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## Highlights

- $\text{Li}_{1.8}\text{Na}_{0.2}\text{TiO}_3:\text{Mn}^{4+}$  was synthesized by the solid-state reaction method in C2/c structure.
- The material was used as a probe for the lifetime-based luminescence thermometry.
- $\text{Mn}^{4+}$  emits around 679 nm due to of  ${}^2\text{E}_g \rightarrow {}^4\text{A}_{2g}$  spin forbidden electron transition.
- The material exhibits a large value of relative sensitivity ( $2.6\% \text{ K}^{-1}$  @ 340 K).
- Excellent  $\text{Mn}^{4+}$  emission repeatability with temperature resolution of 0.1 K.

# **Li<sub>1.8</sub>Na<sub>0.2</sub>TiO<sub>3</sub>:Mn<sup>4+</sup>: the highly sensitive probe for the low-temperature lifetime-based luminescence thermometry**

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## **Abstract:**

In this work, the potential of Li<sub>1.8</sub>Na<sub>0.2</sub>TiO<sub>3</sub>:Mn<sup>4+</sup> for the lifetime-based luminescence thermometry is assessed. The material is prepared by the solid-state reaction of Li<sub>2</sub>CO<sub>3</sub>, Na<sub>2</sub>CO<sub>3</sub>, and nanostructured TiO<sub>2</sub> at 800°C, and its monoclinic structure (space group C2/c) is confirmed by X-ray diffraction analysis. In this host, Mn<sup>4+</sup> provides strong absorption around 330 nm and 500 nm due to <sup>4</sup>A<sub>2g</sub> → <sup>4</sup>T<sub>1g</sub> and <sup>4</sup>A<sub>2g</sub> → <sup>4</sup>T<sub>2g</sub> electric dipole forbidden and spin-allowed electron transitions, respectively, and emits around 679 nm on account of <sup>2</sup>E<sub>g</sub> → <sup>4</sup>A<sub>2g</sub> spin forbidden electron transition. Temperature dependences of emission intensity and emission decay are measured over the 10–350 K range. Due to the low value of energy of <sup>4</sup>T<sub>2g</sub> level (20000 cm<sup>-1</sup>), the strong emission quenching starts at low-temperatures which favors the use of this material for the luminescence thermometry. It is demonstrated that the quite large value of relative sensitivity (2.6% K<sup>-1</sup>@340 K) facilitates temperature measurements with temperature resolution better than 0.1 K, and with the excellent repeatability.

**Keywords:** Luminescence thermometry; phosphor; Mn<sup>4+</sup>; Lithium titanium oxide; Temperature dependence of emission; Temperature sensors

## **1. Introduction**

Much effort is directed to the development of optical methods for temperature measurements these days. Although temperature is one of the most commonly measured physical properties, accurate temperature measurements are still to be realized in some emerging environments, such as at the nanoscale or intracellular environments. [1,2]. Luminescence is a temperature dependent phenomenon. Thus, many materials can be used as luminescence thermometry probes by exploiting temperature induced changes in their excitation and emission band positions and bandwidths [3–5], emission band intensity [6–9], ratio of intensities of two emission bands [10–19], and excited-state lifetimes [20–32].

Temperature readings from the variations the emission intensity ratios and excited-state lifetimes are utilized much more frequently than from other luminescence features since these methods are self-referencing and are robust in respect to fluctuations in excitation power, detector sensitivity, and concentration of the probe material [33].

