

Moreover $m = 0$ gives $\theta = \infty$, and $m = 1$ gives $\theta =$ its value as defined

by the equation
$$\frac{a^2}{f^2 + \theta} + \frac{b^2}{g^2 + \theta} + \frac{c^2}{h^2 + \theta} = 1,$$

so that, reversing the sign, the limits are ∞, θ ; or, finally, writing under the integral sign ϕ in place of θ , the formula is

Resolved Attraction \div Mass of Ellipsoid

$$= \frac{3}{2} a \int_0^\infty \frac{d\phi}{(f^2 + \phi) \sqrt{(f^2 + \phi)(g^2 + \phi)(h^2 + \phi)}},$$

which is a known formula.

On the Solution of Linear Differential Equations in Series.

By J. HAMMOND.

[Read January 14th, 1875.]

By Leibnitz's theorem,

$$D^m \{ \phi(x) y \} = \left\{ \phi(x) D^m + m \phi'(x) D^{m-1} + \frac{m(m-1)}{2} \phi''(x) D^{m-2} + \dots \right\} y$$

$$= \phi(x+d) D^m y,$$

where $dD^m = mD^{m-1}$, $d^2D^m = m(m-1)D^{m-2}$,

and d operates on D only.

Thus the equation

$$\{ \phi_0(x) D^n + \phi_1(x) D^{n-1} + \dots + \phi_n(x) \} y = \phi(x) \dots \dots \dots (1),$$

when differentiated m times, gives

$$\{ [\phi_0(x+d) D^m] D^n + [\phi_1(x+d) D^m] D^{n-1} + \dots + \phi_n(x+d) D^m \} y = \phi^m(x).$$

Now suppose $y = y_0 + y_1 x + y_2 \frac{x^2}{2} + \dots$;

then y_0, y_1, \dots, y_{n-1} are arbitrary constants, and y_n, y_{n+1}, \dots are found from the equation

$$\{ D^n \phi_0(d) + D^{n-1} \phi_1(d) + \dots + \phi_n(d) \} D^m y_0 = \phi^m(0) \dots \dots \dots (2)$$

by putting $m = 0, 1, 2, \dots$, and $D^k y_0 = y_k$.

Now $\phi(d) = \phi + d\phi' + \frac{d^2}{2} \phi'' + \dots$,

where $\phi, \phi', \phi'' \dots$ are written instead of $\phi(0), \phi'(0), \phi''(0) \dots$ for shortness.

Thus (2) becomes

$$m_0 y_{m+n} + m_1 y_{m+n-1} + \dots + m_{m,n} y_0 = \phi^m \dots \dots \dots (3).$$

The general coefficient in (3) is

$$m_\kappa = \phi_\kappa + m\phi'_{\kappa-1} + \frac{m(m-1)}{|2} \phi''_{\kappa-2} + \dots,$$

$\phi_\kappa, \phi'_\kappa, \phi''_\kappa \dots$ being all zero for all values of κ not included among $\kappa = 0, 1, 2, \dots n$.

Thus, if $p\phi_\kappa = \phi'_{\kappa-1}, p^2\phi_\kappa = \phi''_{\kappa-2} \dots,$

$$m_\kappa = (1+p)^m \phi_\kappa,$$

and (3) becomes

$$\phi_0 y_{m+n} + (1+p)^m (\phi_1 y_{m+n-1} + \phi_2 y_{m+n-2} + \dots + \phi_{m+n} y_0) = \phi^m \dots \dots (4),$$

p operating on ϕ only.

