

Coherent structure extraction in turbulent channel flow using boundary adapted wavelets[†]

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ABSTRACT

We introduce boundary adapted wavelets, which are orthogonal and have the same scale in the three spatial directions. The construction thus yields a multiresolution analysis. We analyse direct numerical simulation data of turbulent channel flow computed at a friction Reynolds number of 395, and investigate the role of coherent vorticity. Thresholding of the vorticity wavelet coefficients allows us to split the flow into two parts, coherent and incoherent flows. The coherent vorticity is reconstructed from its few intense wavelet coefficients and the coherent velocity is reconstructed using Biot–Savart's law. The statistics of the coherent flow, i.e. energy and enstrophy spectra, are close to the statistics of the total flow, and moreover, the non-linear energy budgets of the total flow are very well preserved. The remaining incoherent part, represented by the large majority of the weak wavelet coefficients, corresponds to a structureless, i.e. noise-like, background flow whose energy is equidistributed.

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1. Introduction

Wall-bounded turbulent shear flows are of general interest in many engineering applications. Three-dimensional (3D) turbulent channel flow, bounded by two parallel walls, is one of the canonical flows considered for direct numerical simulation (DNS). Starting with the seminal work of Kim et al. [1], many DNS have been performed for increasingly higher Reynolds number, taking advantage of the growing power of supercomputers (see, e.g. the review article [2]). Currently, the DNS with the highest friction-based Reynolds number, $Re_\tau = 5200$, has been carried out by Lee and Moser [3]. The influence of rough walls has been reviewed in Ref. [4].

Turbulent flows are typically characterised by the excitation of a multitude of spatial and temporal scales, which involves a large number of degrees of freedom interacting nonlinearly. Self-organisation of the flow into coherent vortices is observed, even at large Reynolds

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[†]This article is dedicated to Professor Javier Jimenez on the occasion of his 70th birthday.

number [5], where these vortices are superimposed to a random background flow [6]. Moreover, turbulence exhibits significant spatial and temporal intermittency, especially in the dissipative range. This implies that the strong contributions become sparser and sparser as scale becomes smaller in space and time. Wavelets being well-localised functions in both space and scale yield efficient multi-scale decompositions, which have been applied to analyse, model and compute turbulent flows since 1988 [7–10]. Decomposing turbulent flows into a wavelet basis yields a sparse representation, namely the most energetic contributions are concentrated in few wavelet coefficients having large intensity, while the large majority of the remaining wavelet coefficients have negligibly small intensity.

The presence of coherent structures superimposed to a random background flow motivated the development of the coherent vorticity extraction (CVE) method. The idea of CVE, proposed by Farge et al. [11,12], defines coherent structures as what remains after denoising the flow vorticity. Since vorticity is better localised in space than velocity, thus more intermittent, its wavelet decomposition is sparser and only few coefficients are necessary to represent the coherent structures. Moreover, in contrast to velocity, vorticity preserves Galilean invariance and has stronger topological properties owed to Helmholtz' and Kelvin's theorems. Numerous applications of CVE can be found for periodic domains in the literature starting with homogeneous isotropic turbulence [11–16], temporally developing mixing layers [17] and homogeneous shear flow with and without rotation [18].

For wall-bounded flows, the situation becomes more complex, because no-slip boundary conditions have to be taken into account. Indeed, no-slip boundary conditions generate vorticity due to the viscous flow interactions with the walls. For turbulent boundary layers, Khujadze et al. [19] obtained an efficient algorithm to extract coherent vorticity, constructing a locally refined grid using wavelets with mirror boundary conditions. However, this construction does not yield a multiresolution analysis, as the basis functions have mixed scales in the different spatial directions. Fröhlich & Uhlmann [20] constructed wavelets based on second kind Chebyshev polynomials and applied them to channel flow data. Scale-wise statistics in the wall-normal direction have thus been performed. However, no fast wavelet transform (FWT) is available for these Chebyshev wavelet bases. Two-dimensional (2D) wavelets have also been applied to wall-parallel planes in channel flows, in order to examine turbulent statistics, in particular statistics of energy transfer [21,22].

The aim of the present work is to examine the role of coherent and incoherent flow contributions in 3D turbulent channel flow. We propose a novel construction of 3D orthogonal wavelets using boundary wavelets in the wall-normal direction and periodic wavelets in the wall-parallel directions. To this end, Cohen-Daubechies-Jawerth-Vial (CDJV) boundary wavelets [23,24] having three vanishing moments, and the periodised Coiflet 30 wavelets [25] having 10 vanishing moments are employed. These wavelets are orthogonal, the FWT can be used while taking into account boundary conditions, and the basis functions have no mixed scales in the different spatial directions. Hence, the basis functions yield a multiresolution analysis with the same scale in the three directions.

DNS computation of the channel flow has been performed, and the data are analysed at different time instants, using the above boundary adapted 3D wavelets. The flow vorticity is decomposed into an orthogonal wavelet series, and we apply a thresholding to split the coefficients into two sets, the coherent and incoherent ones. The coherent vorticity, reconstructed from the few strongest wavelet coefficients, well preserves the turbulent statistics of the total flow, while the incoherent vorticity, reconstructed from the remaining large majority of the coefficients that are very weak, corresponds to a noise-like background flow. The

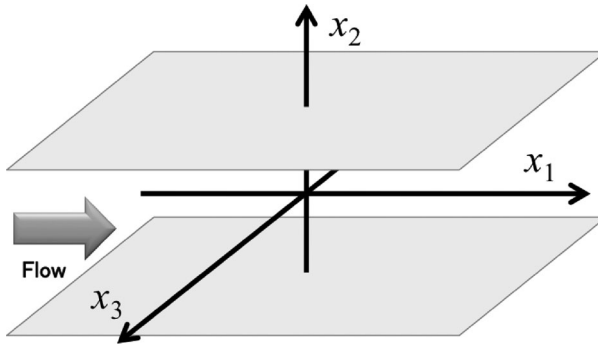


Figure 1. Flow configuration for the turbulent channel flow.

corresponding coherent and incoherent velocity fields are reconstructed from the coherent and incoherent vorticity fields, respectively, using the Biot–Savart relation satisfying the no-slip conditions at the walls. Thus, we can efficiently examine the role of coherent vorticity in turbulent channel flows. Other conventional methods, such as the Q -criterion and the λ_2 method [26,27], could be used to identify coherent vortices in physical space, as regions for which Q or λ_2 is above a given threshold. Here, Q is the second-invariant of the 3D velocity gradient tensor, and λ_2 is the second largest eigenvalue of $S_{ij}S_{jk} + A_{ij}A_{jk}$, where S_{ij} and A_{ij} are, respectively, the symmetric and antisymmetric tensor of the velocity gradient tensor. It should be noticed that these quantities do not preserve the scale information about the vortices, as the smoothness of the flow field is not preserved due to the clipping of vorticity in physical space. In contrast, the proposed wavelet filtering does preserve the smoothness of the coherent vorticity field and the multiscale properties of the coherent structures.

The paper is organised as follows: [Section 2](#) presents the DNS computation and the data we analyse, including the methodology. The construction of wavelets is described, and the CVE method is summarised. Numerical results are shown in [Section 3](#). Conclusions and perspectives are given in [Section 4](#).

2. DNS and methodology

2.1. Direct numerical simulation

We consider 3D incompressible fluid flow in a channel bounded by two parallel walls subjected to a streamwise mean pressure gradient, which is a canonical flow configuration. It is illustrated in [Figure 1](#) together with the Cartesian coordinate system $\mathbf{x} = (x_1, x_2, x_3)$, where the walls are at $x_2 = \pm h$, x_2 being the wall-normal direction and h the half width of the channel. The domain size in the streamwise x_1 -direction is $2\pi h$, and the size in the spanwise x_3 -direction is πh . Periodic boundary conditions are, respectively, imposed in x_1 - and x_3 -directions, while in the x_2 -direction no-slip boundary conditions are satisfied at the walls.

The fluid flow motion obeys the Navier–Stokes equations with the incompressibility condition,

$$\partial_t v_i + \partial_j (v_j v_i) = -\partial_i p + G\delta_{i1} + \nu \partial_j \partial_j v_i, \quad (1)$$

$$\partial_j v_j = 0, \quad (2)$$

where v_i ($i = 1, 2, 3$) is the i -th velocity component, p is the pressure fluctuation, G is the intensity of the mean pressure gradient in the x_1 -direction, δ_{ij} is the Kronecker delta, ν is the kinematic viscosity, t is time, and $\partial_t = \partial/\partial t$ and $\partial_i = \partial/\partial x_i$. Einstein's summation convention is used for repeated indices.