Luminescence thermometry based on excited-state lifetime measurements requires monitoring of just one emission band and, therefore, has an advantage over emission intensity ratiometric method which needs monitoring of two distinctively separated emissions to achieve self-referencing. On the other hand, excited-state lifetime methods show lower temperature sensitivities than ratiometric methods. However, measurements of temporal emission changes are much more accurate than emission intensity measurements, and lower measurement uncertainties compensate for lower sensitivities so that lifetime methods with sensitive probes may provide better temperature resolution than ratiometric methods [17]. Materials of choice are usually lanthanide and transition metal activated phosphors because they exhibit large emission quantum efficiencies, are chemically and thermally stable, show no photobleaching, and have excited state lifetimes on order of  $\mu\text{s}$  and longer. Lanthanide activated phosphors are suitable for high-temperature measurements since their lifetime has an almost constant value in the low- and mid-temperature ranges. In transition metal activated phosphors, excited state lifetime starts to decrease at lower temperatures than in lanthanide phosphors so that luminescence thermometry can be employed from cryogenic temperatures to room temperatures or even slightly higher. Good candidates are  $\text{Mn}^{4+}$  activated oxides since these phosphors have a strong excitation band in the blue spectral region so that emissions can be excited with a cheap LED source, and have low value of  ${}^4\text{T}_{2g}$  energy level which favors emission quenching even at low temperatures. Also, emission from  ${}^2\text{E}_g$  level in the 650-730 nm range is less affected by the environmental (parasitic) luminescence than shorter-wavelength emissions. The materials have optimal values of lifetimes (few tens of  $\mu\text{s}$  to ms) which provide good temporal resolution of measurements (lifetimes greater than 1 ms are not suitable for fast measurements).

Here, we aimed to investigate the potential of  $\text{Li}_{1.8}\text{Na}_{0.2}\text{TiO}_3:\text{Mn}^{4+}$  as a lifetime-based luminescence thermometry probe since the recent report on  $\text{Li}_2\text{TiO}_3:\text{Mn}^{4+}$  [34] showed that it provides favorable luminescence properties for this application.

## 2. Experimental section

### 2.1. Synthesis

$\text{Li}_{1.8}\text{Na}_{0.2}\text{TiO}_3:0.5\%\text{Mn}^{4+}$  sample was prepared by traditional solid state synthesis in air.  $\text{Li}_2\text{CO}_3$  (Alfa Aesar, 99 %),  $\text{Na}_2\text{CO}_3$  (Alfa Aesar, 98 %),  $\text{TiO}_2$  (nanostructured anatase  $\text{TiO}_2$  prepared by titanium isopropoxide hydrolysis [35]) and  $\text{MnO}_2$  (Alfa Aesar, 98 %) were used as starting materials. The precursor powder was grounded in pestle and mortar with ethyl alcohol as a binder for 2 h. Then the mixture is gradually heated to  $670^\circ\text{C}$  and calcined for 3 h. The sample was cold-pressed under 700 MPa pressure and then sintered for 72 h at  $800^\circ\text{C}$ . The pellets were cooled in furnace to room temperature to form pure-phase.

### 2.2. Instruments and measurements

X-ray diffraction (XRD) measurements were performed on a Rigaku SmartLab system operating with  $\text{Cu K}\alpha$  radiation (30 mA, 40 kV) in the  $2\theta$  range from  $10^\circ$  to  $90^\circ$ . Optical Parametric Oscillator (EXPLA NT342) was used as an excitation source for luminescence measurements and emissions were recorded using FHR1000 monochromator (Horiba Jobin Yvon) and iCCD camera (Horiba Jobin Yvon 3771).  $\text{Mn}^{4+}$  emission decay times were measured over the 10–350 K temperature range using a closed-cycle cryostat (Advanced Research System

DE202-AE) equipped with a Lakeshore model 331 controller. Diffuse reflectance measurements were performed with Shimadzu UV-2600 (Shimadzu Corporation, Tokyo, Japan) spectrophotometer equipped with an integrated sphere (ISR-2600).

### 3. Results and Discussions

This material has a monoclinic structure (space group C2/c) [36,37] as evidenced by the synthesized powder X-ray diffraction pattern shown in Fig. 1 (diffractions are indexed according to ICDD: 01-077-8280).

