

# Optics Letters

## Combined photoacoustic imaging to delineate the internal structure of paintings

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**1** In this Letter, we present a combined photoacoustic imaging method, based on consecutive excitation using either the fundamental or the second-harmonic wavelength of a pulsed Nd:YAG laser for the stratigraphy of painted artworks. **2** Near-infrared excitation was employed for the imaging of hidden underdrawings in mock-up samples, whereas visible light for the thickness mapping of the overlying paint through the detection of photoacoustic signal attenuation. The proposed methodology was proven effective in measuring thick and strongly absorbing layers, which would not be possible by means of other pure optical techniques, while also enabling the visualization of features underneath the painted surface. Such an implementation expands significantly the applicability of the previously presented photoacoustic technique, which was limited to point-measurements, and paves the way for novel application in historical and technical studies, as well as in documenting restoring operations. © 2019 Optical Society of America

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Laser-based optical techniques are becoming more and more significant in the field of art diagnostics, thanks to their capability to obtain chemical, structural, and morphological information in a non-invasive way. In regard to paintings, hidden features underneath the painted surface, i.e., underdrawings, signatures, *pentimenti*, ancient and modern materials, overpaintings, and varnishes can be measured without sampling or damaging artwork integrity. More specifically, restoring operations involving the selective removal of altered varnishes, such as the cleaning process, definitely benefit from the micrometric in-depth evaluation of the aged materials to be removed. In the last decades, several optical techniques have been applied for the non-invasive stratigraphic analysis of paintings [1–12]. Among others, optical coherence tomography (OCT) [1,2] provides high-resolution cross-sectional images for the visualization of low scattering varnish and semi-transparent paints, sometimes enabling the visualization of underdrawings [2]. More recently, nonlinear optical microscopy modalities [8–12]

have been tested in painting diagnostics with the aim to extend maximum achievable imaging depth in highly scattering and semi-opaque materials.

In a previous study [13], we introduced a point-wise novel method called photoacoustic signal attenuation analysis (PACSA), which takes advantage of the superficial visible light absorption by the pigment particles for the estimation of paint layer local thickness. The technique is based on the photoacoustic effect, occurring when the absorption of intensity-modulated optical radiation by a medium induces the generation of broadband acoustic waves. Compared to visible and infrared radiation, the transmission of photoacoustic signals through scattering media is considerably higher [14], and characterized by a frequency-dependent exponential attenuation [15–18], which is significantly stronger for higher frequency components [15–17]. By analyzing the photoacoustic signal attenuation through a given material and by estimating the respective attenuation coefficients, it is possible to determine the thickness of an object [18,19]. In our last work, we employed PACSA for local thickness measurements on purposely prepared paint samples. To this aim, we shifted the detected photoacoustic signals from the time to frequency domain and used the amplitude spectrum for the estimation of the average transmitted frequency (ATF) as a characteristic measure of the overall acoustic attenuation during the signal propagation through the paint layer. ATF values were correlated with the thickness of different pigmented layers by fitting exponential decay models which, subsequently, were used to perform spot measurements on a painted canvas sample with high accuracy.

In this Letter, we present a combined photoacoustic imaging methodology employing consecutive excitation through the fundamental and second-harmonic wavelengths emitted by a pulsed Nd:YAG laser source, with the aim of delineating the stratigraphy of paintings. Therefore, we have introduced a novel imaging version of PACSA, called photoacoustic signal attenuation imaging (PACSAI), with optimized irradiation, acquisition, and scanning parameters, enabling the recording of accurate stratigraphic data in opaque paint layers. In principle, a stratigraphy of an object can be obtained by tuning the wavelength according to each layer and knowing the attenuation

coefficients of the materials. In fact, by selecting proper excitation wavelengths, it is possible to confine the generation of photoacoustic waves to a particular layer, while enabling the transmission of the radiation through the other layers. In addition, the visualization of underdrawings can be also obtained by combining the area-wise PACSAI modality with the near-infrared photoacoustic imaging [14], as shown in this Letter. Such in-depth and areal information may definitely turn useful to plan or monitor restoring operations, i.e., the cleaning process, and for historical and technical studies through the disclosure of painting internal features. The dual-content data have been compared with OCT images, validating the capabilities of the proposed diagnostic approach.

