

# Verifying Ice-Sheet Models

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- 1 Why to verify and how to verify a model?
- 2 Verification of the SIA models
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# Model verification and validation

Numerical computer codes for ice sheet flow emerge from two stages of effort:

specification of a continuum model (nonlinear PDEs)

modeling errors arise from not solving the right equations → assessment of modeling errors is called **model validation**.

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the numerical approximation of the model (because of the difficulty of solving the above PDEs exactly)

numerical errors arise from not solving the equations right → assessment of numerical errors is called **model verification**.

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# Ways to verify the ice sheet models

**intercomparison of models** - measuring differences among various models' results on the sets of simplified geometry benchmark tests.

- compares models on variety of tests that resemble real ice sheet geometries;
- provides a set of standards for a modelers;
- examples: EISMINT I, EISMINT II, Ross Ice shelf, ISMIP-HOM, ISMIP-HEINO, ISMIP-POLICE, MISMIP.

**verification by exact solution** - measuring differences between model results and (may be artificially constructed) exact solutions.

- allows modelers check correctness of a code and to *estimate magnitude of numerical errors on a given grid*;
- *allows to measure convergence of numerical methods*;
- *allows tests codes for a variety of cases including different boundary conditions.*

## Example of building a manufactured exact solution (Bueler, 2006)

- completely made-up PDE:

$$\frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2} + u^2$$

is hard to find any exact solutions

- but one can find such for a slightly more general PDE:

$$\frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2} + u^2 + f(x, t) \quad (1)$$

- for *example*, let  $u(x, t) = x^3 + t$ ; compute

$$f = \frac{\partial u}{\partial t} - \frac{\partial^2 u}{\partial x^2} - u^2 = 1 - 6x - (x^3 + t)^2$$

- with this  $f$ , equation (1) has  $u = x^3 + t$  as solution.

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# Verification of the SIA models

## intercomparison of models

- EISMINT I - isothermal nonsliding ice flow (Huybrechts et al., 1996)
  - EISMINT II - thermocoupled nonsliding ice flow (Payne et al., 2000)
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## verification by manufactured exact solutions

- isothermal nonsliding SI (Halfar 1983, Bueler et al 2005)
  - thermocoupled nonsliding SI (Bueler et al 2005)
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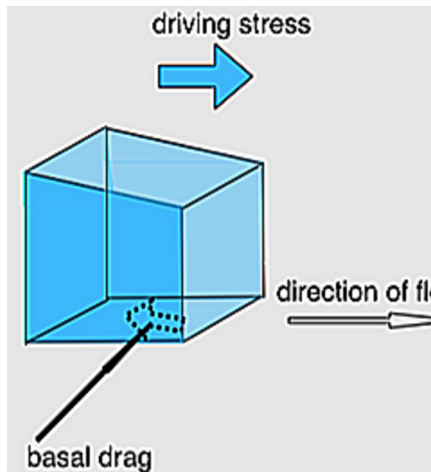
# Shallow Ice Approximation (SIA)

**Reference:** K. Hutter, Theoretical Glaciology. Dordrecht, Kluwer Academic Publishers, 1983.

**Assumptions:** Longitudinal and transverse stresses are neglected.

**Numerics:** Quasi-2D model – 1 degree of freedom per node.

**Conclusion:** Valid only for an ice mass with a small aspect ratio (*ice thickness*  $\ll$  *ice horizontal dimensions*)



**Figure:** Force Balance for Shallow-Ice Approximation

# Ice Sheet Equations of the SIA

conservation of mass

$$h_t = M - \nabla \cdot (\bar{U}h) \quad (2)$$

get velocity in SIA by vertically-integrating this:

$$\bar{U}(z) = -2(\rho g)^n |\nabla h|^{n-1} \nabla h \int_b^z \left( \frac{h - \xi}{A(T^*)} \right)^n d\xi + \bar{U}(b) \quad (3)$$

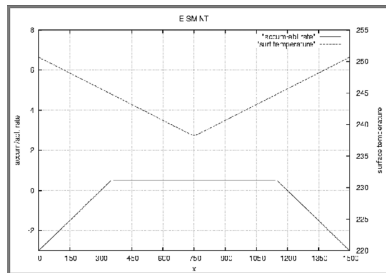
conservation of energy

$$\rho c_p (T_t + \bar{U} \cdot \nabla T) = k \nabla^2 T + (\sigma_{xz}, \sigma_{yz}) \cdot \frac{\partial \bar{U}}{\partial z} \quad (4)$$

where  $\bar{U}$  is vertically averaged horizontal velocity.

