

THE RELATION BETWEEN THE CONVERGENCE OF
SERIES AND OF INTEGRALS

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It has long been known that, when $f(x)$ is a positive function which steadily decreases as x tends to infinity, the series $\sum_{n=1}^{\infty} f(n)$ and the integral $\int_{\mu}^{\infty} f(x) dx$ converge or diverge together: and in case of divergence, the difference

$$\int_{\mu}^{\nu} f(x) dx - \sum_{n=\mu}^{\nu} f(n)$$

tends to a definite limit as ν tends to infinity.*

But, when the series contains terms of both signs (although steadily decreasing in numerical value), it is not easy to make a corresponding statement with reference to the relation between the convergence of the series and of the integral.†

It is known, of course, that, when $f(x)$ decreases steadily to zero as a limit, the two series and the two integrals

$$\sum_{n=1}^{\infty} f(n) \sin na, \quad \int_0^{\infty} f(x) \sin ax dx, \quad \sum_{n=1}^{\infty} f(n) \cos na, \quad \int_0^{\infty} f(x) \cos ax dx$$

are all convergent.‡ And these results can be easily extended to cases in which the periodic factors are of the forms

$$\sin ax P(\sin^2 ax), \quad \cos ax P(\sin^2 ax),$$

where P is a *polynomial*. But, if the polynomial P is replaced by

* The main part of the theorem goes back to Maclaurin: for a proof of the latter part, see my *Infinite Series*, Art. 11. When the monotonic condition is removed from $f(x)$, the theorem is no longer true; for examples, see p. 423 of my book.

† Of course, in a large number of interesting cases, the terms decrease fast enough to ensure *absolute* convergence. This case is covered, from the practical point of view, by Maclaurin's rule, and we shall suppose that absolute convergence is excluded from the cases discussed here.

‡ See, for example, *Infinite Series*, Arts. 19, 20, 169.

a general continuous function, the cosine-series is known to diverge for certain values of α , although the integral is always convergent; however, with the same values of α the sine-series is convergent.*

If the periodic factor in the series is of the form

$$\sin \phi(x) \quad \text{or} \quad \cos \phi(x),$$

where $\phi(x)$ tends to infinity *more rapidly than* x , it is practically certain that the convergence of the integral gives no information with regard to the nature of the series.

Thus, for instance, the integrals

$$\int_0^{\infty} f(x) \sin(ax^p) dx, \quad \int_0^{\infty} f(x) \cos(ax^p) dx,$$

where p is positive, will converge if $f(x)$ steadily decreases to any finite limit (not necessarily zero)†: but, on the other hand, the two series

$$\sum f(n) \sin(an^p), \quad \sum f(n) \cos(an^p)$$

have only been considered for rational values of α/π and integral values of p ; they are then known to diverge [even if $f(x)$ tends to zero] unless a certain condition is satisfied; and this condition is certainly broken even in the simplest case (when p is 2) except for special values of α .‡

The object of the following note is to prove that [with certain restrictions on the functions, stated in (a)-(d) on p. 329], *when $\phi(x)$ tends steadily to infinity, but MORE SLOWLY THAN x , the behaviour of the integrals*

$$\int_0^{\infty} f(x) \sin \phi(x) dx, \quad \int_0^{\infty} f(x) \cos \phi(x) dx$$

entirely settles the character of the series

$$\sum_0^{\infty} f(n) \sin \phi(n), \quad \sum_0^{\infty} f(n) \cos \phi(n).$$

This theorem is then applied, in § 2, to extend (and simplify the proofs of) certain known theorems, the simplest of which is that if

$$A_{\nu} + iB_{\nu} = \sum_1^{\nu} \frac{1}{n^{1+\alpha}} \quad (\alpha \text{ real}),$$

then $A_{\nu} + iB_{\nu} - i/\alpha\nu^{\alpha}$ tends to a definite limit, so that both A_{ν} and B_{ν} have a range of oscillation $2/\alpha$ as ν tends to infinity.