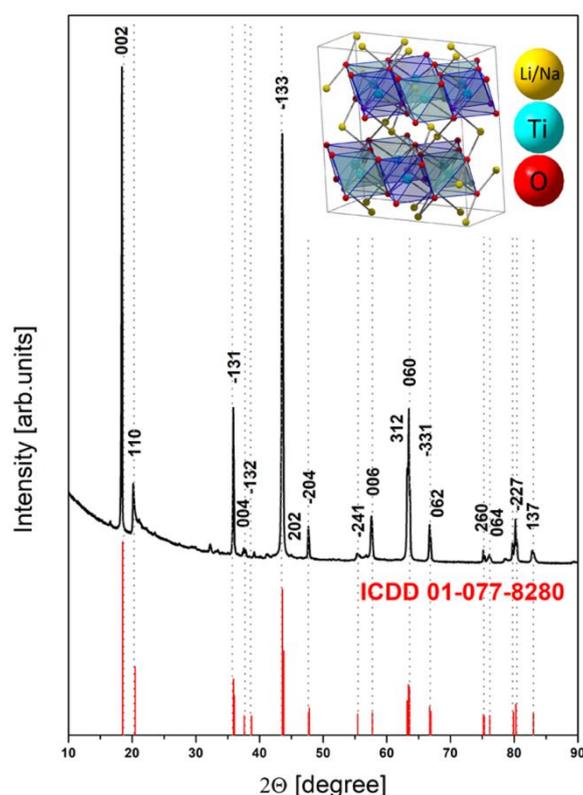


Figure 1. XRD pattern of  $\text{Li}_{1.8}\text{Na}_{0.2}\text{TiO}_3:0.5\%\text{Mn}$  powder indexed according to ICDD: 01-077-8280; inset illustrates the crystal structure of  $\text{Li}_{1.8}\text{Na}_{0.2}\text{TiO}_3$

$\text{Ti}^{4+}$  ions are situated in two nearly identical crystallographic sites (4e symmetry), and in both of them Ti is in octahedral oxygen coordination.

$\text{Mn}^{4+}$  substitutes  $\text{Ti}^{4+}$  in  $\text{Li}_{1.8}\text{Na}_{0.2}\text{TiO}_3$ , so that local Ti site symmetries are reflected in the splitting of the  $\text{Mn}^{4+}$  energy levels and diffuse reflection and photoluminescence emission spectra, Fig. 2.

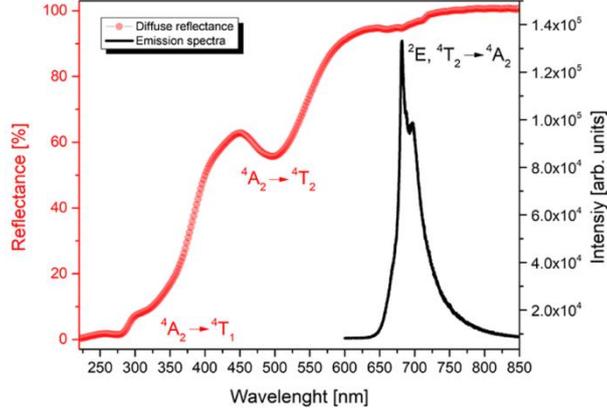


Figure 2. Room temperature diffuse reflectance (red symbols) and emission (black line; excited by 490 nm radiation) spectra of  $\text{Li}_{1.8}\text{Na}_{0.2}\text{TiO}_3:0.5\%\text{Mn}$ .

The diffuse reflection spectrum, Fig. 2 (red symbols), of  $\text{Li}_{1.8}\text{Na}_{0.2}\text{TiO}_3:\text{Mn}^{4+}$  is characterized by absorptions of the  $\text{O}^{2-}-\text{Mn}^{4+}$  and  $\text{O}^{2-}-\text{Ti}^{4+}$  charge transfer bands (at wavelengths < 300 nm) and the  $\text{Mn}^{4+}$  broad intra-configurational ( $3d \rightarrow 3d$ ) absorption bands originating from  ${}^4\text{A}_{2g} \rightarrow {}^4\text{T}_{2g}$  (around 500 nm) and  ${}^4\text{A}_{2g} \rightarrow {}^4\text{T}_{1g}$  (around 330 nm) electric dipole (parity) forbidden and spin-allowed electron transitions. The emission spectrum, Fig. 2 (black line), is due to the  ${}^2\text{E}_g \rightarrow {}^4\text{A}_{2g}$  spin forbidden electron transition (around 679 nm), with the contribution of the associated vibronic sidebands and some contribution from the  ${}^2\text{T}_{1g} \rightarrow {}^4\text{A}_{2g}$  spin-allowed transition.

