

(eV) sterile neutrinos: the global picture

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- I. Oscillation anomalies: ν_e disappearance
 - II. Oscillation anomalies: $\nu_\mu \rightarrow \nu_e$ appearance
 - III. Sterile neutrino models and ν_μ disappearance

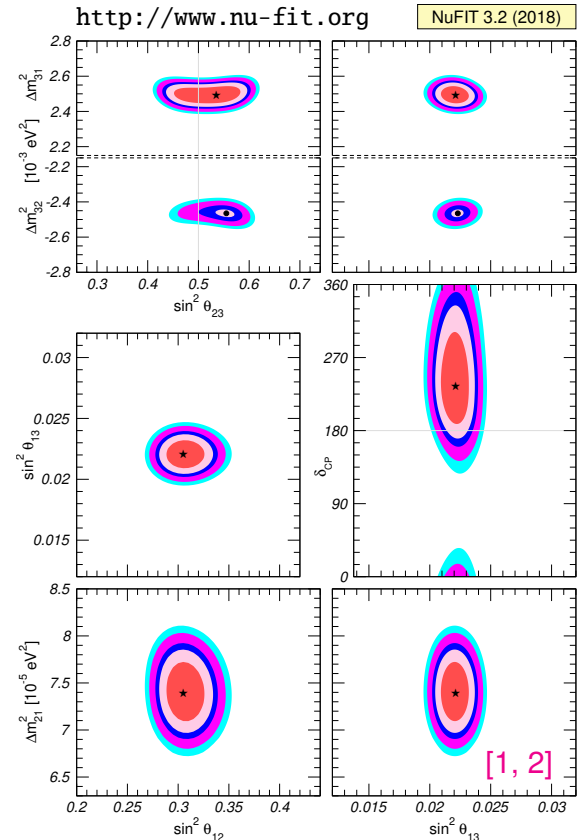
3ν oscillations (until last week)

- Global 6-parameter fit (including δ_{CP}):
 - Solar**: Cl + Ga + SK(1–4) + SNO-full (I+II+III) + Borexino;
 - Atmospheric**: DeepCore;
 - Reactor**: KamLAND + Double-Chooz + Daya-Bay + Reno;
 - Accelerator**: Minos + T2K + NOvA;
- best-fit point and 1σ (3σ) ranges:

$$\begin{aligned} \theta_{12} &= 33.62^{+0.78}_{-0.76} \left({}^{+2.44}_{-2.20} \right), & \Delta m_{21}^2 &= 7.40^{+0.21}_{-0.20} \left({}^{+0.62}_{-0.60} \right) \times 10^{-5} \text{ eV}^2, \\ \theta_{23} &= \begin{cases} 47.2^{+1.9}_{-3.9} \left({}^{+4.4}_{-6.9} \right), \\ 48.1^{+1.4}_{-1.9} \left({}^{+3.6}_{-6.8} \right), \end{cases} & \Delta m_{31}^2 &= \begin{cases} +2.494^{+0.033}_{-0.031} \left({}^{+0.099}_{-0.095} \right) \times 10^{-3} \text{ eV}^2, \\ -2.465^{+0.032}_{-0.031} \left({}^{+0.096}_{-0.097} \right) \times 10^{-3} \text{ eV}^2, \end{cases} \\ \theta_{13} &= 8.54^{+0.15}_{-0.15} \left({}^{+0.44}_{-0.45} \right), & \delta_{\text{CP}} &= 234^{+43}_{-31} \left({}^{+141}_{-90} \right); \end{aligned}$$

- neutrino mixing matrix:

$$|U|_{3\sigma} = \begin{pmatrix} 0.799 \rightarrow 0.844 & 0.516 \rightarrow 0.582 & 0.141 \rightarrow 0.156 \\ 0.242 \rightarrow 0.494 & 0.467 \rightarrow 0.678 & 0.639 \rightarrow 0.774 \\ 0.284 \rightarrow 0.521 & 0.490 \rightarrow 0.695 & 0.615 \rightarrow 0.754 \end{pmatrix}.$$

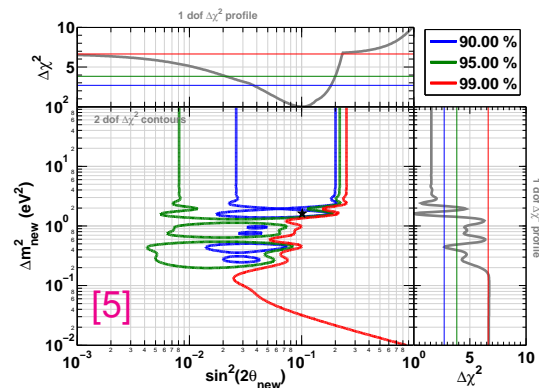
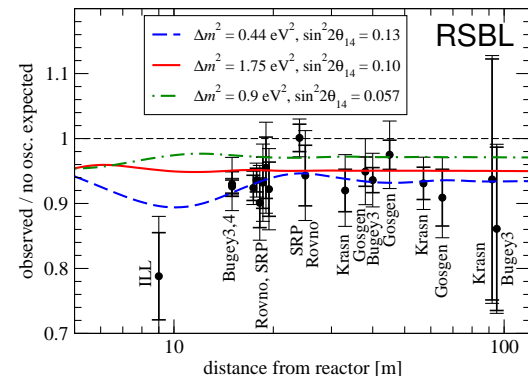


[1] I. Esteban *et al.*, JHEP **01** (2017) 087 [arXiv:1611.01514].

[2] I. Esteban *et al.*, NuFIT 3.2 (2018), <http://www.nu-fit.org>.

$\bar{\nu}_e$ disappearance: the reactor anomaly

- In [3, 4] the reactor $\bar{\nu}$ fluxes were reevaluated;
 - the new calculations result in a small increase of the flux by about **3.5%**;
 - hence, **all** reactor short-baseline (RSBL) finding **no evidence** are actually **observing a deficit**;
 - this deficit **could** be interpreted as being due to SBL neutrino oscillations;
 - no visible dependence on $L \Rightarrow \Delta m^2 \gtrsim 1 \text{ eV}^2$;
 - global data (3σ):
$$\begin{cases} \Delta m_{\text{SOL}}^2 \simeq [6.8 \rightarrow 8.0] \times 10^{-5} \text{ eV}^2, \\ |\Delta m_{\text{ATM}}^2| \simeq [2.4 \rightarrow 2.6] \times 10^{-3} \text{ eV}^2; \end{cases}$$
- \Rightarrow solutions: **add new neutrinos** or **revise fluxes**.



[3] T.A. Mueller *et al.*, Phys. Rev. **C83** (2011) 054615 [arXiv:1101.2663].

[4] P. Huber, Phys. Rev. C **84** (2011) 024617 [arXiv:1106.0687].

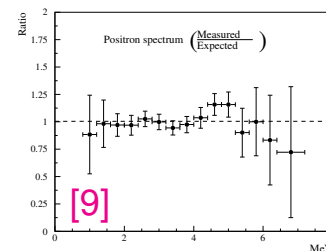
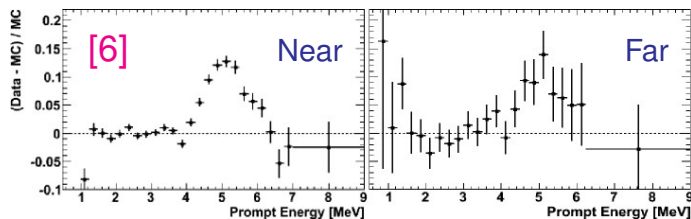
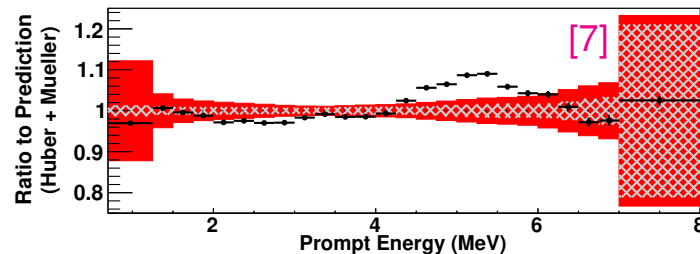
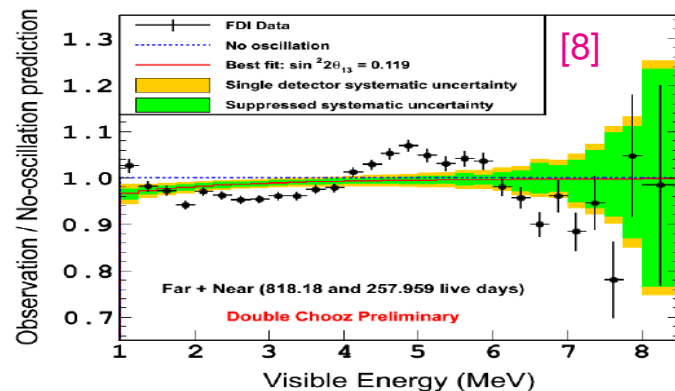
[5] G. Mention *et al.*, Phys. Rev. **D83** (2011) 073006 [arXiv:1101.2755].