We performed DNS of turbulent channel flow at Re_τ of 395 using $N_1 N_2 N_3$ grid points, where $Re_\tau \equiv u_\tau h/\nu$, u_τ is the friction velocity defined by $[\nu dU_1(x_2)/dx_2]^{1/2}$ at $x_2 = -h$. The velocity field v_j is decomposed as $v_j = U_j(x_2) + u_j$ with U_j being the mean velocity defined as $U_j = \langle v_j \rangle$, and u_j are the velocity fluctuations. Here, $\langle \cdot \rangle$ denotes the x_2 -dependent spatial average of \cdot over the x_1 - x_3 plane, and N_i is the number of the grid points in the x_i -direction, $N_1 = N_3 = 256$ and $N_2 = 192$. The toroidal and poloidal representation of Equations (1) and (2) is employed, in order to satisfy the incompressibility constraint as done by Kim et al. [1]. We used the Fourier pseudo-spectral method in the x_1 - x_3 planes, and the Chebyshev-tau method in the x_2 -direction. The Chebyshev collocation points are given by $x_2 = h \cos \{\pi(2j+1)/(2N_2)\}$ ($j = 0, 1, \dots, N_2 - 1$) (see, e.g. Appendix B in Ref. [28]). The aliasing errors are removed by the 3/2 rule in the x_1 - x_3 planes, and by the 2/3 rule in the x_2 -direction. Time advancement is carried out using the first-order implicit Euler method for the viscous terms, and a third-order Runge-Kutta method for the nonlinear terms and the mean pressure gradient term G , whose value is determined so that the total flow rate is kept constant. The DNS code has been developed in Ref. [29].

Statistical quantities shown in this paper are obtained by time averaging over 40 DNS snapshots with intervals of 0.5 washout times, defined by $2\pi h/U_1$. The averaging starts after the total Reynolds stress $-\langle u_1 u_2 \rangle + \nu dU_1/dx_2$ has become quasi-stationary.

2.2. Wavelets

In this subsection, we briefly summarise one-dimensional (1D) orthogonal periodised wavelets and 1D orthogonal boundary wavelets. Then, we propose a 3D orthogonal wavelet transform with one scale in the three spatial directions constructed by tensor product of these 1D wavelets. The CVE based on orthogonal wavelets to extract coherent vorticity out of turbulent channel flow is described in Section 2.3. In Figure 2, we present the flowchart of the CVE method used here.

We first consider 1-periodic wavelets $\psi^P(x)$ and their corresponding scaling function $\phi^P(x)$, and orthogonal boundary adaptive wavelets $\psi^B(x)$ and their scaling function $\phi^B(x)$, with the boundaries at $x = (0, 1)$. Wavelets at scale j are obtained by dilation, so that $\psi_{j,0}^\gamma(x) = 2^{j/2} \psi^\gamma(2^j x)$ and $\phi_{j,0}^\gamma(x) = 2^{j/2} \phi^\gamma(2^j x)$, where $\gamma = P, B$. The periodised orthogonal wavelets are also self-similar with respect to translation. Then the scaling function ϕ^γ and wavelet function ψ^γ at scale 2^{-j} ($j \geq 0$) and position $2^{-j}i$ ($i = 0, 1, \dots, 2^j - 1$), $\phi_{j,i}^P(x)$ and $\psi_{j,i}^P(x)$, are defined as $\phi_{j,i}^P(x) = 2^{j/2} \phi^P(2^j x - i)$ and $\psi_{j,i}^P(x) = 2^{j/2} \psi^P(2^j x - i)$. In contrast, $\psi_{j,0}^B$ and $\phi_{j,0}^B$ are no more translation invariant due to the boundary conditions, which modify the wavelets as position i changes. Readers interested in the details of boundary adapted wavelets may refer to the textbook [30], as the construction of $\psi_{j,i}^B$ and $\phi_{j,i}^B$ is rather technical. All wavelets used here are orthonormal, i.e. $\int_0^1 \psi_{j,i}^\gamma(x) \psi_{j',i'}^\gamma(x) dx = \delta_{ii'} \delta_{jj'}$, $\int_0^1 \phi_{j,i}^\gamma(x) \phi_{j',i'}^\gamma(x) dx = \delta_{ii'}$, and $\int_0^1 \psi_{j,i}^\gamma(x) \phi_{j',i'}^\gamma(x) dx = 0$.

In this paper, we use Coiflet 30 wavelets [25] in the x_1 - and x_3 -directions, and the CDJV wavelets having three vanishing moments [23,24] in the x_2 -direction, both wavelets being

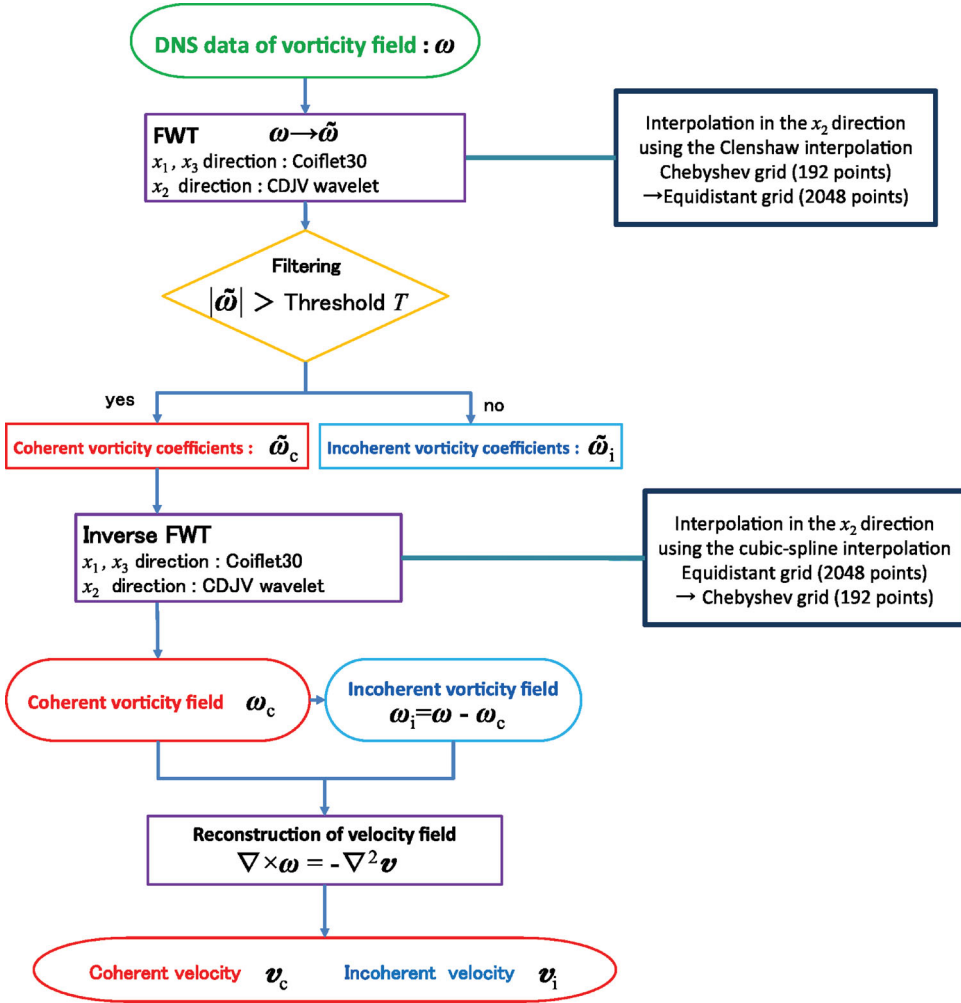


Figure 2. Flowchart of the CVE procedure.

compactly supported. The Coiflet 30 wavelets are quasi-symmetric and have 10 vanishing moments. The largest scale 2^{-j_0} of the CDJV wavelets satisfies $2^{j_0-1} \geq 3$ [24]. The illustrations of these wavelet functions are shown in Figures 3 and 4.

