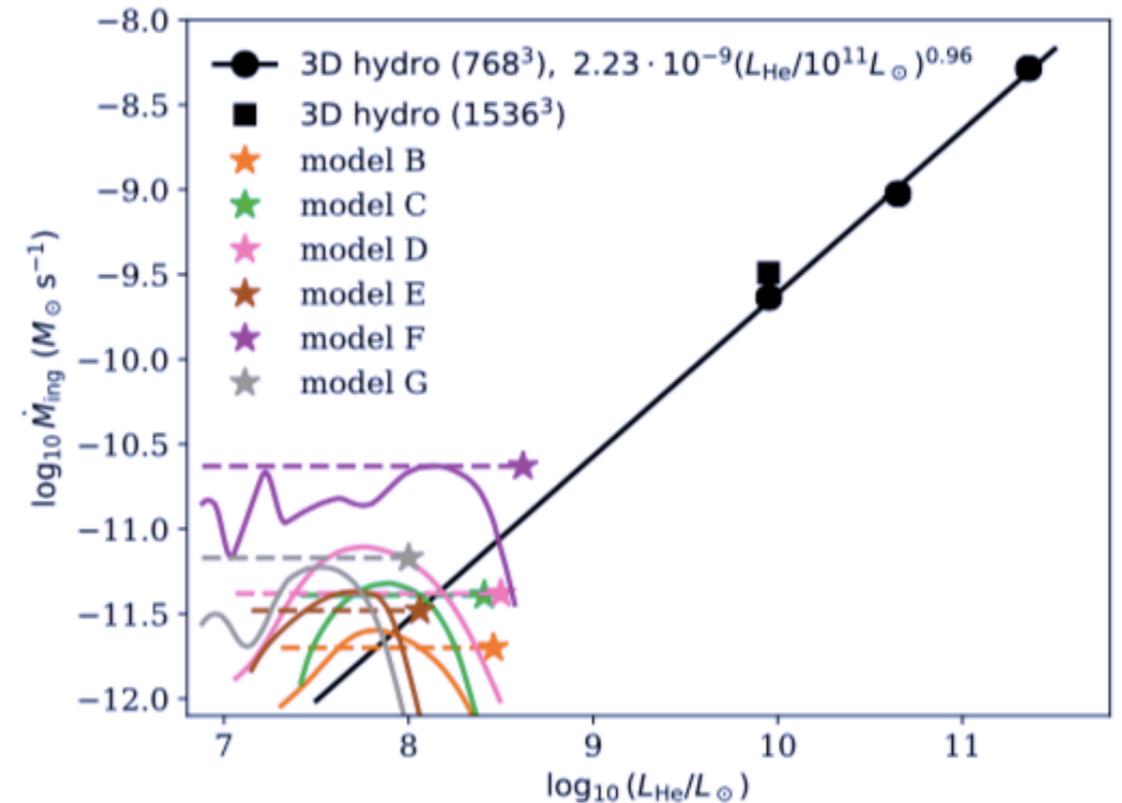
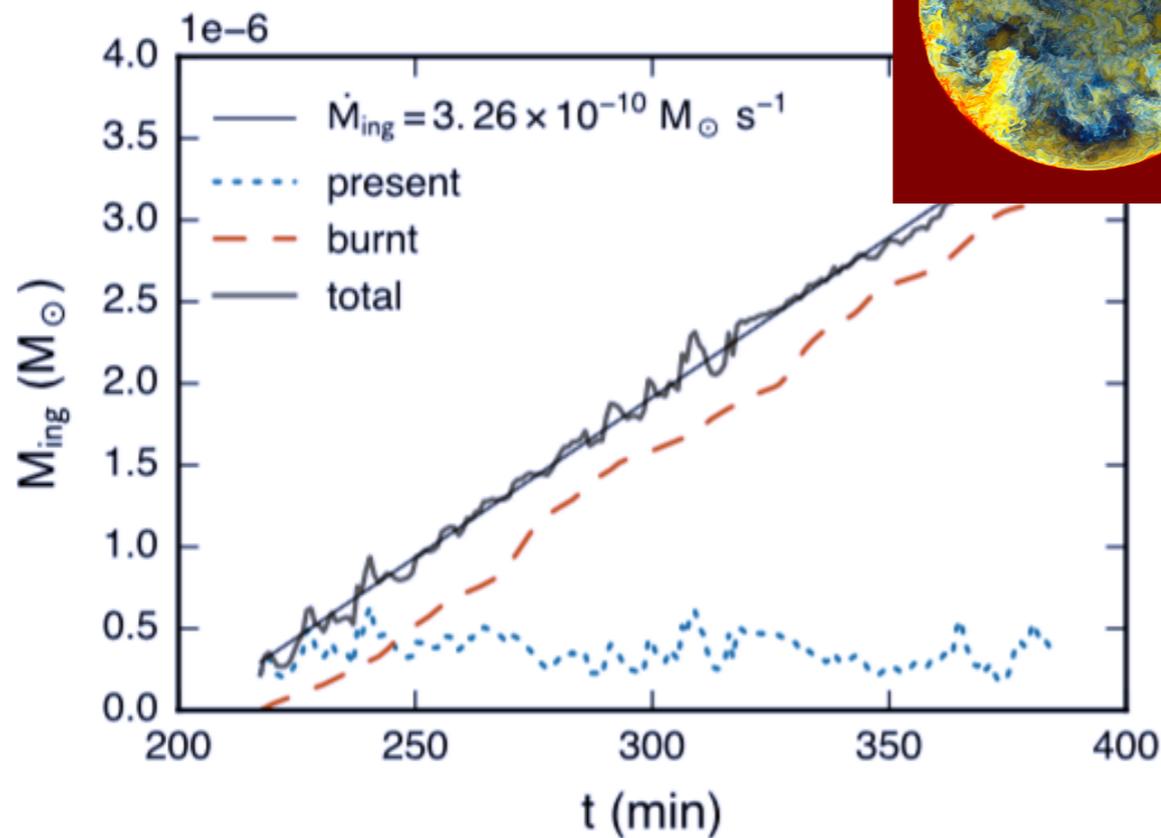
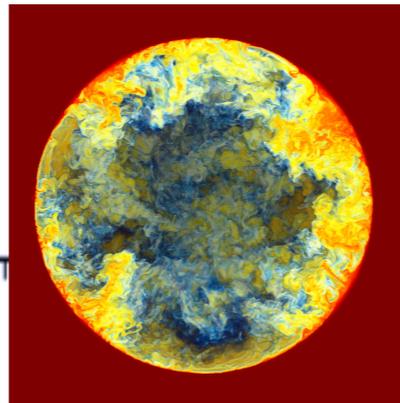
$\chi^2 = 25.064$



# The i-process yields of rapidly accreting white dwarfs from multicycle He-shell flash stellar evolution models with mixing parametrizations from 3D hydrodynamics simulations

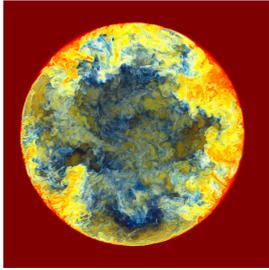
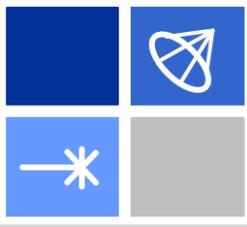
Pavel A. Denissenkov<sup>1,2</sup>\*, Falk Herwig<sup>1,2</sup>\*, Paul Woodward<sup>2,3</sup>, Robert Andrassy<sup>1,2</sup>, Marco Pignatari<sup>2,4,5</sup>† and Samuel Jones<sup>6</sup>†

Rate of H entrainment into He-shell convection zone from 3D hydro sims



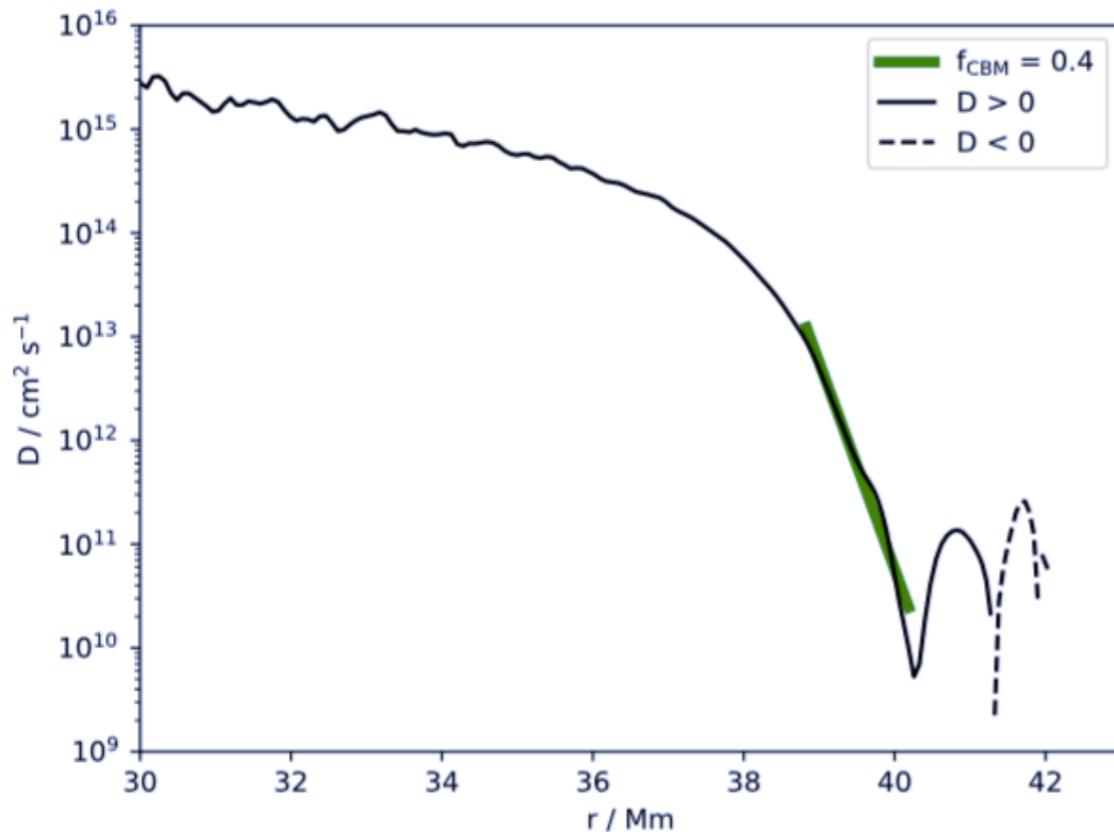
Rate of H entrainment as a function of driving luminosity.