Now write (m, κ) for $(1+p)^m \phi_\kappa$;

and A_m for $(1+p)^m (\phi_{m+1} y_{n-1} + \phi_{m+2} y_{n-2} + \dots + \phi_{m+n} y_0) - \phi^m$;

therefore, from (4),

$$A_m + (m, m) y_n + (m, m-1) y_{n+1} + \dots + (m, 1) y_{m+n-1} + \phi_0 y_{m+n} = 0.$$

This is true for all positive integral values of m . Thus, putting $m = 0, 1, 2, \dots$ in succession,

$$A_0 + \phi_0 y_n = 0,$$

$$A_1 + (1, 1) y_n + \phi_0 y_{n+1} = 0,$$

$$A_2 + (2, 2) y_n + (2, 1) y_{n+1} + \phi_0 y_{n+2} = 0,$$

... ..

$$A_m + (m, m) y_n + (m, m-1) y_{n+1} + \dots + \phi_0 y_{m+n} = 0,$$

solving these equations,

$$\begin{array}{c|c} y_{m+n} & (-1)^{m+1} \\ \hline \begin{array}{l} A_0, \phi_0, 0, 0, \dots\dots 0 \\ A_1, (1, 1), \phi_0, 0, \dots\dots 0 \\ A_2, (2, 2), (2, 1), \phi_0, \dots\dots 0 \\ A_3, (3, 3), (3, 2), (3, 1), \dots\dots 0 \\ \dots \dots \dots \dots \dots \dots \dots \\ A_m, (m, m), (m, m-1), (m, m-2), \dots (m, 1) \end{array} & = \begin{array}{l} \phi_0, 0, 0, \dots\dots 0 \\ (1, 1), \phi_0, 0, \dots\dots 0 \\ (2, 2), (2, 1), \phi_0, \dots\dots 0 \\ \dots \dots \dots \dots \dots \dots \dots \\ (m, m), (m, m-1), (m, m-2), \dots \phi_0 \end{array} \\ & = \left(-\frac{1}{\phi_0}\right)^{m+1}. \end{array}$$

Now $A_0 = \phi_n y_0 + \phi_{n-1} y_1 + \dots + \phi_1 y_{n-1} - \phi,$

$$A_1 = (1, n+1) y_0 + (1, n) y_1 + \dots + (1, 2) y_{n-1} - \phi',$$

$$A_2 = (2, n+2) y_0 + (2, n+1) y_1 + \dots + (2, 3) y_{n-1} - \phi'',$$

... ..

and $(m, \kappa) = \phi_\kappa + m\phi'_{\kappa-1} + \frac{m(m-1)}{|2} \phi''_{\kappa-2} + \dots$

therefore $(1, n+1) = \phi'_n, (2, n+2) = \phi''_n, \dots$

Thus, the expansion of y is

$$\begin{aligned}
 y = & y_0 \left\{ 1 - \frac{1}{\phi_0} \cdot \phi_n \frac{x^n}{n} + \left(\frac{1}{\phi_0}\right)^2 \left| \begin{matrix} \phi_n & \phi_0 \\ \phi'_n(1,1) \end{matrix} \right| \frac{x^{n+1}}{n+1} - \left(\frac{1}{\phi_0}\right)^3 \left| \begin{matrix} \phi_n & \phi_0 & 0 \\ \phi'_n(1,1) & \phi_0 \\ \phi''_n(2,2)(2,1) \end{matrix} \right| \frac{x^{n+2}}{n+2} + \dots \right\} \\
 + & y_1 \left\{ x - \frac{1}{\phi_0} \cdot \phi_{n-1} \frac{x^n}{n} + \left(\frac{1}{\phi_0}\right)^2 \left| \begin{matrix} \phi_{n-1} & \phi_0 \\ (1,n)(1,1) \end{matrix} \right| \frac{x^{n+1}}{n+1} - \left(\frac{1}{\phi_0}\right)^3 \left| \begin{matrix} \phi_{n-1} & \phi_0 & 0 \\ (1,n) & (1,1) & \phi_0 \\ (2,n+1)(2,2)(2,1) \end{matrix} \right| \frac{x^{n+2}}{n+2} + \dots \right\} \\
 + & \dots \dots \dots \dots \dots \dots \dots \dots \dots \dots \dots \dots \dots \\
 + & y_{n-1} \left\{ \frac{x^{n-1}}{n-1} - \frac{1}{\phi_0} \cdot \phi_1 \frac{x^n}{n} + \left(\frac{1}{\phi_0}\right)^2 \left| \begin{matrix} \phi_1 & \phi_0 \\ (1,2)(1,1) \end{matrix} \right| \frac{x^{n+1}}{n+1} - \left(\frac{1}{\phi_0}\right)^3 \left| \begin{matrix} \phi_1 & \phi_0 & 0 \\ (1,2)(1,1) & \phi_0 \\ (2,3)(2,2)(2,1) \end{matrix} \right| \frac{x^{n+2}}{n+2} + \dots \right\} \\
 + & \frac{1}{\phi_0} \cdot \phi \frac{x^n}{n} - \left(\frac{1}{\phi_0}\right)^2 \left| \begin{matrix} \phi & \phi_0 \\ \phi'(1,1) \end{matrix} \right| \frac{x^{n+1}}{n+1} + \left(\frac{1}{\phi_0}\right)^3 \left| \begin{matrix} \phi & \phi_0 & 0 \\ \phi'(1,1) & \phi_0 \\ \phi''(2,2)(2,1) \end{matrix} \right| \frac{x^{n+2}}{n+2} - \dots \dots \dots (5).
 \end{aligned}$$