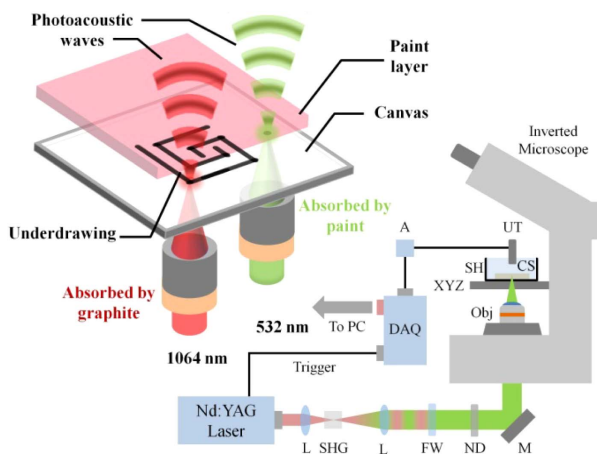
The photoacoustic imaging setup (Fig. 1) is based on a diode-pumped Q-switched Nd:YAG laser at 1064 nm (QIR-1064-200-S, CrystaLaser LC, Reno, Nevada; pulse energy, 29.4  $\mu$ J; pulse duration,  $\sim$ 8 ns; selected repetition rate, 5 kHz;  $M^2$  factor, 1.2) whose second harmonic (at 532 nm) is generated by focusing the beam on a lithium triborate crystal. To switch between infrared and visible excitation, two band-pass filters (FL1064-10, Thorlabs, Newton, New Jersey and FF01-531/40-25, Semrock, Rochester, New York) are used, which are mounted on a filter wheel.

The beam was focused on the back side of the sample, at the interface between the paint layer and the substrate, by an air immersion objective lens (Achromat 8 $\times$ , LOMO, St. Petersburg, Russia; NA: 0.2), after proper attenuation (pulse energy on the sample  $\sim$ 200 nJ). The sample was placed into an optically transparent Petri dish filled with a water-based gel of cellulose ether (carboxymethyl cellulose [CMC] gel-3%) acting as a coupling medium to ensure an efficient transmission of the laser-induced ultrasound signal. CMC is well suited for artwork diagnostics: being an inert material, it is widely used for the cleaning of the painted surfaces due to its safety features, wide availability, and low cost. The Petri dish was fixed on a high-precision motorized XY micrometric stage (8MTF-75LS05, Standa, Vilnius, Lithuania) enabling the point-by-point scanning over the region of interest, whereas the focal plane

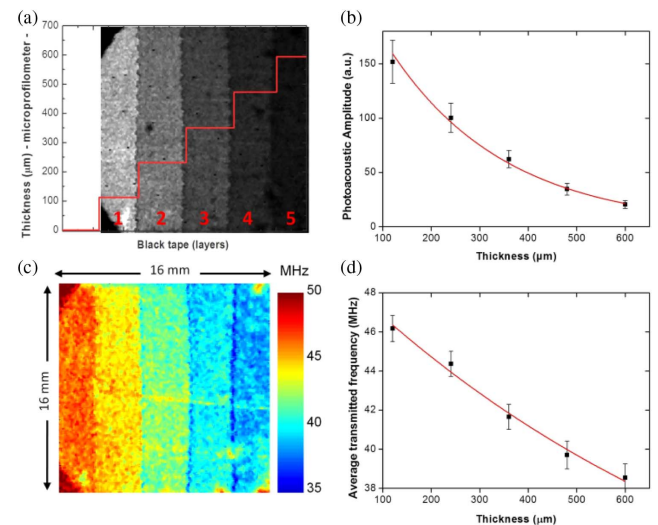
positioning was performed by the built-in manual Z-control of the microscope. For the detection of the generated photoacoustic waves, a single element spherically focused broadband ultrasonic transducer (HFM28, SONAXIS, Besancon, France; central frequency, 73 MHz; effective bandwidth,  $\sim$ 90 MHz at  $-6$  dB; focal distance, 4.53 mm) was immersed into CMC in a confocal and coaxial configuration with respect to the optical axis.