Higher WD mass → higher  $L_{\text{He}}$  → higher entrainment rate → higher  $N_n$  → higher [Ba/La]



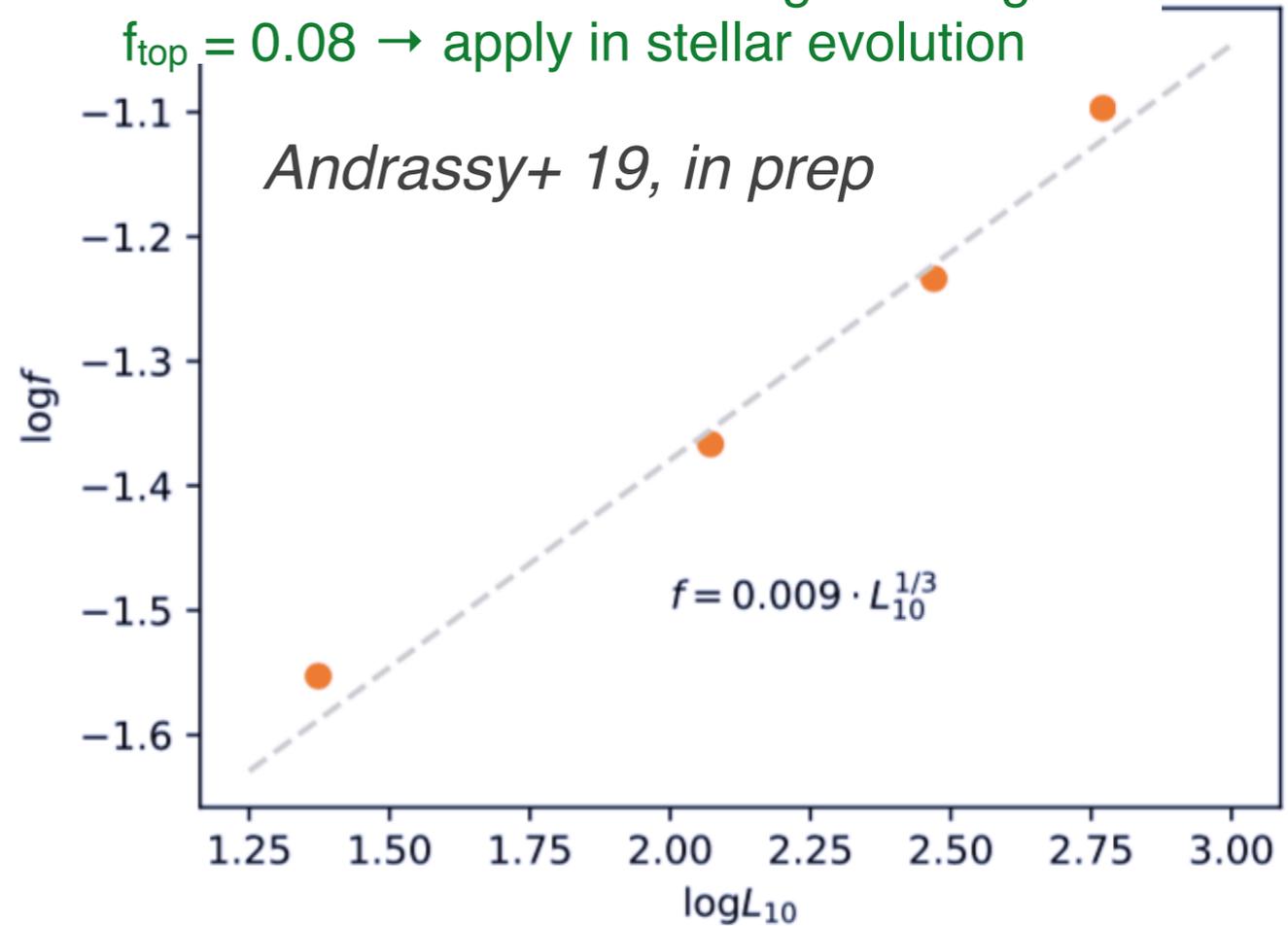
# The i-process yields of rapidly accreting white dwarfs from multicycle He-shell flash stellar evolution models with mixing parametrizations from 3D hydrodynamics simulations

Determination of 1D convective diffusion coefficient from 3D hydro sims (cf. Jones+ 17 for 3D hydro simulations in O-shell convection in a massive star) and fitting of convective boundary mixing parameter  $f$ .

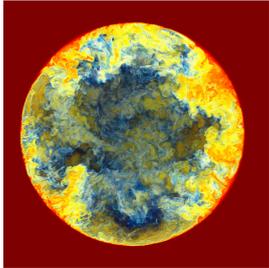
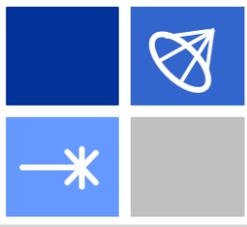


**Figure 6.** Determination of the  $f_{\text{top}}$  mixing parameter from the high-grid-resolution 3D hydrodynamics simulation E10.

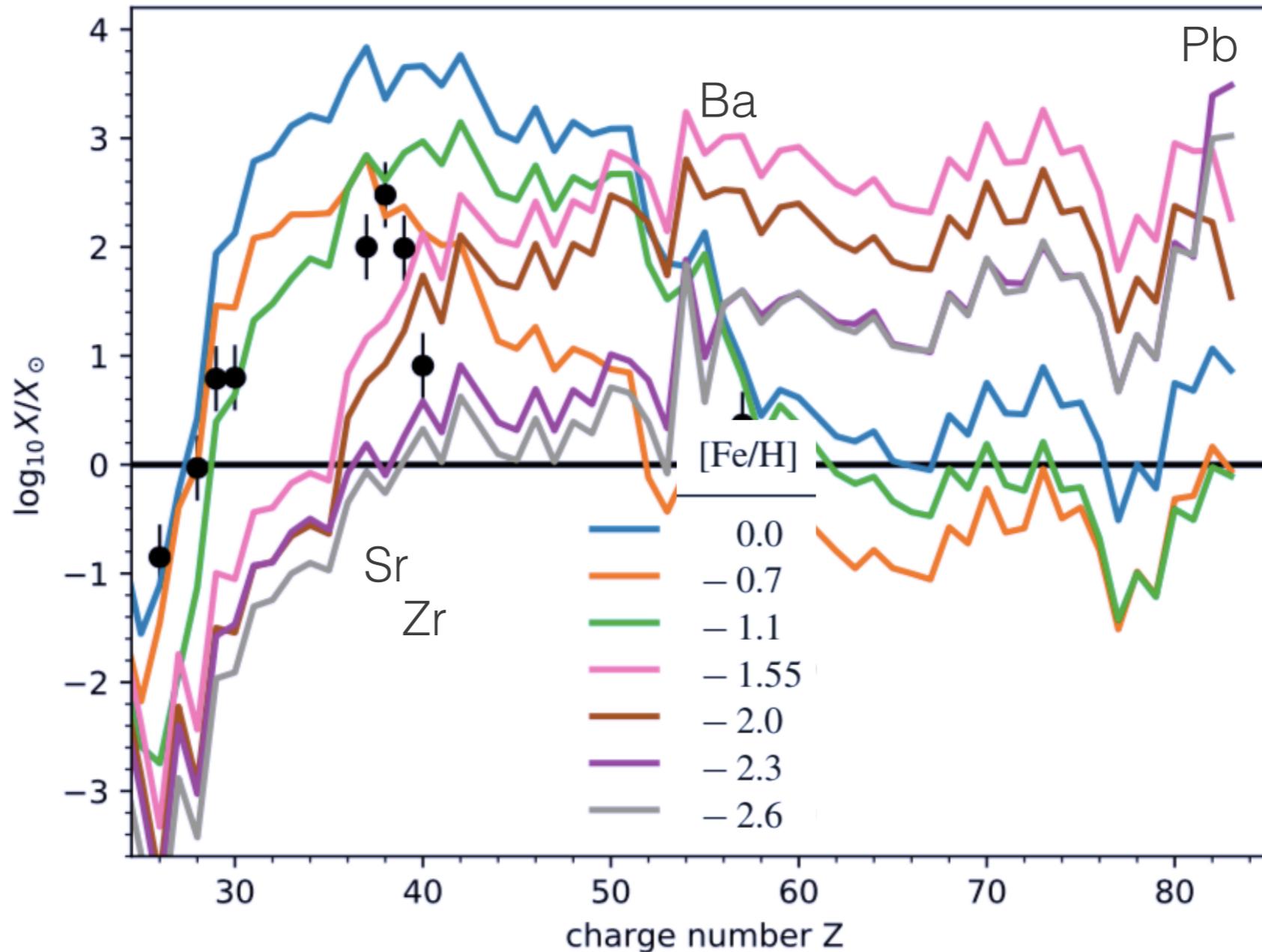
Extrapolate RAWD hydro  $f$  to nominal heating of stellar evolution model using  $f$ -scaling relation  $\rightarrow f_{\text{top}} = 0.08 \rightarrow$  apply in stellar evolution

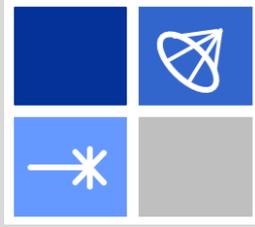


**Figure 7.** Scaling relation of the convective boundary mixing parameter  $f_{\text{top}}$  versus the driving luminosity based on a series of 1536-grid simulations of O-shell convection with the same setup as in Jones et al. (2017).