The 3D orthogonal wavelets Ψ^μ ($\mu = 1, 2, \dots, 7$) are obtained by tensor product such that

$$\Psi_{j,i}^\mu(x_1, x_2, x_3) = \begin{cases} \psi_{j,i_1}^P(x_1)\phi_{j',i_2}^B(x_2)\phi_{j,i_3}^P(x_3) & \text{for } \mu = 1, \\ \phi_{j,i_1}^P(x_1)\psi_{j',i_2}^B(x_2)\phi_{j,i_3}^P(x_3) & \text{for } \mu = 2, \\ \phi_{j,i_1}^P(x_1)\phi_{j',i_2}^B(x_2)\psi_{j,i_3}^P(x_3) & \text{for } \mu = 3, \\ \psi_{j,i_1}^P(x_1)\phi_{j',i_2}^B(x_2)\psi_{j,i_3}^P(x_3) & \text{for } \mu = 4, \\ \psi_{j,i_1}^P(x_1)\psi_{j',i_2}^B(x_2)\phi_{j,i_3}^P(x_3) & \text{for } \mu = 5, \\ \phi_{j,i_1}^P(x_1)\psi_{j',i_2}^B(x_2)\psi_{j,i_3}^P(x_3) & \text{for } \mu = 6, \\ \psi_{j,i_1}^P(x_1)\psi_{j',i_2}^B(x_2)\psi_{j,i_3}^P(x_3) & \text{for } \mu = 7, \end{cases} \quad (3)$$

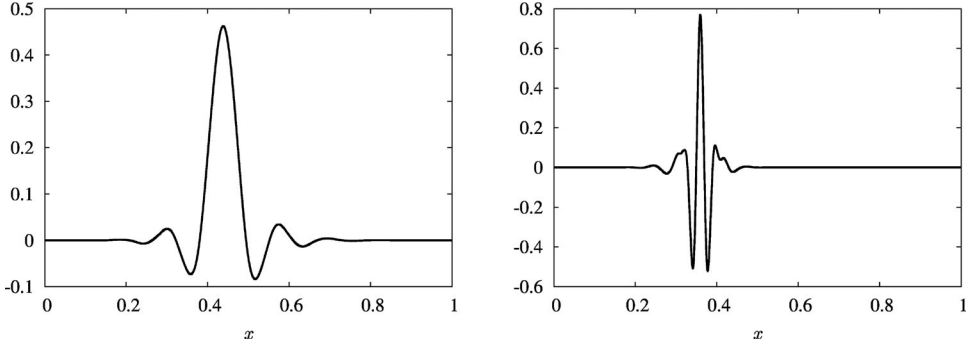


Figure 3. Coiflet 30 wavelet on the periodic domain: scaling function $\phi_{8,i}^P(x)$ (left) and corresponding wavelet $\psi_{8,i}^P(x)$ (right), both at scale $j = 8$.

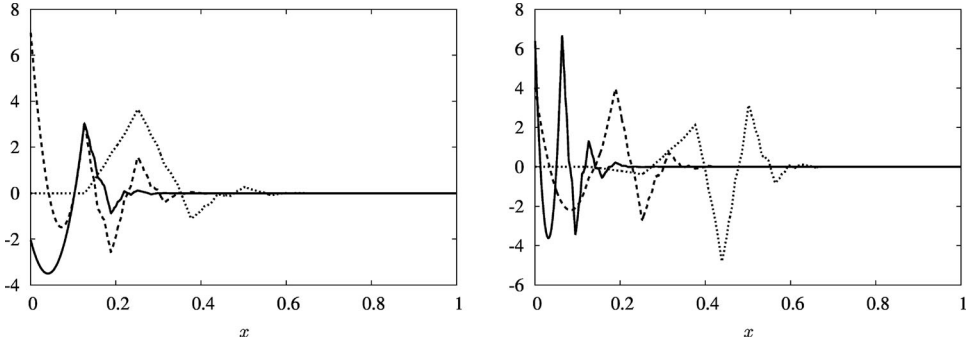


Figure 4. CDJV wavelet on the interval: three scaling functions $\phi_{8,i}^B(x)$ (left) and three wavelets $\psi_{8,i}^B(x)$ (right) at scale $j = 8$ and position $i = 0$ (solid line), 63 (dashed line) and 127 (dotted line) are shown.

where $\mathbf{i} = (i_1, i_2, i_3)$, $j' = j_0 + j$ and $j = 0, \dots, J-1$. The corresponding scaling function is defined as $\Phi(x_1, x_2, x_3) = \phi^P(x_1)\phi^B(x_2)\phi^P(x_3)$.

Now, let us consider a 3D vector field $\mathbf{w}(\mathbf{x}) = (w_1, w_2, w_3)$ in the computational domain D , where $D = \{x_1, x_2, x_3 | 0 \leq x_1 \leq 2\pi h, -h \leq x_2 \leq h, 0 \leq x_3 \leq \pi h\}$. Before applying this wavelet decomposition, we interpolate $\mathbf{w}(\mathbf{x})$ on an equidistant grid in the x_2 -direction from the DNS data non-uniformly sampled on N_2 Chebyshev grid points in the wall-normal direction. We thus get $\mathbf{w}(\mathbf{x})$ uniformly sampled on N'_2 equidistant grid points at $x_{2,n} = h\{-1 + 2n/(N'_2 - 1)\}$ ($n = 0, \dots, N'_2 - 1$) using the Chebyshev interpolation [31]. We choose N'_2 to be equal to 2048 so that the flow field near the walls is kept well-resolved. We have $2/N'_2 \sim 8\pi^2/N_2^2$, which shows that the grid width after the interpolation is comparable to the minimum grid width of the Chebyshev grid. In the x_1 - and x_3 -directions, we keep $\mathbf{w}(\mathbf{x})$ uniformly sampled on N_1 and $N_3 (= N_1)$ equidistant grid points, respectively.

The field $\mathbf{w}(\mathbf{x})$, now sampled on $N_1 \times N'_2 \times N_3$ equidistant grid points, can then be decomposed into an orthogonal wavelet series as follows;

$$\mathbf{w}(\mathbf{x}) = \bar{\mathbf{w}} + \sum_{\mu=1}^7 \sum_{j=0}^{J-1} \sum_{i_1, i_2, i_3=0}^{2^j-1} \tilde{\mathbf{w}}_{j,i}^{\mu} \Psi_{j,i}^{\mu}(\mathbf{x}), \quad (4)$$

with wavelet coefficients computed with wavelets $\Psi_{j,i}^\mu$

$$\tilde{w}_{j,i}^\mu = \frac{1}{\mathcal{V}} \int_0^{L_1} dx_1 \int_{-h}^h dx_2 \int_0^{L_3} dx_3 w(\mathbf{x}) \Psi_{j,i}^\mu \left(\frac{x_1}{2\pi h}, \frac{x_2 + h}{2h}, \frac{x_3}{\pi h} \right), \quad (5)$$

and the mean value computed with the scaling function Φ

$$\bar{w} = \frac{1}{\mathcal{V}} \int_0^{2\pi h} dx_1 \int_{-h}^h dx_2 \int_0^{\pi h} dx_3 w(\mathbf{x}) \Phi \left(\frac{x_1}{2\pi h}, \frac{x_2 + h}{2h}, \frac{x_3}{\pi h} \right), \quad (6)$$

where $J = \log_2 N_1$ and $\mathcal{V} = 4\pi^2 h^3$.

2.3. Coherent vorticity extraction

We extract coherent vorticity out of turbulent channel flow data using the CVE method based on the wavelet decomposition of vorticity $\boldsymbol{\omega} = \nabla \times \mathbf{v}$. In the following, we summarise our method. Since coherent structures do not have a universal definition yet, we define them as what remains after denoising. As first guess we consider the simplest type of noise, namely additive, Gaussian and white, i.e. uncorrelated noise. Readers interested in details of this ansatz may refer to the original articles, e.g. Refs. [11,12,16].

The CVE method is based on nonlinear thresholding of the orthogonal wavelet coefficients of vorticity. To this end, the vorticity $\boldsymbol{\omega}$, interpolated on a sufficiently fine equidistant grid, is decomposed into an orthogonal wavelet series using the FWT. Applying thresholding to the wavelet coefficients, we split the flow into coherent and incoherent contributions. The corresponding coherent and incoherent vorticity fields are then obtained by inverse wavelet transform.

In previous work, we used Donoho's threshold [32] to determine the value of the threshold and estimate the variance of the incoherent vorticity using an iterative scheme. Azzaoui et al. [33] investigated the convergence of the iterative scheme and for isotropic turbulence Okamoto et al. [16] found that, depending on the Reynolds number, 8.7% and 6.0% of the wavelet coefficients are retained as coherent for $Re_\lambda = 167$ and $Re_\lambda = 732$, respectively. In Ref. [11], Farge et al. used one iteration only, which was sufficient to get good compression while preserving the statistics of the total flow. For the turbulent channel flow studied here, we tried Donoho's threshold and found that very few wavelet coefficients keep almost the whole enstrophy of the flow, which is illustrated in the compression curve, shown in Figure 5. The flow visualisation in Figure 6 shows tube-like coherent vortex structures whose intensity is very strong close to the wall and much weaker in the centre of the channel. In the current work, we propose, instead of Donoho's threshold, an *ad hoc* criterion for the threshold defined by $T = \langle |\tilde{\boldsymbol{\omega}}_{j,i}^\mu| \rangle_w + \alpha \langle (|\tilde{\boldsymbol{\omega}}_{j,i}^\mu| - \langle |\tilde{\boldsymbol{\omega}}_{j,i}^\mu| \rangle)^2 \rangle_w^{1/2}$, where $\langle |\tilde{\boldsymbol{\omega}}_{j,i}^\mu| \rangle_w = \sum_{\mu=1}^7 \sum_{j=0}^{J-1} \sum_{i_1, i_2, i_3=0}^{2^j-1} |\tilde{\boldsymbol{\omega}}_{j,i}^\mu| / (N_1 N_2 N_3)$. The field is first decomposed into an orthogonal wavelet series and split into two orthogonal contributions using wavelet thresholding. Hereafter, they are called coherent (wavelet coefficients whose modulus is above the threshold) and incoherent (the remaining weaker wavelet coefficients), in order to be consistent with the terminology of the pioneering work [11]. Our aim is to retain only those wavelet coefficients which are responsible for the nonlinear dynamics of the flow, even if the fully developed turbulent regime has not been yet reached. We set $\alpha = 0.75$ in the threshold

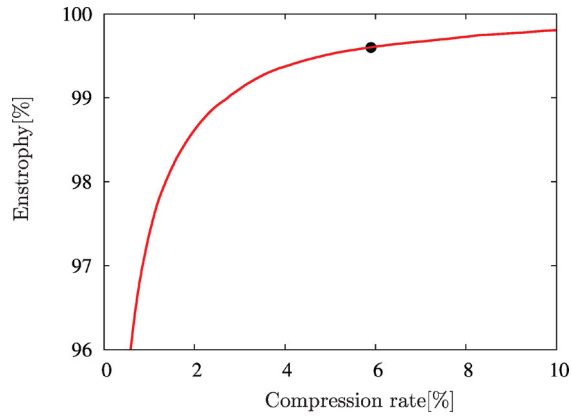


Figure 5. Compression curve for CVE: % of retained enstrophy per unit volume vs. % of retained wavelet coefficients. The circle corresponds to the threshold T used for CVE in the following.

