

NOTE ON THE EVALUATION OF A CERTAIN INTEGRAL  
CONTAINING BESSEL'S FUNCTIONS

By H. M. MACDONALD.

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A PARTICULAR case of the following integral occurs in Prof. H. Lamb's paper "On the Theory of Waves propagated vertically in the Atmosphere" (p. 136 *supra*):—

$$\int_0^{\infty} J_n(ax) J_n(bx) J_m(cx) x^{1-m} dx.$$

Writing  $W = \int_0^{\infty} J_n(ax) J_n(bx) J_m(cx) x^{1-m} dx$   
and using the relation

$$J_m(cx) = \frac{1}{2\pi i} c^m x^m \int_{\gamma-\infty i}^{\gamma+\infty i} e^{xt-c^2x^2/2t} t^{-m-1} dt,$$

where  $\gamma$  is a real positive quantity and  $m > -1$ , it follows that

$$W = \frac{1}{2\pi i} c^m \int_0^{\infty} \int_{\gamma-\infty i}^{\gamma+\infty i} e^{xt-c^2x^2/2t} J_n(ax) J_n(bx) x t^{-m-1} dx dt,$$

whence,  $n > -1$ ,

$$W = \frac{1}{2\pi i} c^{m-2} \int_{\gamma-\infty i}^{\gamma+\infty i} e^{xt-(a^2+b^2)t/2c^2} I_n(abt/c^2) t^{-m} dt,$$

or  $W = \frac{1}{2\pi i} a^{m-1} b^{m-1} c^{-m} \int_{\gamma-\infty i}^{\gamma+\infty i} e^{(c^2-a^2-b^2)t/2c^2} I_n(t) t^{-m} dt.$

When  $c^2 - a^2 - b^2 > 2ab$ , writing  $c^2 - a^2 - b^2 = 2ab\mu$ ,

$$\frac{1}{2\pi i} \int_{\gamma-\infty i}^{\gamma+\infty i} e^{\mu t} I_n(t) t^{-m} dt = \sqrt{\frac{2}{\pi}} \frac{\sin(m-n)\pi}{\pi} e^{(m-\frac{1}{2})\pi} (\mu^2 - 1)^{\frac{1}{2}(2m-1)} Q_{n-\frac{1}{2}}^{\frac{1}{2}-m}(\mu)^* ;$$

therefore, when  $c^2 > (a+b)^2$ ,  $m > -1$ ,  $n > -1$ ,

$$W = a^{m-1} b^{m-1} c^{-m} \sqrt{\frac{2}{\pi}} \frac{\sin(m-n)\pi}{\pi} e^{(m-\frac{1}{2})\pi} (\mu^2 - 1)^{\frac{1}{2}(2m-1)} Q_{n-\frac{1}{2}}^{\frac{1}{2}-m}(\mu),$$

where

$$\mu = (c^2 - a^2 - b^2)/2ab.$$

\* Macdonald, *Proc. London Math. Soc.*, Vol. xxxv., p. 436.

It remains to evaluate,  $W$  when  $(a+b)^2 > c^2$ , for this purpose using the relation

$$I_n(t) = \frac{1}{\sqrt{(2\pi)}} t \int_{-1}^1 e^{t\nu} (1-\nu^2)^{\frac{1}{2}(2l-1)} P_{n-\frac{1}{2}}^{\frac{1}{2}-l}(\nu) d\nu, *$$

where  $n-l$  is an integer greater than  $-1$ .

