ON A THEOREM OF CLEBSCH'S.

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CLEBSCH has shown¹ that the integrals of the equations of motion of a vibrating elastic solid can be put into a very simple and elegant form. His demonstration, while by no means long or difficult, seems, nevertheless, to admit of some simplification.

The equations in question are:2

$$\frac{\partial^{2} u}{\partial t^{2}} = a^{2} \Delta u + (b^{2} - a^{2}) \frac{\partial \sigma}{\partial x}$$

$$\frac{d^{2} v}{\partial t^{2}} = a^{2} \Delta v + (b^{2} - a^{2}) \frac{\partial \sigma}{\partial y}$$

$$\frac{\partial^{2} w}{\partial t^{2}} = a^{2} \Delta w + (b^{2} - a^{2}) \frac{\partial \sigma}{\partial z}$$

$$\Delta = \frac{\partial^{2}}{\partial x^{2}} + \frac{\partial^{2}}{\partial y^{2}} + \frac{\partial^{2}}{\partial z^{2}}$$

$$\sigma = \frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z}$$

in which

.

and

From the above equations, it follows without difficulty that

$$\frac{\partial^2 \sigma}{\partial t^2} = b^2 \Delta \sigma.$$
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Let u, v, w be any given solution of the above equations and let us write

$$u = \frac{\partial P}{\partial x} + u'$$

$$v = \frac{\partial P}{\partial y} + v'$$

$$w + \frac{\partial P}{\partial z} + w'$$
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¹ Clebsch, Ueber die Reflexion an einer Kugelfläche. Borchardt's Journal für die reine u. angew Math., Bd. 61.

² Kirchhoff, Vorlesungen über Mechanik, XI Vorlesung.

From these we obtain

$$\sigma = \Delta P + \sigma'$$

in which

$$\sigma' = \frac{\partial u'}{\partial x} + \frac{\partial v'}{\partial y} + \frac{\partial w'}{\partial z} .$$

If we substitute the above values of u, v, w in the equations of motion, we get

$$\frac{\partial^{2}u'}{\partial t^{2}} + \frac{\partial}{\partial x} \left[\frac{\partial^{2}P}{\partial t^{2}} - b^{2}\Delta P \right] = a^{2}\Delta u' + (b^{2} - a^{2}) \frac{\partial \sigma'}{\partial x}
\frac{\partial^{2}v'}{\partial t^{2}} + \frac{\partial}{\partial y} \left[\frac{\partial^{2}P}{\partial t^{2}} - b^{2}\Delta P \right] = a^{2}\Delta v' + (b^{2} - a^{2}) \frac{\partial \sigma'}{\partial y}
\frac{\partial^{2}vv'}{\partial t^{2}} + \frac{\partial}{\partial z} \left[\frac{\partial^{2}P}{\partial t^{2}} - b^{2}\Delta P \right] = a^{2}\Delta v' + (b^{2} - a^{2}) \frac{\partial \sigma'}{\partial z}$$
5.

Suppose that

$$' = 0.$$
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The necessary and sufficient condition for this, is (equation 4) that

$$\Delta P = \sigma$$
.

Suppose moreover that P satisfies the equation

$$\frac{\partial^2 P}{\partial t^2} = b^2 \Delta P.$$
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Equations 5 then become

$$\frac{\partial^2 u'}{\partial t^2} = a^2 \Delta u'$$

$$\frac{\partial^2 v'}{\partial t^2} = a^2 \Delta v'$$

$$\frac{\partial^2 w'}{\partial t^2} = a^2 \Delta w'$$

If, then, it be possible to choose P so that equations 7 and 8 are simultaneously verified, equations 3 will determine u', v', w'. We shall thus have established that every solution of the given equations can be written in the form of equations 3, P being a common integral of equations 7 and 8, and u', v' w' being three solutions of the equation

$$\frac{\partial^2 \varphi}{\partial t^2} = a^2 \Delta \varphi, \qquad 10$$

these solutions, moreover, being connected by equation 6:

$$\frac{\partial u'}{\partial x} + \frac{\partial v'}{\partial y} + \frac{\partial w'}{\partial z} = 0.$$

In order to show that we can determine P so as to satisfy equations 7 and 8, let us write the equation obtained from them:

$$\frac{\partial^2 P}{\partial t^2} = b^2 \sigma. {12}$$

The pair of equations 7 and 8 may be replaced by the pair of equivalent equations 7 and 12. The integral of equation 12 is

$$P = b^2 \int dt \int \sigma dt + A(x, y, z) \cdot t + B(x, y, z),$$

A and B being two arbitrary functions. From this equation we obtain

$$\Delta P = b^2 \int dt \int \Delta \sigma dt + \Delta A(x, y, z) \cdot t + \Delta B(z, y, z).$$

Choosing the arbitrary functions so that

$$\Delta A(x, y, z) = 0, \quad \Delta B(x, y, z) = 0$$

and making use of equation 2, we obtain

$$\Delta P = \int dt \int \frac{\partial^2 \sigma}{\partial t^2} dt = \sigma.$$

Thus equation 7 is also verified and the theorem stated above is established.

This result may be put into a more striking form if we make use of a theorem tacitly admitted by Clebsch and rigorously established by Duhem.¹ The theorem in question is the following: If u', v', w' are three integrals of equation 10, which also verify equation 11, we may write

$$u' = \frac{\partial W}{\partial y} - \frac{\partial V}{\partial z},$$

$$v' = \frac{\partial U}{\partial z} - \frac{\partial W}{\partial x},$$

$$w' = \frac{\partial V}{\partial x} - \frac{\partial U}{\partial y},$$

in which U, V, W are three solutions of equation 10.

¹ Duhem, Sur l'intégrale des équations des petits mouvements d'un solide isotrope. Mémoires de la Société des Sciences de Bordeaux. 5° Série, t. III., p. 325, 1899.

Thus finally we obtain Clebsch's theorem: Every solution of the equations of motion of a vibrating elastic solid can be written in the form:

$$u = \frac{\partial P}{\partial x} + \frac{\partial W}{\partial y} - \frac{\partial V}{\partial z}$$

$$v = \frac{\partial P}{\partial y} + \frac{\partial U}{\partial z} - \frac{\partial W}{\partial x}$$

$$w = \frac{\partial P}{\partial z} + \frac{\partial V}{\partial x} - \frac{\partial U}{\partial y}$$
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U, V, W being three solutions of the equation

$$\frac{\partial^2 \varphi}{\partial t^2} = a^2 \, \varDelta \, \varphi \tag{14}$$

and P being a solution common to the equations

$$\frac{\partial^2 P}{\partial t^2} = b^2 \Delta P$$

$$\Delta P = \sigma.$$
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Conversely, if U, V, W are any three solutions of equation 14, and P any solution of equation 15, equations 13 furnish a solution of the equations of motion. This is very easily verified by substituting in the original equations the values of u, v, w given by equations 13.

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