The combined photoacoustic imaging method was first tested on a multilayer sample consisting of a number of superimposed black tape layers producing an increasing thickness on a glass support (coverslip  $2.5 \times 2.5$  cm; 150  $\mu$ m thick). This sample of well-defined structural, optical, and acoustic properties was purposely produced to test the performance of PACSAI before proceeding to the analysis of more realistic cases, while clearly highlighting the capability of PACSAI in probing highly turbid and thick materials, in contrast to OCT imaging. The thickness of each area (1–5 tape layers) was preliminarily assessed by a micro-profilometer (Perthometer S5P, Mahr, Göttingen, Germany) with one single tape measured at  $110 \pm 5$   $\mu$ m [Fig. 2(a)]. The sample was irradiated from the back side with the second harmonic of the Nd:YAG laser beam. The generated time-domain photoacoustic signal was recorded point by point over a scanning area of 16 mm<sup>2</sup> resulting in a 200 by 200 pixels image.

Data were then processed in MATLAB and ImageJ environments for the reconstruction of the photoacoustic amplitude image [Fig. 2(a)] where, for each pixel, the brightness value corresponds to the maximum amplitude projection (MAP) obtained by averaging 128 waveforms, which are recorded for S/N enhancement. The amplitude values were calculated as an average on selected areas of  $\sim$ 4600 pixels at the center of each region of different thickness, and then plotted as a function of thickness. The best fitting curve demonstrates an exponential decay of the photoacoustic amplitude with the propagation length through the sample's volume (decay



**Fig. 1.** Photoacoustic setup for the successive imaging of painting samples using two irradiation wavelengths. L, lens; SHG, second-harmonic generation crystal; FW, filter wheel; ND, neutral density filter; M, mirror; Obj, objective lens; XYZ, 3D translational stage; SH, sample holder; CS, canvas sample; UT, ultrasonic transducer; A, amplifier; DAQ, data acquisition card.



**Fig. 2.** PACSAI of a black tape sample. (a) MAP photoacoustic image in grayscale showing the different layers, with the respective thickness values obtained with a microprofilometer (red line); (b) plot of the average photoacoustic amplitude as a function of thickness; (c) ATF image displayed within the range 35–50 MHz; and (d) plot of the average ATF values as a function of total layer thickness.

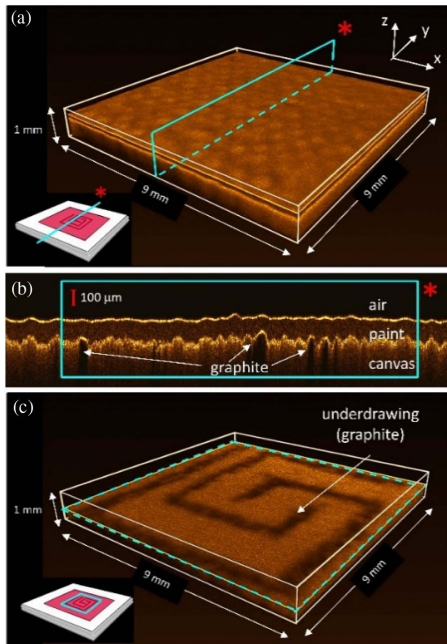
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161 constant:  $4.18 \times 10^{-3}$ ,  $R^2 = 0.992$ ) [Fig. 2(b)]. Finally, a  
 162 respective ATF color-scale image was generated within the  
 163 acoustic band of 10–100 MHz, following a fast Fourier trans-  
 164 form of the acquired time-domain photoacoustic waveforms.  
 165 Similar to the amplitude image, the ATF values were calculated  
 166 as the average on the previously selected areas for each tape layer  
 167 and plotted together with an exponential fitting curve (decay  
 168 constant:  $1.08 \times 10^{-3}$ ,  $R^2 = 0.982$ ) as a function of layer thick-  
 169 ness [Fig. 2(d)].

170 On the basis of these preliminary results, the effectiveness of  
 171 PACSAI was evaluated on an ad-hoc prepared sample simulating  
 172 a real painting on canvas. A geometric pattern was drawn on the  
 173 preparation layer covering the canvas support, using a graphite  
 174 pencil. For the paint layer, we used primary red magenta  
 175 PV19–73900 (Quinacridone [ $C_2OH_{12}N_2O_2$ ], organic) extra-  
 176 fine acrylic color (Maimeri Brera, IT), as in our previous  
 177 work [13].