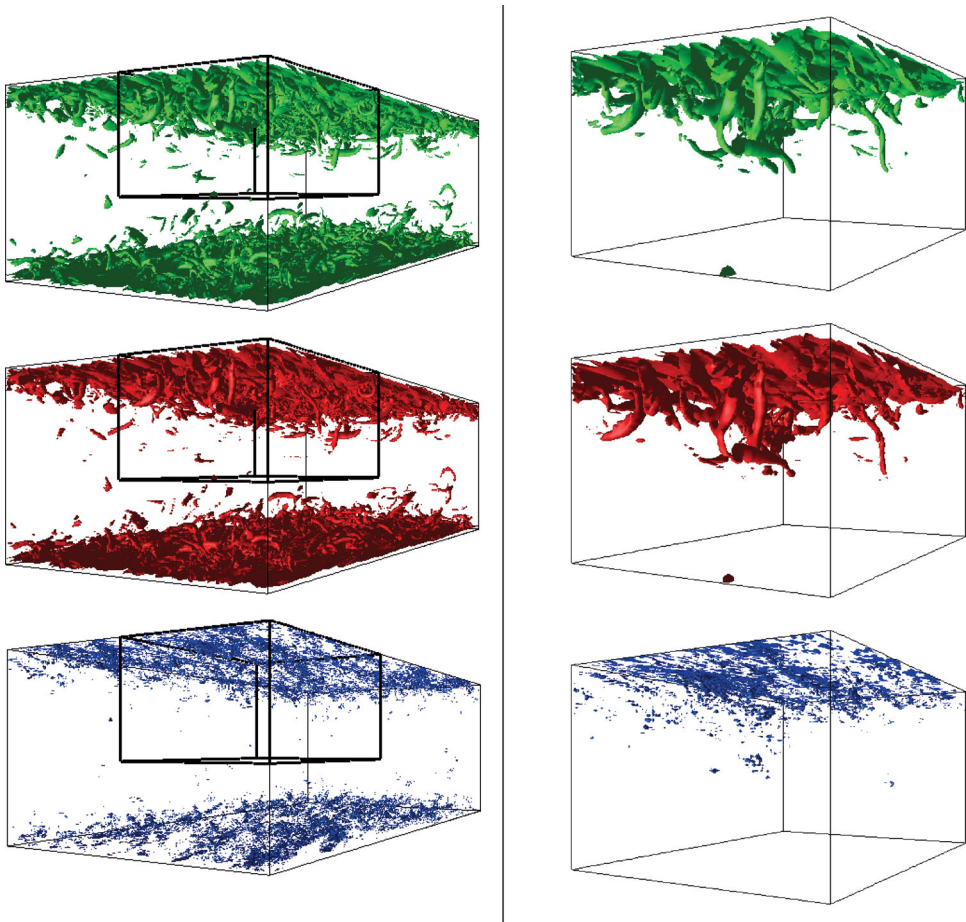


Figure 6. Visualisation of total vorticity ω (top), coherent vorticity ω_c (middle) and incoherent vorticity ω_i (bottom). The left column presents isosurfaces $|\omega^+| = |\omega_c^+| = 0.3$ and $|\omega_i^+| = 0.1$. The right column shows their zooms in the wall region where $0 \leq x_1 \leq 0.71\pi h$, $-0.12h \leq x_2 \leq h$, and $0.5\pi h \leq x_3 \leq \pi h$.

value T such that both velocity and vorticity statistics (as a function of x_2), together with the nonlinear dynamics and structures, are well preserved by the coherent flow. The *ad hoc* criterion of the threshold could be further improved.

Using the inverse FWT, the coherent vorticity ω_c is reconstructed from the wavelet coefficients whose intensity is larger than the threshold value T . The incoherent vorticity ω_i is then obtained using $\omega_i = \omega - \omega_c$. To get ω_c and ω_i sampled on the Chebyshev grid points, which are useful and efficient for the data analysis presented in Section 3, we perform a cubic spline interpolation in the x_2 -direction.

Owing to the orthogonality of the wavelet decomposition, ω_c is orthogonal to ω_i and thus $Z_t = Z_c + Z_i$, where Z_t , Z_c and Z_i are respectively the total, coherent and incoherent enstrophy per unit volume, defined as $Z_\alpha = \iiint_D |\omega_\alpha|^2 d\mathbf{x} / (2V)$ ($\alpha = t, c, i$). The coherent velocity \mathbf{v}_c and the incoherent velocity \mathbf{v}_i are computed from ω_c and ω_i by solving Biot–Savart’s relation, $\nabla^2 \mathbf{v} = -\nabla \times \omega$, respectively. It is noted that \mathbf{v}_i and \mathbf{v}_c are weakly non-orthogonal, i.e. the cross term $\int \mathbf{v}_i \cdot \mathbf{v}_c d\mathbf{x}$ is below 0.4% of the total energy. The extraction method could also be applied to the fluctuating vorticity instead of the total one used here. We checked that the results thus obtained are indeed very similar as those presented in Section 3.

3. Numerical results

Now we analyse 40 snapshots of DNS data for the turbulent channel flow with intervals of 0.5 washout times, and we ensemble-average over those 40 snapshots to guarantee well-converged statistical results. We examine contributions of coherent and incoherent flows obtained with the previously described CVE method. Quantities with the superscript $+$ are expressed in wall units, i.e. they are non-dimensionalised by u_τ and ν . We define the distance from the wall $y = x_2/h + 1$.

3.1. Visualisation

Visualisations of isosurface values of the modulus of vorticity for the total, coherent and incoherent flows given at the same time instant are shown in Figures 6 and 7. Corresponding zooms are also presented to see flow structures more clearly. Figure 6 shows that the most intense vorticity structures are near the walls. Since the incoherent vorticity is much weaker than the total and coherent vorticities in Figure 6, the isosurface value for the incoherent vorticity ω_i is reduced by a factor 3 compared to the coherent and total vorticities. On the other hand, Figure 7 visualises vorticity structures in the core region, using y^+ -dependent isosurface values, $|\omega^+| = |\omega_c^+| = \langle |\omega^+| \rangle + 3\langle (|\omega^+| - \langle |\omega^+| \rangle)^2 \rangle^{1/2}$ and $|\omega_i^+| = \langle |\omega_i^+| \rangle + 3\langle (|\omega_i^+| - \langle |\omega_i^+| \rangle)^2 \rangle^{1/2}$, recalling $\langle \cdot \rangle$ denotes the y^+ -dependent spatial average of \cdot over each wall-parallel plane.