When $\phi_0 = 0$, this expansion fails.

In this case, put $x = z + h$, where h is not a root of $\phi_0(h) = 0$.

Then (1) becomes $\phi_0(z+h) \frac{d^n y}{dz^n} + \dots = \phi(z+h)$,

and y can be expanded in powers of z .

The equation $\phi(x)y = \psi(x) \dots \dots \dots (6)$, when treated in the same way as (1), gives

$$\frac{\psi(x)}{\phi(x)} = \frac{\psi}{\phi} - \left(\frac{1}{\phi}\right)^2 \left| \begin{matrix} \psi & \phi \\ \psi' & \phi' \end{matrix} \right| x + \left(\frac{1}{\phi}\right)^3 \left| \begin{matrix} \psi & \phi & 0 \\ \psi' & \phi' & \phi \\ \psi'' & \phi'' & 2\phi' \end{matrix} \right| \frac{x^2}{2} - \left(\frac{1}{\phi}\right)^4 \left| \begin{matrix} \psi & \phi & 0 & 0 \\ \psi' & \phi' & \phi & 0 \\ \psi'' & \phi'' & 2\phi' & \phi \\ \psi''' & \phi''' & 3\phi'' & 3\phi' \end{matrix} \right| \frac{x^3}{3} + \dots (7).$$

This may be deduced from (5) by putting $n = 0$,

$$\begin{aligned}
 \phi &= \psi, & \phi' &= \psi' \dots \dots \\
 \phi_0 &= \phi, & \phi'_0 &= \phi' \dots \dots
 \end{aligned}$$

and $(n, \kappa) = \frac{|n|}{|\kappa|} \frac{|n|}{|n-\kappa|} \phi^\kappa$.

Again, putting $\psi(x) = 1$ in (6) and (7),

$$\frac{1}{\phi(x)} = \frac{1}{\phi} - \left(\frac{1}{\phi}\right)^2 \phi' \cdot x + \left(\frac{1}{\phi}\right)^3 \left| \begin{matrix} \phi' & \phi \\ \phi'' & 2\phi' \end{matrix} \right| \frac{x^2}{2} - \left(\frac{1}{\phi}\right)^4 \left| \begin{matrix} \phi' & \phi & 0 \\ \phi'' & 2\phi' & \phi \\ \phi''' & 3\phi'' & 3\phi' \end{matrix} \right| \frac{x^3}{3} + \dots (8).$$

The coefficients of (7) suggest that

$$D^n \left\{ \frac{\psi(x)}{\phi(x)} \right\} = \frac{(-1)^n}{\{\phi(x)\}^{n+1}} \left| \begin{matrix} \psi(x) & \phi(x) & 0 & \dots \\ \psi'(x) & \phi'(x) & \phi(x) & \dots \\ \psi''(x) & \phi''(x) & 2\phi'(x) & \dots \\ \dots & \dots & \dots & \dots \\ \dots & \dots & \dots & \dots \end{matrix} \right| \dots \dots \dots (9),$$

there being $(m + 1)$ rows in the determinant, and the last row being

$$\psi^m(x), \phi^m(x), m\phi^{m-1}(x), \frac{m(m-1)}{2} \phi^{m-2}(x), \dots, m\phi'(x).$$

Here (9) may be proved in exactly the same way, except that x is not made $= 0$ in the course of the work.