\Rightarrow Talk: Hayes

\Rightarrow Talk: Suhonen

$\bar{\nu}_e$ disapp: 5 MeV excess

- Neutrino 2014: RENO [6] reported an **excess** of events around 5 MeV;
- excess (not deficit) & independent of $L \Rightarrow$ **flux feature**, not **sterile oscillations**;
- seen by Daya-Bay [7], Double-Chooz [8], and many others (also old Chooz [9]);
- more info: \Rightarrow **Talks: Hayes**



[6] S.H. Seo [RENO collab], talk at Neutrino 2014, Boston, USA, June 2-7, 2014.

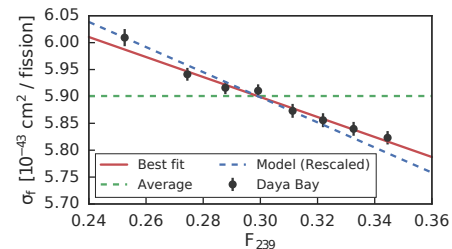
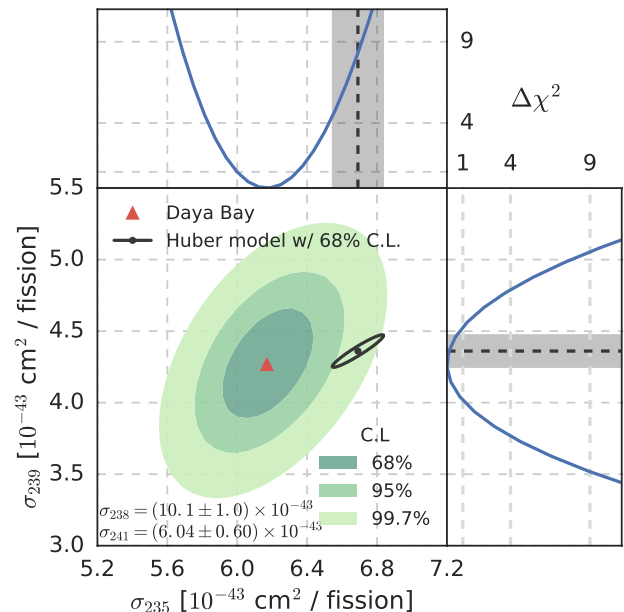
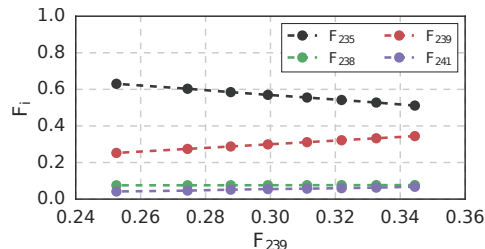
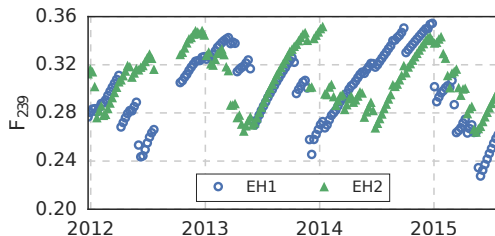
[7] F. P. An *et al.* [Daya-Bay collab], CPC **41** (2017) [arXiv:1607.05378].

[8] I.G. Botella [Double-Chooz collab], talk at EPS 2017, Venice, Italy.

[9] M. Apollonio *et al.* [Chooz collab], PLB **466** (1999) 415 [hep-ex/9907037].

Reactor $\bar{\nu}_e$ flux at the Daya-Bay

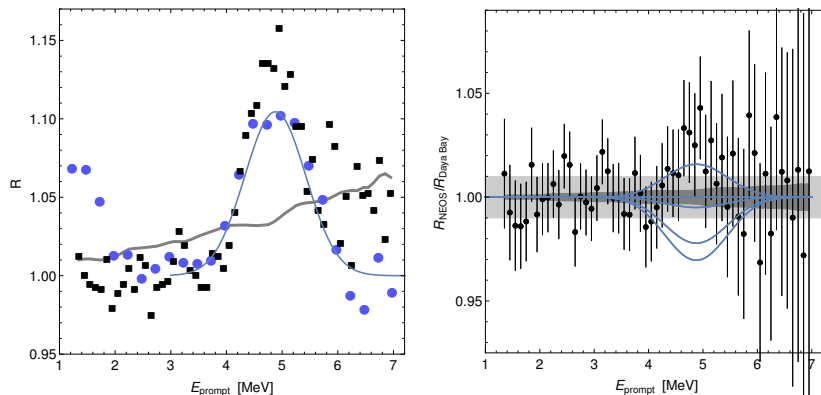
- $\bar{\nu}_e$ from [^{235}U , ^{238}U , ^{239}Pu , ^{241}Pu] fission chains;
- Daya-Bay: varying fuel composition \Rightarrow reconstruct contribution of individual isotopes;
- Results on the overall IBD yield σ_f :
 - (A) 1.7σ below expectation (known);
 - (B) dependence on fuel composition (described by ^{239}Pu) is 3.1σ away from model;
- sterile neutrinos can explain (A) but not (B);
- isotopes: measured ^{239}Pu OK, ^{235}U not OK.



[10] F.P. An et al. [Daya-Bay collab], Phys. Rev. Lett. **118** (2017) 251801 [arXiv:1704.01082].

The 5 MeV excess at Daya-Bay and NEOS

- In Ref. [11], the Daya-Bay spectrum from Ref. [7] was combined with the NEOS spectrum presented at ICHEP 2016 (later published in Ref. [12]) to infer which isotope is most likely to contribute to the 5 MeV excess;
- idea: fission fractions in Daya-Bay ($^{235}\text{U} = 56.1\%$, $^{239}\text{Pu} = 30.7\%$) sizably different from NEOS ($^{235}\text{U} = 65.5\%$, $^{239}\text{Pu} = 23.5\%$) \Rightarrow data are complementary;
- method:
 - construct a double ratio to cancel systematics;
 - add a “bump” to a given isotope and fit its amplitude;
- result: fit is best when the bump is ascribed to ^{235}U .



[7] F. P. An *et al.* [Daya-Bay collab], Chin. Phys. C **41** (2017) [arXiv:1607.05378].

[11] P. Huber, Phys. Rev. Lett. **118** (2017) 042502 [arXiv:1609.03910].

[12] Y. J. Ko *et al.* [NEOS collab], Phys. Rev. Lett. **118** (2017) 121802 [arXiv:1610.05134].

$\bar{\nu}_e$ disapp: new fluxes or new neutrinos?

- ^{235}U flux prediction is in tension with reactor rates [13], with Daya-Bay time evolution [10], and with DB/NEOS spectral data [11];
- however, reactor rates favor sterile solution [14];
- combined fit [14]: flux-only and sterile-only solutions give similar fits, but flux+sterile best;

Sample	235	235+239	OSC	235+OSC	239+OSC
Rates	20.7	17.7	12.8	12.6	12.7
Daya-Bay	3.8	3.6	9.5	3.6	3.8
Combined	25.3	24.8	23.0	20.2	17.5

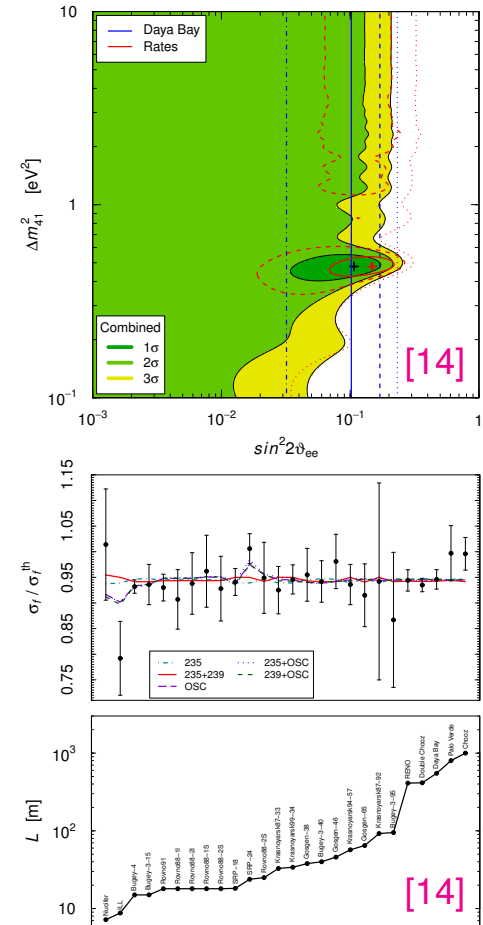
- anyway, free fluxes reduce significance of sterile hint.