178 The sample was also characterized by a spectral-domain  
 179 OCT (Thorlabs Telesto-II, center wavelength, 1300 nm;  
 180 imaging depth, 3.5 mm; axial resolution in air, 5.5  $\mu\text{m}$ ; lateral  
 181 resolution, 7  $\mu\text{m}$ ); tomographic cubes 9 mm  $\times$  9 mm  $\times$  1 mm  
 182 were acquired with sampling step of 5  $\mu\text{m}$  on x and y, and  
 183 3.55  $\mu\text{m}$  on z, enabling the visualization of the graphite under-  
 184 drawing [Figs. 3(a) and 3(c)], the latter for comparison with the  
 185 near-infrared photoacoustic imaging session. A clearer cross-  
 186 sectional visualization of the paint layer was obtained with a  
 187 time-domain confocal-OCT prototype operating at 1550 nm  
 188 (axial resolution, 10  $\mu\text{m}$  in air; lateral resolution, 2.5  $\mu\text{m}$ ), de-  
 189 veloped at the Istituto Nazionale di Ottica (CNR-INO) [2]  
 190 scanning the surface along a line in its middle [light blue rectan-  
 191 gle in Fig. 3(a)]. The image was acquired with 5 and 1  $\mu\text{m}$   
 192 sampling step in y and z directions (image size 25 mm in length

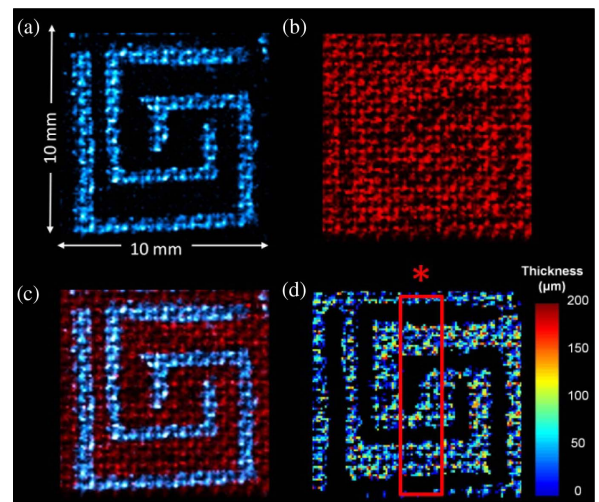


F3:1 **Fig. 3.** OCT reconstruction. (a) Tomocube showing the position of  
 F3:2 one tomographic image (light blue rectangle) acquired for the evalua-  
 F3:3 tion paint's thickness; (b) tomographic image (average paint thick-  
 F3:4 ness:  $82 \pm 9 \mu\text{m}$ ); and (c) x-y section of the tomocube showing  
 F3:5 the underdrawing.

and 1 mm in depth). The paint thickness was computed by  
 averaging 12 measurements along the selected profile, resulting  
 in  $82 \pm 12 \mu\text{m}$  [Fig. 3(b)].

Photoacoustic imaging with excitation at 1064 nm was  
 applied for the visualization of the underdrawing. The sample  
 was irradiated from its back side, and the raster scanning of a  
 10 mm  $\times$  10 mm area (200 by 200 pixels, 750 averaging mea-  
 surements per point) was performed. The photoacoustic waves  
 are generated in correspondence of the underdrawing pattern,  
 due to the high absorption of graphite in the near-infrared,  
 and transmitted through the paint, which is transparent in  
 this spectral region, thus revealing exclusively the hidden sketch  
 [Fig. 4(a)].