$$W = \frac{1}{2\pi i} a^{m-1} b^{m-1} c^{-m} \times \frac{1}{\sqrt{(2\pi)}} \int_{\gamma-\infty}^{\gamma+\infty} \int_{-1}^1 e^{(c^2-a^2-b^2+2ab\nu)t/2ab} (1-\nu^2)^{\frac{1}{2}(2l-1)} P_{n-\frac{1}{2}}^{\frac{1}{2}-l}(\nu) t^{-m} d\nu dt,$$

where  $l$  is chosen less than  $m$ . Now

$$\frac{1}{2\pi i} \int_{\gamma-\infty}^{\gamma+\infty} e^{\kappa t} t^{-m} dt \quad [m > l]$$

vanishes when  $\kappa$  is negative, and has the value  $\kappa^{m-l-1}/\Gamma(m-l-1)$  when  $\kappa$  is positive; therefore, when  $(a-b)^2 > c^2$ ,  $W = 0$  and, when

$$(a+b)^2 > c^2 > (a-b)^2,$$

$$W = a^{m-1} b^{m-1} c^{-m} \frac{1}{\sqrt{(2\pi)}} \frac{1}{\Gamma(m-l-1)} \int_{\mu}^1 (\nu-\mu)^{m-l-1} (1-\nu^2)^{\frac{1}{2}(2l-1)} P_{n-\frac{1}{2}}^{\frac{1}{2}-l}(\nu) d\nu,$$

where  $\mu = (a^2 + b^2 - c^2)/2ab$ , and, using the relation

$$P_{n-\frac{1}{2}}^{\frac{1}{2}-l}(\nu) = \frac{1}{\Gamma(l-\frac{1}{2})} \left(\frac{1+\nu}{1-\nu}\right)^{\frac{1}{2}(1-2l)} F[-n+\frac{1}{2}, n+\frac{1}{2}, l+\frac{1}{2}, \frac{1}{2}(1-\nu)],$$

$$W = a^{m-1} b^{m-1} c^{-m} \frac{1}{\sqrt{(2\pi)}} \frac{1}{\Gamma(m-l-1) \Gamma(l-\frac{1}{2})} \times \int_{\mu}^1 (\nu-\mu)^{m-l-1} (1-\nu)^{l-\frac{1}{2}} F[-n+\frac{1}{2}, n+\frac{1}{2}, l+\frac{1}{2}, \frac{1}{2}(1-\nu)] d\nu,$$

that is,

$$W = a^{m-1} b^{m-1} c^{-m} \frac{1}{\sqrt{(2\pi)}} \frac{1}{\Gamma(m-l-1) \Gamma(-n-\frac{1}{2}) \Gamma(n-\frac{1}{2})} \times \sum_0^{\infty} \frac{\Gamma(r-n-\frac{1}{2}) \Gamma(r+n-\frac{1}{2})}{\Gamma(l+r-\frac{1}{2}) \Gamma(r)} 2^{-r} \int_{\mu}^1 (\nu-\mu)^{m-l-1} (1-\nu)^{l+r-\frac{1}{2}} d\nu.$$

Making the substitution  $\nu-\mu = (1-\mu)\xi$ ,

$$\int_{\mu}^1 (\nu-\mu)^{m-l-1} (1-\nu)^{l+r-\frac{1}{2}} d\nu = (1-\mu)^{m+r-\frac{1}{2}} \int_0^1 \xi^{m-l-1} (1-\xi)^{r+l-\frac{1}{2}} d\xi,$$

that is, the value of this integral is

$$(1-\mu)^{m+r-\frac{1}{2}} \frac{\Gamma(m-l-1) \Gamma(l+r-\frac{1}{2})}{\Gamma(m+r-\frac{1}{2})};$$

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\* *Proc. London Math. Soc.*, Vol. xxxv., p. 431.

therefore

$$W = a^{m-1} b^{m-1} c^{-m} \times \frac{1}{\sqrt{(2\pi)}} \frac{1}{\Gamma(n-\frac{1}{2}) \Gamma(-n-\frac{1}{2})} \sum_0^\infty \frac{\Gamma(r+n-\frac{1}{2}) \Gamma(r-n-\frac{1}{2})}{\Gamma(r+m-\frac{1}{2}) \Gamma(r)} 2^{-r} (1-\mu)^{m+r-\frac{1}{2}}$$

or  $W = a^{m-1} b^{m-1} c^{-m} \frac{1}{\sqrt{(2\pi)}} \frac{(1-\mu)^{m-\frac{1}{2}}}{\Gamma(m-\frac{1}{2})} F[-n+\frac{1}{2}, n+\frac{1}{2}, m+\frac{1}{2}, \frac{1}{2}(1-\mu)],$

that is,  $W = \frac{1}{\sqrt{(2\pi)}} a^{m-1} b^{m-1} c^{-m} (1-\mu^2)^{\frac{1}{2}(2m-1)} P_{n-\frac{1}{2}}^{\frac{1}{2}-m}(\mu).$