We observe that the total flow exhibits intense vortex tubes near the walls, as in previous DNS (e.g. Ref. [34]), but we also see them in the core region, however they are less intense. Looking at the coherent flow, we find that these tubes are well preserved by ω_c , which is reconstructed from only 0.55% of the $256^2 \times 2048 (\simeq 13 \times 10^7)$ wavelet coefficients, i.e. 5.9% of the original $256^2 \times 192 (\simeq 1.2 \times 10^7)$ grid points. The coherent flow retains almost all of the total energy and enstrophy, 99.9% of the total energy and 99.7% of the total enstrophy. In contrast, the incoherent vorticity ω_i looks less organised without exhibiting vortex tubes

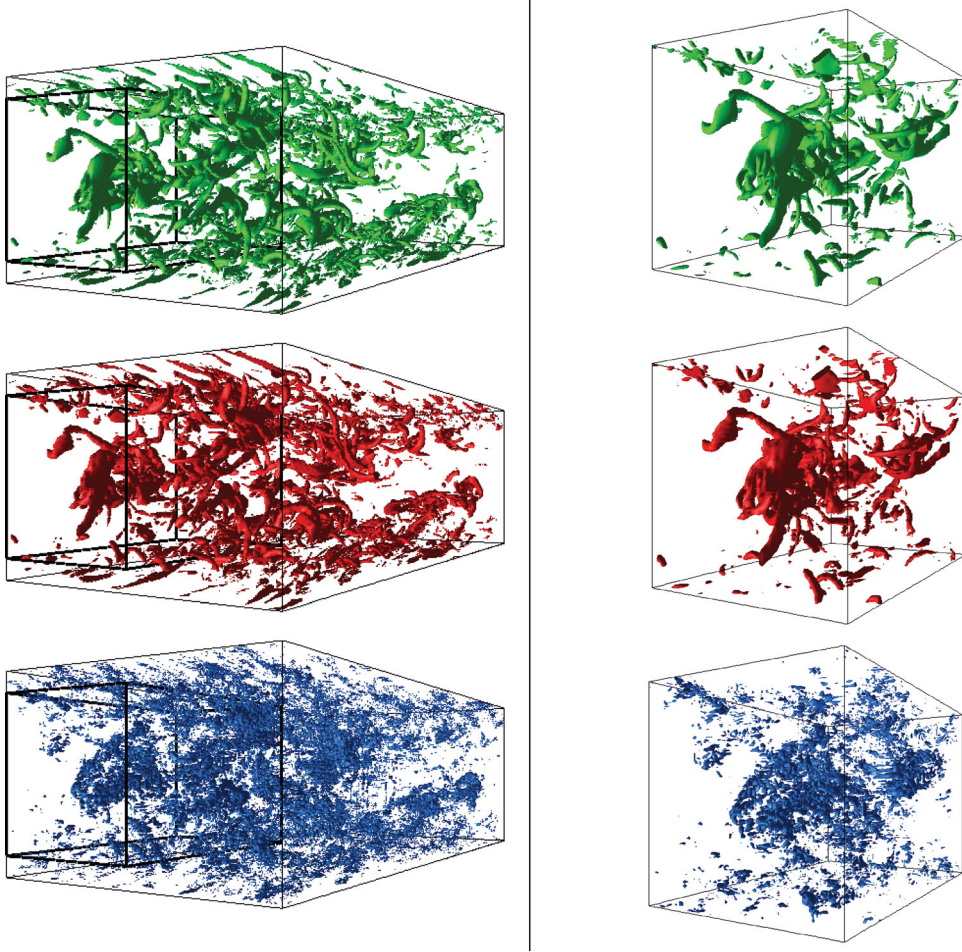


Figure 7. Visualisation of total vorticity ω (top), coherent vorticity ω_c (middle) and incoherent vorticity ω_i (bottom). Isosurfaces $|\omega^+| = |\omega_c^+| = \langle |\omega^+| \rangle + 3(\langle |\omega^+| \rangle - \langle |\omega^+| \rangle)^{1/2}$ and $|\omega_i^+| = \langle |\omega_i^+| \rangle + 3(\langle |\omega_i^+| \rangle - \langle |\omega_i^+| \rangle)^{1/2}$ are shown. The right column presents corresponding zooms in the core where $0 \leq x_1 \leq 0.79\pi h$, $-0.78h \leq x_2 \leq 0.78h$ and $0 \leq x_3 \leq 0.5\pi h$.

near the walls and in the core region. Although the incoherent flow is represented by the remaining majority of wavelet coefficients, it retains a negligible amount of energy, namely $2.3 \times 10^{-3}\%$ of the total energy, and only 0.5% of total enstrophy.

3.2. Mean and root-mean-square velocity and vorticity statistics

We analyse the statistics of the mean velocity and vorticity profiles of the coherent and incoherent flows, and compare them with the total flow. The results are averaged over 40 snapshots. Figure 8 shows the y^+ -dependence of the streamwise mean velocity $U_1^+(y^+)$ and of the spanwise mean vorticity $\Omega_3^+(y^+)$, averaged over x_1 - x_3 planes, for the total, coherent and incoherent flows. It is observed that the coherent flow perfectly preserves $U_1^+(y^+)$ and $\Omega_3^+(y^+)$, while both incoherent contributions are very weak. It can be noted that $U_2^+(y^+)$

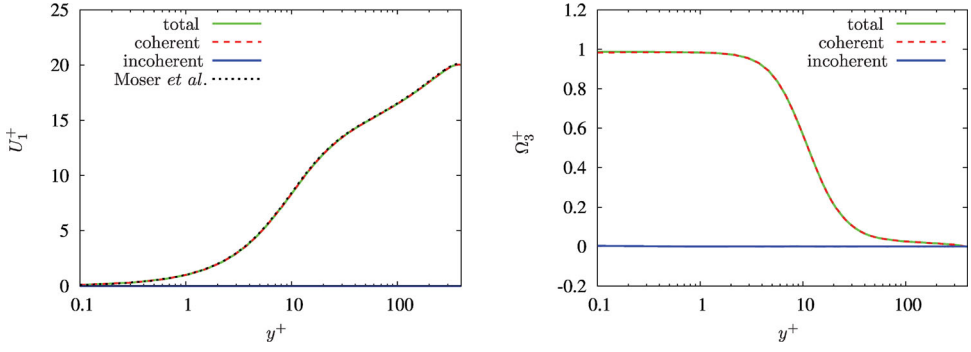


Figure 8. Streamwise mean velocity $U_1^+(y^+)$ (left) and spanwise mean vorticity $\Omega_3^+(y^+)$ (right) for total, coherent and incoherent flows, together with the DNS results of Ref. [35] in the lin-log representation.

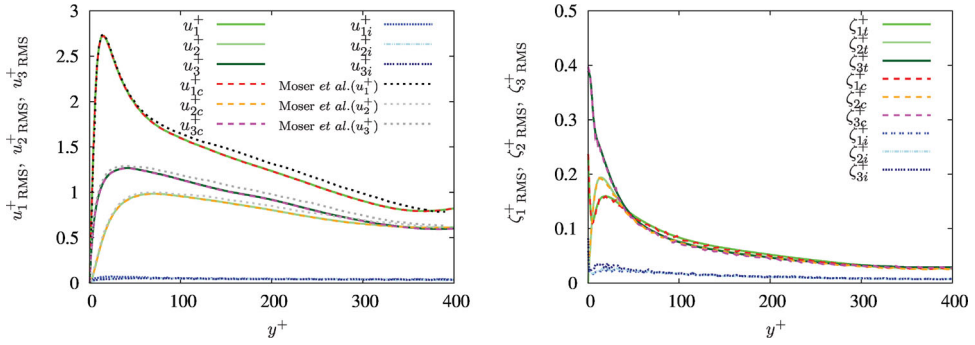


Figure 9. RMS of u_j^+ (left) and RMS of ζ_j^+ (right) for total, coherent and incoherent flows.

vanishes identically and that $U_3^+(y^+)$ almost vanishes for the total, coherent and incoherent flows. This implies that $\Omega_1^+(y^+)$ is almost zero, and $\Omega_2^+(y^+)$ is identically zero. The comparison of U_1^+ with the DNS data at $Re_\tau = 395$ in Moser et al. [35] confirms the validity of the present DNS.

The root-mean square (RMS) of the velocity fluctuations u_j^+ ($j = 1, 2, 3$) as a function of y^+ are shown in Figure 9 (left). Again, we find an excellent agreement between the total and the coherent flow for all values of y^+ , while the incoherent contribution is negligibly small. For the RMS of the vorticity fluctuations ζ_j^+ shown in Figure 9 (right), the coherent contributions well preserve the total RMS of ζ_j^+ . The fluctuations are defined by $\zeta_j = \omega_j - \Omega_j$, where Ω_j is the mean vorticity averaged over the $x_1 - x_3$ plane. The vorticity RMS of the incoherent flow is much smaller than that of the total flow.