Following the near-infrared photoacoustic imaging session,  
 the evaluation of the paint's thickness was performed with  
 PACSAI [Fig. 4(b)]. The sample was scanned in the previously  
 analyzed region using the 532 nm wavelength, by averaging  
 over the same number of measurements for S/N improvement.  
 Since light at 532 nm is absorbed by both graphite and red  
 paint, the photoacoustic waves are generated over the entire  
 scanned area and are attenuated according to the local thick-  
 ness of the paint layer. However, the attenuation is greater in  
 correspondence of the underdrawing, due to the presence of two  
 superimposed absorbing materials, making the underdrawing  
 slightly visible in the final image. Differently, by irradiating  
 at 1064 nm, a much better image contrast is obtained due  
 to the greater difference in the materials absorption. The  
 two amplitude images acquired at 1064 and 532 nm were  
 merged in one single image as blue and red channels, respec-  
 tively [Fig. 4(c)]. Finally, to provide a map of paint thickness  
 over the scanned region, we shifted the time-domain wave-  
 forms recorded at 532 nm to the frequency domain, and  
 estimated the ATF values for each point within the range of  
 10–130 MHz. Acoustic signals originated by the underdrawing  
 regions are excluded from the analysis, since they are gener-  
 ated through the absorption of light by the graphite, rather than the



F4:1 **Fig. 4.** Combined photoacoustic imaging results: (a) MAP photo-  
 F4:2 acoustic image at 1064 nm (blue channel); (b) MAP image at 532 nm  
 F4:3 (red channel); (c) combined image (MAP at 1064 and 532 nm); and  
 F4:4 (d) PACSAI image with estimated paint thickness values on the color  
 F4:5 scale. The considered area for mean thickness measurement is high-  
 F4:6 lighted by the red rectangle\*.



229 overlying paint. We used the exponential decay fitting  
 230 parameters previously defined for primary red magenta [13]  
 231 to relate the estimated ATF values (megahertz) with layer  
 232 thickness  $x$  ( $\mu\text{m}$ ) according to the equation

$$x = -529 \times \ln\left(\frac{\text{ATF} + 0.03}{59.82}\right), \quad (1)$$

233 generating a PACSAI map corresponding to the overlying paint  
 234 layer thickness [Fig. 4(d)].

235 To validate the capabilities of PACSAI for thickness mea-  
 236 surements in paint layers, we have estimated the mean thick-  
 237 ness on a representative area [outlined by the red rectangle  
 238 in Fig. 4(d)] providing results with adequate statistical signifi-  
 239 cance, and compared it with the recorded OCT profiles. It has  
 240 to be noted that all extracted values smaller than the theoretical  
 241 axial resolution of the transducer ( $\sim 15 \mu\text{m}$ ) [20] were excluded  
 242 from the analysis as measurement artifacts. The obtained thick-  
 243 ness for the selected area (686 pixels) was  $77.8 \pm 1.9 \mu\text{m}$ ,  
 244 where measurement uncertainty corresponds to the standard  
 245 error of the mean. PACSAI results are in good agreement with  
 246 OCT thickness ( $\sim 82 \mu\text{m}$ ), demonstrating the reliability of the  
 247 proposed technique.

248 The combined photoacoustic imaging methodology has  
 249 proven to represent a promising diagnostic tool for the analysis  
 250 of painted artworks. The possibility of high-resolution imaging  
 251 to disclose the underdrawing, as well as the stratigraphic  
 252 information of paint layers, using the same experimental setup,  
 253 turns definitely useful in the case of conservation studies and  
 254 restoring operations. The fundamental and second-harmonic  
 255 wavelengths of a single laser source enable the application of  
 256 the presented method to a variety of drawing materials and  
 257 paint layers containing different pigments and binding media.  
 258 Moreover, displaying the results as images eases their interpre-  
 259 tation by final users (restorers, conservators, art historians) di-  
 260 rectly involved in conservation interventions. In view of future  
 261 applications on more complex cases, i.e., real objects with  
 262 multiple paint layers of different pigments or paint mixtures,  
 263 both the experimental apparatus and data processing have to  
 264 be implemented. For example, a multispectral photoacoustic  
 265 imaging approach could be used to excite efficiently a specific  
 266 paint layer by tuning an appropriate irradiation wavelength. For  
 267 paint mixtures, the exponential decay parameters could be  
 268 estimated as the weighted average of the acoustic attenuation  
 269 coefficients of individual paints, according to their light absorp-  
 270 tion properties. Furthermore, the system's performance could  
 271 be enhanced in terms of spatial and temporal resolution  
 272 (e.g., high frequency transducers, fast repetition rate lasers,  
 273 galvo scanners), by taking into account the total time of flight  
 274 of the recorded photoacoustic signal for each point. Finally, a  
 275 pure acoustic imaging technique (e.g., a pulse-echo modality),

using the same ultrasonic transducer, could provide additional 276  
 contrast to the combined photoacoustic imaging modality. 277  
 These upgrades well represent the potential of the proposed 278  
 technique and constitute objects of future research work. 279