The temperature dependence of  $\text{Mn}^{4+}$  emission is recorded over the 10–350 K range is shown in Fig. 3. At low temperatures, the non-radiative deexcitation rate is negligible so that emission intensity and lifetime vary only slightly with temperature (due to variation of the radiative rate). At higher temperatures, both the emission intensity (Fig. 3(a)) and emission lifetime (Fig. 3(b)) are temperature quenched through the  ${}^4\text{T}_{2g}$  level [38].

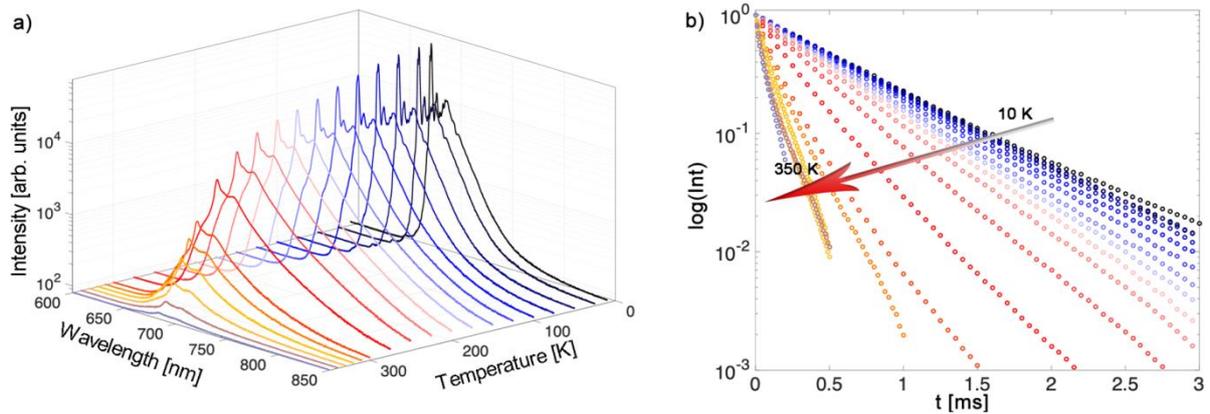


Figure 3. The temperature dependence of  $\text{Li}_{1.8}\text{Na}_{0.2}\text{TiO}_3:0.5\%\text{Mn}^{4+}$  emission in the 10–350 K range: (a) emission spectra measured with 490 nm excitation, (b) decay of 679 nm emissions.

The temperature dependence of the excited state lifetime can be described by the following formula [38]:

$$\tau(T) = \frac{\tau_R(0) \cdot \tanh(h\nu/2k_B T)}{1 + (\tau_R(0) \cdot \tanh(h\nu/2k_B T) / \tau_{NR}) \cdot \exp(-\Delta E/k_B T)}, \quad (1)$$

where  $\tau_R(0)$  is the radiative lifetime at  $T = 0$  K,  $k_B = 0.69503476$  cm<sup>-1</sup>/K is the Boltzmann constant,  $h\nu$  is the average energy of phonons coupled to the  ${}^2E_g \rightarrow {}^4A_{2g}$  transition,  $1/\tau_{NR}$  is the non-radiative decay rate,  $\Delta E$  is the activation energy of the cross-over process via the  ${}^4T_{2g}$  level, and  $T$  represents the temperature. Figure 4(a) shows the experimental data of the lifetime dependence on temperature (symbols) and the fit of Eq. (1) to experimental data (red line;  $\tau_R(0) = 0.6858$  ms,  $h\nu = 257.6$  cm<sup>-1</sup>,  $1/\tau_{NR} = 6 \cdot 10^8$  s<sup>-1</sup>,  $\Delta E = 2574$  cm<sup>-1</sup>,  $R^2=0.998$ ).

