# Current problems in stellar pulsation theory

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The last decade lead to major progress in asteroseismology and stellar physics with the advent of space missions. Thanks to the richness and precision of current oscillation spectra, sophisticated seismic probing techniques allow us now to pinpoint the limits of our current models of stellar structure and evolution. However, the accuracy of the seismic diagnosis depends on the accuracy of the pulsation models. In solar-like oscillations, the main source of inaccuracy comes from the near-surface layers where the oscillations are non-adiabatic and strongly coupled with turbulent convection. Some pulsating stars rotate fast and this must be accurately taken into account in the modeling of their pulsations. In others, the magnetic field or the dynamic tides could play some role. I propose here an overview of the great a
hievements and urrent limitation of asteroseismology.

# $\mathbf{1}$

techniques and talk about future prospects. (2017).

sure the ore rotation of numerous red giants (see Gehan et al., 2018, for the last measurements) with e.g. the method of Goupil et al. (2013). This revealed unexpe
ted slow core rotation due to unknown braking processes 2.2 g-modes (see e.g. Eggenberger et al., 2012; Marques et al., 2013; Eggenberger  $et$   $al., 2017$ ). Differential rotation could be giants rotation rates seem in agreement with models of

1 Introdu
tion angular momentum transport by plume-indu
ed internal  $\alpha$  is golden years of mixed modes allow us now to easily distinguish be $t_{\rm{max}}$  ,  $t_{\rm$ where we have  $\frac{1}{2}$  is the motion of  $\frac{1}{2}$  and  $\frac{1}{2}$  are  $\frac{1}{2}$  and  $\frac{1}{2}$  an  $\frac{1}{2}$  of asterosemologists extra on  $\frac{1}{2}$  ,  $\frac{1}{2}$  ,  $\frac{1}{2}$  as the masses of their helium cores (Montalban et al., 2013; tion spe
tra of unequaled ri
hness. This already lead to Bossini et al., 2015). Some of these great a
hievements major of the discussed in several papers of these proceedings by view of stellar interiors. As a general introdu
tion to the S. Deheuvels, D. Stello, M. Vrard and M. Takata. For  $\mathbf{P}$  summarize here some of the art of red giants asteroseismology as it was in 2017, see Hekker & Christensen-Dalsgaard as it was in 2017, see Hekker & Christensen-Dalsgaard gravity waves (Pinçon *et al.*, 2017). The period spacings  $\mathcal{L}^{\text{max}}$  , and the contract of the cont

 $\sim$  2016), stretching methods now enable us to transform 2.1 Red giants the non-regular os
illation patterns of mixed modes into The distribution of the species of the species of the contracted indicate). Three families of seismic indicators arise:  $\frac{1}{2}$  in red given the  $\frac{1}{2}$  reduced to the classical ones associated to the acoustic cavity (like in the most important and main-sequence solar-like stars), those related to the  $\Gamma$ ,  $\sim$   $\Gamma$  or  $\Gamma$  in the set of the set of the set of the space of the space of the space of the set of the set of the space of the space of the set of the set of the space of the space of the space of the space of  $\frac{1}{2}$  in the mass of  $\frac{1}{2}$  and  $\frac{1}{2}$  and agrees of  $\frac{1}{2}$  in the mass and agent ing on the case, the coupling  $\alpha$  is the canonical corresponding the evanesce and Ke-Canonical and Ke-Canonical corresponding the evanesce in zone. We can already pler were presented the possible was highly well-comed by galaxies. I learn a lot about stellar interiors with that ! Beyond ar
haeologists, opening an entirely new interdis
iplinary that, buoyan
y glit
hes an be looked for and interpreted  $\epsilon$  and  $\epsilon$  is the rest to inonne
tion. For more detail, I refer to its paper in these lude mixed-modes ould be developed. The future prospe
ts for red giants's seismology are great. Based on an improved asymptoti theory (Takata, regular ones (see Mosser et al., 2018, for the most re-

procedure the new step is now to the new step is the new step is the future prospects for subgiants and young red gimeasurements with CAIA. Legfor to the paper of Marchas are also great. With the method developed by De-Pinsonnogult (these proceedings) for paper detail  $\alpha$  heuvels et al. (2017), it is now possible to probe the course the future is even brighter with PLATO (Miglie). Core rotation of red giants from their asymmetric split- $\frac{1}{2}$   $\epsilon$  al., 2017). The avoided to the avoided consistings with mixed  $S_{\text{S}}$  ond, the discrete modes in redship modes in redship and inversion techniques  $\mathbb{S}$  and the peared to be a good gate reality. And this is not a simple den ore. Rotational splittings of mixed modes were rst problem, be
ause ontrary to more evolved red giants, over the set al. (2012). They were used to modes are out of the asymptotic regime in the g-cavity, so that the stret
hing is not easy.

