

# Supplementary Information for “From creep to flow: Granular materials under cyclic shear”

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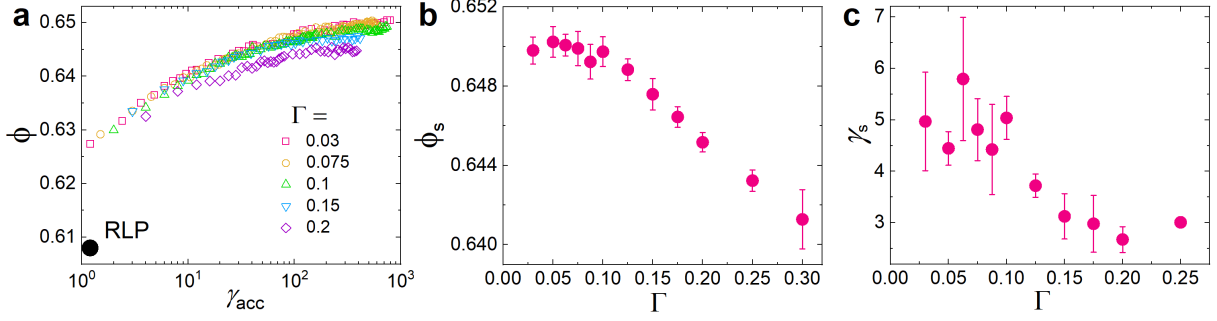
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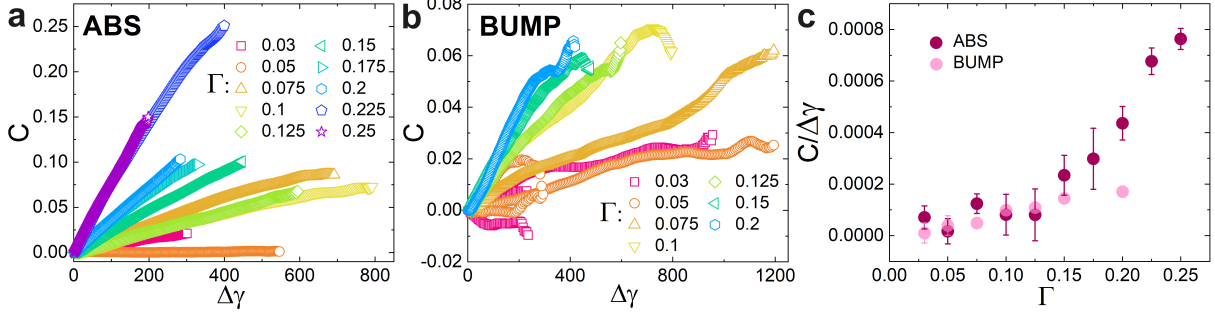
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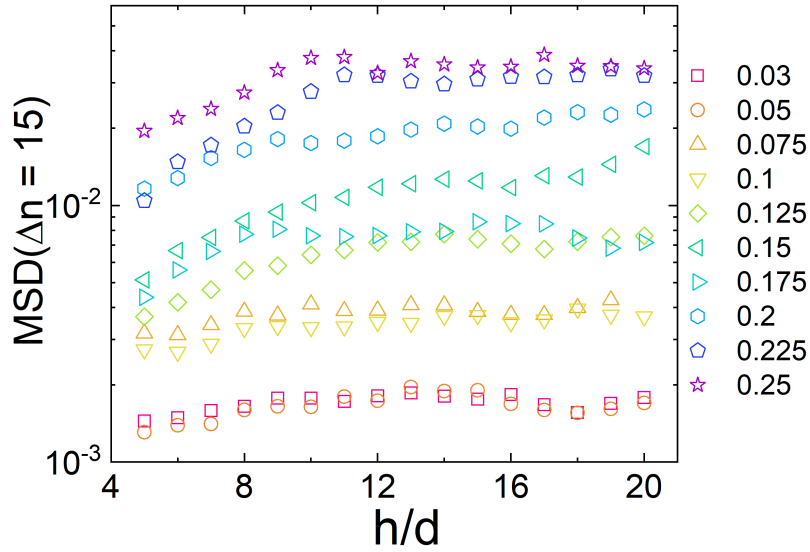
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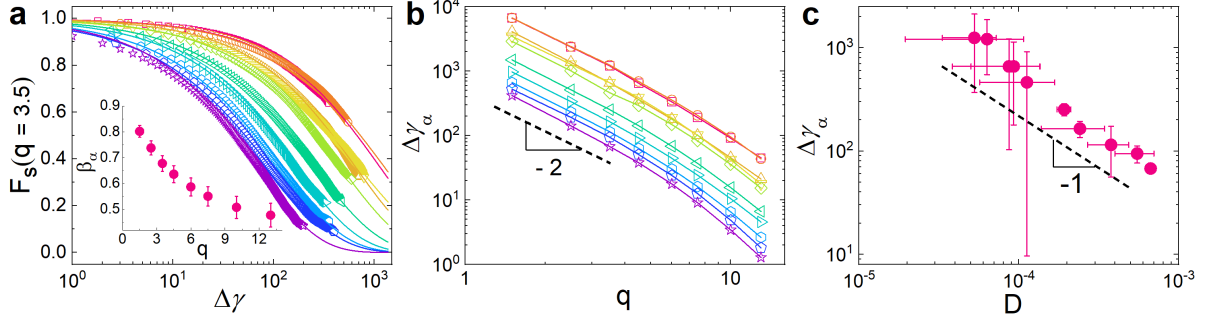
Supplementary Figure 1. **Transient process of compaction for the ABS system.** (a) Packing fraction  $\phi$  as a function of accumulated strain  $\gamma_{acc}$  (defined similarly to  $\Delta\gamma$  but now for a dynamics that is not in steady state) for different  $\Gamma$ , evolving from initially deposited packings, i.e., random loose packings (RLP, marked by black dot) with  $\phi_{RLP} \approx 0.608$ , towards the steady state. This process can be characterized by a stretched exponential law  $\phi = \phi_s + (\phi_{RLP} - \phi_s) \exp[-(\gamma_{acc}/\gamma_s)^\beta]$ , where  $\beta = 0.35$ ,  $\phi_s$  is the steady-state packing fraction, and  $\gamma_s$  is the compaction strain scale. (b) and (c) show, respectively,  $\phi_s$  and  $\gamma_s$  versus  $\Gamma$  and one observes a crossover at  $\Gamma \approx 0.1$ . Error bars represent the standard deviations from different realizations.