There are some particular cases worth notice.

CASE I.—Equation (4) reduces to $\phi_0 y_{m+n} + (1+p)^m y_{m+n-r} = 0$; therefore $(1+p)^m \phi_\kappa = 0$ when κ is not 0 or r .

Putting $m = 0, 1, 2, \dots$ in succession,

$$\left. \begin{aligned} \phi_\kappa &= 0 \\ \phi'_{\kappa-1} &= -\phi_\kappa = 0 \\ \phi''_{\kappa-2} &= \phi_\kappa = 0 \\ \dots &\dots \dots \dots \\ \phi_\kappa &= (-1)^\kappa \phi_\kappa = 0 \end{aligned} \right\} \text{when } \kappa \text{ is not 0 or } r.$$

Thus

$$\begin{aligned} \phi_0(x) &= \phi_0 + \frac{x^r}{r} \phi_0^r = c \left(1 + \alpha_0 \frac{x^r}{r} \right), \\ \phi_1(x) &= \frac{x^{r-1}}{r-1} \phi_1^{r-1} = c\alpha_1 \frac{x^{r-1}}{r-1}, \\ \phi_2(x) &= \frac{x^{r-2}}{r-2} \phi_2^{r-2} = c\alpha_2 \frac{x^{r-2}}{r-2}, \\ &\dots \dots \dots \dots \dots \\ \phi_r(x) &= \text{const.} = c\alpha_r, \\ \phi_{r+1}(x) &= 0, \quad \phi_{r+2}(x) = 0 \dots \dots \end{aligned}$$

Equation (1) then becomes, when $r < n$,

$$\left\{ \left(1 + \alpha_0 \frac{x^r}{r} \right) D^n + \alpha_1 \frac{x^{r-1}}{r-1} D^{n-1} + \dots + \alpha_r D^{n-r} \right\} y = 0 \dots \dots (10);$$

and when $r > n$,

$$\left\{ \left(1 + \alpha_0 \frac{x^r}{r} \right) D^n + \alpha_1 \frac{x^{r-1}}{r-1} D^{n-1} + \dots + \alpha_n \frac{x^{r-n}}{r-n} \right\} y = 0 \dots \dots (11).$$

It is easily seen that the solution of (10) is obtained from that of (11) by solving the equation (11) when n is put $= r$, and then integrating the result $n - r$ times.

Comparing (11) with (1),

$$\begin{aligned} \phi_0 &= 1, \\ \phi_0^r &= \alpha_0, \quad \phi_1^{r-1} = \alpha_1, \quad \phi_2^{r-2} = \alpha_2, \dots \dots \end{aligned}$$

Also $(m, \kappa) = 0$ except when $\kappa = r$. And

$$\begin{aligned} (m, r) &= \phi_r + m\phi'_{r-1} + \frac{m(m-1)}{|2} \phi''_{r-2} + \&c. \\ &= a_r + ma_{r-1} + \frac{m(m-1)}{|2} a_{r-2} + \dots + \frac{m(m-1)\dots(m-r+1)}{|r} a_0. \end{aligned}$$

And since r is to be taken not less than n , $a_r, a_{r-1}, \dots, a_{n+1}$ are all zero; therefore

$$(m, r) = \left\{ a_0 \frac{|m}{|r} \frac{|m-r}{|m-r} + a_1 \frac{|m}{|r-1} \frac{|m-r+1}{|m-r+1} + \dots + a_n \frac{|m}{|r-n} \frac{|m-r+n}{|m-r+n} \right\}.$$