* See Bromwich and Hardy, *Quarterly Journal*, Vol. xxxix., May, 1908, pp. 232, 236, 240, and also below, p. 338.

† See my *Infinite Series*, Art. 169 and Ex. 8, p. 468.

‡ See below, § 3, p. 338; and Genocchi, *Atti di Torino*, t. x., 1875, p. 991.

1. *Proof of the Theorem.*

Let us write $F(x) = f(x) \sin \phi(x)$, where we suppose that

- $$\left. \begin{array}{l} (\alpha) \quad f(x) \text{ tends steadily to zero } * \\ (\beta) \quad \phi(x) \text{ tends steadily to infinity} \\ (\gamma) \quad \phi'(x) \text{ tends steadily to zero} \end{array} \right\} \text{ as } x \rightarrow \infty.$$

From these conditions it follows that

$$(1) \quad f'(x) \leq 0, \quad \phi'(x) \geq 0.$$

Now, if we write $X_n = \int_n^{n+1} F(x) dx - F(n)$, we have the equations

$$(2) \quad X_n = \int_n^{n+1} \{F(x) - F(n)\} dx = \int_0^1 \{F(n+t) - F(n)\} dt,$$

$$(3) \quad F(n+t) - F(n) = \int_0^t F'(n+v) dv.$$

But $F'(x) = f'(x) \sin \phi(x) + f(x) \phi'(x) \cos \phi(x)$;

and so $|F'(n+v)| \leq |f'(n+v)| + |f(n+v)| \cdot |\phi'(n+v)|$.

Now, from conditions (α) , (β) , (γ) and from the inequalities (1), we see that

$$|f'(n+v)| = -f'(n+v),$$

and $|f(n+v)| \cdot |\phi'(n+v)| \leq f(n) \phi'(n)$.

It follows that $|F'(n+v)| \leq f(n) \phi'(n) - f'(n+v)$.

Thus, making use of the last inequality in the equation (3), we find that

$$(3') \quad \begin{aligned} |F(n+t) - F(n)| &\leq t f(n) \phi'(n) + f(n) - f(n+t) \\ &\leq f(n) \phi'(n) + f(n) - f(n+1), \end{aligned}$$

provided that t belongs to the interval $(0, 1)$. Hence, combining (3') and (2), we see that

$$(4) \quad |X_n| \leq f(n) \phi'(n) + f(n) - f(n+1).$$

We now introduce the condition that

$$(\delta) \quad \text{the integral } \int_0^\infty f(x) \phi'(x) dx \text{ is convergent.}$$

* As remarked above (foot-note, p. 327), we suppose absolute convergence excluded, so that $\int_0^\infty f(x) dx$ is divergent. When this integral is convergent, the discussion given here is quite superfluous.

Consequently, by Maclaurin's theorem quoted in the introduction to this paper (p. 327), we see that the series

$$(5) \quad \sum_{\mu}^{\infty} f(n) \phi'(n)$$

is also convergent because $f(x) \phi'(x)$ tends steadily to zero, in virtue of conditions (α), (γ). Further,

$$\sum_{\mu}^{\nu} \{f(n) - f(n+1)\} = f(\mu) - f(\nu+1),$$

and this tends to the limit $f(\mu)$ as ν tends to infinity; thus the series

$$(6) \quad \sum_{\mu}^{\infty} \{f(n) - f(n+1)\}$$

is also convergent.

It follows from (4), (5), and (6) that the series $\sum |X_n|$ converges: and consequently the series $\sum X_n$ is absolutely convergent. But

$$\sum_{\mu}^{\nu} X_n = \int_{\mu}^{\nu} F(x) dx - \sum_{\mu}^{\nu} F(n) + \int_{\nu}^{\nu+1} F(x) dx$$

and
$$\left| \int_{\nu}^{\nu+1} F(x) dx \right| < \int_{\nu}^{\nu+1} f(x) dx < f(\nu),$$

so that
$$\lim_{\nu \rightarrow \infty} \int_{\nu}^{\nu+1} F(x) dx = 0.$$

Thus

$$(7) \quad \lim_{\nu \rightarrow \infty} \left\{ \int_{\mu}^{\nu} F(x) dx - \sum_{\mu}^{\nu} F(n) \right\} = \sum_{\mu}^{\infty} X_n;$$

and accordingly, since the series on the right has been proved to converge, the limit on the left is also definite.