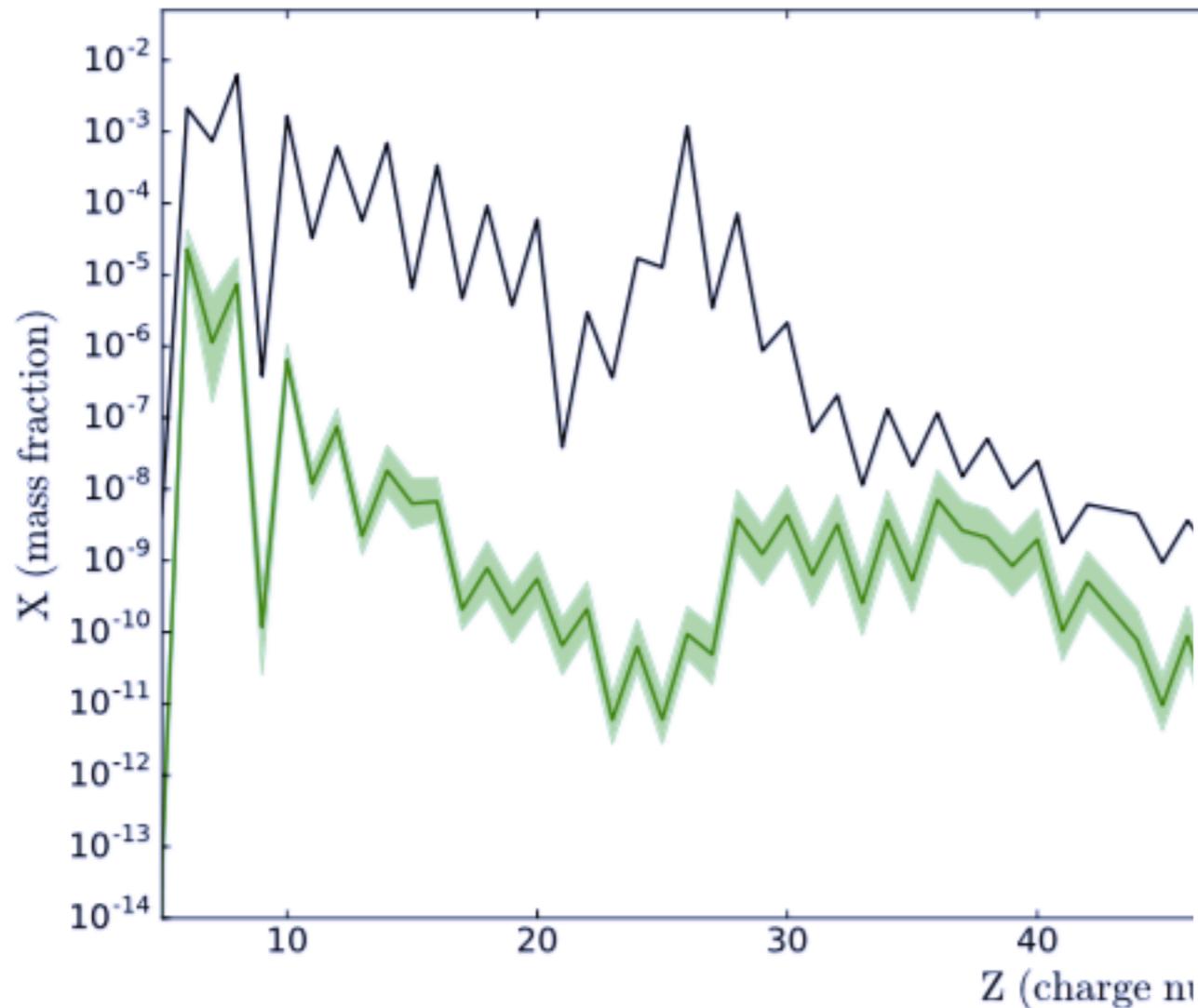
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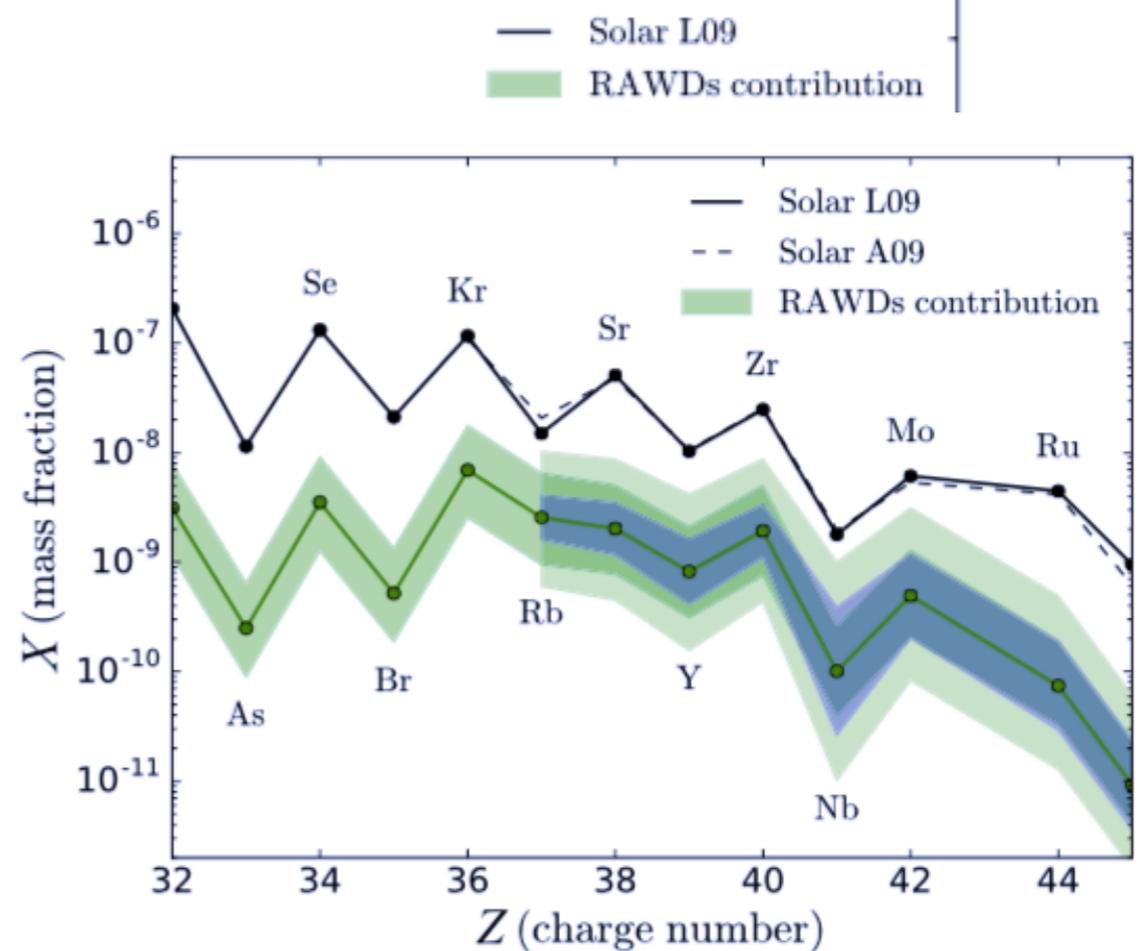


# *i*-process Contribution of Rapidly Accreting White Dwarfs to the Solar Composition of First-peak Neutron-capture Elements

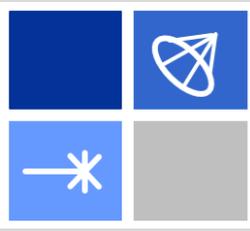
Benoit Côté<sup>1,2,3,10</sup> , Pavel Denissenkov<sup>1,3,10</sup> , Falk Herwig<sup>1,3,10</sup> , Ashley J. Ruiter<sup>4,5,6</sup> , Christian Ritter<sup>1,3,7,10</sup>,  
Marco Pignatari<sup>3,8,10</sup> , and Krzysztof Belczynski<sup>9</sup>



**Figure 8.** Predicted contribution of rapidly accreting white dwarfs (green, RAWDs) to the fiducial model while the green shaded area shows the range of solutions generated by un-



**Figure 9.** Same as in Figure 8, but zoomed on first-peak elements. The dashed black line shows the solar composition of Asplund et al. (2009, A09). The blue shaded area shows the uncertainties generated by nuclear reaction rates (see Section 6.4). The larger lighter-green shaded area shows the combined uncertainties generated by different chemical evolution paths, different ejecta masses for each RAWD, and by nuclear reaction rate uncertainties.



# Nuclear-physics impact studies for i process

IOP Publishing

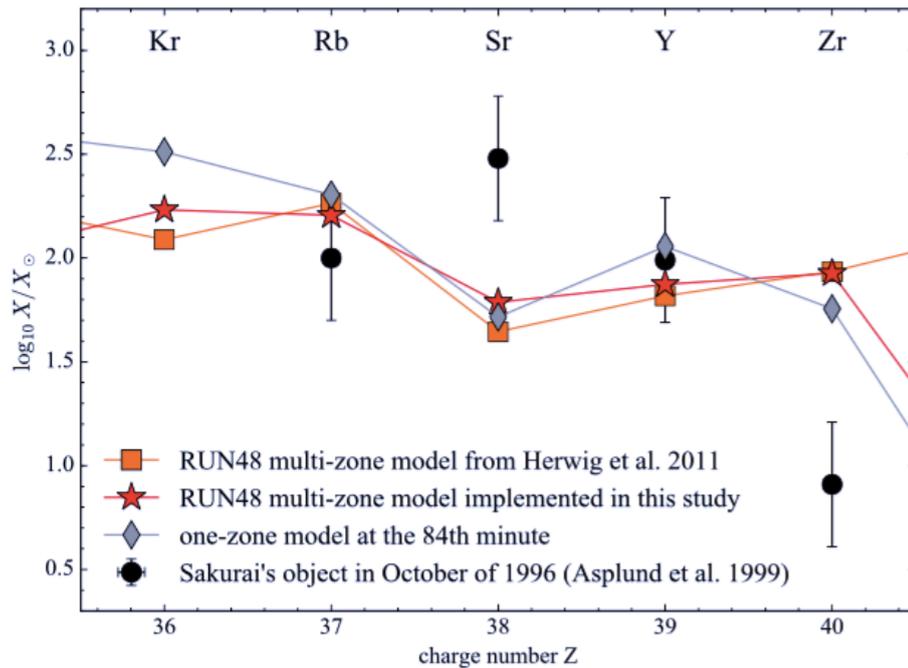
Journal of Physics G: Nuclear and Particle Physics