## 2.2 g-modes

measured in subgiants (Deheuvels et al., 2014) and ore ent types ofstars is a se
ond great a
hievement. Series helium burning stars (Deheuvels *et al.*, 2015). The sub- of consecutive modes are now identified in  $\gamma$  Dor (see The discovery of dense spectra of g-modes in differ-Van Reeth et al., 2015) and SPB stars (see e.g.

although some consider it as "fragile" (Schunker et al.,

stars (Buldgen et al., 2015b). Internal mixing (Buldgen et al., 2015a) and ore overshooting (Deheuvels et al., 2016) an now be seismi
ally measured. Internal rotation ould be probed, revealing nearly uniform internal

methods ould also be envisioned.

The dis
overy of rotation related modes is also an important recent gift. Particularly interesting are the global Rossby modes whi
h appear to be dete
ted in many types of stars:  $\gamma$  Dor stars (Van Reeth et al., 2016), spotted A and B stars, bursting Be stars and the heart beat stars (Saio et al., 2018b). The latter are particularly interesting for the study of the oupling between os
illations and tidal for
es, as detailed in Guo (these pro
eedings). More detail on the high potential of these modes for seismic probing and their interesting properties is given in Saio et al. (2018b) and Saio (these pro-

forward modeling, but inversion begins to be also possi-

Pápi
s et al., 2017), allowing us to measurea

urately ble. The main spe
i limitation of the forward modeling their ore rotation and the extension of their onve
tive approa
h is that it redu
es the ri
hness and omplexity ore (Ouazzani et al., 2017; Christophe et al., 2018; Van of stellar evolution to a small number of parameters to Reeth et al., 2016). More detail about that is given by be determined. On the one hand, there are the physi- Bedding and Ouazzani (these pro
eedings). Dense spe
- al parameters: mass, age, X, Z and, if in
luded in the tra of g-modes are also observed in extreme horizontal models, the rotation rate Ω. On the other hand, there branch stars. Although the interpretation of these spec- are parameters such as the mixing-length parameter  $\alpha,$ tra is a matter of debate (Reed *et al.*, 2011; Østensen be overshooting parameter  $\alpha_{ov}$ , turbulent diffusion coefet al., 2014; Charles the al., 2014; the holds the hope of the holds the holds the holds the second the last parameters are associated the last  $\mu$ a very determined seisming of the models of the near the models of the MLT (typically the MLT) and  $\alpha$ future (see Charpinet, these pro
eedings). Detailed seis-and hemi
al transport. The results obtained by this apmi probing of the internal omposition of white dwarfs proa
h are thus intrinsi
ally limited. They are also model is now possible as shown re
ently by Giammi
hele et al. dependent sin
e they depend on the hoi
e of the opa
- (2018), see also the review of Hermes (these pro
eed-ity table, the equation of state, the onve
tion treatment ings). Finally, there is also the possible dis
overy of (MLT versus FST, instantaneous versus diusive overhigh-order g-modes in the Sun (Fossat et al., 2017; Fossat shooting, . . . ), the initial hemi
al mixture, . . . A rst & S
hmider, 2018, see the paper in these pro
eedings), limitation of inversion te
hniques is that they are linear, 2018). 2017a, for the ina

ura
ies introdu
ed by non-linearity  $\blacksquare$ very commenced are now to perform the normal section are not pentity only the case with the longest Repler Hghtcurves  $t_{\text{total}}$  is solar-type stars. Of particularly  $t_{\text{total}}$  or solar-like stars. However, Buldgen et al. (2019) showed is the soalled Kepler LEGACY Sample of stars, for re
ently that seismi inversion of the mean density of red which are seen to receive the sets of radii, masses and giants is also possible, based on their radial modes only.  $\alpha$  in a set  $\alpha$  in a single set al.,  $\alpha$  and  $\alpha$  common limitations of forward modeling and seismic in-The pre
ision of the frequen
ies and thenumber of de-version are the surfa
e ee
ts problem and the standard ted modes makes makes it possible to extend seismile to extend a seismile to extend a seismile to extend a s  $\mathcal{L}$  include the Sun to the Su which requires a good reference model (see Buldgen *et al.*, in seismic inversions) and complicates their application that they require many identified modes, which is cur-