It follows at once that, if $f(x)$ and $\phi(x)$ are subject to the conditions (α), (β), (γ), and (δ), the series

$$\sum_{\mu}^{\infty} f(n) \sin \phi(n)$$

and the integral
$$\int_{\mu}^{\infty} f(x) \sin \phi(x) dx$$

converge, diverge, or oscillate together.

Further, the equation (7) shews that, in case both oscillate, the amplitude of oscillation is the same for the series as for the integral; and, in case of divergence, the limit on the left of (7) is still finite.

Under the same conditions (α)-(δ), the same results apply to the series

$$\sum_{\mu}^{\infty} f(n) \cos \phi(n)$$

and the integral $\int^{\infty} f(x) \cos \phi(x) dx$;

and, consequently, the same conclusions apply also to the series

$$\sum^{\infty} f(n) \exp \{ \pm i\phi(n) \}$$

and the integrals $\int^{\infty} f(x) \exp \{ \pm i\phi(x) \} dx$.

We note as a special case that the conditions (a)-(d) are certainly satisfied by the function

$$\{M(x)\}^{-\kappa} \quad (\kappa = \beta + i\gamma, \beta > 0),$$

provided that $M(x)$ tends steadily to infinity in such a way that $M'(x)/M(x)$ tends steadily to zero. For we have then to take

$$f(x) e^{\pm i\phi(x)} = \{M(x)\}^{-\kappa}, \text{ or } f(x) = \{M(x)\}^{-\beta}, \text{ and } \phi(x) = |\gamma| \log \{M(x)\},$$

from which it is evident that the first three conditions are satisfied; as to the fourth condition, we must examine the integral

$$\int^{\infty} f(x) \phi'(x) dx = \int^{\infty} \frac{|\gamma| M'(x)}{\{M(x)\}^{(1+\beta)}} dx = |\gamma| \int^{\infty} \frac{dy}{y^{1+\beta}},$$

which is clearly convergent, so that all the four conditions are satisfied.

Again, if the two functions

$$F(x) = f(x) e^{\pm i\phi(x)}, \quad G(x) = g(x) e^{\pm i\psi(x)}$$

satisfy the prescribed conditions, their product $F(x)G(x)$ will also satisfy the conditions, provided that*

$$(\beta') \quad \phi(x) - \psi(x) \text{ tends steadily to infinity,}$$

$$(\gamma') \quad \phi'(x) - \psi'(x) \text{ tends steadily to zero.}$$

For then the product $f(x)g(x)$ will obviously tend steadily to zero, and each of the integrals

$$\int^{\infty} f(x) g(x) \phi'(x) dx, \quad \int^{\infty} f(x) g(x) \psi'(x) dx$$

is convergent, because $\lim f(x) = 0, \lim g(x) = 0$.

* These additional conditions are superfluous when the signs in the two exponential functions are the same.

2. *Examination of certain Special Series.*

1. Consider first the simple case *

$$\sum \frac{1}{n^{1+\alpha i}} \quad (\alpha \text{ real}).$$

Here we can take $M(x) = x$, $\beta = 1$, in the result at the end of § 1; then $M(x)$ tends steadily to infinity, while $M'(x)/M(x)$ tends steadily to zero.

Consequently, the behaviour of the series is determined by that of the integral

$$\int_1^{\nu} \frac{dx}{x^{1+\alpha i}} = \frac{1}{\alpha i} \left\{ 1 - \frac{1}{\nu^{\alpha i}} \right\}.$$

Now, as $\nu \rightarrow \infty$, this integral oscillates, the amplitude (both for real and for imaginary parts) being $2/\alpha$; and so the same is true of the series. The theorem of § 1 can also be applied to cases such as

$$M(x) = \log x, \quad \log(\log x), \quad \dots,$$

but the range of oscillation for the corresponding integrals and series is infinite.