Thus, expanding y by means of the relation

$$\begin{aligned} y_{m+n} &= -(m, r) y_{m+n-r}, \\ y &= y_0 \left\{ 1 - \frac{(r-n, r)}{|r} x^r + \frac{(2r-n, r)(r-n, r)}{|2r} x^{2r} - \dots \right\} \\ &+ y_1 \left\{ x - \frac{(r-n+1, r)}{|r+1} x^{r+1} + \frac{(2r-n+1, r)(r-n+1, r)}{|2r+1} x^{2r+1} - \dots \right\} \\ &+ y_2 \left\{ \frac{x^2}{|2} - \frac{(r-n+2, r)}{|r+2} x^{r+2} + \frac{(2r-n+2, r)(r-n+2, r)}{|2r+2} x^{2r+2} - \dots \right\} \\ &+ \dots \dots \dots \dots \dots \dots \\ &+ y_{n-1} \left\{ \frac{x^{n-1}}{|n-1} - \frac{(r-1, r)}{|n+r-1} x^{n+r-1} + \frac{(2r-1, r)(r-1, r)}{|n+2r-1} x^{n+2r-1} - \dots \right\}. \end{aligned}$$

Many well known expansions are particular cases of the solution of (10) and (11).

Thus

$\{(1+x)D-n\}y = 0$	gives the expansion of $(1+x)^n$,
$\{(1+x)D^2+D\}y = 0$	" " $\log(1+x)$,
$\{(1+x^2)D^2+2xD\}y = 0$	" " $\tan^{-1}x$,
$\{(1-x^2)D^2-xD\}y = 0$	" " $\sin^{-1}x$,
$\{(1-x^2)D^2-xD+m^2\}y = 0$	" " $A \sin(m \sin^{-1}x)$ $+ B \cos(m \sin^{-1}x)$.

CASE II.—The general coefficient of equation (4) is

$$(1+p)^m \phi_\kappa = a_\kappa F(m+n) F(m+n-1) \dots F(m+n-\kappa+1);$$

therefore, putting $m=0$,

$$\phi_\kappa = a_\kappa F(n) F(n-1) \dots F(n-\kappa+1).$$

And (4) reduces to

$$a_0 u_{m+n} + a_1 u_{m+n-1} + \dots + a_n u_m = 0 \dots\dots\dots (12),$$

where
$$u_{m,n} = \frac{y_{m+n}}{\Gamma(m+n)\Gamma(m+n-1)\dots}$$

Here $(1+p)^m = E_n^m$, and $p = \Delta_n$.

Now
$$\phi_{\kappa-r}^r = p^r \phi_\kappa = \Delta_n^r \phi_\kappa;$$

therefore
$$\phi_\kappa^r = E_n^r \Delta_n^r \phi_\kappa.$$

Thus
$$\phi_\kappa(x) = e^{x E_n \Delta_n} \phi_\kappa \dots\dots\dots (13).$$

When $F(n) = 1$, this is the case of linear differential equations with constant coefficients; and the general coefficient in the solution of

$$(a_0 D^n + a_1 D^{n-1} + \dots + a_n) y = 0$$

is the general solution of the difference equation

$$a_0 y_{m+n} + a_1 y_{m+n-1} + \dots + a_n y_m = 0.$$

When $F(n) = n + b - 1$,

$$\phi_\kappa = a_\kappa (n+b-1)(n+b-2)\dots(n+b-\kappa) = a_\kappa \frac{\Gamma(n+b)}{\Gamma(n+b-\kappa)},$$

and $E_n^r \Delta_n^r \phi_\kappa = a_{\kappa+r} (\kappa+r)(\kappa+r-1)\dots(\kappa+1) \frac{\Gamma(n+b)}{\Gamma(n+b-\kappa)};$