J. Phys. G: Nucl. Part. Phys. 45 (2018) 055203 (23pp)

<https://doi.org/10.1088/1361-6471/aabb6e>

## The impact of $(n, \gamma)$ reaction rate uncertainties of unstable isotopes near $N = 50$ on the i-process nucleosynthesis in He-shell flash white dwarfs

Pavel Denissenkov<sup>1,2,10</sup>, Georgios Perdikakis<sup>2,3,4</sup>,  
 Falk Herwig<sup>1,2,10</sup>, Hendrik Schatz<sup>2,4,5,10</sup>,  
 Christian Ritter<sup>1,2,10</sup>, Marco Pignatari<sup>6,7,10</sup>,  
 Samuel Jones<sup>8,9,10</sup>, Stylianos Nikas<sup>2,3</sup> and  
 Artemis Spyrou<sup>2,4,5</sup>



Step 1: observed abundance pattern and a nuclear network simulation (could be 1D multi-zone)

MNRAS submitted, [arXiv:1909.07011](https://arxiv.org/abs/1909.07011)

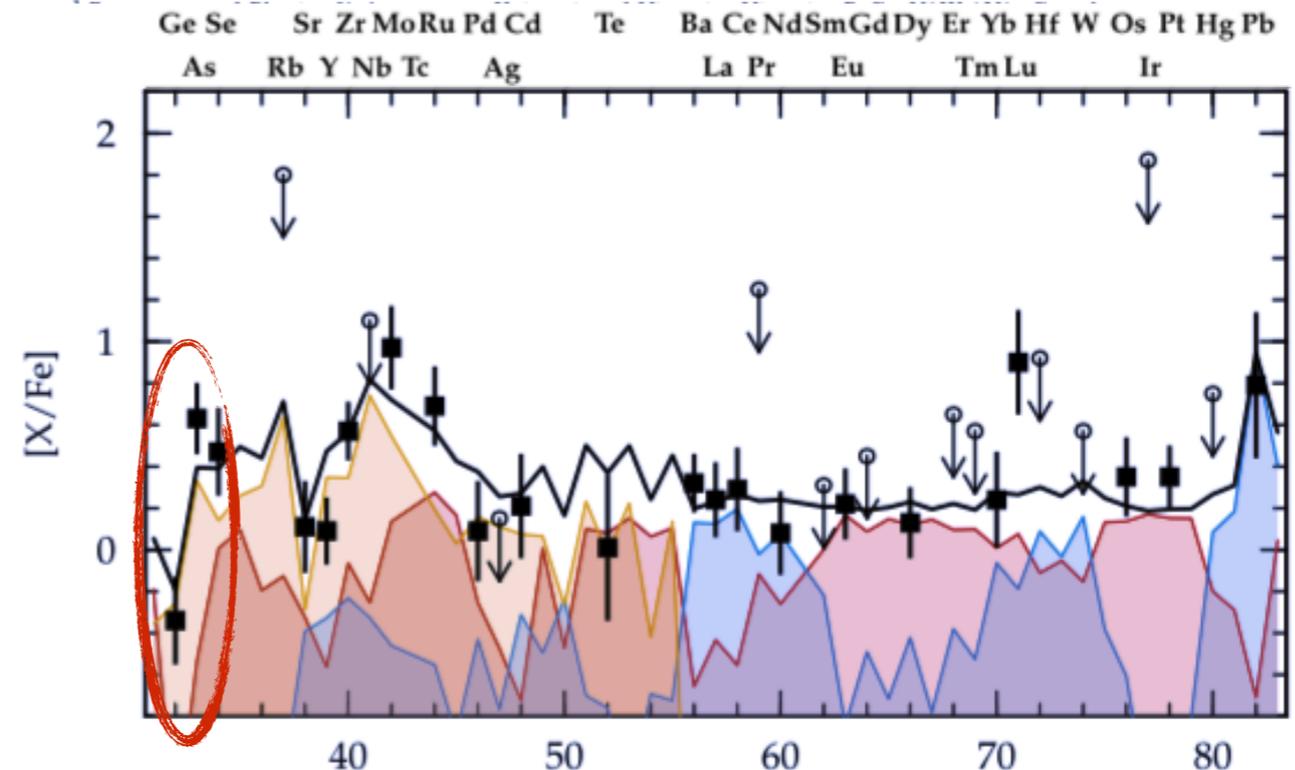
MNRAS 000, 1–10 (2019)

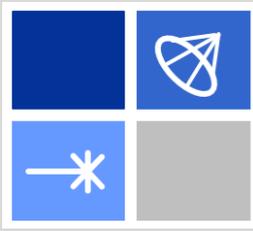
Preprint 13 September 2019

Compiled using MNRAS L<sup>A</sup>T<sub>E</sub>X style file

The impact of  $(n, \gamma)$  reaction rate uncertainties on the predicted abundances of i-process elements with  $32 \leq Z \leq 48$  in the metal-poor star HD94028

John E. McKay<sup>1,2,†</sup>, Pavel A. Denissenkov<sup>1,3,†\*</sup>, Falk Herwig<sup>1,3,†</sup>, Georgios Perdikakis<sup>3,4,5</sup> and Hendrik Schatz<sup>3,5,6,†</sup>





# Step 2: Determine nuclear physics uncertainties

## Step 3: Perform Monte-Carlo simulations

- Determine variation factors from all relevant unstable (n,γ) cross section
- 20 combinations of HF model assumptions listed below
- 10,000 Monte-Carlo runs with randomly drawn factors for each of the included (n,γ) cross sections

**Table 1.** List of models used to describe the nuclear level density and the  $\gamma$  strength in the Hauser–Feshbach calculations. Calculations were performed with all 20 possible combinations of these models.

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### Nuclear level density (NLD) models used in this work

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Constant temperature matched to the Fermi Gas (CT+BSFG) (Dilg *et al* 1973)  
 Back-shifted Fermi Gas (BSFG) (Dilg *et al* 1973, Gilbert and Cameron 1965)  
 Generalized super fluid (GSM) (Ignatyuk *et al* 1979, 1993)  
 Hartree–Fock using Skyrme force (HFS) (Goriely *et al* 2001)  
 Hartree–Fock–Bogoliubov + combinatorial (HFBS-C) (Goriely *et al* 2008)

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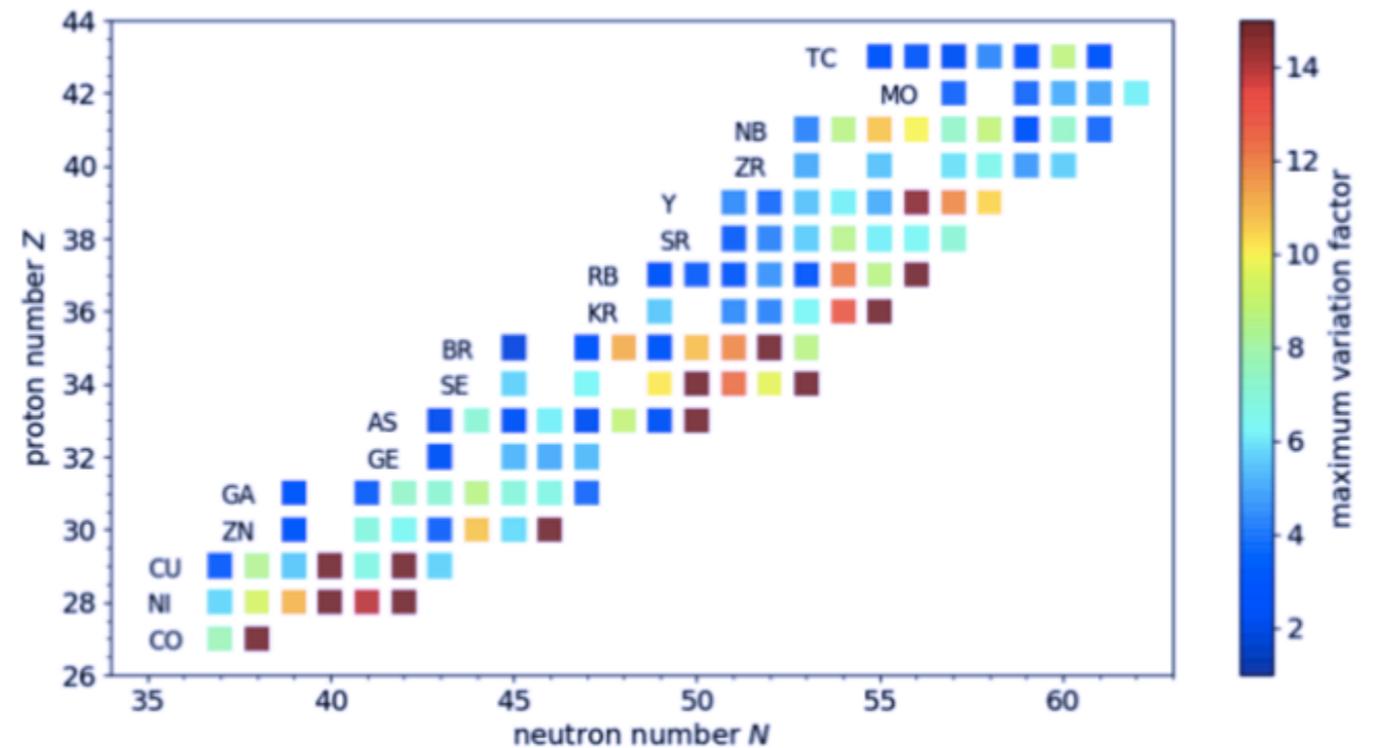
### $\gamma$ ray strength function ( $\gamma$ SF) models used in this work