2. Secondly, let us consider the type of series which is obtained by introducing a complex index in the general logarithmic series: that is, we consider the series

$$\sum \frac{1}{n \cdot l_1 n \cdot l_2 n \dots l_{k-1} n (l_k n)^{1+\alpha i}},$$

where † $l_k x = |\log(l_{k-1} x)|$, $l_1 x = \log x$.

Here we can take $\phi(x) = \alpha l_{k+1} x$, $f(x) = \phi'(x)/\alpha$

or ‡ $f(x) = 1 / \{x \cdot l_1 x \cdot l_2 x \dots l_k x\}$.

Then we can find a constant K , so that

$$f(x) \phi'(x) < K/x^2,$$

and so condition (δ) of § 1 is satisfied and the other conditions are

* This series can be discussed by Weierstrass's rule depending on the quotient of two consecutive terms in the series (see my *Infinite Series*, p. 204). The particular case of the rule which is needed here is, however, rather troublesome to establish; and it would be almost impossible to use a similar method in the other cases given below.

† Of course, after a certain stage, the logarithms are all positive, and the sign of the modulus may then be omitted; it is often simpler to suppose that the earlier terms are left out from the series so as to avoid this complication.

‡ Because, using accents for differential coefficients, we have

$$l'_{k+1} = \frac{l'_k}{l_k}, \quad l'_k = \frac{l'_{k-1}}{l_{k-1}}, \quad \dots, \quad l'_1 = \frac{1}{x}.$$

evidently satisfied. Thus we have to consider the integral

$$\int_{\mu}^{\nu} f(x) e^{-i\phi(x)} dx = \frac{1}{i\alpha} (e^{-i\phi(\mu)} - e^{-i\phi(\nu)}),$$

which again oscillates (as $\nu \rightarrow \infty$) with an amplitude $2/\alpha$; and so the series has the same range of oscillation.

3. It follows without further proof that, if $\psi(x)$ is any function tending steadily to zero (as $x \rightarrow \infty$) the series

$$\sum \frac{\psi(n)}{n^{1+\alpha}}, \quad \sum \frac{\psi(n)}{n \cdot l_1 n \cdot l_2 n \dots l_{k-1} n \cdot (l_k n)^{1+\alpha}}$$

are both convergent, in virtue of Dirichlet's test of convergence* and the results obtained in (1), (2) above.

Thus, as a simple example, we may note that the series

$$\sum \frac{1}{n^{1+\alpha} (\log n)^{\beta}}$$

is convergent if β is positive.

We can generalise these results still further by supposing that $\psi(x)$ is complex, but tends to zero as x tends to infinity in such a way that

$$\int^{\infty} |\psi'(x)| dx$$

is convergent. For then we have the inequality †

$$\left| \int_{X_1}^{X_2} F(x) \psi(x) dx \right| < HV \quad (X_2 > X_1),$$

where $V = \int_{X_1}^{\infty} |\psi'(x)| dx,$

and H is the upper limit to the integral

$$\left| \int_{X_1}^X F(x) dx \right| \quad (X_1 \leq X \leq X_2).$$

* See my *Infinite Series*, Art. 20.

† See *Proc. London Math. Soc.*, Vol. 6, 1907, p. 65; it is perhaps worth while to remark that there the integral is proved to be less than HV' , where

$$V' = \int_{X_1}^{X_2} |\psi'(x)| dx + |\psi(X_2)|.$$

But $\psi(X_2) = -\int_{X_2}^{\infty} \psi'(x) dx;$

and so $|\psi(X_2)| \leq \int_{X_2}^{\infty} |\psi'(x)| dx.$

Thus $V' \leq \int_{X_1}^{\infty} |\psi'(x)| dx,$

and so the value of V given in the text is larger than V' .