Collecting the results,

$W = 0 \quad [(a-b)^2 > c^2],$

$$W = \frac{1}{\sqrt{(2\pi)}} a^{m-1} b^{m-1} c^{-m} (1-\mu^2)^{\frac{1}{2}(2m-1)} P_{n-\frac{1}{2}}^{\frac{1}{2}-m}(\mu),$$

where

$$\mu = (a^2 + b^2 - c^2)/2ab \quad [(a+b)^2 > c^2 > (a-b)^2],$$

$$W = a^{m-1} b^{m-1} c^{-m} \sqrt{\frac{2}{\pi}} \frac{\sin(m-n)\pi}{\pi} e^{(m-\frac{1}{2})\pi i} (\mu^2 - 1)^{\frac{1}{2}(2m-1)} Q_{n-\frac{1}{2}}^{\frac{1}{2}-m}(\mu),$$

where

$$\mu = (c^2 - a^2 - b^2)/2ab \quad [c^2 > (a+b)^2],$$

and

$$n > -1, \quad m > -\frac{1}{2}.$$

When  $n-m$  is an integer greater than  $-1$  these results become

$W = 0 \quad [(a-b)^2 > c^2],$

$$W = \frac{1}{\sqrt{(2\pi)}} a^{m-1} b^{m-1} c^{-m} (1-\mu^2)^{\frac{1}{2}(2m-1)} P_{n-\frac{1}{2}}^{\frac{1}{2}-m}(\mu) \quad [(a+b)^2 > c^2 > (a-b)^2],$$

$W = 0 \quad [c^2 > (a+b)^2],$

where  $\mu = (a^2 + b^2 - c^2)/2ab$ , and in this case  $P_{n-\frac{1}{2}}^{\frac{1}{2}-m}(\mu)$  can be expressed in finite terms. The result in the particular case when  $n = m$  is, remembering that

$$P_{n-\frac{1}{2}}^{\frac{1}{2}-n}(\mu) = \frac{2^{\frac{1}{2}-n}}{\Gamma(n-\frac{1}{2})} (1-\mu^2)^{\frac{1}{2}(2n-1)},$$

$W = 0 \quad [(a-b)^2 > c^2],$

$$W = \frac{a^{-m} b^{-m} c^{-m}}{2^{2m-1} \sqrt{\pi} \Gamma(m-\frac{1}{2})} \{ (a+b+c)(a+b-c)(a-b+c)(-a+b+c) \}^{m-\frac{1}{2}} \quad [(a+b)^2 > c^2 > (a-b)^2],$$

$W = 0 \quad [c^2 > (a+b)^2];$

this particular result has been given by Sonine.\*

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\* *Math. Ann.*, Bd. xvi.

When  $m = \frac{1}{2}$ , remembering that

$$J_{\frac{1}{2}}(cx) = \left(\frac{2}{\pi cx}\right)^{\frac{1}{2}} \sin cx$$

and writing 
$$V = \int_0^{\infty} J_n(ax) J_n(bx) \sin cx \, dx,$$

the above results become

$$V = 0 \quad [(a-b)^2 > c^2],$$

$$V = \frac{1}{2} a^{-\frac{1}{2}} b^{-\frac{1}{2}} P_{n-\frac{1}{2}}(\mu) \quad [(a+b)^2 > c^2 > (a-b)^2],$$

where

$$\mu = (a^2 + b^2 - c^2)/2ab,$$

$$V = \frac{\cos n\pi}{\pi} a^{-\frac{1}{2}} b^{-\frac{1}{2}} Q_{n-\frac{1}{2}}(\mu) \quad [c^2 > (a+b)^2],$$