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## REFERENCES

1. P. Targowski, M. Göra, and M. Wojtkowski, *Laser Chem.* **2006**, 35373 286  
(2006). 287
2. R. Fontana, A. Dal Fovo, J. Striova, L. Pezzati, E. Pampaloni, M. 288  
Raffaelli, and M. Barucci, *Appl. Phys. A* **121**, 957 (2015). 289
3. B. Blümich, J. Perlo, and F. Casanova, *Prog. Nucl. Magn. Reson.* 290  
*Spectrosc.* **52**, 197 (2008). 291
4. N. Proietti, D. Capitani, and V. Di Tullio, *Sensors* **14**, 6977 (2014). 292
5. A. J. L. Adam, P. C. M. Plancken, S. Meloni, and J. Dik, *Opt. Express* 293  
**17**, 3407 (2009). 294
6. G. Filippidis, M. Massaouti, A. Selimis, E. J. Gualda, J. M. Manceau, 295  
and S. Tzortzakis, *Appl. Phys. A* **106**, 257 (2012). 296
7. H. Liang, M. Mari, C. S. Cheung, S. Kogou, P. Johnson, and G. 297  
Filippidis, *Opt. Express* **25**, 19640 (2017). 298
8. G. Latour, J. P. Echard, M. Didier, and M. C. Schanne-Klein, *Opt.* 299  
*Express* **20**, 24623 (2012). 300
9. G. Filippidis, G. T. Tserevelakis, A. Selimis, and C. Fotakis, *Appl.* 301  
*Phys. A* **118**, 417 (2015). 302
10. M. Oujja, S. Psilodimitrakopoulos, E. Carrasco, M. Sanz, A. 303  
Philippidis, A. Selimis, P. Pouli, G. Filippidis, and M. Castillejo, 304  
*Phys. Chem. Chem. Phys.* **19**, 22836 (2017). 305
11. T. E. Villafana, W. P. Brown, J. K. Delaney, M. Palmer, W. S. Warren, 306  
and M. C. Fischer, *Proc. Natl. Acad. Sci. USA* **111**, 1708 (2014). 307
12. A. Dal Fovo, R. Fontana, J. Striova, E. Pampaloni, M. Barucci, M. 308  
Raffaelli, R. Mercatelli, L. Pezzati, and R. Cicchi, in *Proceedings of* 309  
*the International Conference LACONA XI* (2017). 310
13. G. J. Tserevelakis, A. Dal Fovo, K. Melessanaki, R. Fontana, and G. 311  
Zacharakis, *J. Appl. Phys.* **123**, 123102 (2018). 312
14. G. J. Tserevelakis, I. Vrouvaki, P. Siozos, K. Melessanaki, K. 313  
Hatzigiannakis, C. Fotakis, and G. Zacharakis, *Sci. Rep.* **7**, 747 314  
(2017). 315
15. B. E. Treeby and B. T. Cox, *Proc. SPIE* **7177**, 717716 (2009). 316
16. S. R. Cherry, R. D. Badawi, and J. Qi, *Essentials of In Vivo Biomedical* 317  
*Imaging* (CRC Press, 2016), p. 105. 318
17. X. L. Deán-Ben, D. Razansky, and V. Ntziachristos, *Phys. Med. Biol.* 319  
**56**, 6129 (2011). 320
18. H. Ammari, E. Bretin, V. Jugnon, and A. Wahab, *Mathematical* 321  
*Modeling in Biomedical Imaging II* (Springer, 2012). 322
19. H. Ammari, E. Bretin, J. Garnier, and A. Wahab, *Mathematical and* 323  
*Statistical Methods for Imaging* (American Mathematical Society, 324  
2011). 325
20. J. Yao and L. V. Wang, *Laser Photonics Rev.* **7**, 758 (2013). 326

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