### 3.1 Model dependen
e

 $\frac{1}{2}$ nasimetely tyjes as fact as their midletitudes (Dependix small number of parameters defining standard models.  $p_{\text{rel}}(s)$  and these proceedings and  $\sum_{n=1}^{\infty}$  This problem is coupled with the small number of avail- $\epsilon$   $\epsilon$   $\epsilon$   $\epsilon$   $\epsilon$   $\epsilon$   $\epsilon$  able independent seismic indicators in many cases such as in the seismilling developed for the setsmining and giants. Fitting a red for the seismology of red giants. Fitting a red  $\circ$  , indicate stars is to be  $\circ$  beyond the giant with two parameters (its age and mass) is less than usual seismic indicators and introduce new ones. One the 4 parameters of the von Neumann's elephant! ...It  $\frac{1}{2}$  by  $\frac{1}{2}$  be phase math is the phase math  $\frac{1}{2}$  be should also not be forgot that asterosemology probes the  $P$  Rombing (2016). And interior one was reported that  $P$  interior or a star as it is now, not its evolutions and its  $\alpha$  Farnir et al. (2018) (and these processes such as atomic difof work is also done on the development of new kinds fusion, nu
lear burning and ma
ros
opi transport pro of optimization algorithms for forward seismi modeling. esses. Many papers in these pro
eedings are devoted to For the future, the development of non-linear inversion stellar physi
s. I just summarize here the main sour
es It is useful to onsider with a little more attention the of un
ertainty.

2.4 Processing and the modes of the union of diffusion. Neglecting it in forward seismic Concerning first the microphysics, we have the ubiqmodeling introduces systematic inaccuracies in e.g. age measurements. It is still treated approximately in most stellar evolution odes: partial ionization is generally negle
ted, metals are treated as a whole and radiative levitation is negle
ted. Some stellar evolution odes treat the diffusion element by element, include radiative forces and couple microscopic transport with macroscropic mixing (turbulen
e, thermoaline onve
tion, . . . ). This is important but the ost in term of omputation time is huge. I refer to Deal (these pro
eedings) for more detail.

 $\mathbf{r}_1$ . It is not possible yet to intermediate  $\mathbf{r}_1$  and  $\mathbf{r}_2$  $\overline{3}$  directly affect the temperature gradient and thus oscilla-However, we must not forget the limitations of present tion frequen
ies and ages (see e.g. Lebreton et al., 2014, asteroseismi te
hniques. The main urrent approa
h is g. 18). The new abundan
e determinations by Asplund Opa
ity omputations are still approximate. Indeed, and take into account the coupling between all states into account; a compromise is unavoidable. Opacities et al. (2009) lead to significant discrepancy with the seis-

mically inverted sound speed profile, the so-called solar problem. A local increase of the opacity just below the convective envelope is the most probable path to solve this problem (Basu  $&$  Antia, 2008). However, the recent new opacity computations by the Los Alamos National Laboratory (Colgan et al., 2016) and by the CEA (Mondet *et al.*, 2015) did not allow to solve the problem (Buldgen et  $al.$ , 2017b). This problem also illustrates the impact of the chemical mixture on the opacities. Erroneous assumptions on the chemical mixture and, in particular, assuming homogeneity of metal abundances is a source of errors. So, it is clear that opacities are still a source of unknown systematic inaccuracy in forward seismic modeling, either due to intrinsic inaccuracies in present opacity computations or due to inaccurate chemical mixture's assumptions. We must not forget also the uncertainties related to the equation of state. They are particularly important in brown dwarfs and probably also in white dwarfs. For more detail, I refer to Pain (these proceedings).

Macroscopic processes are subject to even larger uncertainties. The so-called rotational mixing hides in reality a complex interplay between angular momentum transport, chemical mixing, magnetism, tidal effects, mass loss, ... A state of the art of the problems associated to the modeling of these processes can be found in e.g Buldgen et al. (these proceedings) and their impact is discussed in e.g. Meynet  $et$  al. (2016). The transport of angular momentum by waves and modes is still very difficult to quantify. However, a new model of waves generation by penetrative convection recently proposed by Pinçon et al. (2016) could explain the internal rotation of subgiants (Pincon *et al.*, 2017).