# EISMINT II (European Ice Sheet Modeling Initiative)

- intercomparison of 10 ice-sheet models on a series of experiments.
- a circular ice sheet is used and steady states and responses to stepped changes in climate are investigated.
- Exp. A: starting from zero ice, ice accumulation/ablation rate and ice-surface temperature are fixed as functions of geographical position:



$$M(x, y) = \min \left[ M_{max}, S_b \left( R_{el} - \sqrt{(x - x_{summit})^2 + (y - y_{summit})^2} \right) \right]$$
$$T_{surface}(x, y) = T_{min} + S_T \sqrt{(x - x_{summit})^2 + (y - y_{summit})^2}$$

where  $M_{max}$  is the maximum accumulation rate;  $R_{el}$  is a distance from the summit  $(x_{summit}, y_{summit})$  where the accumulation rate becomes zero;  $S_b$  is the gradient of accum. rate change with horizontal distance;  $T_{min}$  is the minimum surface temperature;  $S_T$  is the gradient of air temperature change with horizontal distance.

# EISMINT: Experiment A

- wide range of results;
- how to estimate magnitude of numerical errors for a particular model?

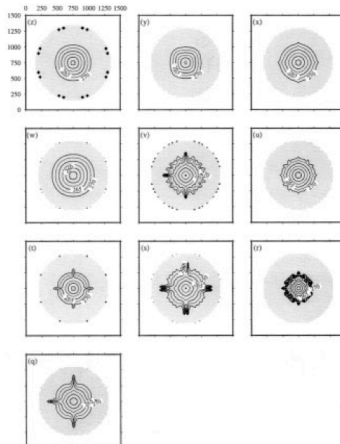


Figure: Predicted Steady-state basal temperatures in Exp. A (from EISMINT II)

# EISMINT: Experiment B

- Use as initial condition the final steady-state ice sheet of Exp. A (constant  $T_{min} = 238.15K$ ) and
- surface temperature experiences  $5K$  warming ( $T_{min} = 243.15K$ ).

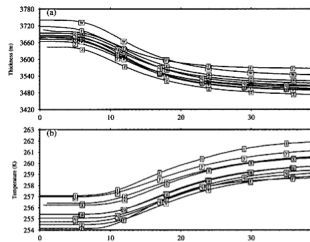


Figure: from EISMINT II: Time series of thickness and basal

# EISMINT: Experiment F

- Use as initial condition the final steady-state ice sheet of Exp. A (constant  $T_{min} = 238.15K$ ) and
- surface temperature experiences  $15K$  cooling ( $T_{min} = 223.15K$ )
- Are the spokes (in exp. F) just numerical errors? No, they reflect a sensitivity of the continuum equations to perturbation in some geometry/temperature regimes (Bueler, Hindmarsh).

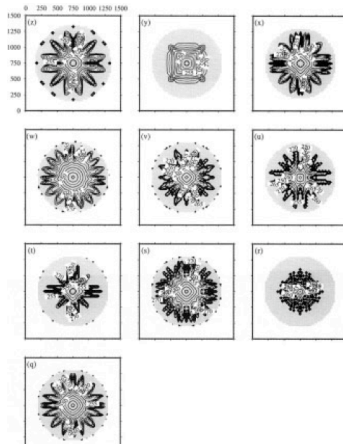
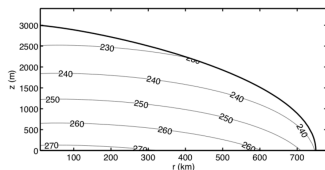


Figure: Predicted Steady-state basal temperatures in Exp. F (from EISMINT)

# Verification of the SIA model using manufactured solution (Bueler)

## Exact solution to thermocoupled SIA



$h, T$  chosen (circular ice caps like EISMINT)  $\rightarrow$  compute accumulation, velocity, etc. which satisfy all equations.