3.3. Probability density functions of velocity and vorticity

Figure 10 (left) shows the probability density functions (PDFs), estimated using histograms with 200 bins, of the streamwise velocity fluctuations u_1^+ for the total, coherent and incoherent flows at three representative positions y^+ : in the viscous sublayer, the log region and near the centre of the channel. In all cases, we observe that the PDFs for the total and coherent velocity fluctuations perfectly superimpose, which indicates that high order statistics are

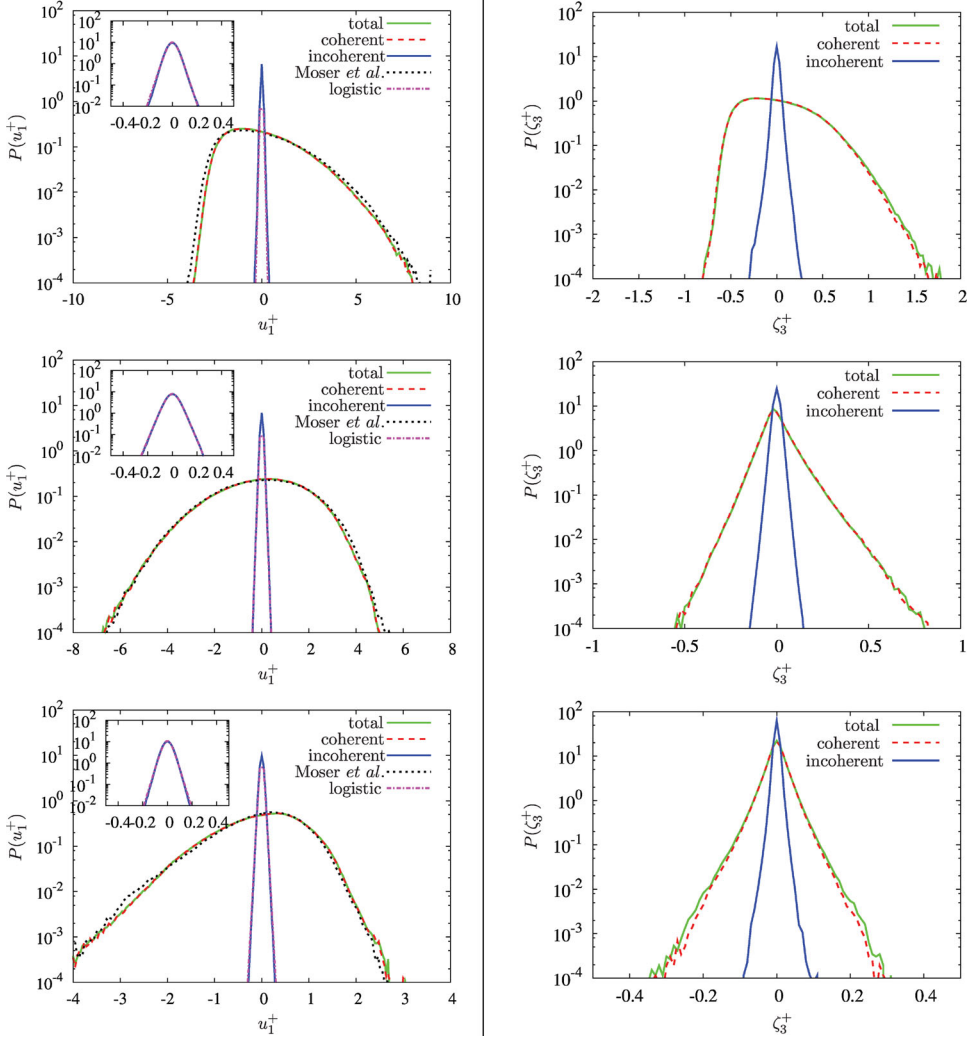


Figure 10. PDFs of u_1^+ (left) and ζ_3^+ (right) ; (top) $y^+ = 4.6$ viscous sublayer, (middle) $y^+ = 96.8$ around the log region and (bottom) $y^+ = 378.8$ around the centre of the channel.

well preserved by the coherent flow. We also find that the velocity PDFs remain skewed in the different regions and agree well with the data of Ref. [35], using appropriate renormalisation. In contrast, the PDFs of the incoherent velocity fluctuations are symmetric, and have strongly reduced variances. For the incoherent velocity, we analysed y^+ -dependent flatness, and found values around 4 in the viscous sublayer and in the log region, which decrease to 3.6 near the centre of the channel. For the y^+ -dependent skewness, fluctuations around zero are observed with an amplitude below 0.05. The PDFs of the incoherent velocity well superimpose the logistic distribution with zero mean and the variances $\sigma^2(y^+)$ of the incoherent velocity, though their flatness is 1.2, which is much smaller than the PDFs of the incoherent velocity. The logistic distributions $P(u_1^+)$ are given by $\exp(-\pi u_1^+/s)/\{s(1 + \exp(-u_1^+/s))\}$, where $s = 3^{1/2}\sigma(y^+)/\pi$.

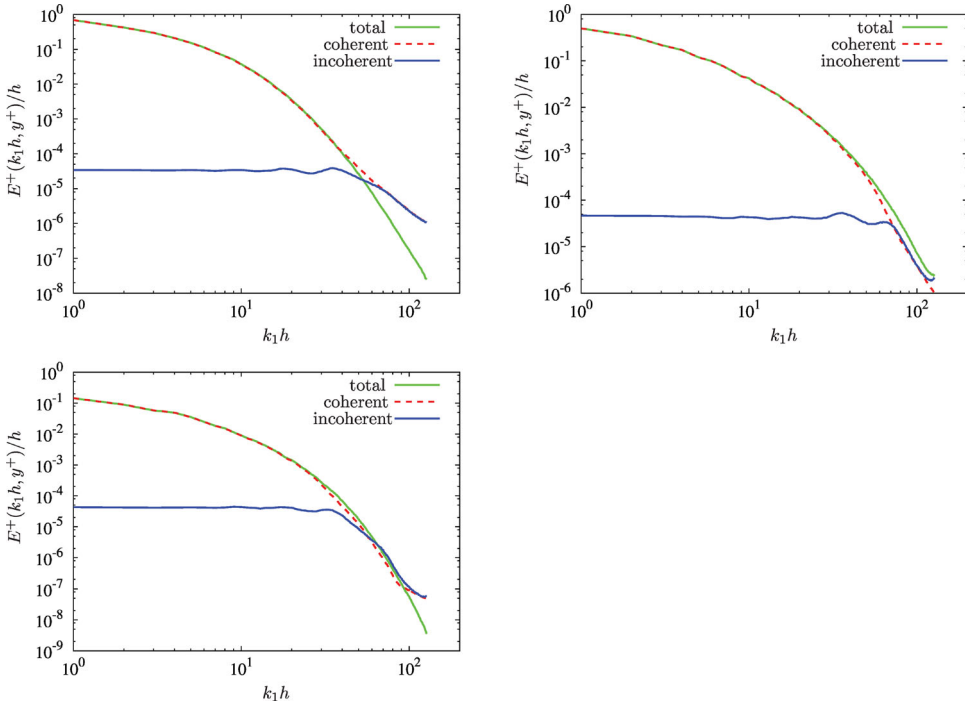


Figure 11. Dimensionless energy spectra of u_1^+ in the x_1 -direction at three representative y^+ : (top left) $y^+ = 4.6$ viscous sublayer, (top right) $y^+ = 96.8$ around the log region, and (bottom) $y^+ = 378.8$ around the centre of the channel.

In [Figure 10](#) (right), we illustrate the PDFs of the vorticity fluctuations ζ_3^+ at three representative positions y^+ : in the viscous sublayer, the log region and near the centre of the channel. The coherent vorticity fluctuations well represent the total vorticity PDFs which are skewed in all cases, while the corresponding incoherent PDFs are symmetric. The variances of these incoherent PDFs are, respectively, significantly weaker than the variances of the total and coherent PDFs.

3.4. Energy spectra

To get insight into the scale distribution of turbulent kinetic energy, we analyse the 1D energy spectra of the streamwise velocity u_1^+ in the streamwise direction $E^+(k_1h, y^+)$, which is defined as $E^+(k_1h, y^+) = \sum' |\hat{u}_i(k_1h, y^+, k_3h)|^2/2$, where $\hat{u}_i(k_1h, y, k_3h)$ is the Fourier transform of $\mathbf{u}^+(\mathbf{x})$ in the x_1 - x_3 planes, \sum' denotes the summation in terms of all k_3 . The results are shown in [Figure 11](#), again for the total, coherent and incoherent flows at three representative positions; in the viscous sublayer, the log layer and near the centre of the channel. The dimensionless wavenumber in the x_1 -direction is denoted by k_1h . [Figure 11](#) shows that the spectral distribution of turbulent kinetic energy is well preserved by the coherent flow, from the viscous sublayer to the centre of the channel. In contrast, the incoherent energy exhibits an almost flat spectrum, which corresponds to equipartition of incoherent energy, i.e. decorrelation of the incoherent flow in physical space.

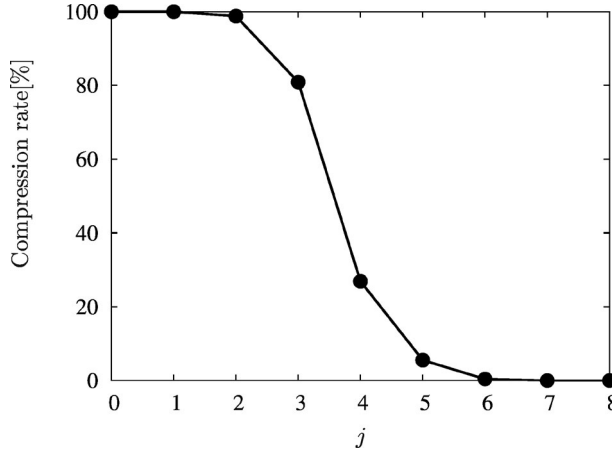


Figure 12. Scale-dependent compression rate.