Thus, since $H \leq 2/a$ [see (1), (2) above], we can choose X_1 so that

$$\left| \int_{X_1}^{X_2} F(x) \psi(x) dx \right| < \epsilon \quad (\text{if } X_2 > X_1),$$

because, by proper choice of X_1 , the integral V can be made as small as we please. Consequently the integral

$$\int F(x) \psi(x) dx$$

is convergent: and so the same is true of the series

$$\sum F(n) \psi(n),$$

provided that the conditions of § 1 are satisfied by the function $F(x) \psi(x)$.

As an illustration of the last result, we may take

$$F(x) = x^{-(1+ai)}, \quad \psi(x) = (\log x)^{-\kappa} \quad (\kappa = \beta + i\gamma, \beta > 0).$$

For then

$$|\psi'(x)| = \frac{|\kappa|}{x(\log x)^{1+\beta}},$$

and so

$$\int |\psi'(x)| dx = \int \frac{|\kappa| dx}{x(\log x)^{1+\beta}},$$

which converges when β is positive, because the indefinite integral is

$$\frac{|\kappa|}{\beta} (\log x)^{-\beta}.$$

Further, as was pointed out in (1) above, the conditions of § 1 are satisfied by the functions $F(x), \psi(x)$; and so the conditions are satisfied also by their product, since the function in the exponential is here

$$a \log x + \gamma \log (\log x).$$

Thus we see that the series

$$\sum \frac{1}{n^{1+ai} (\log n)^{\beta+\gamma i}} \quad (\beta > 0)$$

is convergent; and the same method can be easily extended to more complicated cases. In this way it can be proved that the series*

$$\sum \frac{(l_2 n)^{\kappa_2} (l_3 n)^{\kappa_3} \dots (l_k n)^{\kappa_k}}{n^{1+ai} (l_1 n)^{\beta+\gamma i}} \quad (\beta > 0)$$

is convergent, whatever the indices $\kappa_2, \kappa_3, \dots, \kappa_k$ may be (real or complex).

* These examples were suggested to me by Mr. Hardy.

4. As a last example, we shall find an asymptotic formula for the series

$$\sum_1^{\nu} \frac{(\log n)^{\kappa}}{n^{1+\alpha i}} \quad (\kappa = \beta + i\gamma, \beta > 0),$$

which has been discussed in the special case $\kappa = 1$, by Mertens.*

Here we have to put

$$f(x) = \frac{(\log x)^{\beta}}{x}, \quad \phi(x) = \alpha \log x - \gamma \log(\log x),$$

and so we find
$$f'(x) = -\frac{(\log x)^{\beta}}{x^2} \left(1 - \frac{\beta}{\log x}\right),$$

$$\phi'(x) = \frac{1}{x} \left(\alpha - \frac{\gamma}{\log x}\right).$$

Thus the first three conditions of § 1 will be satisfied, as soon as $\log x$ is greater than both β and γ/α ; further, the integral (δ) of § 1 will converge, provided that

$$\int \frac{(\log x)^{\beta}}{x^2} dx$$

is convergent; but this reduces to the known integral

$$\int \xi^{\beta} e^{-\xi} d\xi \quad (\text{if } \xi = \log x),$$

and so all the conditions of § 1 are satisfied here.

Thus the asymptotic formulæ for the series

$$\sum_1^{\nu} \frac{(\log n)^{\kappa}}{n^{1+\alpha i}}$$

is given by the asymptotic formula for the integral

$$\int_1^{\nu} \frac{(\log x)^{\kappa}}{x^{1+\alpha i}} dx = \int_0^{\log \nu} \xi^{\kappa} e^{-\alpha i \xi} d\xi.$$

Thus, on integrating by parts, we get the formula

$$\begin{aligned} \int_1^{\nu} \frac{(\log x)^{\kappa}}{x^{1+\alpha i}} dx = & -\frac{1}{\nu^{\alpha i}} \left[\frac{(\log \nu)^{\kappa}}{\alpha i} + \kappa \frac{(\log \nu)^{\kappa-1}}{(\alpha i)^2} + \kappa(\kappa-1) \frac{(\log \nu)^{\kappa-2}}{(\alpha i)^3} + \dots \right. \\ & \left. + \kappa(\kappa-1) \dots (\kappa-m+1) \frac{(\log \nu)^{\kappa-m}}{(\alpha i)^{m+1}} \right] \\ & + \int_0^{\log \nu} \frac{\kappa(\kappa-1) \dots (\kappa-m)}{(\alpha i)^{m+1}} \xi^{\kappa-m-1} e^{-\alpha i \xi} d\xi, \end{aligned}$$

provided that $\beta - m$ is positive.