## Exact solution formulas

$$h(r, t) = h(r) + \phi(r)\gamma(t), \quad T(r, t, z) = T(r) \frac{\nu(t, r) + h(r, t)}{\nu(t, r) + z}, \quad \text{where}$$

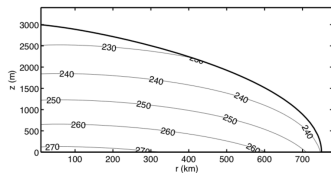
$$h(r) = \frac{h_0}{\left(1 - \frac{1}{n}\right)^{\frac{n}{2n+2}}} \left[ \left(1 + \frac{1}{n}\right) s - \frac{1}{n} + (1-s)^{1+\frac{1}{n}} - s^{1+\frac{1}{n}} \right]^{\frac{n}{2n+2}},$$

$$\phi(r) = \cos^2 \left( \frac{\pi(r - 0.6L)}{0.6L} \right), \quad \gamma(r) = A_r \sin \left( \frac{2\pi t}{t_p} \right),$$

$$\nu(t, r) = \frac{kT(r)}{2G} \left( 1 + \sqrt{1 + 4 \frac{h(t, r)G}{kT(r)}} \right), \quad s = r/L.$$

# Verification of the SIA model using manufactured solution (Bueler)

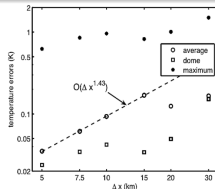
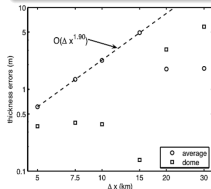
## Exact solution to thermocoupled SIA



$h$ ,  $T$  chosen (circular ice caps like EISMINT)  $\rightarrow$  compute accumulation, velocity, etc. which satisfy all equations.

## Useful for

- calculation of numerical errors ( $h$  and  $T$ ) and
- measuring convergence rate under grid refinement





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# Higher-order (HO) and full-Stokes (FS) 3-D models

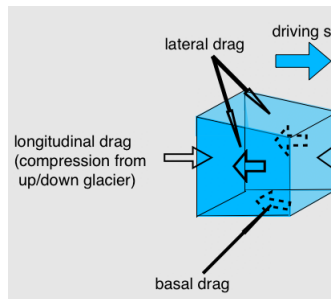
**Reference:** F. Pattyn, Investigating the stability of subglacial lakes with a full-Stokes ice-sheet model, ... 2008.

R.C.A.Hindmarsh, A numerical comparison of approximation to the Stokes equations used in ice sheet and glacier modeling, ... 2004.

**Assumptions:** Higher-order – variational stresses are neglected,  
full-Stokes – all stresses are included.

**Numerics:** Higher-order – 2D models;  
full-Stokes – 3D models.

**Conclusion:** Better predictions but computationally intensive.



**Figure:** Force Balance

# Verification of the full-Stokes and HO models

## intercomparison of models

- ISMIP-HOM - isothermal flow (Pattyn et al, 2008)
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## verification by exact solution

- quasi-analytical solutions for the 1-st order approximation equations (Blatter 1995)
  - analytical solutions for transient 2-D flow (Hutter 1980, Hutter 1983)
  - 3-D solution of the linearized 0-th order problem (Gudmundsson 2003)
  - manufactured solutions of a steady-state isothermal 2-D and 3-D flow (Fastook and Sargent)
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# ISMIP-HOM

- Intercomparison of 28 full-Stokes and HO models.
- 6 experiments (2-D and 3-D geometries): 5 - steady-state (Glen-type flow law), 1 - time-dependent (constant viscosity);
- 1 experiment with data from Haut Glacier d'Arolla;
- Isothermal ice mass;
- Periodic lateral boundary conditions.

# ISMIP-HOM: Experiment B

- 2-D: steady-state ice flow over a rippled bed;
- Boundary conditions: frozen bed, stress-free surface, periodic lateral;
- The surface elevation and the bed topography are defined as:

$$s(x, y) = -x \cdot \tan \alpha,$$

$$b(x, y) = s(x, y) - 1000 + 500 \sin(\omega x),$$

where  $\omega = 2\pi/L$ ,  $L$  is the ice length.