[10] F.P. An *et al.* [DB], PRL **118** (2017) 251801 [arXiv:1704.01082].

[11] P. Huber, PRL **118** (2017) 042502 [arXiv:1609.03910].

[13] C. Giunti, PLB **764** (2017) 145 [arXiv:1608.04096].

[14] C. Giunti *et al.*, JHEP **10** (2017) 143 [arXiv:1708.01133].

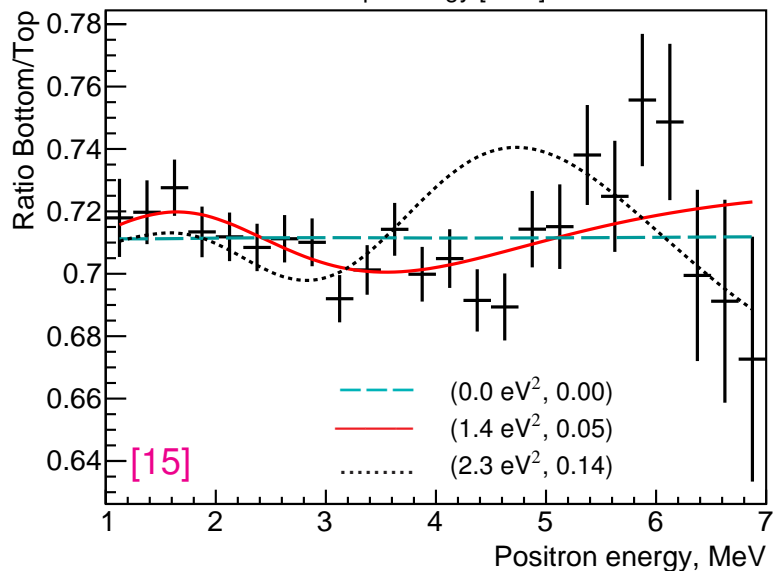
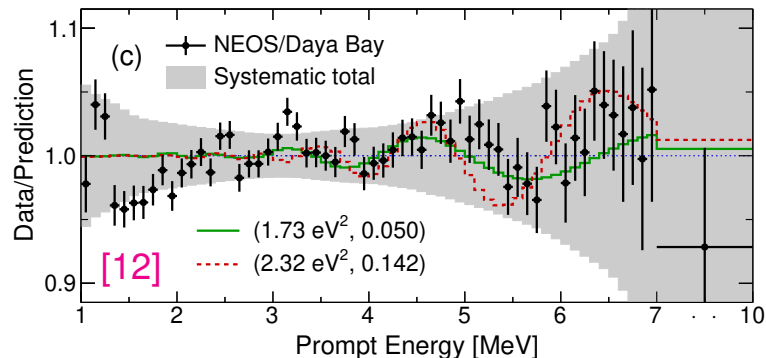


NEOS and DANSS results

- Both detectors have measured reactor neutrinos at very short baseline:
 - NEOS [12]: 24 m;
 - DANSS [15]: 10.7 m \rightarrow 12.7 m;
- data: near/far spectral ratios \Rightarrow insensitive to flux shape & normalization:
 - NEOS: normalized to Daya-Bay;
 - DANSS: movable detector;
- both detectors observe small energy modulations \Rightarrow hints of sterile ν .

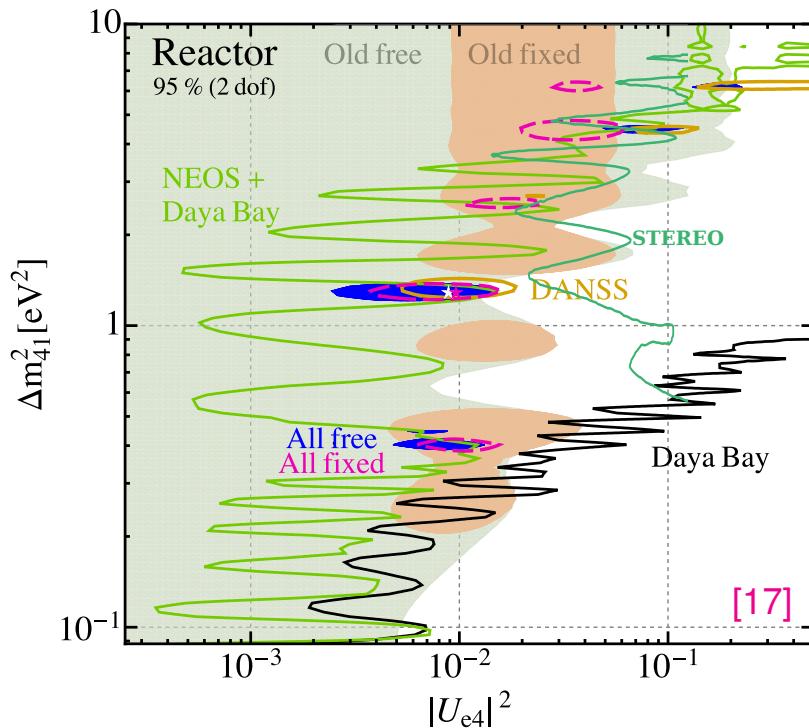
[12] Y. J. Ko *et al.* [NEOS collab], PRL **118** (2017) 121802 [arXiv:1610.05134].

[15] I. Alekseev *et al.* [DANSS collaboration], arXiv:1804.04046.



Global analysis of all reactor $\bar{\nu}_e$ disappearance data

- Total rates only \Rightarrow two models [16]:
 - “free”: unconstrained normalizations for [^{235}U , ^{238}U , ^{239}Pu , ^{241}Pu];
 - “fixed”: assumes Huber fluxes;
- spectral data as near/far ratios \Rightarrow independent of flux assumptions;
- results [17]:
 - 2.9σ (3.5σ) hint for sterile ν from analysis with free (fixed) fluxes;
 - fit dominated by DANSS+NEOS;
 - DANSS osc. agree with NEOS;
 - DANSS osc. in tension with fixed.



[16] M. Dentler *et al.*, JHEP **11** (2017) 099 [arXiv:1709.04294].

[17] M. Dentler *et al.*, arXiv:1803.10661.

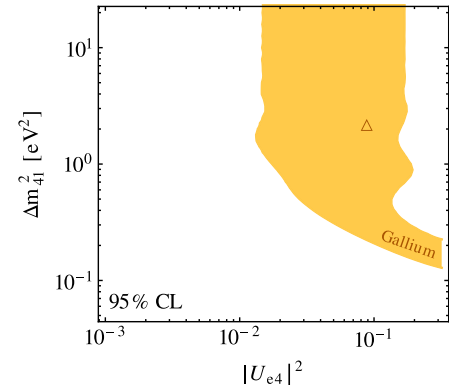
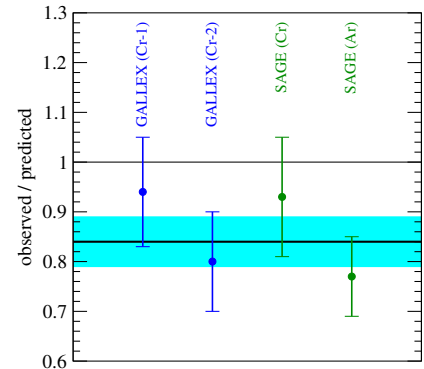
\Rightarrow Poster #M/147: Á. Hernández-Cabezudo

ν_e disappearance: the gallium anomaly

- The $^{71}\text{Ga} \rightarrow ^{71}\text{Ge}$ neutrino capture cross-section, relevant for the **GALLEX** and **SAGE** solar neutrino experiments, was calibrated with intense ^{51}Cr and ^{37}Ar neutrino sources;
- these measurements show a significant deficit with respect to the predicted values:

$$\begin{array}{l} \text{GALLEX:} \left\{ \begin{array}{l} R_1(\text{Cr}) = 0.94 \pm 0.11 \text{ [18]} \\ R_2(\text{Cr}) = 0.80 \pm 0.10 \text{ [18]} \end{array} \right\} \\ \text{SAGE:} \left\{ \begin{array}{l} R_3(\text{Cr}) = 0.93 \pm 0.12 \text{ [19]} \\ R_4(\text{Ar}) = 0.77 \pm 0.08 \text{ [20]} \end{array} \right\} \end{array} \Rightarrow 0.84 \pm 0.05$$

- such 3σ deficit can be interpreted in terms of ν oscillations;
- once again, data suggests $\Delta m^2 \gtrsim 1 \text{ eV}^2$.