where

$$\mu = (c^2 - a^2 - b^2)/2ab.$$

When  $n$  is half a positive odd integer these results become

$$V = 0 \quad [(a+b)^2 > c^2],$$

$$V = \frac{1}{2} a^{-\frac{1}{2}} b^{-\frac{1}{2}} P_{n-\frac{1}{2}}(\mu) \quad [(a+b)^2 > c^2 > (a-b)^2].$$

$$V = 0 \quad [c^2 > (a+b)^2].$$

The various known series for the spherical harmonics enable the values of  $W$  or  $V$  to be calculated to any degree of accuracy; the values in the neighbourhood of the discontinuities will be given below.

When  $c^2 - (a-b)^2$  is small and tends to zero, it follows from the series for  $P_{n-\frac{1}{2}}^{\frac{1}{2}-m}(\mu)$  in powers of  $(1-\mu)$  that  $W$  tends to

$$\frac{1}{\sqrt{(2\pi)}} a^{m-1} b^{m-1} c^{-m} \frac{(1-\mu)^{m-\frac{1}{2}}}{\Gamma(m-\frac{1}{2})},$$

that is, to zero when  $m > \frac{1}{2}$ ; also, that  $V$  tends to the value  $\frac{1}{2} a^{-\frac{1}{2}} b^{-\frac{1}{2}}$  when  $c^2 - (a-b)^2$  tends to zero. When  $(a+b)^2 - c^2$  is small  $\mu$  is nearly equal to  $-1$ , and it is convenient to replace  $P_{n-\frac{1}{2}}^{\frac{1}{2}-m}(\mu)$  by harmonics of  $-\mu$ . If  $m - \frac{1}{2}$  is not an integer,

$$P_{n-\frac{1}{2}}^{\frac{1}{2}-m}(\mu) = \frac{\cos n\pi}{\cos m\pi} P_{n-\frac{1}{2}}^{\frac{1}{2}-m}(-\mu) + \frac{\sin(n-m)\pi}{\cos m\pi} \frac{\Gamma(n-m)}{\Gamma(n+m-1)} P_{n-\frac{1}{2}}^{m-\frac{1}{2}}(-\mu),$$

and therefore, when  $n-m$  is not an integer,  $W$  tends to the value

$$\frac{1}{\sqrt{\pi}} 2^{m-1} a^{m-1} b^{m-1} c^{-m} \frac{\sin(n-m)\pi}{\cos m\pi} \frac{\Gamma(n-m)}{\Gamma(n+m-1) \Gamma(\frac{1}{2}-m)},$$

that is, to the value

$$-\frac{1}{\pi \sqrt{\pi}} 2^{m-1} a^{m-1} b^{m-1} c^{-m} \sin(n-m)\pi \frac{\Gamma(n-m) \Gamma(m-\frac{3}{2})}{\Gamma(n+m-1)}.$$

If  $m - \frac{1}{2}$  is an integer,

$$P_{n-\frac{1}{2}}^{\frac{1}{2}-m}(\mu) = \cos(n-m)\pi P_{n-\frac{1}{2}}^{\frac{1}{2}-m}(-\mu) - \frac{2}{\pi} \sin(n-m)\pi Q_{n-\frac{1}{2}}^{\frac{1}{2}-m}(-\mu),$$

and, making use of the appropriate expression for  $Q_{n-\frac{1}{2}}^{\frac{1}{2}-m}(-\mu)$ ,\* it follows that, when  $n-m$  is not an integer,  $W$  tends to the value

$$-\frac{1}{\pi\sqrt{\pi}} 2^{m-1} a^{m-1} b^{m-1} c^{-m} \sin(n-m)\pi \frac{\Pi(n-m)\Pi(m-\frac{3}{2})}{\Pi(n+m-1)},$$

when  $m > \frac{1}{2}$ .

