

PRISMS PhaseField

Faceted Anisotropy (with Coupled CH-AC Dynamics)

This application is essentially a specialization of the CHAC_anisotropyRegularized application with a particular choice of interfacial energy anisotropy $\gamma(\mathbf{n})$. In this document, we repeat the formulation for that model for completeness and then describe the anisotropy used in this application. Consider a free energy expression of the form:

$$\Pi(c, \eta, \nabla\eta) = \int_{\Omega} (f_{\alpha}(1-H) + f_{\beta}H) + \frac{1}{2}|\gamma(\mathbf{n})\nabla\eta|^2 + \frac{\delta^2}{2}(\Delta\eta)^2 dV \quad (1)$$

where f_{α} and f_{β} are the free energy densities corresponding to α and β phases, respectively, and are functions of composition c . H is a function of the structural order parameter η . δ is a scalar regularization parameter. The interface normal vector \mathbf{n} is given by

$$\mathbf{n} = \frac{\nabla\eta}{|\nabla\eta|} \quad (2)$$

for $\nabla\eta \neq \mathbf{0}$, and $\mathbf{n} = \mathbf{0}$ when $\nabla\eta = \mathbf{0}$.

1 Variational treatment

Following standard variational arguments (see Cahn-Hilliard formulation), we obtain the chemical potentials:

$$\mu_c = (f_{\alpha,c}(1-H) + f_{\beta,c}H) \quad (3)$$

$$\mu_{\eta} = (f_{\beta,c} - f_{\alpha,c})H_{,\eta} - \nabla \cdot \mathbf{m} + \delta^2\Delta(\Delta\eta) \quad (4)$$

The components of the anisotropic gradient \mathbf{m} are given by

$$m_i = \gamma(\mathbf{n}) \left(\nabla\eta + |\nabla\eta|(\delta_{ij} - n_i n_j) \frac{\partial\gamma(\mathbf{n})}{n_j} \right), \quad (5)$$

where δ_{ij} is the Kronecker delta.

2 Kinetics

Now the PDE for Cahn-Hilliard dynamics is given by:

$$\frac{\partial c}{\partial t} = -\nabla \cdot (-M_c \nabla\mu_c) \quad (6)$$

$$= M_c \nabla \cdot (\nabla(f_{\alpha,c}(1-H) + f_{\beta,c}H)) \quad (7)$$

and the PDE for Allen-Cahn dynamics is given by:

$$\frac{\partial \eta}{\partial t} = -M_{\eta} \mu_{\eta} \quad (8)$$

$$= -M_{\eta} [(f_{\beta,c} - f_{\alpha,c})H_{,\eta} - \nabla \cdot \mathbf{m} + \delta^2\Delta(\Delta\eta)] \quad (9)$$

where M_c and M_η are the constant mobilities. In order that the formulation only includes second order derivatives, an auxiliary field ϕ is introduced to break up the biharmonic term:

$$\phi = \Delta\eta \quad (10)$$

and the PDE for Allen-Cahn dynamics becomes

$$\frac{\partial\eta}{\partial t} = -M_\eta ((f_{\beta,c} - f_{\alpha,c})H_{,\eta} - \nabla \cdot \mathbf{m}) + \delta^2 \Delta\phi. \quad (11)$$

3 Time discretization

Considering forward Euler explicit time stepping, we have the time discretized kinetics equation:

$$\phi^{n+1} = \Delta\eta^n \quad (12)$$

$$\eta^{n+1} = \eta^n - \Delta t M_\eta ((f_{\beta,c}^n - f_{\alpha,c}^n)H_{,\eta}^n - \nabla \cdot \mathbf{m}^n + \delta^2 \Delta\phi^n) \quad (13)$$

$$c^{n+1} = c^n + \Delta t M_\eta \nabla \cdot (\nabla(f_{\alpha,c}^n(1 - H^n) + f_{\beta,c}^n H^n)) \quad (14)$$