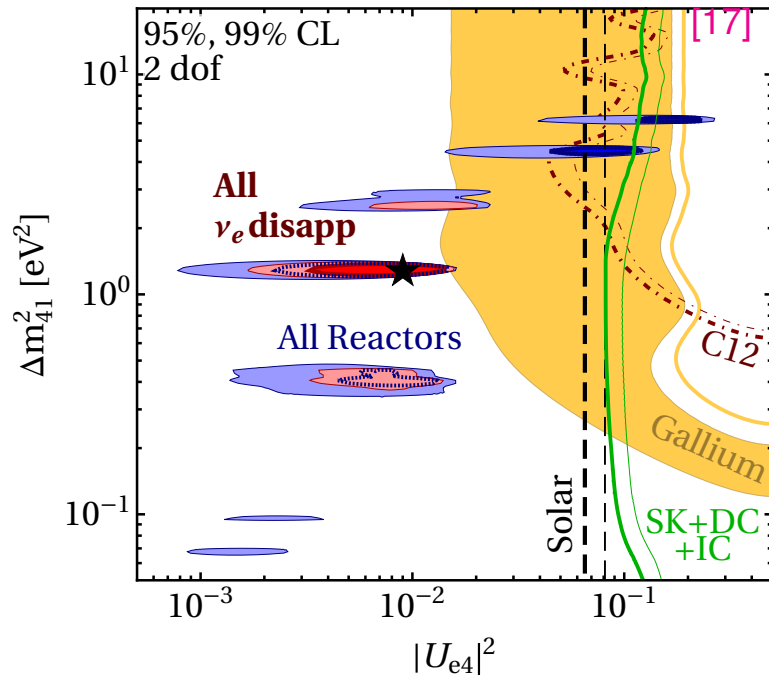
[18] F. Kaether *et al.*, Phys. Lett. **B685** (2010) 47–54 [arXiv:1001.2731].

[19] J. Abdurashitov *et al.* [SAGE collab], Phys. Rev. **C59** (1999) 2246–2263 [hep-ph/9803418].

[20] J. Abdurashitov *et al.* [SAGE collab], Phys. Rev. **C73** (2006) 045805 [nucl-ex/0512041].

Global analysis of ν_e and $\bar{\nu}_e$ disappearance

- In addition to **reactor** ($\bar{\nu}_e$) and **gallium** (ν_e) data, we include:
 - **Karmen** [21] and **LSND** [22] data on $\nu_e + {}^{12}\text{C} \rightarrow e^- + {}^{12}\text{N}$ reaction;
 - **atmospheric** neutrino data ($\nu_e + \bar{\nu}_e$);
 - **solar** neutrino data (ν_e);
- results [17]:
 - small tension (2.2σ) between **gal-**
lium and **reactor** (**free** fluxes) data;
 - hint for **sterile** slightly increases to 3.2σ (3.8σ) for **free** (**fixed**) fluxes;
 - global best-fit at $\Delta m_{41}^2 = 1.3 \text{ eV}^2$.



[17] M. Dentler *et al.*, arXiv:1803.10661.

[21] B. Armbruster *et al.* [Karmen collab], Phys. Rev. C **57** (1998) 3414 [hep-ex/9801007].

[22] L. B. Auerbach *et al.* [LSND collab], Phys. Rev. C **64** (2001) 065501 [hep-ex/0105068].

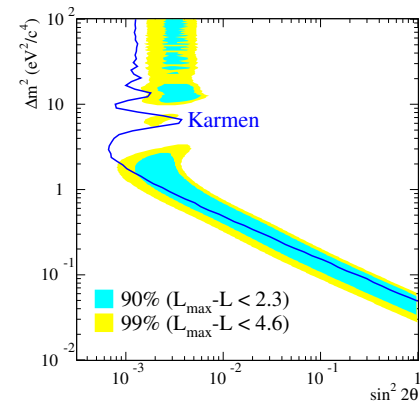
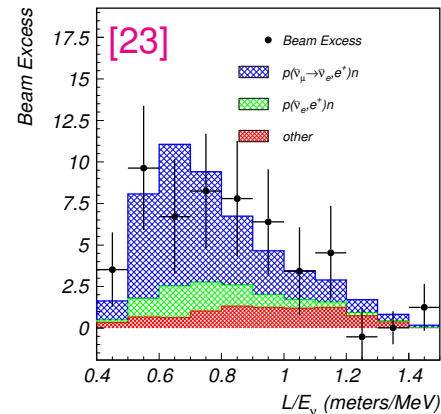
The LSND anomaly

- The **LSND** experiment observed an excess of $\bar{\nu}_e$ events in a $\bar{\nu}_\mu$ beam ($E_\nu \sim 30$ MeV, $L \simeq 35$ m) [23];
- the **Karmen** collaboration did not confirm the claim, but couldn't fully exclude it either [24];
- the signal is compatible with $\bar{\nu}_\mu \rightarrow \bar{\nu}_e$ oscillations provided that $\Delta m^2 \gtrsim 0.1 \text{ eV}^2$;
- on the other hand, global neutrino data give (at 3σ):

$$\Delta m_{\text{SOL}}^2 \simeq [6.8 \rightarrow 8.0] \times 10^{-5} \text{ eV}^2,$$

$$|\Delta m_{\text{ATM}}^2| \simeq [2.4 \rightarrow 2.6] \times 10^{-3} \text{ eV}^2;$$

- again, to explain LSND with mass-induced ν oscillations one needs **new** neutrino mass eigenstates;
- **MiniBooNE**: much larger E_ν and L but similar L/E_ν .

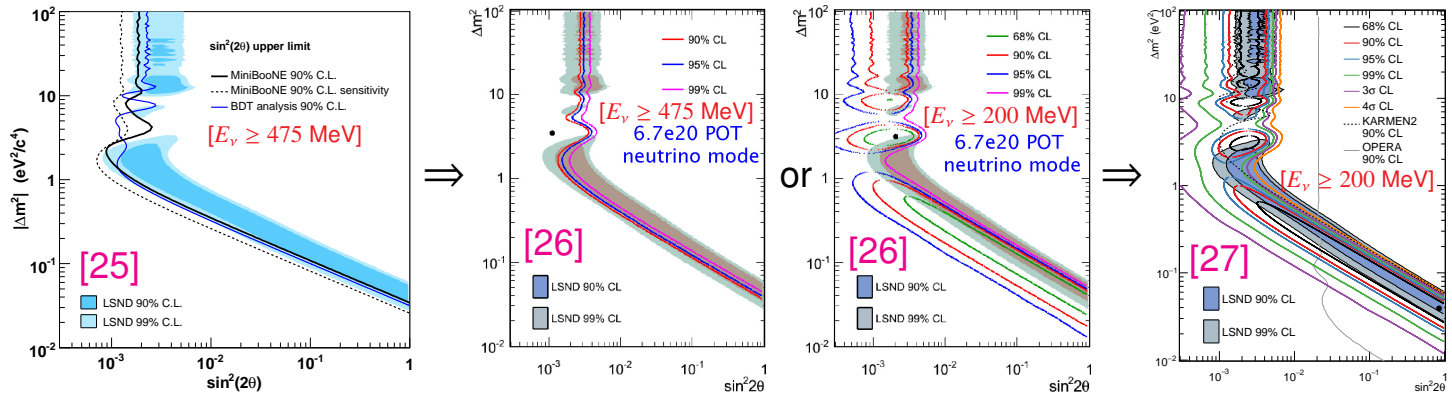


[23] A. Aguilar-Arevalo *et al.* [LSND collab], Phys. Rev. D **64** (2001) 112007 [hep-ex/0104049].