If  $m = \frac{1}{2}$ , and  $n - \frac{1}{2}$  is not an integer,  $V$  tends to the value

$$-\frac{1}{2\pi} \cos n\pi a^{-\frac{1}{2}} b^{-\frac{1}{2}} \log \{(a+b)^2 - c^2\},$$

when  $(a+b)^2 - c^2$  tends to zero. If  $n-m$  is an integer and  $m > \frac{1}{2}$ ,  $W$  tends to zero, when  $(a+b)^2 - c^2$  tends to zero. If  $m = \frac{1}{2}$  and  $n - \frac{1}{2}$  is an integer,  $V$  tends to the value  $\frac{1}{2} \sin n\pi a^{-\frac{1}{2}} b^{-\frac{1}{2}}$ , when  $(a+b)^2 - c^2$  tends to zero.

When  $c^2 > (a+b)^2$  the expression for  $W$  can be replaced by

$$W = \frac{1}{\pi} a^{m-1} b^{m-1} c^{-m} 2^{1-m} \sin(m-n)\pi \times \frac{1}{\Pi(m-1)} \int_{\eta}^{\infty} e^{-n\xi} (2 \cosh \xi - 2 \cosh \eta)^{m-1} d\xi,$$

where  $2ab \cosh \eta = c^2 - a^2 - b^2,$

and  $m$  is positive.

When  $\eta$  tends to zero, this expression, writing  $e^{-\xi} = x$ , tends to

$$\frac{1}{\pi} a^{m-1} b^{m-1} c^{-m} 2^{1-m} \sin(n-m)\pi \frac{1}{\Pi(m-1)} \int_0^1 x^{n-m} (1-x)^{2m-2} dx,$$

provided  $m > \frac{1}{2}$ . Hence, if  $m > \frac{1}{2}$ ,  $W$  tends to the value

$$\frac{1}{\pi\sqrt{\pi}} 2^{m-1} a^{m-1} b^{m-1} c^{-m} \sin(m-n)\pi \frac{\Pi(n-m)\Pi(m-\frac{3}{2})}{\Pi(n+m-1)},$$

when  $c^2 - (a+b)^2$  tends to zero. When  $m = \frac{1}{2}$ ,

$$V = \frac{1}{\pi} \cos n\pi a^{-\frac{1}{2}} b^{-\frac{1}{2}} \int_{\eta}^{\infty} e^{-n\xi} (2 \cosh \xi - 2 \cosh \eta)^{-\frac{1}{2}} d\xi,$$

and the limit to which

$$\int_{\eta}^{\infty} e^{-n\xi} (2 \cosh \xi - 2 \cosh \eta)^{-\frac{1}{2}} d\xi$$

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\* Macdonald, *Proc London Math. Soc.*, Vol. xxxi., p. 276.

tends when  $\eta$  tends to zero is  $-\log \eta$ ; therefore  $V$  tends to

$$-\frac{1}{2\pi} a^{-\frac{1}{2}} b^{-\frac{1}{2}} \cos n\pi \log \{c^2 - (a+b)^2\},$$

when  $c^2 - (a+b)^2$  tends to zero.

[*Added January 19th, 1909.*]

The following integral is closely related to the integral investigated above, but differs from it inasmuch as it is continuous for all real values of  $a, b, c$ . Writing

$$U = \int_0^\infty J_n(ax) J_n(bx) K_m(cx) x^{m+1} dx,$$

and making use of the relation

$$K_m(cx) = \frac{1}{2} \int_0^\infty e^{-\frac{1}{2}cx(s+s^{-1})} s^{m-1} ds,$$

which is equivalent to

$$K_m(cx) = \frac{1}{2} (cx)^{-m} \int_0^\infty e^{-\frac{1}{2}s - \frac{1}{2}c^2x^2s^{-1}} s^{m-1} ds,$$

it follows that

$$U = \frac{1}{2} c^{-m} \int_0^\infty \int_0^\infty e^{-\frac{1}{2}s - \frac{1}{2}c^2x^2s^{-1}} J_n(ax) J_n(bx) s^{m-1} x dx ds;$$

