# Germanium Photovoltaic Cells with MoO<sub>x</sub> Hole Selective Contacts

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# Abstract

12 Very thin, thermally evaporated  $MoO_x$  (x<3) layer has been used as transparent hole-selective contact on an n-type Germanium substrate to effectively demonstrate PV conversion capability. 13 The fabricated  $MoO_x/Ge$  heterojunction PV cell shows a photocurrent density of 44.8 mA/cm<sup>2</sup> 14 15 under AM1.5G illumination, which is comparable to that of conventional Ge PV cells. However, a low open-circuit voltage of 138 mV is obtained, which might be explained by the 16 17 presence of tunnelling mechanisms through the MoO<sub>x</sub>/Ge interface. To our knowledge, this is the first demonstration of a hole-selective contact made of transition metal oxide on an n-type 18 semiconductor different from c-Si. Thus, this work may have important implications toward the 19 20 development of new device architectures, such as novel low-cost Ge PV cells with possible 21 applications in multijunction solar cells and thermophotovoltaics.

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### Main text

23 Historically, the driving force for the use of Ge in photovoltaic (PV) applications has been as a 24 substrate for GaAs space solar cells (Miller and Harris 1980), the main reason being the higher 25 thermal conductivity and the possibility of manufacturing thinner and lighter wafers with Ge 26 than with GaAs. Later on, Ge/GaAs tandem solar cells were pursued to enhance the conversion 27 efficiency (Chand et al. 1986) by using the Ge bottom cell to convert the infrared part of the 28 solar spectrum. This progress eventually derived in the development of the current standard 29 technology for space solar cells that consists of triple junction Ge/GaAs/GaInP structures, with 30 AMO conversion efficiencies in the range of 28-30%. These cells have been also used in 31 terrestrial applications within concentrated-PV (CPV) systems, where they reached AM1.5D conversion efficiencies of 41.6% (R.R. King et al. 2009), just slightly below the current world-32

record for solar-to-electricity conversion efficiency of 46.0% (Dimroth et al. 2014). Apart from 33 34 solar applications, Ge PV cells have been considered as a low-cost replacement for low band gap III-V semiconductors in thermophotovoltaic (TPV) converters, in which thermal radiation is 35 36 directly converted into electricity by infrared sensitive PV devices (Bauer 2011; Chubb 2007). 37 In this context, Ge TPV cells could be used in a broad range of applications such as waste heat 38 recovery (Bauer et al. 2003), solar-thermal power (Ungaro, Gray, and Gupta 2015; Lenert et al. 39 2014; Alejandro Datas and Algora 2013), space power (A. Datas and Martí 2017), and energy 40 storage (Alejandro Datas et al. 2016), among many others.

41 Current state of the art of Ge PV cells consist of p-n junctions created by diffusion of dopants at 42 high temperatures (Bitnar 2003). For instance, p-n junctions in p-Ge have been created by 43 diffusion of V-group atoms (typically P and As) during the first growing step of GaInP or GaAs nucleation layers within a Metal-Organic CVD (MOCVD) reactor at temperatures of ~ 650°C 44 45 (Fernandez et al. 2008; Fernández 2010; Barrigón Montañés 2014). Other groups have used the 46 diffusion of Zn in n-Ge substrates within a LPE reactor (Khvostikov et al. 2002). In an effort to 47 reduce manufacturing costs of standalone Ge PV cells, IMEC reported devices with p-n 48 junctions created by spin-on diffusion of P on p-Ge by rapid thermal annealing at different 49 temperatures (450-700 °C) (Posthuma et al. 2007; van der Heide 2009; van der Heide et al. 50 2009) leading to the best reported 1-sun AM1.5G conversion efficiency for stand-alone Ge PV 51 cells of 7.9% (van der Heide et al. 2009). Surface passivation has been accomplished by 52 forming different kinds of heterojunctions on Ge surface, such as Ge/GaAs (Khvostikov et al. 53 2002) or Ge/GaInP (Fernandez et al. 2008; Fernández 2010; Barrigón Montañés 2014) by 54 MOCVD or LPE (Khvostikov et al. 2002), or Ge/a-Si (Posthuma et al. 2007; van der Heide 55 2009; van der Heide et al. 2009; Posthuma et al. 2005), Ge/SiNx (Nagashima, Okumura, and Yamaguchi 2007) and Ge/a-Si<sub>x</sub>C<sub>1-x</sub> (Fernandez et al. 2008; Fernández 2010) by PECVD. 56

In order to further reduce the fabrication cost of Ge PV cells, it is desirable to eliminate the high temperature diffusion, and complex MOCVD or PECVD processes. In this regard, a particularly appealing option consists of substituting the doping step by carrier-selective coatings with surface passivation properties that could be deposited at low temperatures. For this purpose, high electron-affinity transition metal oxides (TMOs) such as MoO<sub>3</sub>, WO<sub>3</sub>, and V<sub>2</sub>O<sub>5</sub>, are very interesting candidates that have already been found effective to produce hole-selective contacts on both n-type and p-type c-Si (Gerling et al. 2016; Battaglia et al. 2014; Bullock et al. 2014).