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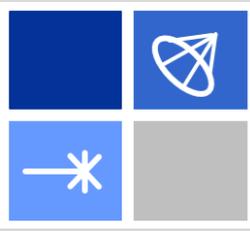
Kopecky–Uhl generalized Lorentzian (KU) (Kopecky and Uhl 1990)  
 Hartree–Fock BCS + QRPA (HF-BCS+QRPA) (Goriely and Khan 2002)  
 Hartree–Fock–Bogolyubov + QRPA (HFB+QRPA) (Goriely *et al* 2004)  
 Modified Lorentzian (Gor-ML) (Goriely 1998)

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## Case weak i process

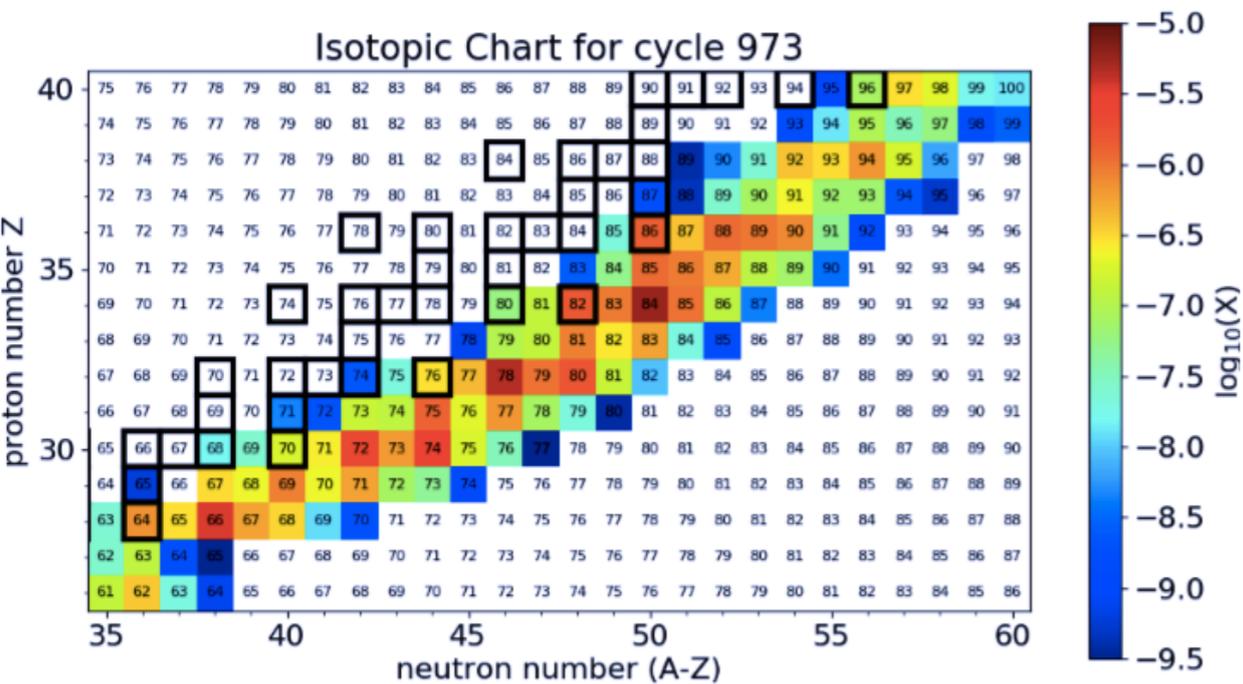


**Figure 1.** The unstable isotopes whose (n,γ) reaction rates were varied in this study and the maximum variation factors used for the rates. The maximum variation factors for  $^{65}\text{Co}$ ,  $^{68}\text{Ni}$ ,  $^{70}\text{Ni}$ ,  $^{69}\text{Cu}$ ,  $^{71}\text{Cu}$ ,  $^{76}\text{Zn}$ ,  $^{83}\text{As}$ ,  $^{84}\text{Se}$ ,  $^{87}\text{Se}$ ,  $^{87}\text{Br}$ ,  $^{91}\text{Kr}$ , and  $^{93}\text{Rb}$  exceed the maximum value of  $v_i^{\text{max}} = 15$  assigned to the color map.



# Step 4: Analyse results (weak i process)

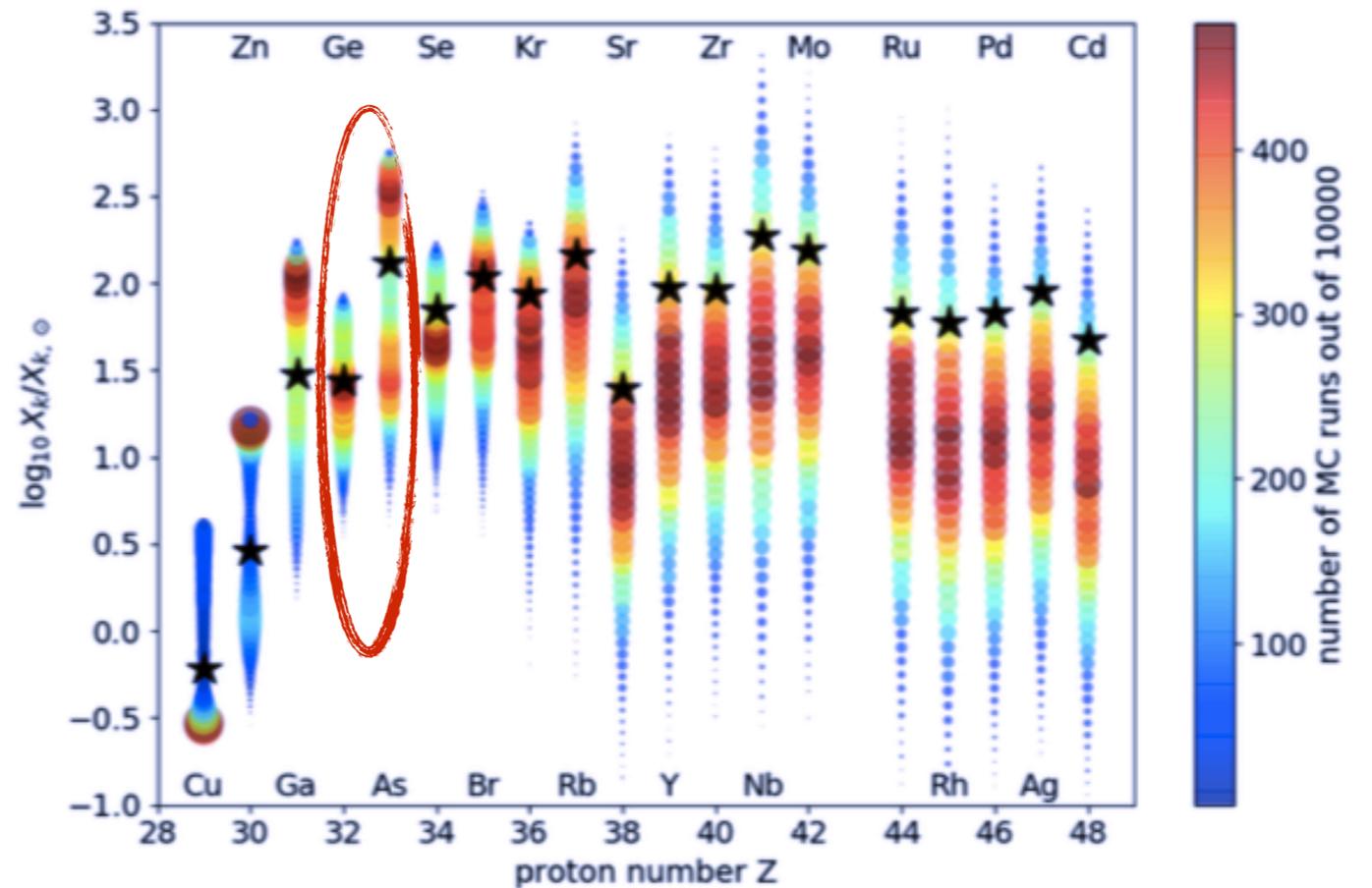
## Network flux in abundance chart



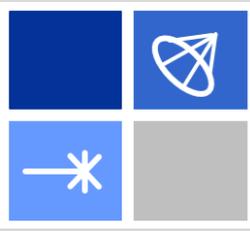
**Figure 5.** Undecayed isotopic abundances (mass fractions  $X$ ) for the best-fit cycle 973 near the i-process nucleosynthesis peak.

- In some cases, like here for As, the MC simulation yields a bimodal distribution corresponding to a specific branch point, in this case  $^{75}\text{Ga}$

## Benchmark case and MC results

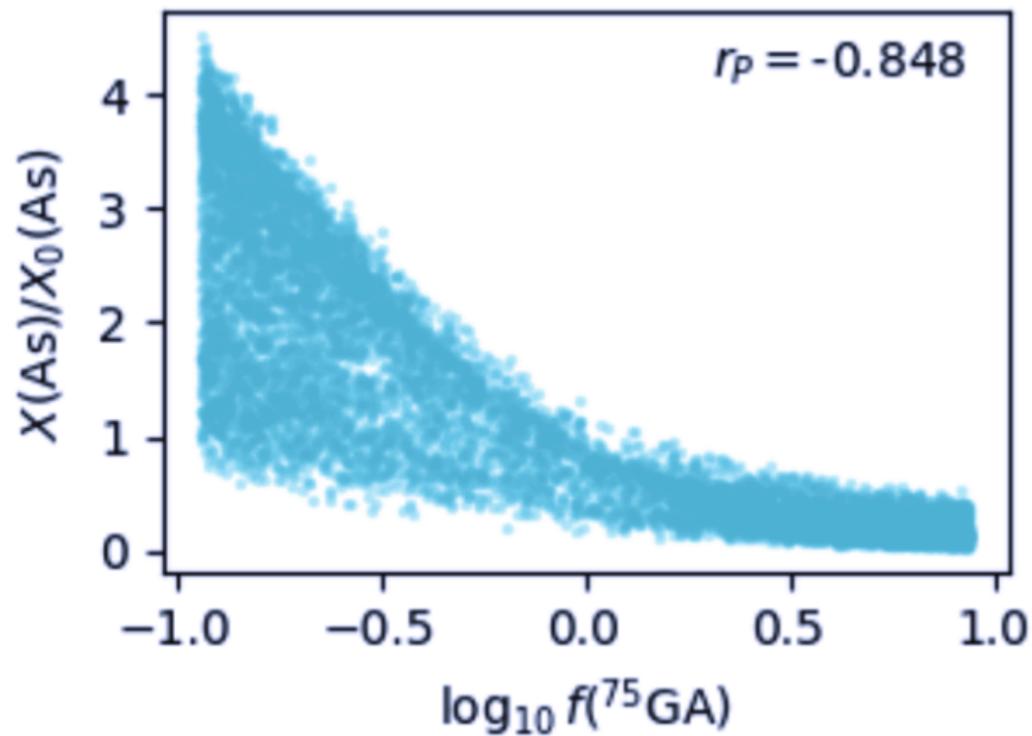


**Figure 6.** The distributions of the i-process elemental abundances generated in the Monte Carlo simulation by randomly varying  $(n, \gamma)$  reaction rates of the unstable isotopes that are displayed in Figure 1. The larger and redder circles correspond to a larger number of MC runs contributing to a given abundance. The black star symbols show the benchmark simulation abundances.