Finally, there is of course the complexity of convection: on the one hand the uncertainties related to overshooting and semi-convection above convective cores (see Buldgen, these proceedings, for more detail), and on the other hand the uncertainties associated to convective envelopes and their coupling with oscillations, which I discuss in the next section.

### $\bf 3.2$ Surface effects

This leads me to consider the so-called surface effects problem. This warrants indeed a special attention. What are surface effects? In a nutshell, inaccurate modeling of the superficial layers affects the frequencies of high-order p-modes and leads thus to inaccurate seismic inferences. It must not be forgot that there are two sources of inaccuracies. On the one hand, the structural inaccuracies mainly associated to the modeling of convection in atmosphere models, and on the other hand the modal inaccuracies associated to the adiabatic approximation (neglecting thus the fact that oscillations are nonadiabatic and the coupling between oscillations and convection is strong in superficial layers).

### $3.2.1$ *Structural inaccuracies*

As detailed in Ludwig (these proceedings), 3D atmosphere models are now on the market. But how to use them appropriately for stellar evolution and asteroseismology is still under development. A first approach, the simplest one, is to use them to calibrate empirical frequencies corrections. The most recent work in this direction was done by Sonoi et al.  $(2015)$ , Ball et al.  $(2016)$ and Trampedach et al. (2017). A second approach is to use the 3D atmospheres to calibrate the convection

parameters of the approximate convection models used in our stellar evolution codes. Most recent work in this direction was done by Trampedach et al. (2014), Magic  $et\ al.$  (2015), Sonoi  $et\ al.$  (2018) and these proceedings. Finally, interpolation in 3D grids can also be envisioned. Preliminary work in this direction was recently done by Jørgensen et al.  $(2018)$ .

### 3.2.2 Modal inaccuracies

Modal inaccuracies are another piece -of cake...Oscillations are totally non-adiabatic near the surface. Moreover, the convective, thermal and oscillation time-scales are of the same order in the outermost layers of solar-like oscillators. Time-Dependent Convection (TDC) models are thus needed. I worked on that and I am strongly convinced that current models are by far too approximate and in many cases do not even catch the real physics of the coherent interaction between convection and oscillations. A few linear non-adiabatic oscillation models of the time-dependent interaction between convection and oscillations have been proposed and implemented. First, there is the model of Balmforth (1992), which is a non-local generalization of the MLT theory of Gough (1977), widely used by G. Houdek and his collaborators. Second, there is the model of Gabriel (1996) and Grigahcène et al. (2005), which is based on the approach originally proposed by Unno (1967). These two MLT perturbative theories are compared in Houdek  $&$  Dupret (2015). Finally, there is the even more complex TDC model developed by Xiong *et al.* (2015). All these models are clearly reaching their limits: on the one hand they encounter difficulties to fit observations (typically the mode line-widths, see below), and on the other hand, their complexity hides crude approximations. It is time to start trying to model this problem in all its 4D (3D space  $+$  time) complexity if we want to go out of this deadlock.

The good point which can help to progress at this level is that there are additional seismic constraints associated to solar-like stochastic excited oscillations: on the one hand the linewidths in the power spectrum, which are directly related to the mode damping rates, and on the other hand the amplitudes. The theoretical damping rates are obtained with non-adiabatic oscillation models including time-dependent convection and the theoretical amplitudes require the use of a stochastic excitation models, too. Confrontation to the observed values constrains thus these models and gives their more weight when they are used to model surface effects. The most recent confrontations with mode linewidths of Kepler stars are presented in Houdek (these proceedings) and Aarslev *et al.*  $(2018)$ .

### 3.2.3 Non-adiabaticity in classical pulsators

In other types of pulsating stars, non-adiabaticity has a negligible impact on the frequencies, so no problem of surface effect for them, which is a big advantage. But that does not mean that nonadiabatic modeling is useless for these stars. It enables to understand and characterize the driving processes at the origin of pulsations. Moreover, the predicted range of excited modes, amplitude ratios and phases can be computed and compared with observations. This provides strong constraints on the opacity in  $\beta$  Cep and SPBs (e.g. Walczak et al., 2013; Salmon et al., 2012; Daszyńska-Daszkiewicz et al., 2005; Dupret et al., 2004), in sdBs and in hot white dwarfs (Quirion *et al.*, 2009). Since the opacities depend on

spheri
al symmetry and transforms the usual 1D eigen value problem into a 2D non-separable one. Codes solvwork have been implemented (Ouazzani et al., 2012; Reese et al., 2006). They provide an entirely new view of fast rotating stars' pulsations, but the pri
e in omputation time is huge, making usual seismic probing methods impractical.