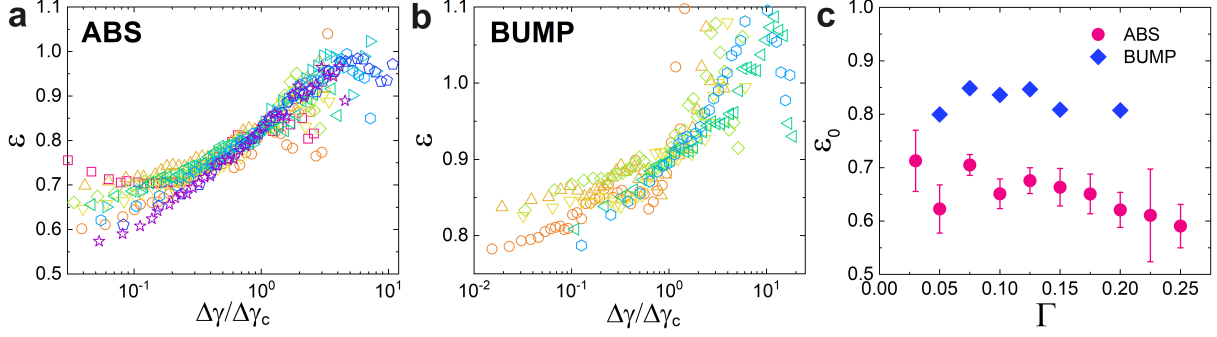
Supplementary Figure 2. **Convection strength in steady state for both systems.** (a) and (b): Up/down asymmetry in the particle motion in the two systems, characterized by  $C(\Delta\gamma) = \langle \delta z \rangle / \sqrt{\langle \delta x^2(\Delta\gamma) \rangle}$ , as a function of  $\Delta\gamma$ .  $C$  grows linearly with  $\Delta\gamma$  with slopes that depend on  $\Gamma$ , indicating an upward motion, i.e., convection. (c) The associated slope  $C/\Delta\gamma$  as a function of  $\Gamma$  shows a crossover at  $\Gamma \approx 0.1$  for the ABS system. This is further evidence that the dynamics changes at  $\Gamma \approx 0.1$  for ABS, allowing the convection to set in for large  $\Gamma$ . Error bars represent the standard deviations from different realizations.