In this letter we report a Ge PV cell formed by a thin sub-stoichiometric  $MoO_x$  (x<3) layer on top of an n-type crystalline Ge (c-Ge) substrate, which behaves as a hole selective contact. To our knowledge, this is the first demonstration of a hole-selective contact made with a TMO on an n-type semiconductor different than c-Si. Thus, it might open the door to new device architectures, not only for PV applications, but also in photonics and CMOS electronics, where
the integration of TMOs is being investigated (Sanchez et al. 2016), along with the use of
different semiconductors having higher carrier mobilities and extended spectral response than cSi, such as Ge (Reboud et al. 2017; Toriumi and Nishimura 2017).

72 The PV cell structure was fabricated on (100) oriented, Czochralski, n-type Ge substrates ( $\rho$ = 73 0.37  $\Omega$ ·cm, 350 µm-thick). The substrate was cleaned by HCl: H<sub>2</sub>O (33%) immediately prior to 74 rear side passivation by PECVD of  $(i/n^+)$  a-SiC<sub>x</sub>:H (4/15nm, x~0.2) and a-SiC (80nm) stack 75 deposited at ~ 300°C. Next, the rear contact was created by laser firing of the a-SiC stack to produce an array of ~ 60 µm diameter local diffusion points, separated by 600 µm pitch. Laser 76 77 firing was accomplished by means of a ~1200 mW,  $\lambda$ = 1064 nm Nd/YAG laser system at a frequency of 4 kHz with 6 pulses per spot, following a similar approach than in (López et al. 78 79 2018) . The rear contact was finalized by means of an e-beam evaporated Ti/Pd/Ag metal stack 80 that provides lateral interconnection between fired points. The hole selective contact was 81 formed at the front side of the device by means of very thin (nominally 20 nm) MoO<sub>x</sub> layer 82 thermally evaporated from powdered MoO<sub>3</sub> sources at  $\sim 8 \cdot 10^{-6}$  mbar and a deposition rate of  $\sim$ 0.2 Å/s. A 75 nm-thick ITO layer was subsequently deposited by RF-Sputtering on top of the 83 84 MoO<sub>x</sub> layer to increase lateral electrical conductivity and minimize optical reflectivity. A sketch 85 of the full PV cell structure and the TEM image of the MoOx/ITO interface are shown in Figure 86 1, where a pronounced inter-diffusivity between the layers is clearly observed. The  $1x1 \text{ cm}^2$ active area of the PV cells was defined by conventional lithographic techniques followed by 87 88 mesa etching of the MoO<sub>x</sub>/ITO layers. Finally, the front Ag grid electrode (2 µm thick) was 89 evaporated through a shadow mask for a 4% contacted area

The current density-voltage (*J-V*) curve under 1-sun illumination is shown in Figure 2. The short-circuit current density ( $J_{SC}$ = 44.8 mA/cm<sup>2</sup>) outperforms that of the best performing state of the art Ge PV cells (43.2 mA/cm<sup>2</sup>) (van der Heide et al. 2009). On the other hand, a much lower open circuit voltage (138 mV) is measured, compared to those reported in (Fernández 2010; van der Heide et al. 2009) (up to 265 mV), which ultimately results in a lower FF (40.9 %), partially due to a non-optimized metal grid that introduces a series resistance of 0.65  $\Omega$ cm<sup>2</sup>. As a result, an AM1.5G conversion efficiency of 2.53 % is obtained.