4 Weak formulation

In the weak formulation, considering an arbitrary variation w , the above equations can be expressed as residual equations.

$$\int_{\Omega} w \phi^{n+1} dV = \int_{\Omega} \nabla w \cdot \underbrace{\nabla \eta^n}_{r_{\phi x}} dV \quad (15)$$

$$\int_{\Omega} w \eta^{n+1} dV = \int_{\Omega} w \eta^n - w \Delta t M_\eta ((f_{\beta,c}^n - f_{\alpha,c}^n)H_{,\eta}^n - \kappa \Delta \eta^n) dV \quad (16)$$

$$= \int_{\Omega} w \left(\underbrace{\eta^n - \Delta t M_\eta ((f_{\beta,c}^n - f_{\alpha,c}^n)H_{,\eta}^n)}_{r_\eta} \right) + \nabla w \cdot \underbrace{(-\Delta t M_\eta)(\mathbf{m}^n - \delta^2 \phi^n)}_{r_{\eta x}} dV \quad (17)$$

and

$$\int_{\Omega} w c^{n+1} dV = \int_{\Omega} w c^n + w \Delta t M_c \nabla \cdot (\nabla(f_{\alpha,c}^n(1 - H^n) + f_{\beta,c}^n H^n)) dV \quad (18)$$

$$= \int_{\Omega} w \underbrace{c^n}_{r_c} + \nabla w \underbrace{(-\Delta t M_c) [(f_{\alpha,c}^n(1 - H^n) + f_{\beta,c}^n H^n) \nabla c + ((f_{\beta,c}^n - f_{\alpha,c}^n)H_{,\eta}^n \nabla \eta^n)]}_{r_{cx}} dV \quad (19)$$

The above values of $r_{\phi x}$, r_η , $r_{\eta x}$, r_c and r_{cx} are used to define the residuals in the following equations file: *applications/anisotropyFacet/equations.h*

5 Anisotropy

The above formulation is generic to any $\gamma(\mathbf{n})$. In this application, we use an anisotropy of the form

$$\gamma(\mathbf{n}) = \gamma_0 \left(1 - \sum_{i=1} \alpha_i (\mathbf{m}_i \cdot \mathbf{n})^{w_i} \Theta(\mathbf{m}_i \cdot \mathbf{n}) \right), \quad (20)$$

where \mathbf{m} is a unit vector corresponding to a crystallographic orientation, γ_0 is a scaling factor for interfacial energy, α_i and w_i are scalar parameters specific to each orientation, and $\Theta(\cdot)$ is the Heaviside function. The derivatives with respect to components of the normal are

$$\frac{\partial \gamma(\mathbf{n})}{\partial n_j} = -\gamma_0 \sum_{i=1} w_i \alpha_i m_{ij} (\mathbf{m}_i \cdot \mathbf{n})^{w_i-1} \Theta(\mathbf{m}_i \cdot \mathbf{n}), \quad (21)$$

Calculation of $\gamma(\mathbf{n})$ and $\partial\gamma(\mathbf{n})/\partial n_j$ is performed in an application-specific function located in *applications/anisotropyFacet/facet_anisotropy.h*.

This anisotropy was developed by M. Salvalaglio et al. (doi: 10.1021/acs.cgd.5b00165), and is extensively documented in their paper. Briefly, we note that α_i determines the interfacial energy at the orientation \mathbf{m}_i , and w_i determines how localized the change interfacial energy is around \mathbf{m}_i . The Heaviside function $\Theta(\mathbf{m}_i \cdot \mathbf{n})$, which returns zero if $\mathbf{m}_i \cdot \mathbf{n} < 0$ and one otherwise, ensures that orientations are considered independently; i.e. there is no change in $\gamma(\mathbf{n})$ around $-\mathbf{m}_1$ unless that corresponds to another listed orientation \mathbf{m}_i . In its intended configuration, with $0 < \alpha_i < 1$ and high w_i (e.g. $w_i = 50$), this anisotropy results in nearly flat facets at the orientations \mathbf{m}_i .