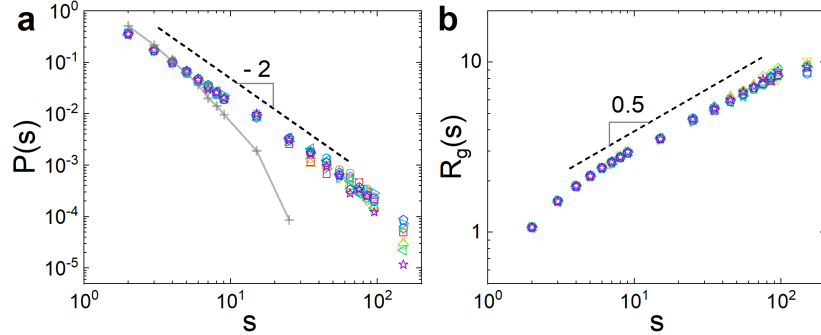
Supplementary Figure 3. **Height profile of MSD at  $\Delta n = 15$  for the ABS system.**  $h/d$  is the height from the cell bottom normalized by the small particle diameter, which starts from  $h/d = 4$  since we exclude the particles at the boundary. One finds only a very mild height dependence of the dynamics, in contrast to the clear  $\Gamma$ -dependence.



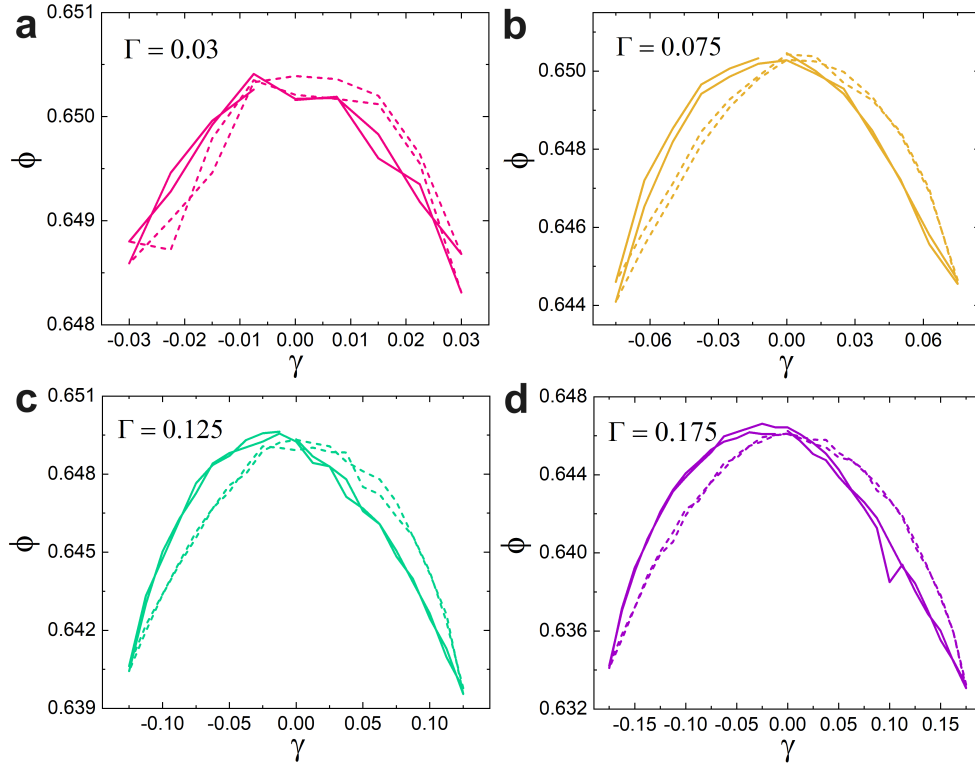
Supplementary Figure 4. **Self-intermediate scattering function  $F_s(q, \Delta\gamma)$  for the ABS system.** (a)  $F_s(q, \Delta\gamma)$  for  $q = 3.5$ , i.e., the first peak in the static structure factor, and different  $\Gamma$ . Solid curves indicate a stretched exponential fit  $F_s(q, \Delta\gamma) = \exp[-(\Delta\gamma/\Delta\gamma_\alpha)^{\beta_\alpha}]$ , where  $\Delta\gamma_\alpha$  is the relaxation strain scale. Inset: Stretching exponent  $\beta_\alpha$  as a function of  $q$ . (b)  $\Delta\gamma_\alpha$  as a function of  $q$  for different  $\Gamma$ . The scaling  $\Delta\gamma_\alpha \propto q^{-2}$  for small  $q$  (dashed line) indicates the Gaussian dynamics for large length scale, in agreement with the limit  $\beta_\alpha \rightarrow 1$  for  $q \rightarrow 0$  shown in the inset of panel (a). The color codes in (a) and (b) are the same as in Supplementary Fig. 2(a). (c)  $\Delta\gamma_\alpha \propto D^{-1}$  (dashed line) for  $q = 3.5$  and different  $\Gamma$ , indicating that the relaxation of  $F_s(q, \Delta\gamma)$  conveys similar information as the MSD presented in the main text. Error bars represent the standard deviations from different realizations.