[24] B. Armbruster *et al.* [KARMEN collab], Phys. Rev. D **65** (2002) 112001 [hep-ex/0203021].

MiniBooNE neutrino data

- Statistics: 5.58 (2007) \rightarrow 6.46 (2008) \rightarrow 12.84 (2018) $\times 10^{20}$ POT;
- is ν signal compatible with 2ν oscillations? $\left\{ \begin{array}{l} 2007: P_{\text{osc}} \simeq 1\% \Rightarrow \text{no it isn't [25];} \\ 2012: P_{\text{osc}} \simeq 6\% \Rightarrow \text{maybe it is [26];} \\ 2018: P_{\text{osc}} \simeq 15\% \Rightarrow \text{yes it is [27];} \end{array} \right.$
- do MB- ν rule out LSND- $\bar{\nu}$ signal? 2007: yes [25]; 2012: not really [26]; 2018: no [27].



[25] A.A. Aguilar-Arevalo *et al.* [MiniBooNE collab], Phys. Rev. Lett. **98** (2007) 231801 [arXiv:0704.1500].

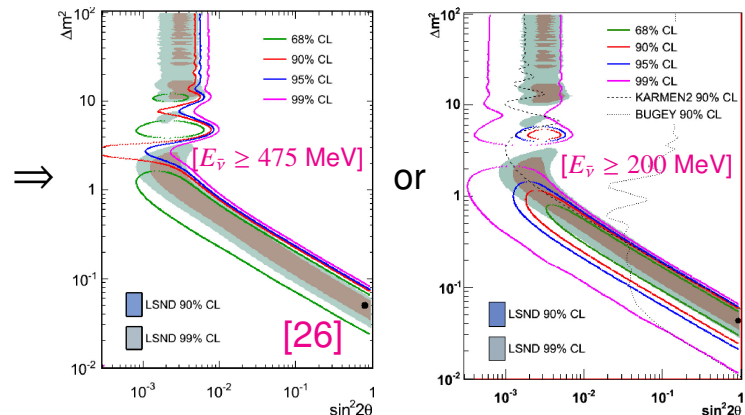
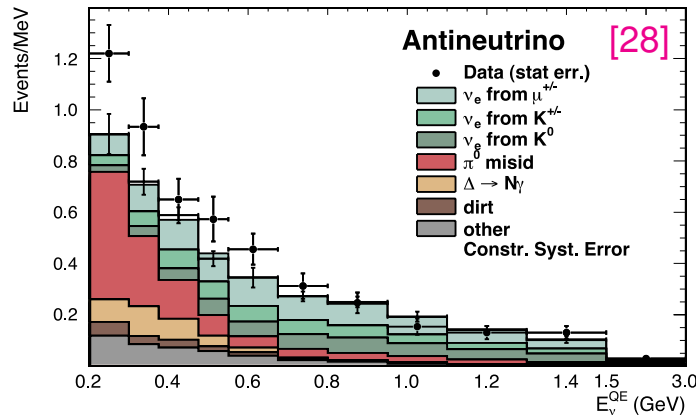
[26] C. Polly, talk at Neutrino 2012, Kyoto, Japan, June 3-9, 2012.

\Rightarrow Talk: Huang

[27] A.A. Aguilar-Arevalo *et al.* [MiniBooNE collab], arXiv:1805.12028.

MiniBooNE antineutrino data

- New data presented at Neutrino 2012, statistics doubled ($\rightarrow 11.27 \times 10^{20}$ POT) [26];
- compatibility with ν data: $\left\{ \begin{array}{l} \text{low-energy excess increased} \Rightarrow \text{better agreement;} \\ \text{mid-energy excess reduced} \Rightarrow \text{better agreement;} \end{array} \right.$
- is $\bar{\nu}$ signal compatible with 2ν oscillations? $P_{\text{osc}} = 66\% \Rightarrow \text{definitely yes}$ [28];
- is MB- $\bar{\nu}$ signal compatible with LSND? **Yes**, irrespective of the energy threshold.

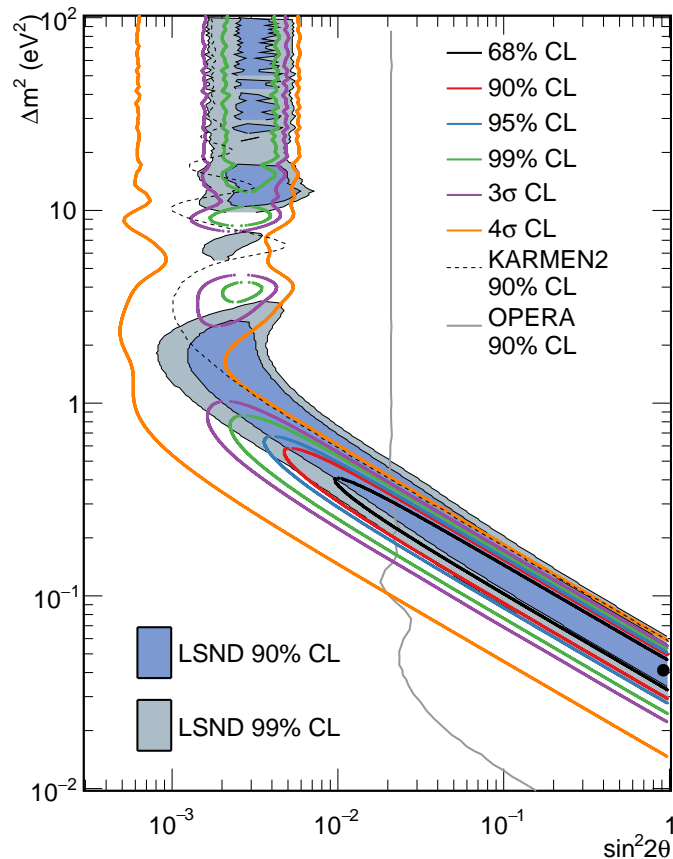
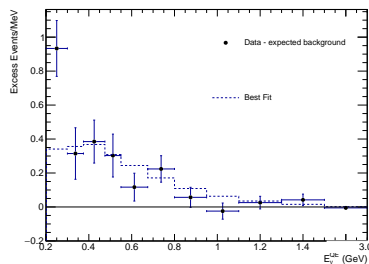
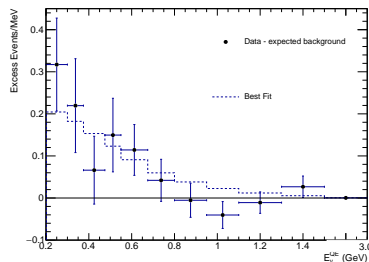
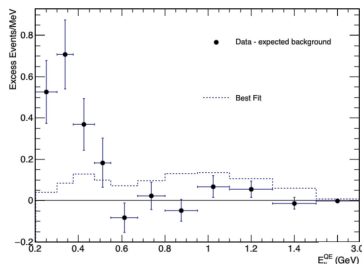


[26] C. Polly, talk at Neutrino 2012, Kyoto, Japan, June 3-9, 2012.

[28] A.A. Aguilar-Arevalo *et al.* [MiniBooNE collab], PRL **110** (2013) 161801 [arXiv:1303.2588].

MiniBooNE combined data

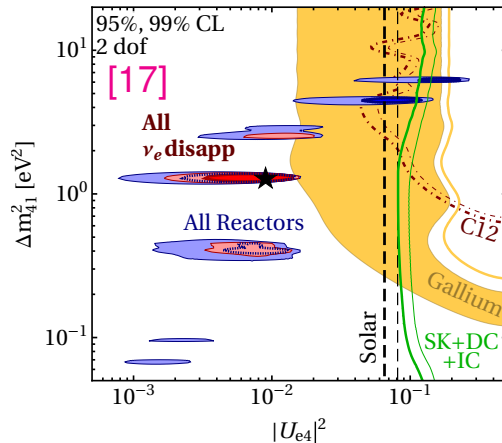
- Low-energy excess in 1st half of ν data absent from 2nd half (except lowest bin) \Rightarrow overall excess is only mild;
- combined $\nu_e + \bar{\nu}_e$ fit: 4.8σ evidence for sterile;
- LSND signal: 3.8σ ;
- global LSND + MB preference for sterile neutrinos: 6.1σ .