* *Göttingen Nachrichten*, 1887, p. 266; the method adopted by Mertens is to differentiate the first series of p. 332 with respect to α .

When β is an integer, the formula just obtained is only true if $m = \beta - 1$; to extend it to $m = \beta$, we divide the remainder integral into two, from 0 to 1 and from 1 to $\log \nu$. The second of these integrals may be again integrated by parts, and so we obtain the term given by $m = \beta$ together with a constant; the new remainder integral is then proportional to

$$\int_1^{\log \nu} \xi^{i\gamma-1} e^{-\alpha i \xi} d\xi,$$

which is easily seen to converge to a definite value as ν tends to infinity.* Consequently, if we write $m = \beta$ in the expression in square brackets at the foot of p. 335, the difference between this formula and the sum of the series will tend to a definite limit as ν tends to infinity.

When β is not an integer, we take m as the integer next less than β , and the remainder integral can then be proved to converge (as $\nu \rightarrow \infty$) by a method similar to that used in the last case.

Summing up, we have now the result that

$$\sum_1^{\nu} \frac{(\log n)^\kappa}{n^{1+\alpha i}} \sim -\frac{1}{\nu^{\alpha i}} \left[\frac{(\log \nu)^\kappa}{\alpha i} + \kappa \frac{(\log \nu)^{\kappa-1}}{(\alpha i)^2} + \kappa(\kappa-1) \frac{(\log \nu)^{\kappa-2}}{(\alpha i)^3} + \dots \right. \\ \left. + \kappa(\kappa-1) \dots (\kappa-m+1) \frac{(\log \nu)^{\kappa-m}}{(\alpha i)^{m+1}} \right],$$

where m is the integral part of β . Thus in particular we have Mertens's result

$$\sum_1^{\nu} \frac{\log n}{n^{1+\alpha i}} \sim -\frac{1}{\nu^{\alpha i}} \left\{ \frac{\log \nu}{\alpha i} + \frac{1}{(\alpha i)^2} \right\}.$$

These asymptotic equations imply that the difference between the expressions on the two sides of the symbol \sim has a finite limit as ν tends to infinity.

It would be easy to multiply examples of this type by introducing more logarithms, but enough has been said to indicate the scope of the method.

3. A Different Test for Series which contain Periodic Factors.

Suppose that we wish to discuss the series

$$\sum f(n) v(n),$$

In fact, the real part of the integral can be written in the form

$$\text{const.} + \int_c^{\log \nu} \left\{ \left(a - \frac{\gamma}{\xi} \right) \cos(\alpha \xi - \gamma \log \xi) \right\} \frac{d\xi}{\alpha \xi - \gamma} \quad (\alpha c > \gamma),$$

to which we can apply Dirichlet's test for convergence (*Infinite Series*, Art. 169). Similarly the imaginary part of the integral can be proved to converge.

where $f(x)$ has the same properties as in § 1, but $v(n)$ has the period ω , so that

$$v(n+\omega) = v(n).$$

Then the necessary and sufficient condition for convergence is that

$$\sum_1^{\omega} v(n) = 0.$$

For, suppose that $\Omega = \sum_1^{\omega} v(n),$

then $v(\omega) = \Omega - v(1) - v(2) - \dots - v(\omega - 1).$

Then, since $v(r\omega + s) = v(s),$ we find that

$$\sum_1^{\lambda\omega} f(n) v(n) = \Omega \sum_{r=1}^{\lambda} f(r\omega) + \sum_{s=1}^{\omega-1} v(s) S,$$

where $S = f(s) - f(\omega) + f(s + \omega) - f(2\omega) + \dots + f[s + (\lambda - 1)\omega] - f(\lambda\omega).$