compromise between fast computation and accuracy  $\frac{1}{2}$ <br>(Ballet *at al* 2012) In particular using it appears to ings). (Ballot et al., 2012). In parti
ular, using it appears to be justified for the interpretation of the wonderful oscillation spectra detected in Kepler gamma Dor stars References the traditional approximation an be used to disentan- $\frac{1}{2}$  i.e.  $\frac{1}{2}$  os  $\frac{1}{2}$  i.e.  $\frac{1}{2}$  i.e.  $\frac{1}{2}$  Asplund, M., Grevesse, N., Sauval, A. J., & Scott, P. Bouabid et al., 2013). It shows that the two main seismic quantities that can be measured are the buoyancy radius  $\Pi = (\int N/r dr)^{-1}$ , directly related to the size  $\Xi_{\text{Ball W H Roock R}}^{4/3}$  Roock B Cameron B H  $\ell$  $\int \Omega N/r dr \Pi$ . Differential rotation could also be detected<br>Rallot I Ligniance E Prat V, Reese D, R, & Rieuof their onve
tive ore and their average ore rotation an we get more than these two measurements? I think the answer is yes. First, it is well known that trapping is possible in the  $\mu$ -gradient region, leading to oscillations of the period spacing (Miglio *et al.*, 2008). It could be  $\frac{588}{389}$ used to constrain the sharpness of the chemical transi- $\text{Basu}, \text{S.} \& \text{All} \& \text{A.} \& \text{All} \& \text{A.} \& \text{All} \& \text{A.} \& \text{All} \&$ and theoreti
al predi
tions, dips are sometimes present Be
k, P. G., Montalban, J., Kallinger, T., De Ridder, J., in the period space in the period contract to the association of the contract and the contract and the contract of the contrac ferential rotation and/or mode coupling not taken into  $R_{\text{odding}}^T T_R$ . Mossor  $R_{\text{dipole}}^T T_L$  Mortalbán, I account by the traditional approximation (Saio et al.,

There are however cases where the variable separation is not justified: the fast rotating  $\delta$  Sct and Be stars Benomar, O., Takata, M., Shibahashi, H., Ceillier, T., & are the clearest example. In  $\delta$  Sct stars, the equiva- $\Omega$  and  $\Omega$  and  $\Omega$  are determining and used an to measure their mean density (García Hernández et al., 2009). Mirouh et al. (2019) (and these pro
eedings) de- Bouabid, M. P., Dupret, M. A., Salmon, S., Montalbán, veloped a very promizing method of mode classification  $\cdots$  and stars based on network in the stars below  $\cdots$   $\cdots$  for mode identification. Important theoretical work was also also at the mode in contract on on our cases on any pured on pured on reese, we appeal the second transformation of the second transformation of the second transformation of the second transformation of the second tra servables, but this remains very difficult (Reese  $et$   $al.$ ,  $\mathcal{Q}$  and  $\mathcal{Q}$  are matrix remains the huge major different matrix  $\mathcal{Q}$  and  $\mathcal{Q}$  and  $\mathcal{Q}$  and  $\mathcal{Q}$  are  $\mathcal{Q}$  and  $\mathcal{Q}$  and  $\mathcal{Q}$  and  $\mathcal{Q}$  are  $\mathcal{Q}$  and  $\mathcal{Q}$  and  $\mathcal{Q}$  are  $\mathcal{$ lenge of omputing realisti evolutionary models of fast  $\alpha$  is the stars near the matrix of work  $\alpha$ ,  $\alpha$  is the stars in  $\alpha$ has been done at this level, e.g. in the frame of the ES-

### 3.4 Magnetic field, tidal effects, non-linearity

and tidal effects. Fortunately, this is mostly justified.