[27] A.A. Aguilar-Arevalo *et al.* [MiniBooNE collab], arXiv:1805.12028.

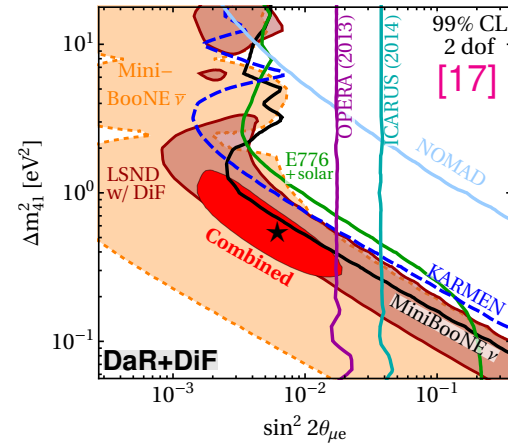
ν_e disappearance

- Relevant experiments:
 - Gallium (ν)
 - SBL reactors ($\bar{\nu}$)
 - LBL reactors ($\bar{\nu}$)
 - KamLAND ($\bar{\nu}$)
 - Atmos ($\nu, \bar{\nu}$)
 - Solar (ν)
 - ^{12}C (ν)



$\nu_\mu \rightarrow \nu_e$ appearance

- Relevant experiments:
 - LSND ($\bar{\nu}$)
 - MiniBooNE ($\nu, \bar{\nu}$)
 - E776 ($\nu, \bar{\nu}$)
 - ICARUS (ν)
 - KARMEN ($\bar{\nu}$)
 - NOMAD (ν)
 - OPERA (ν)

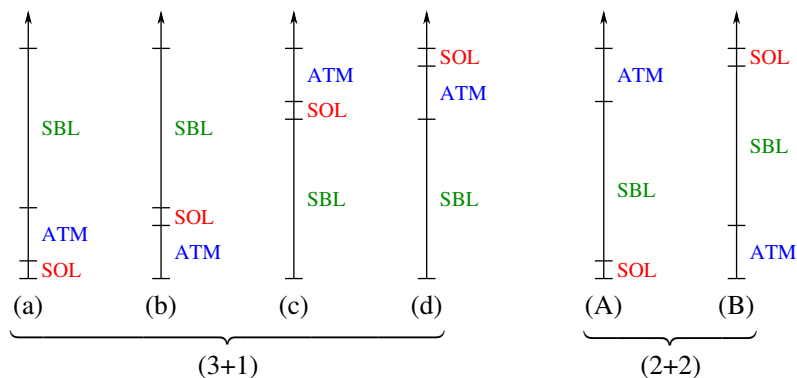


- Note: $\bar{\nu}_e \rightarrow \bar{\nu}_e$ and $\bar{\nu}_\mu \rightarrow \bar{\nu}_e$ probe the same Δm^2 but a different mixing angle \Rightarrow mutual comparison requires embedding them into a **general oscillation model**.

[17] Dentler, Hernández-Cabezudo, Kopp, Machado, MM, Martinez-Soler, Schwetz, arXiv:1803.10661.

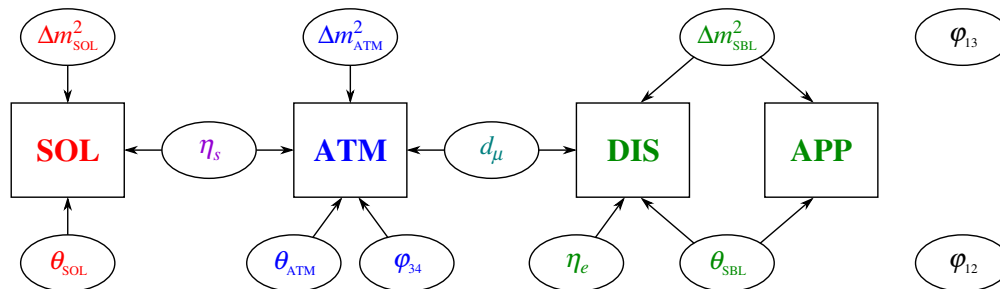
Four neutrino mass models

- Approximation: $\Delta m_{\text{SOL}}^2 \ll \Delta m_{\text{ATM}}^2 \ll \Delta m_{\text{SBL}}^2 \Rightarrow$ 6 different mass schemes:

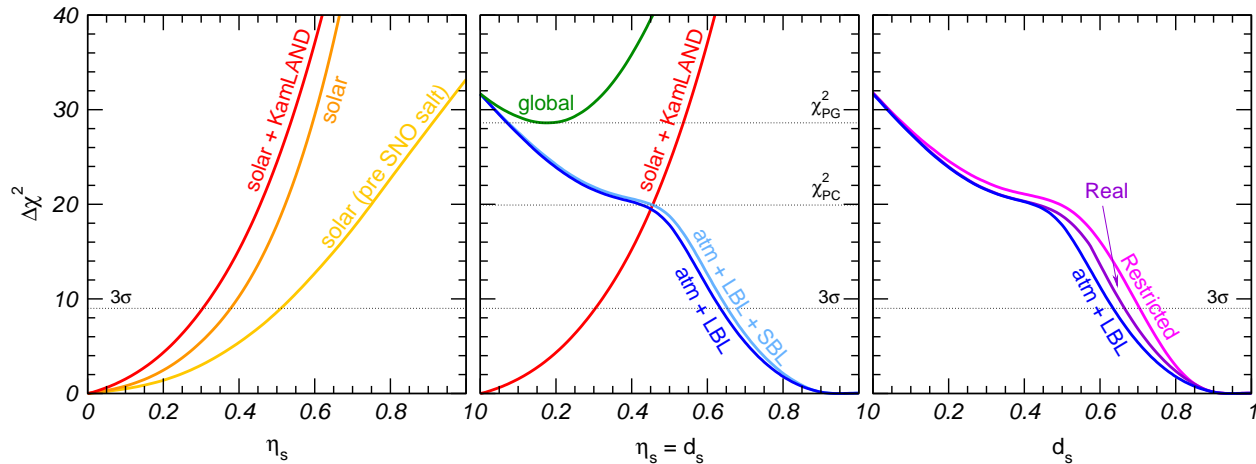


$$U = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} & U_{e4} \\ U_{\mu1} & U_{\mu2} & U_{\mu3} & U_{\mu4} \\ U_{\tau1} & U_{\tau2} & U_{\tau3} & U_{\tau4} \\ U_{s1} & U_{s2} & U_{s3} & U_{s4} \end{pmatrix}$$

- Total: 3 Δm^2 , 6 angles, 3 phases. Different set of experimental data *partially decouple*:



(2+2): ruled out by solar and atmospheric data



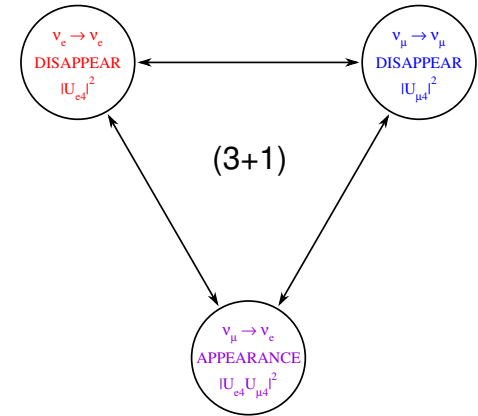
- in (2+2) models, fractions of ν_s in **solar** (η_s) and **atmos** ($1 - d_s$) add to one $\Rightarrow \boxed{\eta_s = d_s}$;
- 3σ allowed regions $\eta_s \leq 0.31$ (**solar**) and $d_s \geq 0.63$ (**atmos**) do not overlap; superposition occurs only above 4.5σ ($\chi_{\text{PC}}^2 = 19.9$);
- the χ^2 increase from the combination of **solar** and **atmos** data is $\chi_{\text{PG}}^2 = 28.6$ (1 dof), corresponding to a $\text{PG} = 9 \times 10^{-8}$ [29].

[29] M. Maltoni, T. Schwetz, M.A. Tortola, J.W.F. Valle, Nucl. Phys. **B643** (2002) 321 [hep-ph/0207157].