therefore, expanding (13) and putting $\kappa=0$,

$$\phi_0(x) = a_0 + a_1 x + a_2 x^2 + \dots + a_n x^n.$$

The general value of $\phi_\kappa(x)$, obtained by expanding (13), is

$$\begin{aligned} \phi_\kappa(x) &= \frac{\Gamma(n+b)}{\Gamma(n+b-\kappa)} \left\{ a_\kappa + a_{\kappa+1} (\kappa+1)x + a_{\kappa+2} \frac{(\kappa+2)(\kappa+1)}{2} x^2 + \dots \right\} \\ &= \frac{\Gamma(n+b)}{\Gamma(n+b-\kappa)} \lfloor \kappa \left\{ a_\kappa \lfloor \kappa + a_{\kappa+1} \lfloor \kappa+1 x + a_{\kappa+2} \frac{\lfloor \kappa+2}{2} x^2 + \dots \right. \\ &\qquad \qquad \qquad \left. \dots + a_n \frac{\lfloor n}{n-\kappa} x^{n-\kappa} \right\} \\ &= \frac{\Gamma(n+b)}{\Gamma(n+b-\kappa)} \lfloor \kappa D^\kappa \phi_0(x). \end{aligned}$$

Thus equation (1) becomes

$$\left\{ \frac{\phi_0(x)}{\Gamma(n+b)} D^n + \frac{\phi_0'(x)}{\Gamma(n+b-1)} D^{n-1} + \frac{\phi_0''(x)}{\Gamma(n+b-2)} \lfloor 2 D^{n-2} + \dots \right. \\ \left. \dots + \frac{\phi_0^n(x)}{\Gamma(b) \lfloor n} \right\} y = 0 \dots\dots\dots (14),$$

where $\phi_0(x)$ is any rational integral function of x of the n th degree. And equation (12) becomes

$$\frac{a_0 y_{m+n}}{\Gamma(m+n+b)} + \frac{a_1 y_{m+n-1}}{\Gamma(m+n+b-1)} + \dots + \frac{a_n y_m}{\Gamma(m+b)} = 0 \dots (15).$$

Thus if α be a single root of the equation $\phi_0\left(\frac{1}{x}\right) = 0$, the corresponding solution of (15) is

$$y_m = \Gamma(m+b) A \alpha^m,$$

and the corresponding expansion of y from (14)

$$y = A \left\{ \Gamma(b) + 1' (b+1) \alpha x + \Gamma(b+2) \frac{\alpha^2 x^2}{2} + \dots \right\} \dots \dots (16).$$

The n arbitrary constants A, \dots are not y_0, y_1, \dots , but are connected with them by linear equations.

When $b=1$, equation (14) reduces to

$$\left\{ \phi_0(x) D^n + n \phi_0'(x) D^{n-1} + \frac{n(n-1)}{2} \phi_0''(x) D^{n-2} + \dots \right\} y = 0,$$

or

$$D^n \{ \phi_0(x) y \} = 0.$$

This gives the expansion for rational fractions; and in the same manner, when b is put $=0$, or any positive or negative integer, the expansion obtained is that of $D^{b-1}(f)$, where f stands for the rational fraction, and $D^{b-1}(f)$ for its $(b-1)^{\text{th}}$ differential coefficient; negative indices of course meaning integrations.

*The Diagonal Scale Principle applied to Angular Measurement
in the Circular Slide Rule.* By JOHN R. CAMPBELL.

[Abstract of Paper, read January 14th, 1875.]

Before entering upon the construction of diagonal scales, having, in place of the usual equidistant parallel lines crossed by a straight diagonal, as many equidistant concentric arcs of circles crossed by a curved diagonal, it will be necessary for me briefly to describe the instrument, or rather the somewhat rough home-made model of one, in which I have introduced them.

It is simply a form of circular slide rule, combining in one arrangement both the ordinary principle of two logometric scales for multiplication and division, and that introduced by the late Dr. Roget (*vide* Phil. Trans., Nov. 17, 1814) for the finding of powers and roots.

Fig. 1 is a plan of the face; fig. 2, a section of the instrument by a vertical plane through the centre.

The face is a circular cardboard surface AA, 12 or 14 inches in diameter, but which might well be made of smaller dimensions. It consists of two parts,—an annular rim AB for the *fixed* scale, and a circular disc BCB corresponding to the *slide*, turning on an axis C in