hence, changing the order of integration and using the relation

$$\int_0^\infty e^{-\frac{1}{2}c^2x^2s^{-1}} J_n(ax) J_n(bx) x dx = \frac{s}{c^2} e^{-(a^2+b^2)s/2c^2} I_n\left(\frac{abs}{c^2}\right),$$

$$U = \frac{1}{2} c^{-m-2} \int_0^\infty e^{-\frac{1}{2}s - \frac{1}{2}(a^2+b^2)sc^{-2}} I_n\left(\frac{abs}{c^2}\right) s^m ds,$$

when  $n > -\frac{1}{2}$ .

Writing  $a^2 + b^2 + c^2 = 2ab\mu, \quad abs = c^2t,$

this becomes 
$$U = \frac{1}{2} c^{2n} (ab)^{-n-1} \int_0^\infty e^{-\mu t} I_n(t) t^n dt,$$

that is, 
$$U = \frac{1}{2} c^{2n} (ab)^{-n-1} \sum_0^\infty \frac{1}{2^{n+2k} \Pi(n+k) \Pi(k)} \int_0^\infty e^{-\mu t} t^{n+m+2k} dt;$$

hence 
$$U = \frac{1}{2} c^{2n} (ab)^{-n-1} \sum_0^\infty \frac{\Pi(n+m+2k)}{2^{n+2k} \Pi(n+k) \Pi(k)} \frac{1}{\mu^{n+m+2k+1}},$$

provided that  $n+m > -1$ , that is,

$$U = \frac{1}{2} c^{2n} (ab)^{-n-1} \sum_0^\infty \frac{2^{n+m+2k} \Pi\left(\frac{n+m}{2} + k\right) \Pi\left(\frac{n+m-1}{2} + k\right)}{2^{n+2k} \Pi(n+k) \Pi(k) \Pi\left(-\frac{1}{2}\right) \mu^{n+m+2k+1}};$$

and therefore

$$U = 2^{m-1} \pi^{-\frac{1}{2}} c^m (ab)^{-m-1} \frac{\Pi\left(\frac{n+m}{2}\right) \Pi\left(\frac{n+m-1}{2}\right)}{\Pi(n)} \frac{1}{\mu^{n+m+1}} \\ \times F\left(\frac{n+m+2}{2}, \frac{n+m+1}{2}, n+1, \frac{1}{\mu^2}\right).$$

Now  $Q_{n-\frac{1}{2}}^{m+\frac{1}{2}}(\mu) = e^{(m+\frac{1}{2})\pi i} \frac{\Pi(n+m) \Pi(-\frac{1}{2})}{2^{n+\frac{1}{2}} \Pi(n)} \frac{(\mu^2-1)^{\frac{1}{2}(2m+1)}}{\mu^{n+m+1}} \\ \times F\left(\frac{n+m+2}{2}, \frac{n+m+1}{2}, n+1, \frac{1}{\mu^2}\right);$

therefore

$$U = 2^{n+m-\frac{1}{2}} \pi^{-\frac{1}{2}} c^m (ab)^{-m-1} e^{-(m+\frac{1}{2})\pi i} \frac{\Pi\left(\frac{n+m}{2}\right) \Pi\left(\frac{n+m-1}{2}\right)}{\Pi(n+m) \Pi(-\frac{1}{2})} \\ \times (\mu^2-1)^{-\frac{1}{2}(2m+1)} Q_{n-\frac{1}{2}}^{m+\frac{1}{2}}(\mu);$$

whence  $U = 2^{-\frac{1}{2}} \pi^{-\frac{1}{2}} e^{-(m+\frac{1}{2})\pi i} c^m (ab)^{-m-1} (\mu^2-1)^{-\frac{1}{2}(2m+1)} Q_{n-\frac{1}{2}}^{m+\frac{1}{2}}(\mu);$