(3+1): appearance versus disappearance

- (3+1): $P_{\nu_\mu \rightarrow \nu_e} \propto |U_{e4}U_{\mu 4}|^2$ with $\begin{cases} |U_{e4}|^2 \propto P_{\nu_e \rightarrow \nu_e}, \\ |U_{\mu 4}|^2 \propto P_{\nu_\mu \rightarrow \nu_\mu}; \end{cases}$
- hence, $P_{\nu_\mu \rightarrow \nu_e} > 0$ requires $\begin{cases} P_{\nu_e \rightarrow \nu_e} > 0, \\ P_{\nu_\mu \rightarrow \nu_\mu} > 0; \end{cases}$

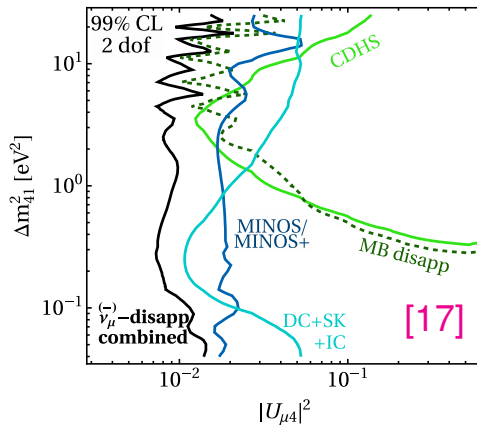
❓ are $\nu_\mu \rightarrow \nu_\mu$ searches compatible with this?



ν_μ disappearance: present status

- Many experiments have been performed:

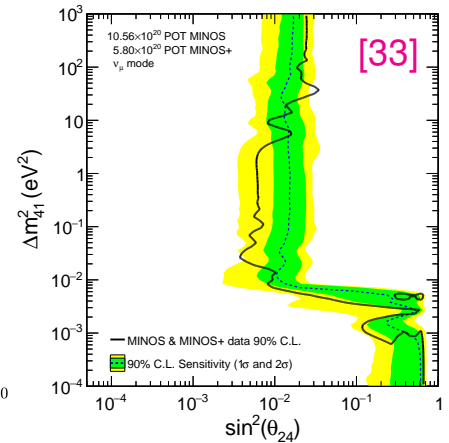
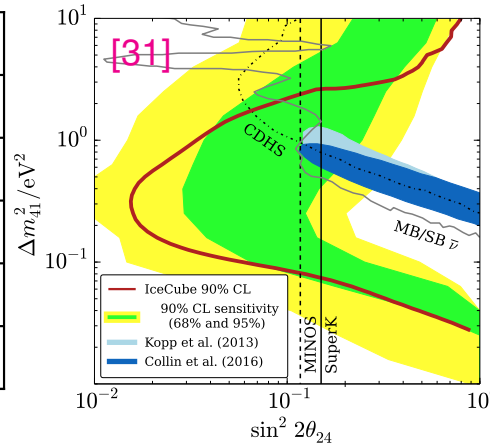
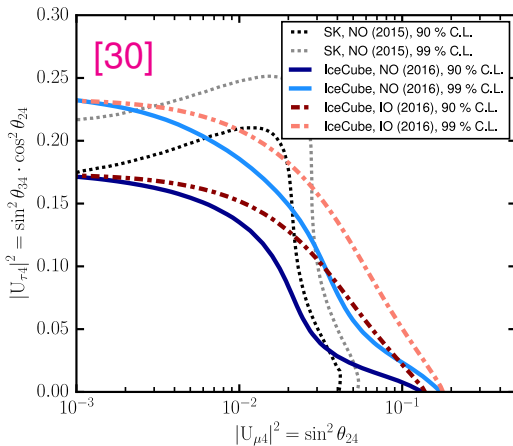
– CDHS (ν)	– MINOS (ν)
– MiniBooNE ($\nu, \bar{\nu}$)	– NO ν A (ν)
– SciBooNE ($\nu, \bar{\nu}$)	– SK atmos ($\nu, \bar{\nu}$)
- no hint of ν_μ disappearance has been observed;
- bound on $|U_{\mu 4}|^2$ may be in tension with other data...



[17] Dentler, Hernández-Cabezudo, Kopp, Machado, MM, Martinez-Soler, Schwetz, arXiv:1803.10661.

ν_μ disappearance: recent atmospheric and SBL data

- new DeepCore [30] and IceCube [31] atmos. data probe ν_μ mixing with heavy state;
- sterile neutrino oscillations also studied by NO ν A using neutral-current data [32];
- recent MINOS/MINOS+ analysis [33] improves bound on ν_μ disapp. \Rightarrow Talk: Aurisano



[30] M.G. Aartsen *et al.* [IceCube collab], Phys. Rev. D **95** (2017) 112002 [arXiv:1702.05160].

[31] M.G. Aartsen *et al.* [IceCube collab], Phys. Rev. Lett. **117** (2016) 071801 [arXiv:1605.01990].

[32] P. Adamson *et al.* [NO ν A collab], Phys. Rev. D **96** (2017) 072006 [arXiv:1706.04592].

[33] P. Adamson *et al.* [MINOS collab], [arXiv:1710.06488].

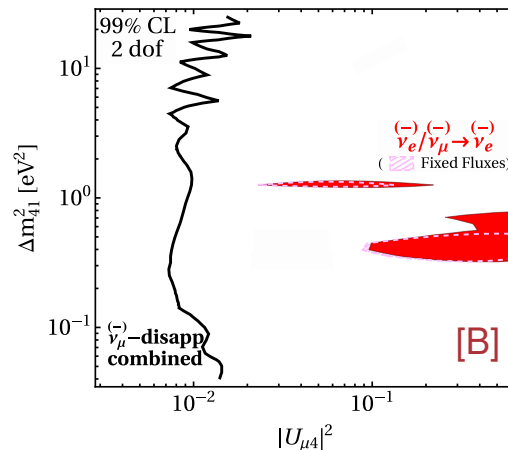
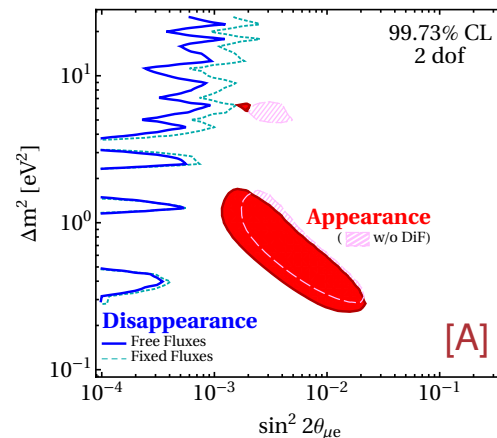
(3+1): tension among data samples

- Limits on $\nu_e \rightarrow \nu_e$ and $\nu_\mu \rightarrow \nu_\mu$ disappearance imply a bound on the $\nu_\mu \rightarrow \nu_e$ appearance probability;
- such bound is stronger than what is required to explain the **LSND** and **MiniBooNe** excesses [A];
- hence, severe tension arises between **APP** and **DIS** data: $\chi^2_{\text{PG}}/\text{dof} = 29.6/2 \Rightarrow \text{PG} = 3.7 \times 10^{-7}$ [17];
- a similar result is visible when comparing “ ν_e -data” ($\nu_e \rightarrow \nu_e$ and $\nu_\mu \rightarrow \nu_e$) and “ ν_μ -data” ($\nu_\mu \rightarrow \nu_\mu$) [B];
- note: tension between **APP** and **DIS** data first pointed out in 1999 [34]. Full global fit in 2001 [35] cornered (3+1) models. No conceptual change since then...

[17] M. Dentler *et al.*, arXiv:1803.10661.