At large wavenumbers close to the cut-off scale, imposed by the resolution of the DNS, we find that the incoherent energy dominates the total energy in the viscous sublayer and the log layer, while it dominates the coherent energy as well as the total one at the large wavenumbers around the centre of the channel. However, this is not surprising since the wavelet decomposition is orthogonal for vorticity but not for velocity, due to the fact that the Biot–Savart operator is not diagonal in wavelet space. Note that $\hat{\omega}_c^+(k_1, y^+, k_3)$ and $\hat{\omega}_i^+(k_1, y^+, k_3)$ are not orthogonal for any fixed (k_1, k_3) at every y^+ . But even though the cross-term $\langle \omega_c^+ \cdot \omega_i^+ \rangle(x_2) \neq 0$, its averaged contribution vanishes, $\int_{-h}^h dx_2 \langle \omega_c^+ \cdot \omega_i^+ \rangle(x_2) = 0$.

The compression is most efficient for small scales, i.e. large j and large wavenumbers (Figure 12). This implies that the scale-by-scale incoherent enstrophy is comparable or larger than the scale-by-scale coherent enstrophy.

3.5. Nonlinear dynamics

To get further insight into the nonlinear dynamics, we consider the energy budget given in the equation for $\langle u_j^+ u_j^+ \rangle / 2$ per unit mass [36]:

$$\frac{1}{2}(\partial_t + U_j^+ \partial_j) \langle u_j^+ u_j^+ \rangle = P(\mathbf{v}^+) + T(\mathbf{u}^+) + \Pi(\mathbf{u}^+, p^+) - \epsilon(\mathbf{u}^+) + V(\mathbf{u}^+), \quad (7)$$

where $P(\mathbf{v}^+) = -\langle u_j^+ u_l^+ \rangle \partial_l U_j^+$, $T(\mathbf{u}^+) = -\partial_l \langle u_j^+ u_j^+ u_l^+ \rangle / 2$, $\Pi(\mathbf{u}^+, p^+) = -\partial_l \langle p^+ u_l^+ \rangle$, $\epsilon(\mathbf{u}^+) = \nu \langle \partial_l u_j^+ \partial_l u_j^+ \rangle$, $V(\mathbf{u}^+) = \nu \partial_l \langle u_j^+ u_j^+ \rangle$. In Figure 13 (left), we see that three nonlinear coherent contributions, corresponding to production $P(\mathbf{v}_c^+)$, turbulent diffusion $T(\mathbf{u}_c^+)$ and pressure diffusion $\Pi(\mathbf{u}_c^+, p_c^+)$, are in good agreement with the corresponding total ones. Hence, the coherent flow almost perfectly preserves the nonlinear dynamics. Thus, we anticipate that the departure of the coherent flow from the total flow is negligibly small. Indeed, the incoherent contribution to the different terms, defined by $P(\mathbf{v}^+) - P(\mathbf{v}_c^+)$, $T(\mathbf{u}^+) - T(\mathbf{u}_c^+)$ and $\Pi(\mathbf{u}^+, p^+) - \Pi(\mathbf{u}_c^+, p_c^+)$, is almost zero. The two viscous contributions, $\epsilon(\mathbf{u}^+)$ and $V(\mathbf{u}^+)$, are also well retained by the coherent flow, $\epsilon(\mathbf{u}_c^+)$ and $V(\mathbf{u}_c^+)$,

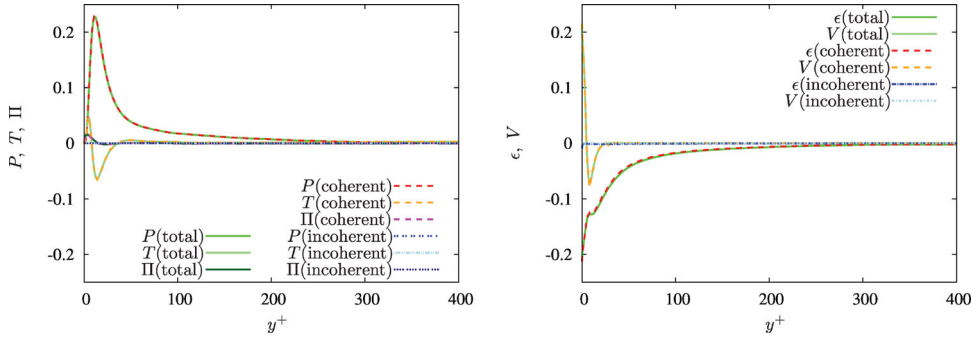


Figure 13. Production term P , turbulent diffusion term T and pressure diffusion term Π vs. y^+ (left). Energy dissipation ϵ and viscous diffusion V vs. y^+ (right). Coherent and incoherent contributions are presented together with the total one.

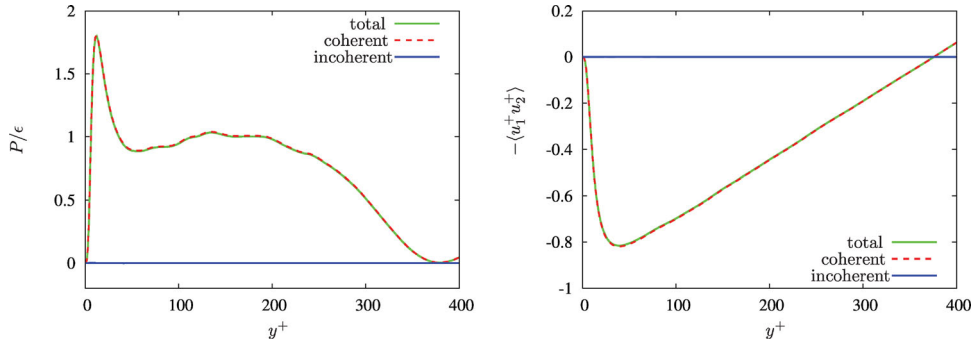


Figure 14. The ratios of total, coherent and incoherent productions P to the dissipation of turbulent kinetic energy for total flow, ϵ , vs. y^+ (left). The Reynolds stress $-\langle u_1^+ u_2^+ \rangle$ vs. y^+ (right).

as confirmed in Figure 13 (right). In the viscous sublayer, the incoherent flow has small contribution on $\epsilon(\mathbf{u}^+)$ and $V(\mathbf{u}^+)$. The incoherent contribution to the viscous terms, respectively measured by $\epsilon(\mathbf{u}^+) - \epsilon(\mathbf{u}_c^+)$ and $V(\mathbf{u}^+) - V(\mathbf{u}_c^+)$, becomes even smaller and more negligible as y^+ increases, a behaviour which is expected.

The ratio between the production and the dissipation yields insight into the equilibrium of the turbulent flow in the log region, as discussed in Ref. [35]. Figure 14 (left) shows this balance. Considering $P(\mathbf{v}^+)/\epsilon(\mathbf{u}^+)$, the coherent contribution perfectly superimposes with the ratio of the total flow, $P(\mathbf{v}^+)/\epsilon(\mathbf{u}^+)$. The corresponding quantity for the incoherent flow, $\{P(\mathbf{v}^+) - P(\mathbf{v}_c^+)\}/\epsilon(\mathbf{u}^+)$, is negligible, as expected from Figure 13 (left).

The Reynolds stress defined by $-\langle u_1^+ u_2^+ \rangle$ measures the fluctuation of turbulent momentum. The analysis of the Reynolds stress provides detailed information on the contribution to the turbulence production from various events occurring in the flow. Figure 14 (right) shows that the coherent Reynolds stress well represents the Reynolds stress for the total flow, while its incoherent contribution is negligibly small. The interaction between the coherent flows is predominant over the stress. In contrast, the remaining interactions play a non-significant role in the stress, not only between the incoherent flows but also between the coherent and the incoherent flows.

4. Conclusions and perspectives

DNS data of turbulent channel flow at moderate Reynolds number have been analysed using the CVE method. Boundary-adapted wavelets have been developed and implemented into a FWT. By thresholding the wavelet coefficients with one scale in three spatial directions, the flow has been decomposed into coherent and incoherent contributions. We found that few percentage of wavelet coefficients, i.e. 6%, are sufficient to represent the coherent structures of the flow. The low order statistics, mean velocity, mean vorticity, RMS velocity and RMS vorticity of the coherent part agree very well with those of the total flow. A spectral decomposition of turbulent kinetic energy confirms that the coherent flow matches the spectrum all along the inertial range. In contrast, the incoherent noise-like flow exhibits energy equipartition, which suggests that filtering it out corresponds to modelling turbulent dissipation. In order to obtain reliable statistical results, averaging over 40 flow snapshots has been performed. To get insight into the flow dynamics, we analysed the energy budget and we found that the coherent flow almost perfectly retains the nonlinear dynamics. The production/dissipation ratio of the coherent flow superimposes well the one of the total flow in the log layer, while the interactions between incoherent–incoherent and coherent–incoherent contributions are negligibly small. Although the coherent and incoherent vorticity fields are not perfectly divergence free, the divergence issue is not crucial as discussed in Appendix 1.