and therefore

$$\int_0^\infty J_n(ax) J_n(bx) K_m(cx) x^{m+1} dx \\ = 2^{-\frac{1}{2}} \pi^{-\frac{1}{2}} e^{-(m+\frac{1}{2})\pi i} c^m (ab)^{-m-1} (\mu^2-1)^{-\frac{1}{2}(2m+1)} Q_{n-\frac{1}{2}}^{m+\frac{1}{2}}(\mu),$$

where  $2ab\mu = a^2 + b^2 + c^2 \quad (n > -\frac{1}{2}, n+m > -1).$

Remembering that  $K_{-m} = K_m$ , it follows that

$$\int_0^\infty J_n(ax) J_n(bx) K_m(cx) x^{1-m} dx \\ = 2^{-\frac{1}{2}} \pi^{-\frac{1}{2}} e^{(m-\frac{1}{2})\pi i} c^{-m} (ab)^{m-1} (\mu^2-1)^{\frac{1}{2}(2m-1)} Q_{n-\frac{1}{2}}^{\frac{1}{2}-m}(\mu),$$

where  $n > -\frac{1}{2}, \quad m < n+1.$

Substituting  $m = -\frac{1}{2}$  in the first result, and writing for  $K_{-\frac{1}{2}}(cx)x^{\frac{1}{2}}$ , its value  $2^{-\frac{1}{2}} \pi^{\frac{1}{2}} c^{-\frac{1}{2}} e^{-cx}$ , it becomes

$$\int_0^\infty e^{-cx} J_n(ax) J_n(bx) dx = \frac{1}{\pi} a^{-\frac{1}{2}} b^{-\frac{1}{2}} Q_{n-\frac{1}{2}}(\mu),$$

a result previously obtained by the writer.\*

\* *Proc. London Math. Soc.*, Vol. xxvi., pp. 160, 161, 165.

Again, substituting  $m = 0$ , the first result becomes

$$\int_0^\infty J_n(ax) J_n(bx) K_0(cx) x dx = 2^{-\frac{1}{2}} \pi^{-\frac{1}{2}} e^{-\frac{1}{2} \pi i} (ab)^{-1} (\mu^2 - 1)^{-\frac{1}{2}} Q_{n-\frac{1}{2}}^{\frac{1}{2}}(\mu),$$

that is, 
$$\int_0^\infty J_n(ax) J_n(bx) K_0(cx) x dx = \frac{1}{2ab \sinh \psi} e^{-n\psi},$$

where 
$$\cosh \psi = \mu = (a^2 + b^2 + c^2) / 2ab.$$

Substituting in the first result  $m = n$ , it becomes

$$\begin{aligned} \int_0^\infty J_n(ax) J_n(bx) K_n(cx) x^{n+1} dx \\ = 2^{-\frac{1}{2}} \pi^{-\frac{1}{2}} c^n (ab)^{-n-1} e^{-(n+\frac{1}{2}) \pi i} (\mu^2 - 1)^{-\frac{1}{2}(2n+1)} Q_{n-\frac{1}{2}}^{n+\frac{1}{2}}(\mu), \end{aligned}$$

that is,

$$\begin{aligned} \int_0^\infty J_n(ax) J_n(bx) K_n(cx) x^{n+1} dx \\ = 2^{-\frac{1}{2}} \pi^{-\frac{1}{2}} c^n (ab)^{-n-1} (\mu^2 - 1)^{-\frac{1}{2}(2n+1)} \frac{\Pi(2n) \Pi(-\frac{1}{2})}{2^{n+\frac{1}{2}} \Pi(n)} (\mu^2 - 1)^{-\frac{1}{2}(2n+1)}; \end{aligned}$$

and therefore

$$\begin{aligned} \int_0^\infty J_n(ax) J_n(bx) K_n(cx) x^{n+1} dx \\ = 2^{n-1} \pi^{-\frac{1}{2}} \Pi(n - \frac{1}{2}) c^n (ab)^{-n-1} (\mu^2 - 1)^{-n-\frac{1}{2}} \quad (n > -\frac{1}{2}). \end{aligned}$$