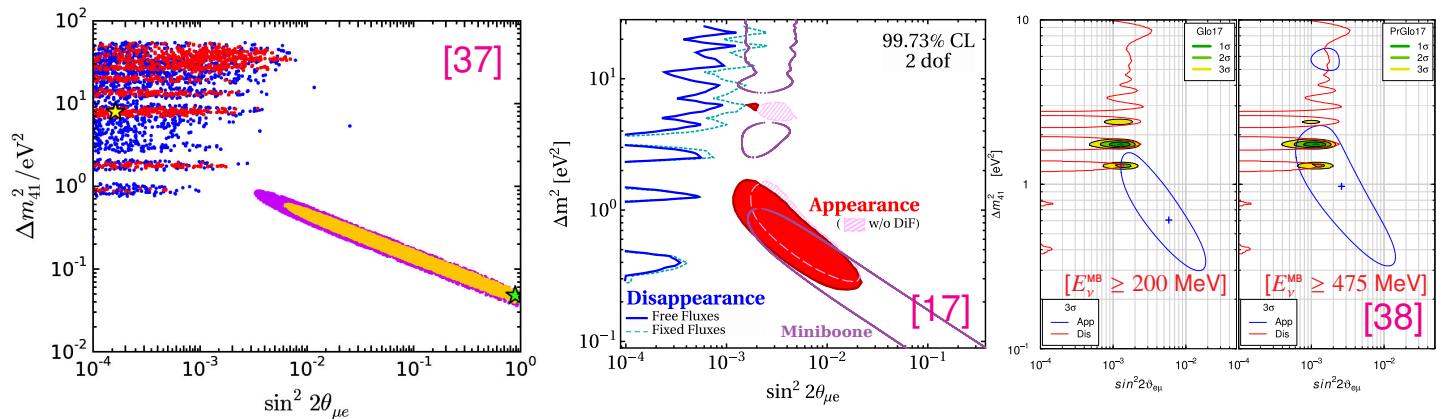
[34] S.M. Bilenky *et al.*, PRD **60** (1999) 073007 [hep-ph/9903454].

[35] MM, Schwetz, Valle, PLB **518** (2001) 252 [hep-ph/0107150].



(3+1): comparison with other analyses

- A few independent analyses of **APP** versus **DIS** compatibility have been presented;
- results of [36, 37] similar to ours \Rightarrow (3+1) ruled out: $\chi^2_{\text{PG}}/\text{dof} = 22.9/2 \Rightarrow \text{PG} = 10^{-5}$;
- results of [38]: $\begin{cases} \text{(3+1) ruled out: } \chi^2_{\text{PG}}/\text{dof} = 17.2/2 \Rightarrow \text{PG} = 0.019\% \text{ (MB } E_\nu > 200 \text{ MeV),} \\ \text{(3+1) very poor: } \chi^2_{\text{PG}}/\text{dof} = 7.2/2 \Rightarrow \text{PG} = 2.7\% \text{ (MB } E_\nu > 475 \text{ MeV).} \end{cases}$



[17] Dentler, Hernández-Cabezudo, Kopp, Machado, MM, Martinez-Soler, Schwetz, arXiv:1803.10661.

[36] J.M. Conrad et al., Adv. High Energy Phys. **2013** (2013) 163897 [arXiv:1207.4765].

[37] G.H. Collin et al., Nucl. Phys. B **908** (2016) 354 [arXiv:1602.00671].

[38] S. Gariazzo, C. Giunti, M. Laveder and Y. F. Li, JHEP **06** (2017) 135 [arXiv:1703.00860].

(3+1): robustness of the result

- The tension cannot be eliminated by discarding any *individual* experiment.

Analysis	$\chi^2_{\min, \text{global}}$	$\chi^2_{\min, \text{app}}$	$\Delta\chi^2_{\text{app}}$	$\chi^2_{\min, \text{disapp}}$	$\Delta\chi^2_{\text{disapp}}$	$\chi^2_{\text{PG}}/\text{dof}$	PG
Global	1120.9	79.1	11.9	1012.2	17.7	29.6/2	3.7×10^{-7}
Removing anomalous data sets							
w/o LSND	1099.2	86.8	12.8	1012.2	0.1	12.9/2	1.6×10^{-3}
w/o MiniBooNE	1012.2	40.7	8.3	947.2	16.1	24.4/2	5.2×10^{-6}
w/o reactors	925.1	79.1	12.2	833.8	8.1	20.3/2	3.8×10^{-5}
w/o gallium	1116.0	79.1	13.8	1003.1	20.1	33.9/2	4.4×10^{-8}
Removing constraints							
w/o IceCube	920.8	79.1	11.9	812.4	17.5	29.4/2	4.2×10^{-7}
w/o MINOS(+)	1052.1	79.1	15.6	948.6	8.94	24.5/2	4.7×10^{-6}
w/o MB disapp	1054.9	79.1	14.7	947.2	13.9	28.7/2	6.0×10^{-7}
w/o CDHS	1104.8	79.1	11.9	997.5	16.3	28.2/2	7.5×10^{-7}
Removing classes of data							
$\bar{\nu}_e$ -dis vs app	628.6	79.1	0.8	542.9	5.8	6.6/2	3.6×10^{-2}
$\bar{\nu}_\mu$ -dis vs app	564.7	79.1	12.0	468.9	4.7	16.7/2	2.3×10^{-4}
$\bar{\nu}_\mu$ -dis+solar vs app	884.4	79.1	13.9	781.7	9.7	23.6/2	7.4×10^{-6}

More sterile neutrinos? The case of (3+2) models

- With **one** extra sterile neutrino, m_4 :

$$P_{\mu e}^{4\nu} = 4|U_{e4}|^2|U_{\mu 4}|^2 \sin^2 \phi_{41} \quad \text{with} \quad \phi_{ij} \equiv \frac{\Delta m_{ij}^2 L}{4E};$$

- for large energy $P_{\mu e}^{4\nu}$ drops as $1/E^2$;
- however, the low-energy MB excess is much sharper ($\sim 1/E^3$);
- On the other hand, with **two** extra neutrinos, m_4 and m_5 :

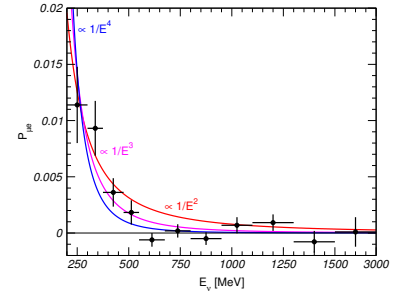
$$P_{\mu e}^{5\nu} = 4|U_{e4}|^2|U_{\mu 4}|^2 \sin^2 \phi_{41} + 4|U_{e5}|^2|U_{\mu 5}|^2 \sin^2 \phi_{51} + 8|U_{e4}U_{e5}U_{\mu 4}U_{\mu 5}| \sin \phi_{41} \sin \phi_{51} \cos(\phi_{54} - \delta);$$

- terms of order $1/E^2$ suppressed if $\delta \approx \pi$ and $|U_{e4}U_{\mu 4}\Delta m_{41}^2| \approx |U_{e5}U_{\mu 5}\Delta m_{51}^2|$;

\Rightarrow two extra sterile states provide a better description of the MB low-energy ν data [39].

- also, $\delta = \arg(U_{e4}^* U_{\mu 4} U_{e5} U_{\mu 5}^*)$ differentiates between ν (MB) from $\bar{\nu}$ (MB/LSND);
- however, (3+2) models suffer from **the same APP/DIS tension** as (3+1) models;
- also, (3+2) models have stronger problems with **cosmology** since $\sum m_\nu$ is larger;

\Rightarrow (3+N) models **do not present substantial advantages** over the simpler (3+1) model.



[39] M. Maltoni, T. Schwetz, Phys. Rev. **D76** (2007) 093005 [arXiv:0705.0107].

- Anomalies in $\nu_e \rightarrow \nu_e$ **disappearance** and $\nu_\mu \rightarrow \nu_e$ **appearance** experiments point towards conversion mechanisms beyond the well-established 3ν oscillation paradigm;
 - each of these anomalies can be **individually** explained by sterile neutrinos;
 - sterile neutrinos still succeed in simultaneously explaining groups of anomalies **sharing the same oscillation channel**. However some problem arises:
 - $\nu_e \rightarrow \nu_e$ **disappearance** data face issues with flux normalization and the 5 MeV bump, as well as small tensions in reactor vs gallium and “rates” vs DANSS/NEOS;
 - $\nu_\mu \rightarrow \nu_e$ **appearance** data show an excess in low-E neutrino data, which is not so manifest in antineutrino data.
 - in contrast, no anomaly is found in any $\nu_\mu \rightarrow \nu_\mu$ **disappearance** data set;
- ⇒ sterile neutrino models **fail to simultaneously account** for **all** the $\nu_e \rightarrow \nu_e$ data, the $\nu_\mu \rightarrow \nu_e$ data and the $\nu_\mu \rightarrow \nu_\mu$ data. This conclusion is robust;
- if the $\nu_e \rightarrow \nu_e$ and $\nu_\mu \rightarrow \nu_e$ anomalies are confirmed, and the $\nu_\mu \rightarrow \nu_\mu$ bounds are not refuted, new physics will be needed. Such new physics may well involve extra sterile neutrinos, but together with something else (or some “unusual” neutrino property).