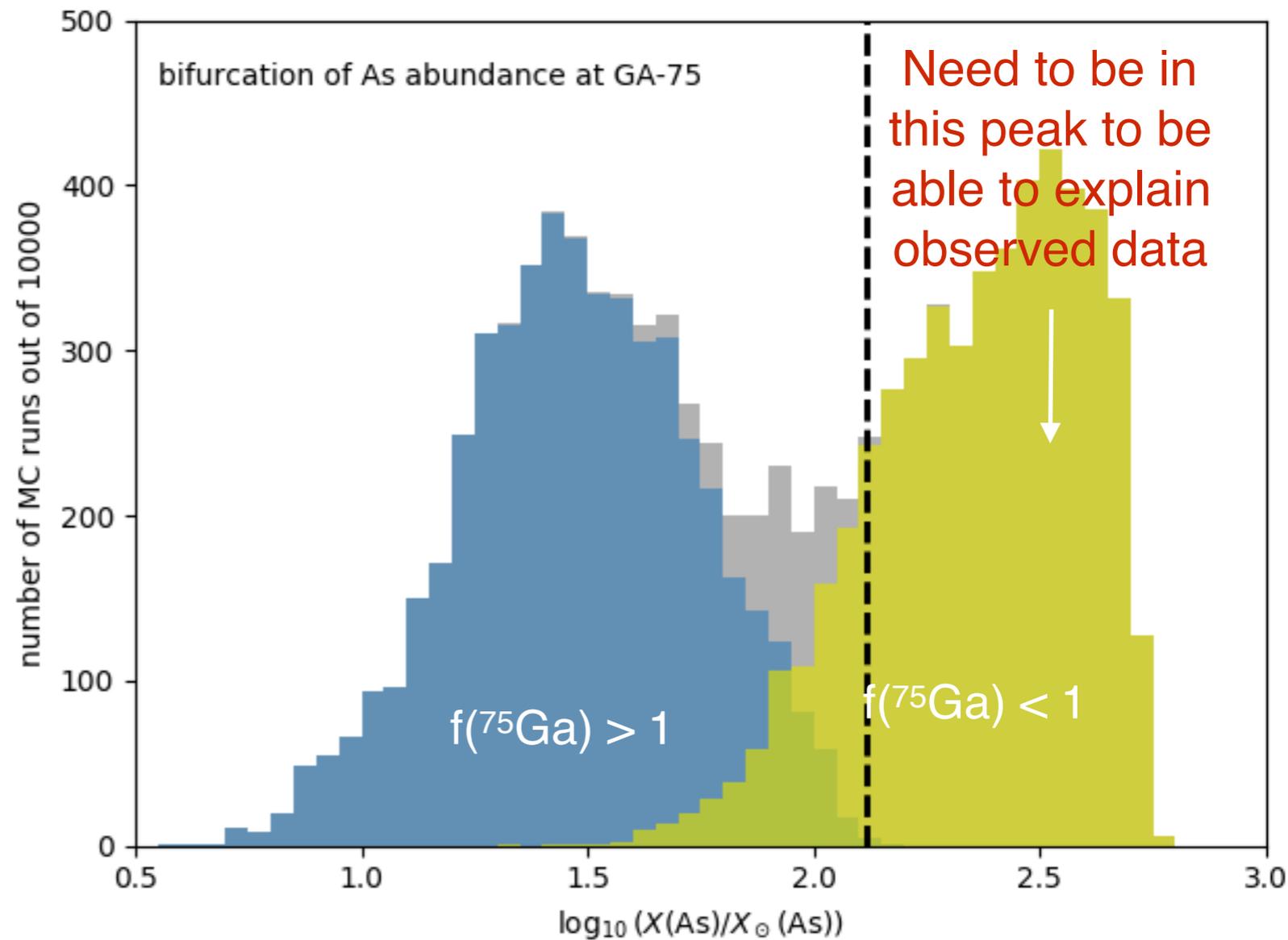
# Key importance of $^{75}\text{Ga}$ for As prediction

Correlation As abundance with  $^{75}\text{Ga}$  cross section

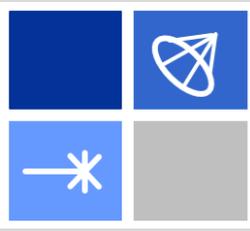


In some cases, like here for As, the MC simulation yields a bimodal distribution corresponding to a specific branch point, in this case  $^{75}\text{Ga}$

Benchmark case and MC results



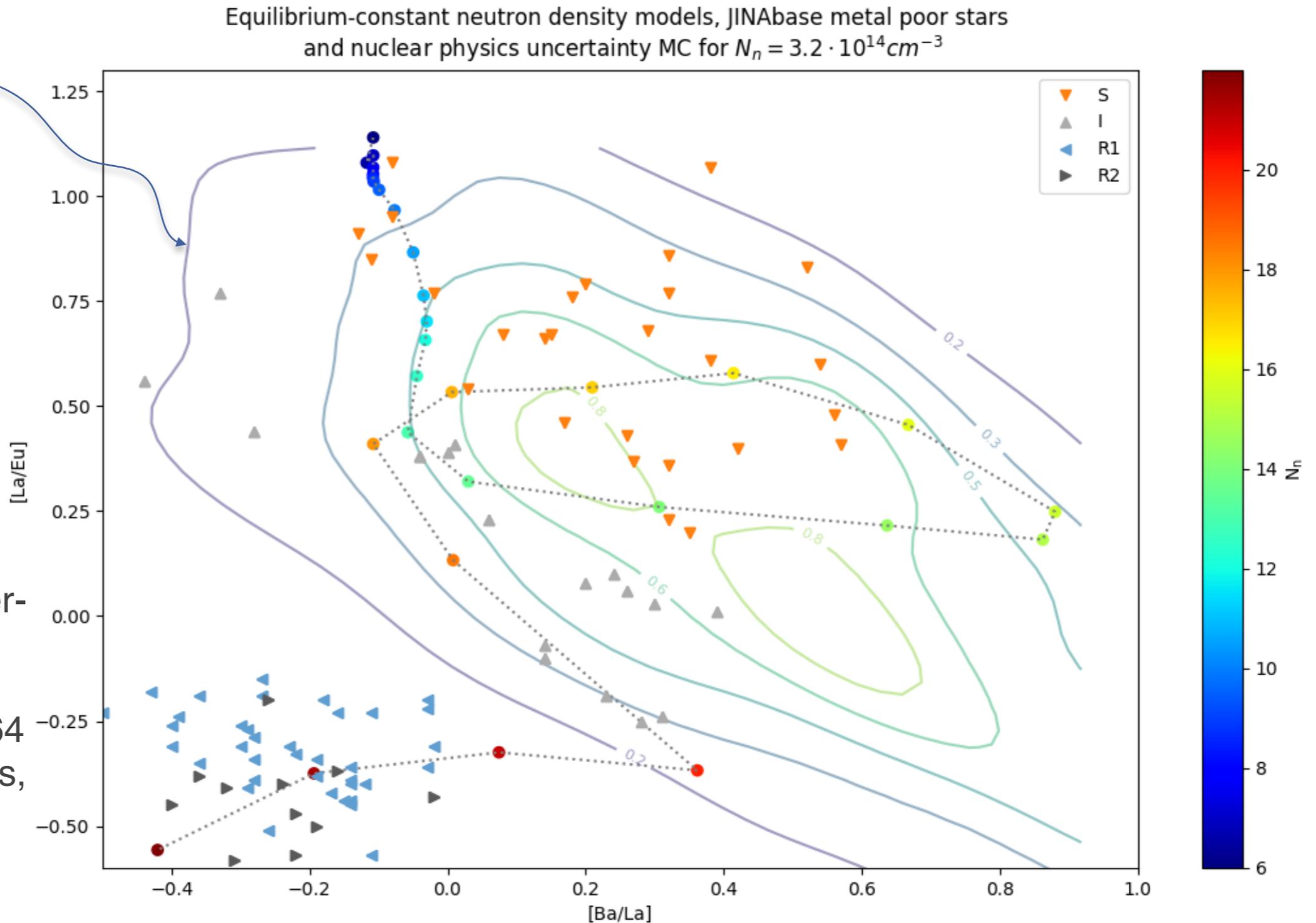
In order for i process to explain the observed  $[\text{As}/\text{Ge}]$  in HD94028 the  $^{75}\text{Ga}(n,\gamma)$  cross section must be in lower part of range.



# Second-peak nuclear physics impact study

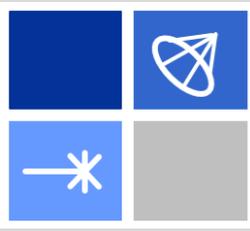
Contour lines:  
probability  
density nuclear  
physics  
uncertainties  
according to  
Monte-Carlo  
nuclear physics  
impact study.

10,000  
simulations  
involving Hauser-  
Feshbach  
uncertainty  
estimates for 164  
unstable species,  
for just one  
neutron density.