the hemi
al omposition, onstraints on hemi
al trans-For the magneti eld, the only major ex
eption is the port pro
esses an also be obtained, as shown in Hu modeling of ro Ap pulsations, in whi
h the Lorentz for
e et al. (2011). In older stars, this provides tests of time-has a signi
ant dynami
al ee
t on pulsations. A nondependent onve
tion models and their urrent limita-perturbative model for axisymmetri p-mode pulsations tions (Dupret et al., 2005a,b). of stars with dipole magneti elds was developed by detail the effects of a strong magnetic field on internal The last limitation of oscillation models I complete Here  $\alpha$  and  $\alpha$ is the usual separation in spherical marmonics, in foods, where  $\mu$  is the distribution  $\mu$  interacted the distribution in the distribution of the distribution of the distribution of the distribution of the distribution is on the error of fast rotation. Fast rotation breaks the solar in  $\lambda$  is and  $\lambda$  is a solar interior, and  $\lambda$ ing rigorously the pulsation equations in this 2D frame- fancy factor  $\theta$  frame- factor  $\theta$  ()  $h$  and  $g$  is gravitor of variables obtained within the widely used for the modeling of high amplitude radial se called traditional approximation appears to be a good with pulsations and could help to explain longstanding probcompromise between fest computation and escureoy. Jems such as the Blazhko effect (Kolláth, these proceede.g. Saio & Gauts
hy (2004). More re
ently, Loi & Papaloizou (2018) (and these proceedings) analysed in ompa
t star os
illations. Studying the impa
t of tidal forces on pulsations in close binaries is still in its indetection of tidally excited oscillations in heartbeat stars observed by Kepler (Guo et al., 2017) (and these proceedings), with first models developed by Fuller  $(2017)$ . The linear approximation is ubiquitous in oscillation models used in asteroseismology. Currently, it seems unavoidable. On the opposite, non-linear pulsation models are ings).