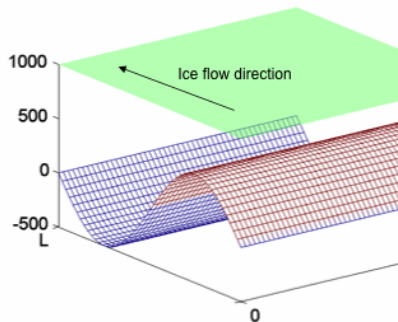
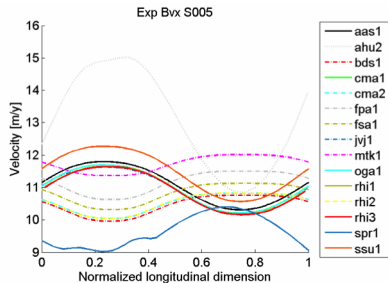


Figure: Experiment B - from Pattyn, 2008.

# ISMIP-HOM: Conclusions (from Pattyn, 2008)

- Benchmarks work better for longer length scales than for smaller;
- However, interesting features appear at smaller length scales ( $L = 5$ ): distinction between FS and HO models;
- Differences between models are due to either physical approximations or numerical problems/inaccuracies.

▶ FS verification



**Figure:** Experiment B: clear distinction in behavior between HO and FS models (Pattyn, 2008).

# Steady-state isothermal 2-D flowline model

conservation of mass

$$\frac{\partial u}{\partial x} + \frac{\partial w}{\partial z} = 0, \quad (5)$$

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kinematic boundary conditions

$$u(x, s(x)) \frac{ds}{dx} - w(x, s(x)) = \dot{a}, \quad (6)$$

$$u(x, b(x)) \frac{db}{dx} - w(x, b(x)) = 0, \quad (7)$$

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# Steady-state isothermal 2-D flowline model

## conservation of momentum

$$\frac{\partial (2\mu \frac{\partial u}{\partial x} + p)}{\partial x} + \frac{\partial (\mu (\frac{\partial u}{\partial z} + \frac{\partial w}{\partial x}))}{\partial z} = 0, \quad (8)$$

$$\frac{\partial (\mu (\frac{\partial w}{\partial x} + \frac{\partial u}{\partial z}))}{\partial x} + \frac{\partial (2\mu \frac{\partial w}{\partial z} + p)}{\partial z} = \rho g, \quad (9)$$

## boundary conditions

- stress-free surface;
- frozen bed or sliding bed;
- lateral bc: periodic or Dirichlet.

## constitutive relation (Glen's ice flow law)

$$\mu = \frac{B}{2} \left( \frac{1}{2} \left( \frac{\partial u}{\partial z} + \frac{\partial w}{\partial x} \right)^2 - \frac{\partial u}{\partial x} \frac{\partial w}{\partial z} \right)^{\frac{1-n}{2n}}, \quad (10)$$



## Derivation of a manufactured exact solution

Let's assume that in the domain  $s(x) > b(x)$ :

$$w(x, z) = u(x, z) \left( \frac{db}{dx} \frac{s-z}{s-b} + \frac{ds}{dx} \frac{z-b}{s-b} \right) - \dot{a} \frac{z-b}{s-b}. \quad (11)$$

then

- 1 kinematic boundary conditions are satisfied and
- 2 conservation of mass equation is reduced to the equation of one variable  $u(x, z)$ :

$$\frac{\partial u}{\partial x} + \frac{\partial u}{\partial z} \left( \frac{db}{dx} \frac{s-z}{s-b} + \frac{ds}{dx} \frac{z-b}{s-b} \right) + u \frac{\frac{ds}{dx} - \frac{db}{dx}}{s-b} - \frac{\dot{a}}{s-b} = 0. \quad (12)$$

This equation has a solution:

$$u(x, z) = \frac{1}{s(x) - b(x)} \vartheta \left( \frac{z - b(x)}{s(x) - b(x)} \right) + \frac{\dot{a}x}{s(x) - b(x)}, \quad (13)$$

where  $\vartheta$  is an arbitrary function of one variable.