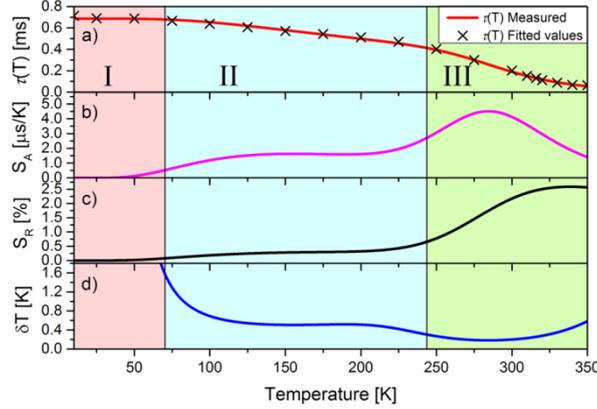


Figure 4. (a) The temperature dependence of the  $\text{Mn}^{4+}$  emission lifetime (symbols); the red line represent fit to Eq. (1), (b,c) the absolute and relative sensitivity respectively, and (d) the temperature resolution of thermometry.

One can observe three temperature regions in which lifetimes change differently with increase in temperature. In the first one, where temperature is lower than 75 K, the lifetime is constant and equals the value of  $\tau_R(0)$ . For temperatures between 75 and 230 K (the second temperature region), the lifetime decreases because of an increased probability of phonon induced transitions between the  $\text{Mn}^{4+} {}^2E_g$  and  ${}^4A_{2g}$  states. In the third region, for temperatures larger than 230 K, the lifetime strongly decreases due to an increase in the rate nonradiative processes. The applicability of this material for luminescence thermometry can be assessed from the change of the lifetime with temperature, i.e. from the absolute ( $S_a$ ) or relative ( $S_R$ ) sensitivity of the measurement, which can be derived from the Eq. (1) as:

$$S_a(T) [s \cdot K^{-1}] = \frac{d\tau(T)}{dT}$$

$$= - \frac{\tau_R(0) \cdot \tanh(h\nu/2k_B T)}{k_B T^2} \frac{h\nu / (2 \cdot \sinh(h\nu/2k_B T) \cdot \cosh(h\nu/2k_B T)) + \Delta E \cdot (\tau_R(0) \cdot \tanh(h\nu/2k_B T) / \tau_{NR}) \cdot \exp(-\Delta E/k_B T)}{[1 + (\tau_R(0) \cdot \tanh(h\nu/2k_B T) / \tau_{NR}) \cdot \exp(-\Delta E/k_B T)]^2},$$

$$S_R(T) [\%] = \left| \frac{1}{\tau(T)} \cdot \frac{d\tau(T)}{dT} \right| \cdot 100\% =$$

$$= \left| \frac{1}{k_B T^2} \frac{h\nu / (2 \cdot \sinh(h\nu/2k_B T) \cdot \cosh(h\nu/2k_B T)) + \Delta E \cdot (\tau_R(0) \cdot \tanh(h\nu/2k_B T) / \tau_{NR}) \cdot \exp(-\Delta E/k_B T)}{1 + (\tau_R(0) \cdot \tanh(h\nu/2k_B T) / \tau_{NR}) \cdot \exp(-\Delta E/k_B T)} \right| \cdot 100\% \quad (2)$$

are given in Figs. 4(b) and (c). The onset of applicability is at 75 K, and the relative sensitivity increases in value with a temperature increase. In the first temperature region, the relative sensitivity has a moderate value, in the 0.1–0.5 % K<sup>-1</sup> range. For temperatures larger than 230 K, in the third temperature region, the relative sensitivity is a quite high, with maximum value

of  $2.59 \text{ \%K}^{-1}$  at 340 K. Such high sensitivity facilitates measurement with temperature resolution which is among the highest ever reported for the luminescence thermometry, Fig. 4(c). The temperature resolution is calculated from the uncertainty in lifetime measurement, Fig. 5(a), as  $\delta T(T) = \sigma/S_a$ , where  $\sigma = 0.8228 \text{ }\mu\text{s}$  represents the value of measurement uncertainty. It has good values (from 1 to 0.2 K) in the 75–230 K temperature range, and exceptional value 0.1K and lower around room temperatures, third region in Fig. 4(c). Cycles of lifetime measurements were performed between two temperatures, see Fig. 5(b), and an excellent repeatability was observed.