- $\lambda_1$  and  $\lambda_2$ ,  $\lambda_3$  are asymptotic model in the asymptotic contract  $\lambda_1$ . Houdek, G., Handberg, R.,  $\&$ Christensen-Dalsgaard, J. 2018, MNRAS, 478, 69.
	- 2009, Annual Review of Astronomy and Astrophysi
	s, 47, 481.
	- $\mathcal{L}_{\text{tot}}$  relation to the size of the 2016, A&A, 592, A159.
- $\frac{1}{2}$  is the function for the future is  $\frac{1}{2}$  for the future in  $\frac{1}{2}$  and  $\frac{1}{2}$  in  $\frac{1}{2}$  for  $\frac{1}{2}$  subsets in  $\frac$ Ballot, J., Lignières, F., Prat, V., Reese, D. R., & Rieuwith Helio- and Asteroseismology, edited by H. Shibahashi, M. Takata, & A. E. Lynas-Gray, Astronomi
al Society of the Pacific Conference Series, vol. 462, p.
	- Balmforth, N. J. 1992, MNRAS, 255, 632.
	-
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	- Bedding, T. R., Mosser, B., Huber, D., Montalbán, J., Be
	k, P., et al. 2011, Nature, 471, 608.
- $\mathcal{P}$  and  $\mathcal{P}$  and  $\mathcal{P}$  benomar, O., Bazot, M., Nielsen, M. B., Gizon, L., Sekii, T., et al. 2018, Science, 361, 1231.
	- Gar
	ía, R. A. 2015, MNRAS, 452, 2654.
	- talbán, J., et al. 2015,MNRAS, 453, 2290.
	- J., Miglio, A., et al. 2013, MNRAS, 429, 2500.
	- 583, A62.
	- 598, A21.
	- 2015b, A&A, 574, A42.
	- A., et al. 2019,MNRAS, 482, 2305.
- $T = T$ ,  $\Gamma$ problem is far from being fully solved. Reese, D. R., et al. 2017b, A&A, 607, A58.
	- Current pulsation models usually negle
	t magneti G. Fontaine, & S. Charpinet, Astronomi
	al So
	iety of Charpinet, S., Brassard, P., Van Grootel, V., & Fontaine, G. 2014, In 6th Meeting on Hot Subdwarf Stars and Related Obje
	ts, edited by V. van Grootel, E. Green, the Pacific Conference Series, vol. 481, p. 179.
- Christophe, S., Ballot, J., Ouazzani, R. M., Antoci, V., & Salmon, S. J. A. J. 2018, A&A, 618, A47.
- Colgan, J., Kilcrease, D. P., Magee, N. H., Sherrill, M. E., Abdallah, J., J., et al. 2016, ApJ, 817, 116.
- Cunha, M. S., Stello, D., Avelino, P. P., Christensen-Dalsgaard, J., & Townsend, R. H. D. 2015, ApJ, 805, 127.
- Daszyńska-Daszkiewicz, J., Dziembowski, W. A., & Pamyatnykh, A. A. 2005, A&A, 441, 641.
- Deheuvels, S., Ballot, J., Beck, P. G., Mosser, B., Østensen, R., et al. 2015, A&A, 580, A96.
- Deheuvels, S., Brandão, I., Silva Aguirre, V., Ballot, J., Michel, E., et al. 2016, A&A, 589, A93.
- Deheuvels, S., Doğan, G., Goupil, M. J., Appourchaux, T., Benomar, O., et al. 2014, A&A, 564, A27.
- Deheuvels, S., Ouazzani, R. M., & Basu, S. 2017, A&A, 605, A75.
- Dupret, M. A., Grigahcène, A., Garrido, R., De Ridder, J., Scuflaire, R., et al. 2005a, MNRAS, 361, 476.
- Dupret, M. A., Grigahcène, A., Garrido, R., De Ridder, J., Scuflaire, R., et al. 2005b, MNRAS, 360, 1143.
- Dupret, M. A., Thoul, A., Scuflaire, R., Daszyńska-Daszkiewicz, J., Aerts, C., et al. 2004, A&A, 415, 251.
- Eggenberger, P., Lagarde, N., Miglio, A., Montalbán, J., Ekström, S., et al. 2017, A&A, 599, A18.
- Eggenberger, P., Montalbán, J., & Miglio, A. 2012, A&A, 544, L4.
- Farnir, M., Dupret, M.-A., Salmon, S. J. A. J., Noels, A., & Buldgen, G. 2018,  $arXiv$  e-prints,  $arXiv:1812.04984$ .
- Fossat, E., Boumier, P., Corbard, T., Provost, J., Salabert, D., et al. 2017, A&A, 604, A40.
- Fossat, E. & Schmider, F. X. 2018, A&A, 612, L1.
- Fuller, J. 2017, MNRAS, 472, 1538
- Gabriel, M. 1996, Bulletin of the Astronomical Society of India, 24, 233.
- García Hernández, A., Moya, A., Michel, E., Garrido, R., Suárez, J. C., et al. 2009, A&A, 506, 79.
- Gehan, C., Mosser, B., Michel, E., Samadi, R., & Kallinger, T. 2018, A&A, 616, A24.
- Giammichele, N., Charpinet, S., Fontaine, G., Brassard, P., Green, E. M., et al. 2018, Nature, 554, 73.
- Gough, D. O. 1977, ApJ, 214, 196.
- Goupil, M. J., Mosser, B., Marques, J. P., Ouazzani, R. M., Belkacem, K., et al. 2013, A&A, 549, A75.
- Grigahcène, A., Dupret, M. A., Gabriel, M., Garrido, R., & Scuflaire, R. 2005, A&A, 434, 1055.
- Guo, Z., Gies, D. R., & Fuller, J. 2017, ApJ, 834, 59.
- Hekker, S. & Christensen-Dalsgaard, J. 2017, Astronomy and Astrophysics Review, 25, 1.