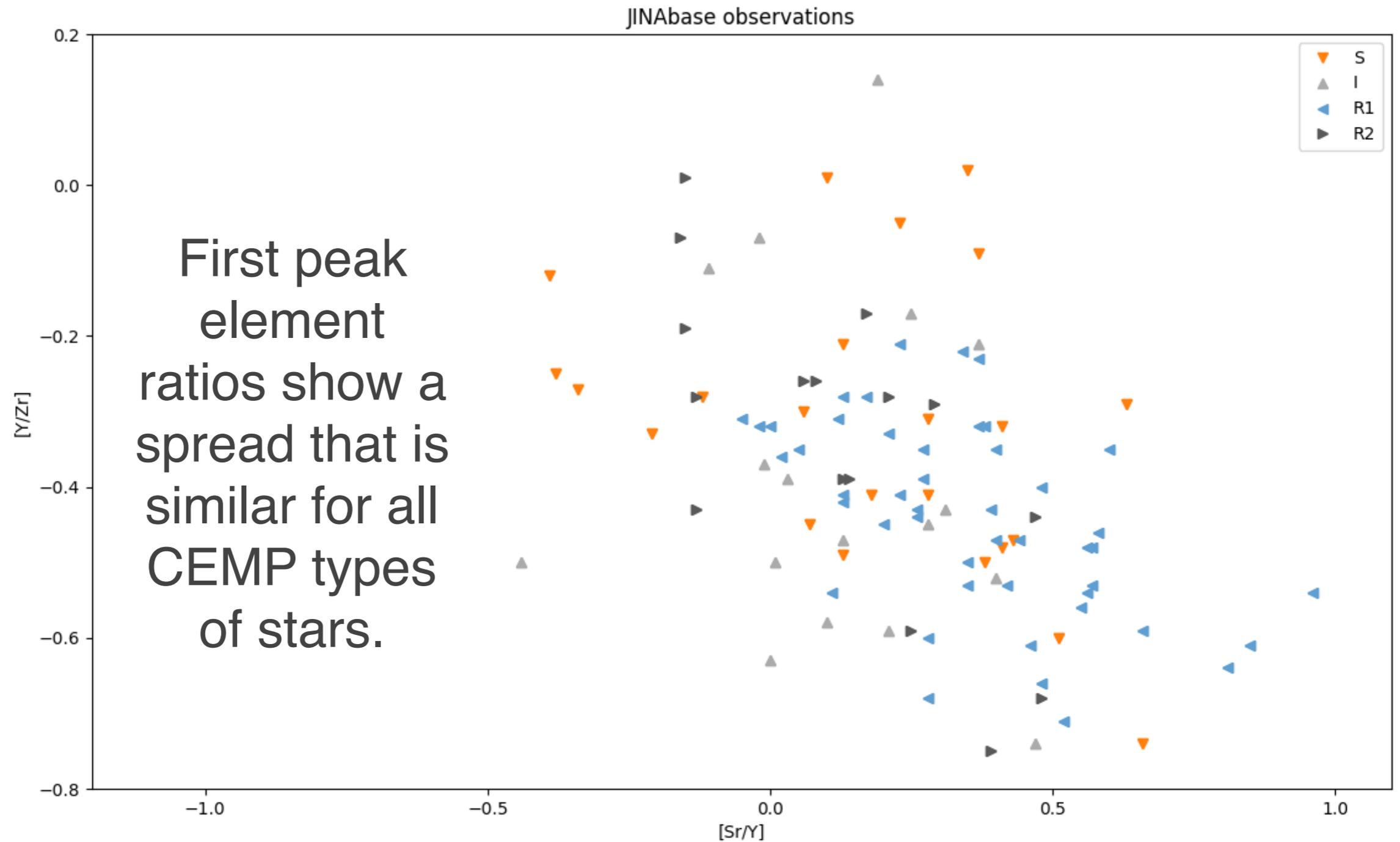


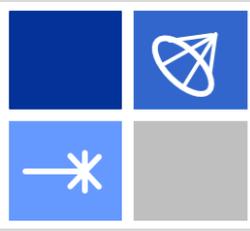
Denissenkov P, Perdikakis G, Herwig F, et al. 2018. *Journal of Physics G Nuclear Physics*. 45(5):055203



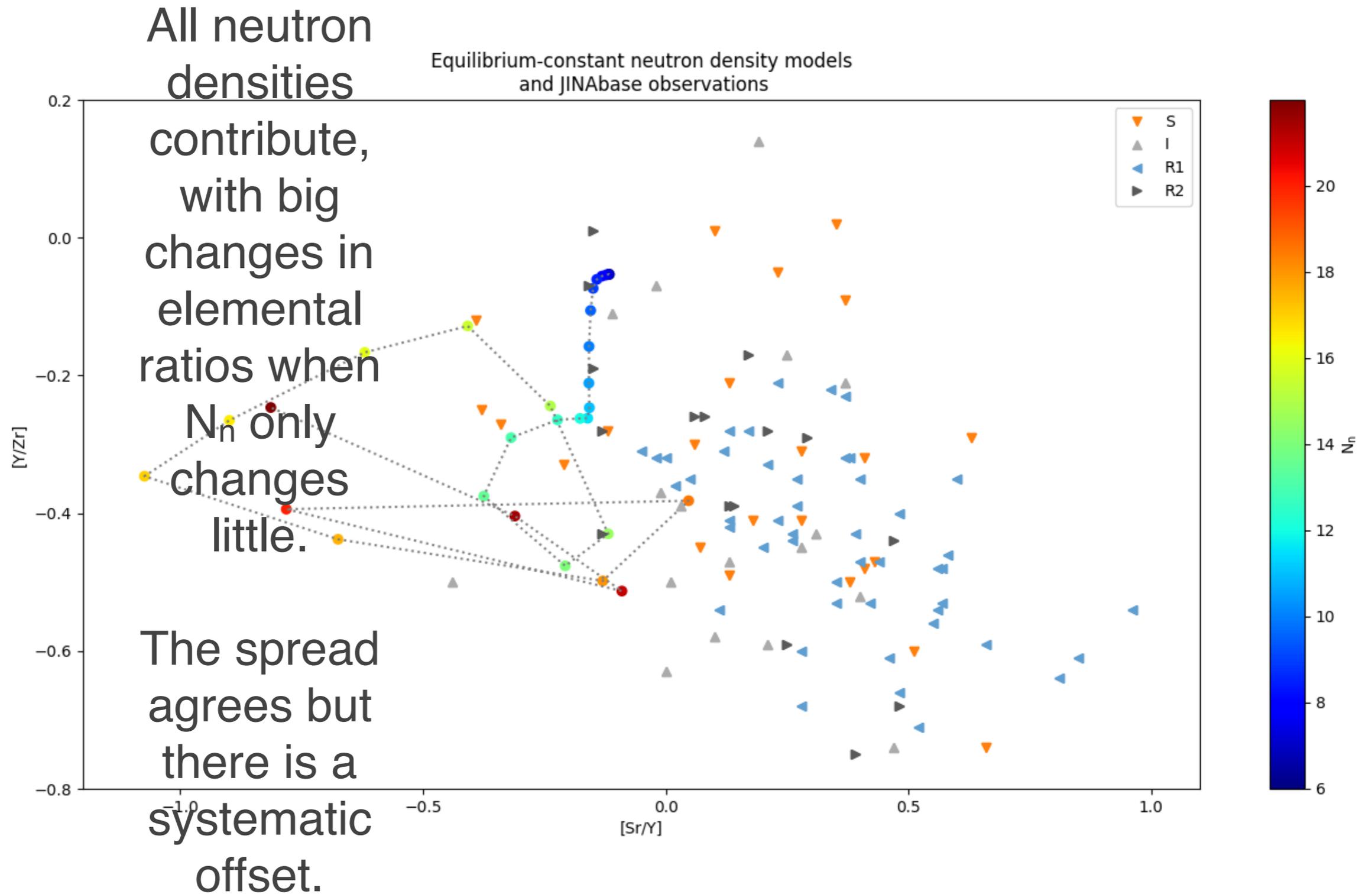


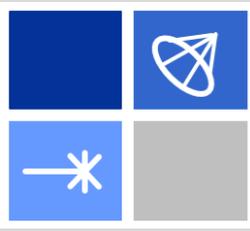
# First-peak elements metal-poor stars





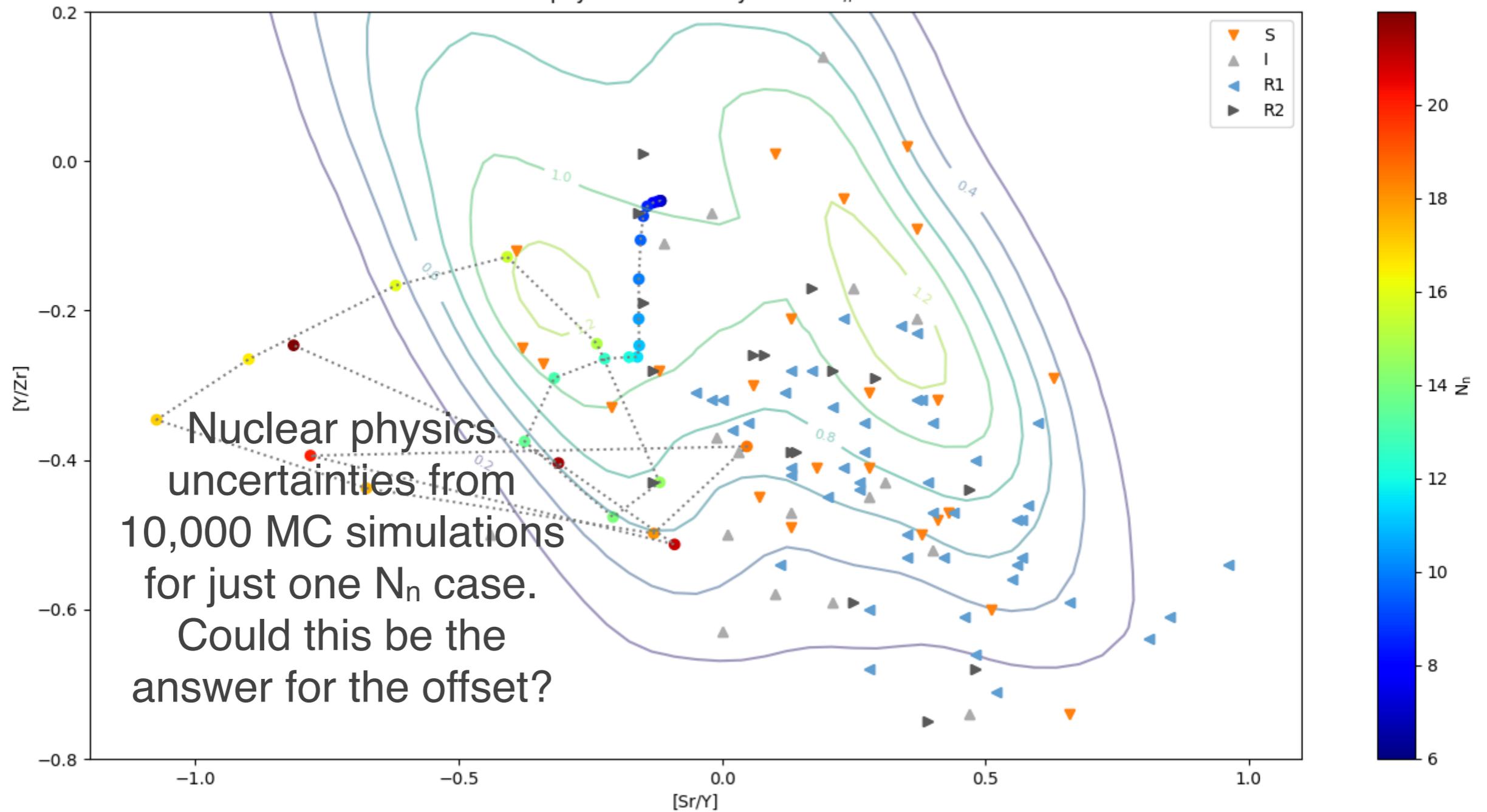
# First-peak & constant neutron density models



