Comparisons of measured and requantized classical molecular dynamics calculated line shape of air-broadened isolated transitions of molecular oxygen

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13th International HITRAN Conference, Cambridge, 23-25 June 2014

## Why molecular oxygen ?

 $\Box$  O<sub>2</sub> absorption bands (A, B,  $\gamma$ ) are used for a number of remote sensing applications as:

- wind measurements,

- determination of surface pressure, clouds and aerosols proprieties, vertical profiles of pressure and temperature,í

□ But the Voigt profile, widely used for the analysis, is **inaccurate** as it neglectes:

- The speed-dependences of the collisionnal widths  $\gamma(v)$  and shifts  $\delta(v)$ 

-The velocity changes induced by collisions (Dicke narrowing effect)





## Line shape model beyond the Voigt profile

□ Recently, a new simplified line shape model (Hartmann-Tran model) has been developed in order to account for the effects neglected by the Voigt profile. This model has been proposed as a standard for the future updates of the spectroscopic databases.

(Tennyson et al., Pure and Applied Chemistry, in press (2014))

□ The parameters involved in this model can be obtained from theoretical spectra obtained from Classical Molecular Dynamics Simulations (CMDS).

□ These simulations provide (among other quantities) the autocorrelation function of the dipole, yielding to the spectra through a Laplace-Fourier Transform.

□ They have been successfully applied for line shape predictions for pure CO<sub>2</sub>, H<sub>2</sub>O, and for one transition (P11P11) of pure O<sub>2</sub>.

• Can we developed a HT model for O<sub>2</sub>-air ?

## **Classical Molecular Dynamics Simulations**

#### $\Box N_M$ (>10<sup>6</sup>) molecules (20% O<sub>2</sub>, 80% N<sub>2</sub>) treated simultaneously

- Placed in a cubic box (size determined from  $N_M$  and molecular density n)
- Periodic boundary (treated box surrounded by 4096 identical other boxes)
- When a molecule gets out of the box, it comes back-in from the opposite box

# □ The state of each molecule *m* (linear, rigid rotor) is parameterized by its:

- Center of mass (CoM) position and velocity
- Molecule orientation
- Molecule rotational speed



**The site-site intermolecular potentials for**  $O_2$ - $O_2$ ,  $O_2$ - $N_2$  et  $N_2$ - $N_2$  have been taken from the literature (as 6-12 Lennard-Jones)

#### **The O<sub>2</sub> electronic spin has been neglected in the calculations**

- $O_2$  transitions are under the form  $\Delta N(N'') \Delta J(J'')$
- Hence õPPö and õPQö transitions types are equivalent in the calculations

## **Classical Molecular Dynamics Simulations**

#### □ Initialization (time t=0)

- Random CoM positions and axis orientations
- CoM velocity and rotation: random orientations, modules from Maxwell-Boltzmann

**Time evolution for all molecules treated sequentially** (with small enough time step *dt*)

- At each time *t* compute force and torque on each molecule from sum of potential gradient of over surrounding neighbors (cut-off sphere of 20 Å)

- Then compute acceleration of CoM and of orientation
- Then compute molecule parameters at t+dt from those at t

□ For each molecule *m*, a requantification procedure is applied based on the correspondence principle: Associates the rotational quantum number  $N_m$ to the rotational speed  $\omega_m$  by

$$\frac{1}{2}I\omega_m^2 \approx \frac{\hbar^2}{2I}N_m(N_m+1)$$

### Spectral line shape

The spectrum  $F(\omega)$  is given by the Laplace-Fourier transform of the autocorrelation function (ACF)  $\Phi(\omega,t)$ 

$$F(\omega) = \operatorname{Re}\left\{\frac{1}{\pi}\int_{0}^{+\infty}\Phi(\omega,t)e^{-i\omega t}dt\right\}$$



## Analysis

□ The effects beyond the Voigt profile can be studied by analyzing the  $\tilde{O}W\ddot{O}$ type signatures in the residuals of the fits. <sup>1.0</sup> O<sub>2</sub>-air. P9P9 P=150 Torr. T=295 K

□ All the studied spectra were adjusted with a Voigt profile:

- The Doppler width  $\Gamma_D$  has been **fixed** to its theoretical value,

- The intensity *S*, the Lorentz width  $\Gamma_L$  have been **adjusted**.



□ In order to validate the calculations, experimental spectra were recorded at NIST for the R1Q2, P9P9, P11P11, P13P13, and P15Q14 transitions of the  $a {}^{1}\Delta_{g} \leftarrow X {}^{3}\Sigma_{g}^{-}(0,0)$  band, using a frequency-stabilized cavity ringdown spectrometer (FS-CRDS)

### Results: Lorentz widths $\Gamma_L$



### Results: Lorentz widths $\Gamma_L$



### Results: study of the residuals

Residuals of the fit of theoretical and measured spectra of the P9P9 transition



#### **Results:**

-Typical õWö-shaped residuals

-Both amplitudes and widths of the õWö are in good agreement between the calculations and the measurements

- As seen for other systems (H<sub>2</sub>O, CO<sub>2</sub>): For  $\Gamma_L / \Gamma_D = 1$ , the amplitude is maximum For  $\Gamma_L / \Gamma_D \rightarrow 0$ , the amplitude goes to 0 For  $\Gamma_L / \Gamma_D \rightarrow +\infty$ , the amplitude goes to an asymptotique value

### Results: study of the residuals



### Conclusion and future work

□ The comparison of theoretical and experimental spectra trhough their fits by a Voigt profile demonstrates that the CMDS can be used as a prediction of the (small) deviations to the Voigt profile.

□ The results presented here validate the first step of this work that aims at understanding the physical processes that affect the isolated line shapes of molecular oxygen.

□ The next step will aim to determinate the parameters describing the velocity changes induced by collisions from the CMDS spectra described here.

*JL is pleased to acknowledge support of this research by the French National Research Agency (ANR) through the project ASGGRS (ANR-12-PDOC-0012-01).*