- Houdek, G. & Dupret, M.-A. 2015, Living Reviews in Solar Physics, 12, 8.
- Hu, H., Tout, C. A., Glebbeek, E., & Dupret, M. A. 2011, MNRAS, 418, 195.
- Jørgensen, A. C. S., Mosumgaard, J. R., Weiss, A., Silva Aguirre, V., & Christensen-Dalsgaard, J. 2018, MN-RAS, 481, L35.
- Lebreton, Y., Goupil, M. J., & Montalbán, J. 2014, In EAS Publications Series, EAS Publications Series, vol. 65, pp. 99-176.
- Loi, S. T. & Papaloizou, J. C. B. 2018, MNRAS, 477, 5338
- Magic, Z., Weiss, A., & Asplund, M. 2015, A&A, 573, A89.
- Marques, J. P., Goupil, M. J., Lebreton, Y., Talon, S., Palacios, A., et al. 2013, A&A, 549, A74.
- Meynet, G., Maeder, A., Eggenberger, P., Ekstrom, S., Georgy, C., et al. 2016, Astronomische Nachrichten, 337, 827.
- Miglio, A., Chiappini, C., Mosser, B., Davies, G. R., Freeman, K., et al. 2017, Astronomische Nachrichten, 338, 644.
- Miglio, A., Montalbán, J., Baudin, F., Eggenberger, P., Noels, A., et al. 2009, A&A, 503, L21.
- Miglio, A., Montalbán, J., Noels, A., & Eggenberger, P. 2008, MNRAS, 386, 1487.
- Mirouh, G. M., Angelou, G. C., Reese, D. R., & Costa, G. 2019, MNRAS, 483, L28.
- Mondet, G., Blancard, C., Cossé, P., & Faussurier, G. 2015, The Astrophysical Journal Supplement Series, 220, 2.
- Montalbán, J., Miglio, A., Noels, A., Dupret, M. A., Scuflaire, R., *et al.* 2013, ApJ, 766, 118.
- Mosser, B., Gehan, C., Belkacem, K., Samadi, R., Michel, E., et al. 2018, A&A, 618, A109.
- Østensen, R. H., Telting, J. H., Reed, M. D., Baran, A. S., Nemeth, P., et al. 2014, A&A, 569, A15.
- Quazzani, R. M., Dupret, M. A., & Reese, D. R. 2012,  $A&A, 547, A75.$
- Ouazzani, R.-M., Salmon, S. J. A. J., Antoci, V., Bedding, T. R., Murphy, S. J., et al. 2017, MNRAS, 465, 2294.
- Pápics, P. I., Tkachenko, A., Van Reeth, T., Aerts, C., Moravveji, E., et al. 2017, A&A, 598, A74.
- Pinçon, C., Belkacem, K., & Goupil, M. J. 2016, A&A, 588, A122.
- Pinçon, C., Belkacem, K., Goupil, M. J., & Marques, J. P. 2017, A&A, 605, A31.
- Quirion, P.-O., Fontaine, G., & Brassard, P. 2009, In Journal of Physics Conference Series, vol. 172, p. 012077.
- Reed, M. D., Baran, A., Quint, A. C., Kawaler, S. D., O'Toole, S. J., et al. 2011, MNRAS, 414, 2885.
- Reese, D., Lignières, F., & Rieutord, M. 2006, A&A, 455, 621.
- Reese, D. R., Lignières, F., Ballot, J., Dupret, M. A., Barban, C., et al. 2017, A&A, 601, A130.
- Rieutord, M. & Espinosa Lara, F. 2013, Ab Initio Modelling of Steady Rotating Stars, p. 49.
- Roxburgh, I. W. 2016, A&A, 585, A63.
- Saio, H., Bedding, T. R., Kurtz, D. W., Murphy, S. J., Antoci, V., et al. 2018a, MNRAS, 477, 2183.
- Saio, H. & Gautschy, A. 2004, MNRAS, 350, 485.
- Saio, H., Kurtz, D. W., Murphy, S. J., Antoci, V. L., & Lee, U. 2018b, MNRAS, 474, 2774.
- Salmon, S., Montalbán, J., Morel, T., Miglio, A., Dupret, M. A., et al. 2012, MNRAS, 422, 3460.
- Schunker, H., Schou, J., Gaulme, P., & Gizon, L. 2018, SoPh, 293, 95.
- Silva Aguirre, V., Lund, M. N., Antia, H. M., Ball, W. H., Basu, S., et al. 2017, ApJ, 835, 173.
- Sonoi, T., Ludwig, H. G., Dupret, M. A., Montalbán, J., Samadi, R.,  $et\;$   $al.\;$  2018, arXiv  $\;$  e-prints, arXiv:1811.05229.
- Sonoi, T., Samadi, R., Belkacem, K., Ludwig, H. G., Caffau, E., et al. 2015, A&A, 583, A112.
- Takata, M. 2016, Publications of the Astronomical Soci $ety$  of Japan, 68, 109.
- Trampedach, R., Aarslev, M. J., Houdek, G., Collet, R., Christensen-Dalsgaard, J., et al. 2017, MNRAS, 466, L43.
- Trampedach, R., Stein, R. F., Christensen-Dalsgaard, J., Nordlund, Å., & Asplund, M. 2014, MNRAS, 445, 4366.
- Unno, W. 1967, Publications of the Astronomical Society of Japan, 19, 140.
- Van Reeth, T., Mombarg, J. S. G., Mathis, S., Tka
henko, A., Fuller, J., et al. 2018, A&A, 618, A24.
- Van Reeth, T., Tka
henko, A., & Aerts, C. 2016, A&A, 593, A120.
- Van Reeth, T., Tka
henko, A., Aerts, C., Pápi
s, P. I., Triana, S. A., et al. 2015, The Astrophysi
al Journal Supplement Series, 218, 27.
- Wal
zak, P., Daszy«ska-Daszkiewi
z, J., Pamyatnykh, A. A., & Zdravkov, T. 2013, MNRAS, 432, 822.
- Xiong, D. R., Deng, L., & Zhang, C. 2015, MNRAS, 451,