The present construction requires that the DNS data be interpolated onto an equidistant grid. This limits the applicability of the current CVE algorithm as higher resolution DNS data may not be handled due to the implied memory requirements. One way to overcome this is the use of Chebyshev wavelets, see e.g. [20,37]. In Appendix 2, we tested this approach and we have shown that similar results in terms of statistics and compression rate are indeed obtained.

The CVE results are encouraging for developing coherent vorticity simulation (CVS) of wall bounded turbulent flows. We anticipate that for higher Reynolds number, the compression rate will further improve, similar to what was found for isotropic turbulence [16]. CVS is based on a deterministic computation of the coherent flow evolution using an adaptive orthogonal wavelet basis [13]. The influence of the incoherent background flow is neglected to model turbulent dissipation. Applications of CVS to turbulent mixing layers and isotropic turbulence can be found in Refs. [17] and [38], respectively.

Finally, let us remark that the wavelet bases are not orthogonal in 2D planes for fixed y^+ . This implies that 2D statistics cannot be done, especially at small scales. In this case, it would be better to apply 2D wavelets in each plane, as done in previous work [21].

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Disclosure statement

No potential conflict of interest was reported by the authors.

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Appendices

Appendix 1. Divergence issues

The vector-valued wavelet basis used here is not divergence-free, since the orthogonal wavelet transform does not commute with the differential operator. Thus, the coherent vorticity, ω_c , and also the incoherent one, ω_i , are not divergence-free. In the following, we quantify the y^+ -dependent contribution of the divergent component $\nabla \xi$ of ω_c on the

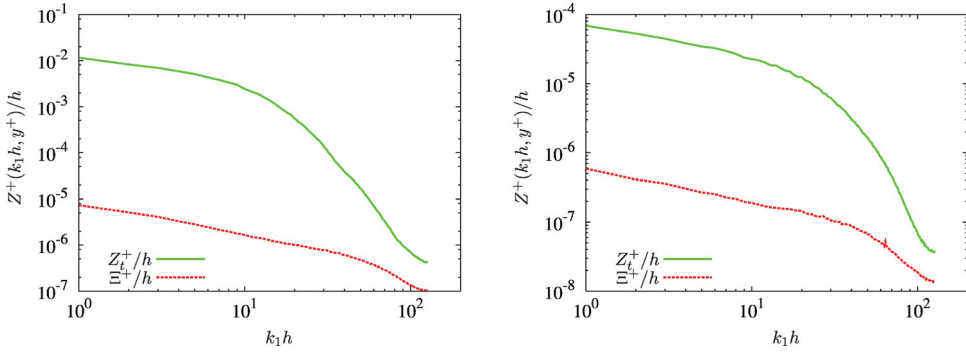


Figure A1. Enstrophy spectra, $Z^+(k_1h, y^+)$, and spectra of the divergence component $\Xi^+(k_1h, y^+)$ in the x_1^+ -direction, at (left) $y^+ = 4.6$ in the viscous sublayer, and (right) at $y^+ = 381.9$ around the centre of the channel.

streamwise vorticity spectra in the x_1 -direction. **Figure A1** shows dimensionless spectra of total streamwise vorticity ω^+ and those of $\nabla\xi^+$ at two representative values of y^+ , which are, respectively, located near the wall and around the centre of the channel. The contribution of $\nabla\xi^+$ appears mostly in the dissipative range, not only in the viscous sublayer but also around the centre of the channel. It can be seen that the contributions of ξ^+ are weak in the lower wavenumber region. The intensity of $\nabla\xi^+$, denoted by $\langle |\nabla\xi|^2 \rangle(y^+)$, is about $2.8 \times 10^{-2}\%$ of the total enstrophy in the viscous sublayer, and about 1.89% around the centre of the channel. Therefore, this divergence issue in ω_c is negligible for the statistics, but also for simulations, since ω_c is almost divergence-free.

Appendix 2. CVE using Chebyshev wavelets

In the following, we briefly summarise Chebyshev wavelets which yield an alternative construction of wavelets on the interval [39]. The idea is to perform a change of variables, similar to what is done for the trigonometric definition of Chebyshev polynomials. The efficient numerical implementation of Chebyshev wavelets is based on the periodic wavelet transform, in analogy with the fast Chebyshev transform which uses the cosine transform. The CVE results presented here use Chebyshev wavelets in the x_2 -direction instead of the CDJV wavelets, while in the x_1 and x_3 -directions, periodic Coiflet 30 wavelets are used.

B1. On Chebyshev wavelets

Using the coordinate transform $x = \cos(\theta)$ we map the interval $x \in [-1, 1]$ onto $\theta \in [0, \pi]$. Then π -periodic orthogonal wavelets $\psi^P(\theta)$ are used to construct wavelets $\psi^B(\theta)$, [20], which are even functions:

$$\psi^B(\theta) = \psi^P(\theta) + \psi^P(\pi - \theta). \quad (\text{B1})$$

The corresponding dilated and translated wavelets are obtained by $\psi_{j,i}^B(\theta) = 2^{j/2} \psi^B(2^j\theta - i)$. Setting $\theta = \arccos x$ we obtain the boundary wavelets $\psi^B(x)$ on the interval $[-1, 1]$

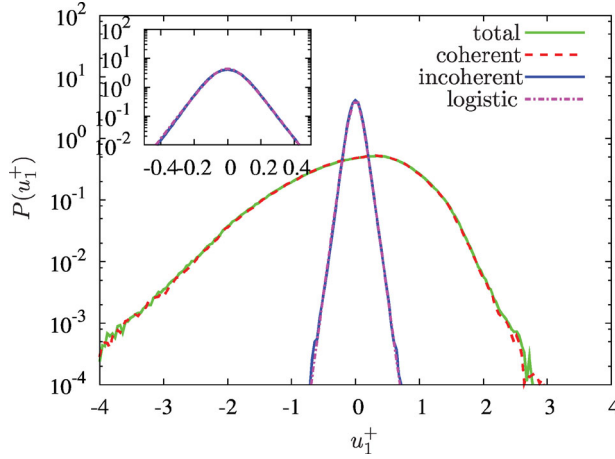


Figure B1. PDFs of u_1^+ ; $x_2^+ = 378.8$ around the centre of the channel with CVE using Chebyshev wavelets. The inset shows a zoom of the PDF of the incoherent velocity around 0.

which yield an orthogonal basis with respect to the weighted scalar product, i.e.

$$\int_{-1}^1 \psi_{j,i}^B(x) \psi_{j',i'}^B(x) \frac{1}{\sqrt{(1-x^2)}} dx = \delta_{jj'} \delta_{ii'}.$$

To compute the Chebyshev wavelet transform efficiently, we use periodic orthonormal Coiflet 30 wavelets with period 2π and extend the vorticity $\omega(x_1, \theta, x_3)$ as an even function $\mathbf{g}(x_1, \theta, x_3)$ for each (x_1, x_3) ,

$$\mathbf{g}(x_1, \theta, x_3) = \begin{cases} \omega(x_1, \theta, x_3) & \text{for } 0 \leq \theta \leq \pi, \\ \omega(x_1, -\theta, x_3) & \text{for } -\pi \leq \theta < 0. \end{cases} \quad (\text{B2})$$

Before applying the extension of ω , we interpolate the vorticity given on 192 Chebyshev grid points onto 256 equidistant grid points in the θ -coordinate. Then we can proceed with the CVE method and apply the FWT to \mathbf{g} using 3D orthogonal wavelets constructed by a tensor product from $\psi^P(x_1)$, $\psi^P(\theta)$ and $\psi^P(x_3)$.

B2. Numerical results

Now we extract coherent vorticity out of the turbulent channel flow at $Re_\tau = 395$, using the previously described Chebyshev wavelets. For the threshold value T we use the coefficient $\alpha = 0.10$. We find that the coherent flow, reconstructed from only 4.8% of the $256^2 \times 512$ wavelet coefficients, i.e. 6.4% of the original $256^2 \times 192$ grid points, retains almost all of the total energy and enstrophy, i.e. 99.9% of the total energy and 99.0% of the total enstrophy. In contrast, the incoherent flow represented by the remaining majority of the wavelet coefficients has little energy and enstrophy, namely 10^{-2} % of the total energy and 1.3% of the total enstrophy.

Inspecting [Figure B1](#) confirms that the PDFs for the total and coherent velocity fluctuations perfectly superimpose, indicating that high-order statistics are well preserved by the coherent flow. In contrast, the PDFs of incoherent velocity fluctuations have strongly

reduced variances and are not skewed, in contrast to what is found for the total and coherent fluctuations. Coherent and incoherent flows exhibit very similar properties as in [Section 3](#), where we used CDJV wavelets instead of the Chebyshev wavelets (figure with flow visualisations is omitted). Thus, Chebyshev wavelets can be more efficient for CVE than CDJV wavelets if the flow data have a large number of grid points, as no interpolation onto a fine equidistant grid is required.