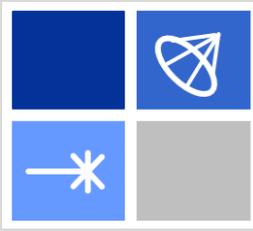


# Nuclear physics impact

Equilibrium-constant neutron density models, JINABase metal poor stars  
and nuclear physics uncertainty MC for  $N_n = 10^{15} \text{cm}^{-3}$



Nuclear physics  
uncertainties from  
10,000 MC simulations  
for just one  $N_n$  case.  
Could this be the  
answer for the offset?



# Nuclear physics impact studies for i process and convective-reactive regimes

First-peak i-process impact study.  $^{88}\text{Kr}(n, \gamma)$   
**Published.**  
Denissenkov+ 18 J. Phys. G

Weak i process proposed by  
Roederer+ 16. Produces  
anomalous [As/Ge]. **Submitted.**  
[arXiv:1909.07011](https://arxiv.org/abs/1909.07011)

Pop III i process (→Edinburgh  
student Sam Lloyd).  
**In progress.**

Odd-Z elements P, Cl, K, Sc in convective-reactive O-  
and C-shell mergers, site explored Ritter+18. **Planned.**

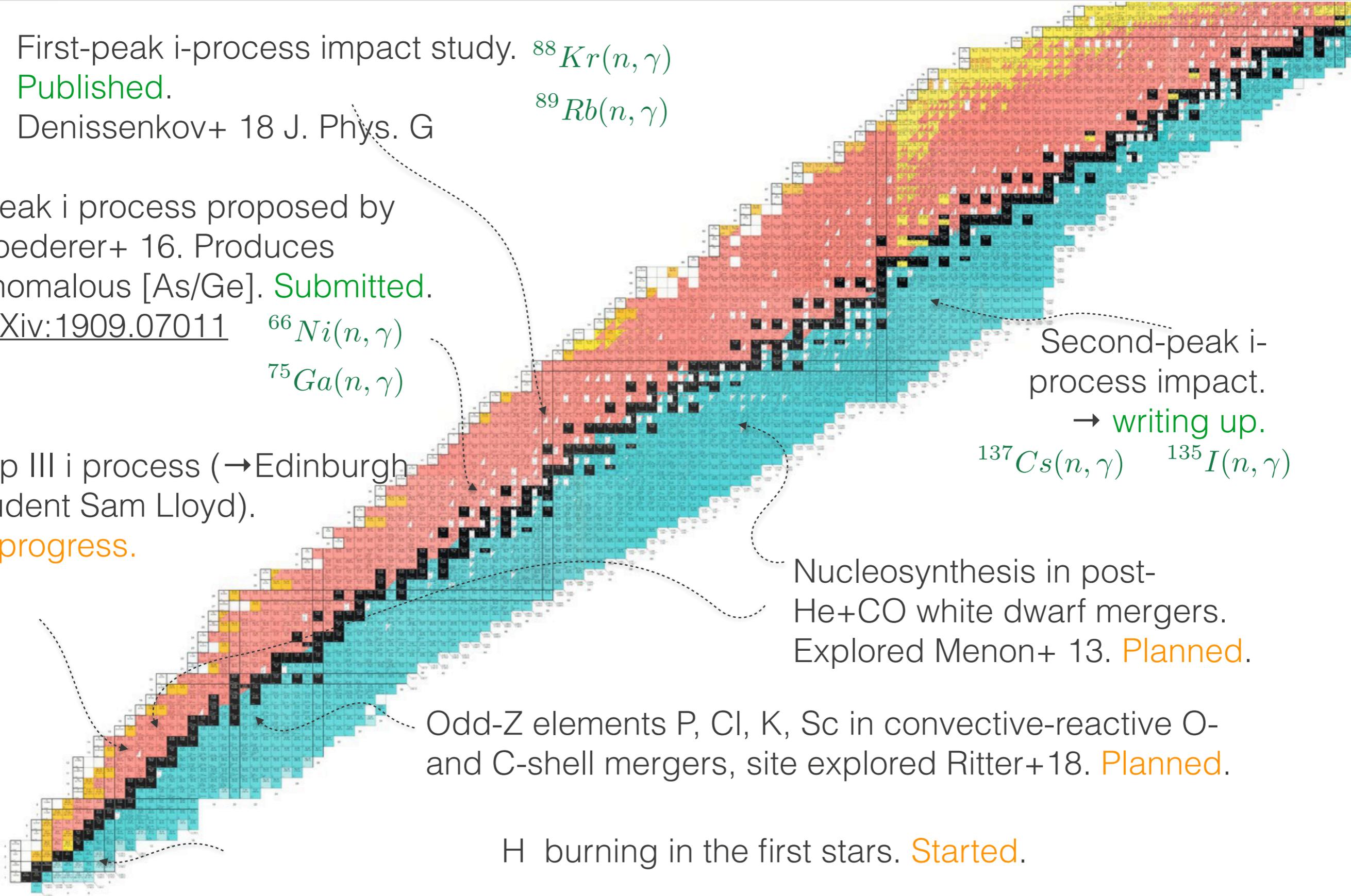
H burning in the first stars. **Started.**

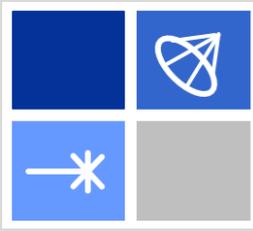


Second-peak i-  
process impact.  
→ **writing up.**



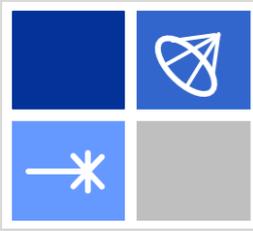
Nucleosynthesis in post-  
He+CO white dwarf mergers.  
Explored Menon+ 13. **Planned.**





# Summary

- CEMP stars show the abundance signature of a wide range of neutron densities, including s and r, but also from intermediate neutron capture conditions
- i process is a unique neutron-capture process induced by convective H mixing into He-shell convection
- There are many potential i-process sites in massive and low-mass stars. Rapidly accreting white dwarfs naturally produce the right conditions, for example, for CEMP-r/s (CEMP-i) stars
- **It is absolutely key to improve nuclear physics data for the i process, detailed impact studies are now available**



# Bibliography i process theory

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- **Malaney** RA. 1986. *MNRAS*. 223(4):683–707, **Herwig** F et al. 2011. *ApJ*. 727(2):89 *investigate i-process nucleosynthesis in post-AGB stars*
- **Campbell** SW, Lugaro M, Karakas AI. 2010. *A&A*. 522:L6, **Cruz** MA, Serenelli A, Weiss A. 2013. *A&A*. 559:A4 *present models of i process in H-ingestion into He-core flash*
- **Bertolli** MG, Herwig F, Pignatari M, Kawano T. 2013. *Phys Rev C submitted, arXiv:1310.4578*, **Dardelet** L, Ritter C, Prado P, Heringer E, Higgs C, et al. 2014. *I process and CEMP-s+r stars*, **Hampel** M, Stancliffe RJ, Lugaro M, Meyer BS. 2016. *ApJ*. 831(2):171 *first make the connection of i-process with CEMP-r/s stars → CEMP-i*
- **Côté** B, Denissenkov P, Herwig F et al. 2018. *ApJ*. 854(2):105 *GCE simulations to demonstrate iRAWd contribution to solar system*
- **Clarkson** O, Herwig F, Pignatari M. 2018. *MNRAS*. 474(1):L37–L41; Clarkson+ 19 *propose Pop III i process for CEMP-no stars*; **Banerjee** P, et al 2018. *ApJ*. 865(2):120 *considers metal-poor stars in more general*
- *Large body of literature on H-ingestion flashes (HIF or PIE) in AGB stars (Iben, Schönberner, Fujimoto, Sudo, Iwamoto, Campbell, Schlattl ....*
- i-process impact study: **Denissenkov** P, Perdikakis G, Herwig F, et al. 2018. (First peak elements) *Journal of Physics G Nuclear Physics*. 45(5):055203; **McKay** J, et al. 2019. (Weak i-process). *MNRAS*, submitted, astro-ph; **Denissenkov** et al. in prep. (Second-peak elements); **Lloyd** S, et al. in prep (Pop III light i process)