97 External quantum efficiency (EQE) of the PV cell is shown in Figure 3 at short-circuit 98 conditions along with the EQE of Ge PV cells reported in (van der Heide 2009) for a direct 99 comparison. The improved EQE for wavelengths shorter than 600 nm might be explained by the 100 reduction of the recombination close to the front surface compared to the one existing in the 101 highly-doped emitters  $(10^{19}-10^{21} \text{ cm}^{-3})$  used in (van der Heide 2009). Such a low recombination 102 does not necessary indicate a good chemical surface passivation, i.e. strong reduction of 103 interface state density, but it could be related to a strong electric field that unbalances carrier 104 densities, i.e. field-effect passivation. In order to measure the electrostatic potential barrier built 105 at the junction  $(V_{bi})$ , capacitance-voltage measurements in reverse bias were performed 106 following the same approach than in (Almora et al. 2017) where similar structures on c-Si substrates are characterized. This data can be obtained by fitting the  $C^{-2}$  vs. V curve, known as 107 Mott-Schottky plot, using the following equation  $1/C^2 = 2(V_{bi} - V - 2k_BT/q)/q\varepsilon_S N_D,$ 108 where symbols have their usual meanings. By applying this model to the experimental data, we 109 get an almost perfect linear fit ( $R^2$ =0.99988) leading to  $V_{bi}$  =317±4 mV. Additionally, the 110 111 doping density  $(N_D)$  can be obtained from the slope of the curves leading to a  $N_D$  value of  $6.9\pm0.1\cdot10^{15}$  cm<sup>-3</sup>, which fully agrees with the Ge substrate specifications. The calculated V<sub>bi</sub> 112 indicates that the surface is highly inverted, i.e. hole density at the surface is even higher than 113 114 the doping density  $N_{\rm D}$ , reducing interface recombination due to the scarce availability of 115 electrons. This might explain the relatively high EQE values measured under short-circuit 116 conditions in the UV-visible range.

117 In order to investigate the origin of the low  $V_{OC}$ , a further understanding of the current 118 mechanisms taking place in the MoO<sub>x</sub>/Ge heterojunction is needed. With this aim, open-circuit 119 voltage ( $V_{OC}$ ) is measured as a function of photogenerated current ( $J_{ph}$ ) by means of a flash lamp. For every flash,  $V_{oc}$  values of the cell are recorded in an oscilloscope, while  $J_{ph}$  is 120 121 estimated from the light intensity measured by a reference Ge PV cell (Kerr, Cuevas, and Sinton 122 2001). It is well known that applying the superposition principle and taking into account that the 123 device is kept under open-circuit conditions,  $J_{ph}$  must be equal to the current that would be 124 measured in the cell at dark conditions and the series resistance has no effect on the 125 measurement. As a consequence, the analysis of  $J_{ph}$ - $V_{OC}$  curves enables the extraction of useful 126 information otherwise hidden by the series resistance effects in conventional dark J-V 127 characteristics. This advantage is crucial in our devices given the combination of relatively high 128 currents with significant series resistance. Figure 4 shows the  $J_{ph}$ - $V_{oc}$  curves measured at 129 temperatures ranging 293-323 K in 5 K steps. The experimental data are fitted to an exponential trend given by  $J_{ph} = J_0(T)[\exp(A(T) \cdot V_{OC}) - 1]$  and two examples for the highest and lowest 130 temperature measurement are also shown in Figure 4. Notice that in this model no series 131 132 resistance is included and consequently we have only two free parameters: the saturation current density,  $J_0(T)$ , and the exponential factor, A(T). In Figure 5 we show the Arrhenius plot of these 133 parameters where a constant value of  $A \approx 34$  V<sup>-1</sup> and an activation energy of 0.462 eV for  $J_0(T)$ 134 135 suggests that tunnelling mechanism dominates at the MoO<sub>x</sub>/Ge interface (Sze and Ng, n.d.). 136 This tunnelling current jeopardizes the electron blocking properties of the junction leading to a 137 high saturation current density and, thus, low  $V_{oc}$  values. A deeper knowledge of the band

- 138 structure and interface characteristics of  $MoO_x/Ge$  junction is needed to fully understand how 139 this tunnel mechanism takes place and to improve the obtained  $V_{oc}$  values.
- In conclusion, we have reported for the first time a heterojunction MoO<sub>x</sub>/Ge PV cell that 140 141 effectively demonstrates the possibility of creating hole selective contacts in n-type c-Ge. 142 Photovoltaic performance of the device shows excellent  $J_{sc}$  values (44.8 mA/cm<sup>2</sup>) mainly related to an enhanced spectral response at short wavelengths. On the other hand, low  $V_{oc}$  values (138) 143 144 mV) might be explained by an excess of tunnel current at the MoO<sub>x</sub>/Ge interface resulting in 145 high saturation currents. With evident room for improvement, these results could eventually open a new route for cost-reduction of Ge-based PV devices, including the development of new 146 147 kind of low cost thermophotovoltaic converters. Eventually, it could also open the door for the 148 integration of transition metal oxides in Ge photonics and CMOS electronics.

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Figure 3. External quantum efficiency of the Ge PV cell manufactured in this work along with
that of the Ge PV cell reported in (van der Heide 2009).



Figure 4.  $J_{ph}$ -V<sub>OC</sub> curves at different temperatures from 293 to 333 K in 5 K steps.





Figure 5. Activation energy and current as a function of temperature