## Satisfying the conservation of momentum equation and the stress-free surface boundary conditions

Substitution of the manufactured solution to the conservation of momentum equations and the surface boundary conditions will result in additional terms in the PDEs:

### conservation of momentum

$$\frac{\partial \left( 2\mu \frac{\partial u}{\partial x} + p \right)}{\partial x} + \frac{\partial \left( \mu \left( \frac{\partial u}{\partial z} + \frac{\partial w}{\partial x} \right) \right)}{\partial z} = \Sigma_x, \quad (14)$$

$$\frac{\partial \left( \mu \left( \frac{\partial w}{\partial x} + \frac{\partial u}{\partial z} \right) \right)}{\partial x} + \frac{\partial \left( 2\mu \frac{\partial w}{\partial z} + p \right)}{\partial z} - \rho g = \Sigma_z, \quad (15)$$

### boundary conditions

- stress-free surface bc (add an artificial term to the RHS);
- frozen bed bc – satisfied automatically;
- periodic side bc – satisfied automatically.

# Verification of the FS model using manufactured solution

Let's, for simplicity, assume that function  $\vartheta$  is as follows:

$$\vartheta(x) = x^\lambda + c_b, \quad (16)$$

where  $\lambda \geq 2$  and  $c_b \geq 0$  are constants;  $c_b = 0$  for frozen-bed solutions; then

Exact solutions are

$$u(x, z) = \frac{1}{s(x) - b(x)} \left( \frac{z - b(x)}{s(x) - b(x)} \right)^\lambda + \frac{c_b}{s(x) - b(x)},$$
$$w(x, z) = u(x, z) \left( \frac{\partial b}{\partial x} \frac{s - z}{s - b} + \frac{\partial s}{\partial x} \frac{z - b}{s - b} \right)$$

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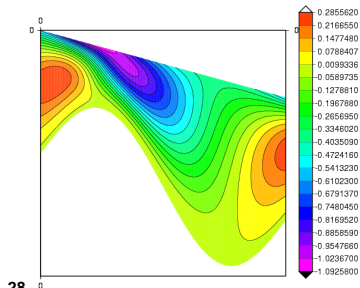
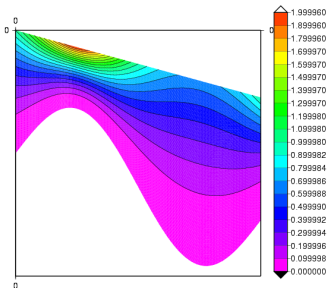
where  $s(x)$  and  $b(x)$  are the ice surface elevation and bed topography curves.

# Verification of the FS model using manufactured solution: EISMINT Experiment B type solution

If surface elevation and bed topography are defined as in Exp. B of ISMIP-HOM:

$$s(x, y) = -x \cdot \tan \alpha,$$
$$b(x, y) = s(x, y) - 1000 + 500 \sin(\omega x),$$

then the horizontal and vertical velocities are as follows:



# Verification of the FS using manufactured solution: How realistic is the solution?

- conservation of mass flux:  $q = hu = 1$  is satisfied:

$$\text{for } z = b, \quad u(x, b) = 0, \quad w(x, b) = 0;$$

$$\text{for } z = s, \quad u(x, s) = \frac{1}{s - b} = \frac{1}{h}, \quad w(x, s) = 0$$

- This anti-correlated relationship between  $u$  and  $h$  is consistent with the simulation of a Exp. B in ISMIP-HOM by *all* flowline full-Stokes models.

▶ Experiment B, EISMINT

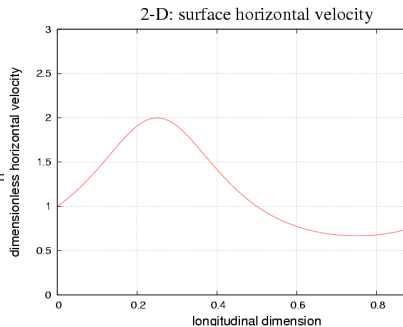











Figure: Surface horizontal velocity.

# Verification of the FS using manufactured solution

## Summary:

- Manufactured analytical solutions for 2-D steady-state isothermal flowline models are derived:
    - variable viscosity;
    - solutions can be specified for different surface and bed geometries;
    - solutions are periodic;
    - solutions are easy to use.
  - Similar manufactured solutions derived for 3-D full-Stokes ice flow model.
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Thank you for your attention!