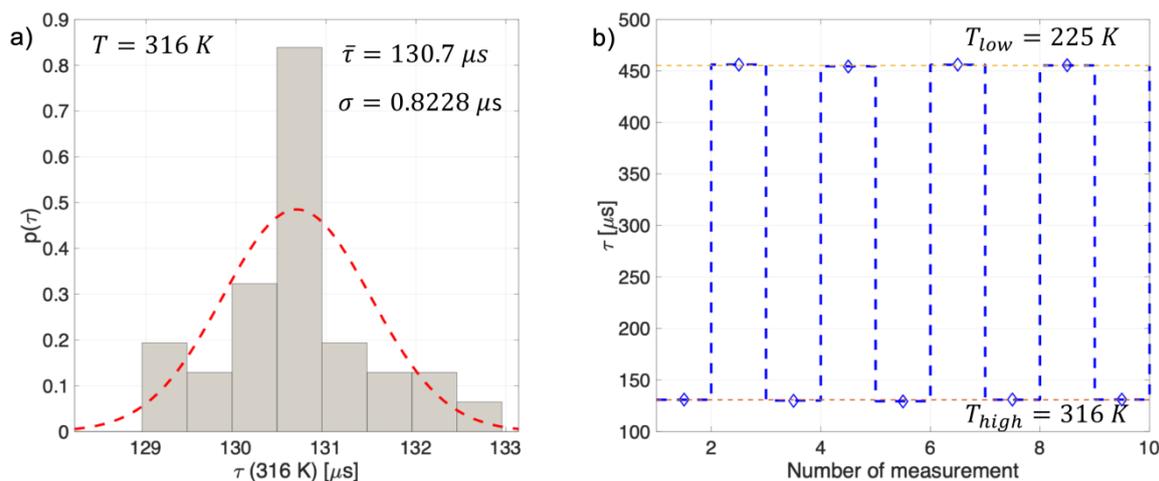


Figure 5. (a) The distribution of lifetime values measured at 316 K, (b) the repeatability of lifetime measurements demonstrated by measuring lifetime while cycling between two temperatures.

## Conclusion

$\text{Li}_{1.8}\text{Na}_{0.2}\text{TiO}_3:\text{Mn}^{4+}$  probe for the lifetime-based luminescence thermometry provides high-sensitivity temperature readings for temperatures higher than 75 K with a quite high temperature resolution, better than 0.1 K around room temperature. This probe can be excited by the blue light and emits in the deep-red spectral region (because of the strong nephelauxetic effect in oxides) where the contribution of parasitic background luminescence is minimal. The strong quenching of emissions at relatively low temperatures is a consequence of the low energy value of  ${}^4\text{T}_{2g}$  level which facilitates the low-energy activation ( $\Delta E = 2574 \text{ cm}^{-1}$ ) of the quenching process via the cross-over from  ${}^2\text{E}_g$  to  ${}^4\text{T}_{2g}$  and finally to the ground state. As for the most of lanthanide and transition metal activated phosphors, the repeatability of measurements is almost perfect. The presented results show that  $\text{Li}_{1.8}\text{Na}_{0.2}\text{TiO}_3:0.5\%\text{Mn}^{4+}$  is a suitable candidate for the development of accurate and self-referencing luminescent temperature sensors. The results also suggest that it is possible to engineer the lifetime-based  $\text{Mn}^{4+}$  activated luminescent probes for targeted applications and temperature ranges by selecting host materials of suitable energy of the  ${}^4\text{T}_{2g}$  state and the energy of phonon coupling to the  ${}^2\text{E}_g \rightarrow {}^4\text{A}_{2g}$  transition.

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