Now, in virtue of the decreasing character of $f(x),$ the sum S has a definite limit $\phi(s)$ less than $f(s),$ as λ tends to infinity, and so

$$\lim_{\lambda \rightarrow \infty} \sum_{s=1}^{\omega-1} v(s) S = \sum_{s=1}^{\omega-1} v(s) \phi(s).$$

But $\sum_{r=1}^{\lambda} f(r\omega) > \int_1^{\lambda} f(\omega\xi) d\xi = \int_1^{\omega\lambda} f(x) dx,$

and so this sum tends to infinity with $\lambda,$ since $\int_1^{\infty} f(x) dx$ is divergent.

It follows that the sum $\sum_1^{\lambda\omega} f(n) v(n)$

also tends to infinity with λ unless Ω is zero; in the latter case, the sum has the finite limit

$$\sum_{s=1}^{\omega-1} v(s) \phi(s).$$

Now, if $0 < \mu < \omega,$ we have

$$\left| \sum_{\lambda\omega+1}^{\lambda\omega+\mu} f(n) v(n) \right| < V_{\mu} f(\lambda\omega),$$

where V denotes the largest of the values $|v(1)|, |v(2)|, \dots, |v(\omega)|.$ Consequently

$$\lim_{\lambda \rightarrow \infty} \sum_{\lambda\omega+1}^{\lambda\omega+\mu} f(n) v(n) = 0,$$

and so the behaviour of the general series

$$\sum_1^{\infty} f(n) v(n)$$

is the same as that of $\lim_{\lambda \rightarrow \infty} \sum_1^{\lambda \omega} f(n) v(n)$.

Thus the series diverges unless Ω is zero, and converges if Ω is zero.

As a simple example, we can infer the convergence of the series

$$\sum f(n) \sin nh \Theta(\sin^2 nh),$$

where $h = p\pi/q$ (p, q being positive integers), and Θ is any continuous function; because here $\omega = 2q$ if p is odd, or $\omega = q$ if p is even, and in either case $\Omega = 0$.

But the series $\sum f(n) \cos nh \Theta(\sin^2 nh)$

can be made to diverge by adjustment of Θ if p is even and q is odd.*

Similarly the series $\sum (-1)^{[nh]} f(n)$,

where $[\xi]$ denotes the integral part of ξ , will diverge if $h = p/q$, where p is even and q is odd,† because again $\omega = q$ and so $\Omega = 1$.

The applications to series of the type

$$\sum f(n) \sin (n^s h), \quad \sum f(n) \cos (n^s h),$$

where s is a positive integer, are equally obvious. Thus, for example, the series

$$\sum f(n) \sin \left(n^2 \frac{2\pi}{q} \right), \quad \sum f(n) \cos \left(n^2 \frac{2\pi}{q} \right)$$

can converge only if

$$\sum_0^{q-1} \sin \left(n^2 \frac{2\pi}{q} \right) = 0, \quad \sum_0^{q-1} \cos \left(n^2 \frac{2\pi}{q} \right) = 0,$$

respectively.

Thus, when q is of the form $4k+1$ the sine-series converges, but the cosine-series diverges; but if q is of the form $4k+3$, the cosine-series converges, while the sine-series is divergent; if q is of the form $4k$, both series diverge, but if q is of the form $4k+2$, both converge. These results follow from the values found by Gauss for $\sum \exp(2\pi i n^2/q)$ in his investigations on quadratic residues.‡

* For instance, if we take $p = 2, q = 3$, the value of Ω is easily seen to be

$$\Theta(0) - \Theta\left(\frac{2}{3}\right),$$

which, of course, may have any value.

† Bromwich and Hardy, *l.c.*, p. 240.

‡ *Werke*, Bd. II., p. 9 (§ 19).