Supplementary Figure 5. **Scaling of the mean squared displacement to determine the yielding strain  $\Delta\gamma_c$  for both systems.** (a) and (b): Local power-law exponent  $\varepsilon = d \log(\langle \delta x^2 \rangle) / d \log(\Delta\gamma)$  as a function of  $\Delta\gamma/\Delta\gamma_c$  for the ABS and BUMP systems, respectively. The color codes are respectively the same as in Figs. 2(a) and (b) of the main text. The yielding strain  $\Delta\gamma_c$  is defined by locating the crossover point  $\varepsilon = 0.825$  (or 0.9) for ABS (or BUMP), i.e., a value in the middle between the sub-diffusive exponent at small strain and 1.0 in the large strain limit. (c) The estimated sub-diffusive exponents from averaging  $\varepsilon(\Delta\gamma/\Delta\gamma_c \lesssim 0.1)$  for both systems, which confirms that they are independent of  $\Gamma$ . Error bars represent the standard deviations from different realizations for the ABS system.



Supplementary Figure 6. **Cluster characterization of the most mobile particles for the ABS system.** We apply the same criterion as in Fig. 4(c) (see main text) to define clusters of the top 10% fastest particles. Color codes are the same as in Supplementary Fig. 2(a). (a) Cluster size distribution  $P(s)$  for different  $\Gamma$  at the yielding point, i.e.,  $\Delta\gamma = \Delta\gamma_c$ . The dashed line indicates an approximately universal scaling  $P(s) \propto s^{-2}$ .  $P(s)$  from the randomly chosen 10% (line with crosses) deviates strongly from this master curve. (b) Also the average radius of gyration  $R_g(s)$  of the clusters as a function of  $s$  shows a universal scaling  $R_g(s) \propto s^{0.5}$  (dashed line).



Supplementary Figure 7. **Evolution of the packing fraction during the shear cycle for the ABS system.** Data for  $\Gamma = 0.03, 0.075, 0.125,$  and  $0.175$  are shown in (a-d), and each panel covers two consecutive cycles after the system has reached the steady state. For a better visualization, we use a solid curve for the forward direction of shear (increasing  $\gamma$ ) and a dashed curve for the backward direction. The persistent hysteresis found here resembles the findings of Ref. [45], indicating the absence of an elastic regime.