# Preliminary Sketch of Biquaternions. By Prof. Clifford, M.A. 

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I. -

The vectors of Hamilton are quantities having magnitude and direction, bat no particular position; the vector $A B$ being regarded as identical with the vector $C D$ when $A B$ is equal and parallel to $C D$ and in the same sense. The translation of a rigid body is an example of such a quantity; for since all particles of the body move through equal distances along parallel straight lines in the same sense, the motion is entirely specified by a straight line of the given length and direction drawn through any point whatever. A couple, again, may be adequately represented by a vector; since the axis of a couple is any line of length proportional to its moment drawn perpendicular from a given face of its plane.

For many parposes, however, it is necessary to consider quantities which have not only magnitude and direction, but position also. The rotational velocity of a rigid body is about a certain definite axis, and equal rotations about two parallel axes are not equivalent to one another. A force acting upon a solid has a definite line of action, and equal forces acting along parallel lines differ by a certain couple. The difference between the two kinds of quantities is clearly seen when we consider the geometric calculus which is used for the study of each. In studying the motions of a particle or the composition of couples, the only construction required is that of the "force-polygon," and the theory involved is that of the addition of vectors; but in the static or kinematic of solids we require in addition the construction of the "linkpolygon," and there is involved the theory of the involution of lines in space, or of the linear complex.

The name vector may be conveniently associated with a velocity of translation, as the simplest type of the quantity denoted by it. In analogy with this, I propose to use the name rotor (short for rotator) to mean a quantity having magnitude, direction, and position, of which the simplest type is a velocity of rotation about a certain axis. A rotor will be geometrically represented by a length proportional to its magnitude measured upon its axis in a certain sense. The rotor AB will be identical with CD if they are in the same straight line, of the same length, and in the same sense; i.e., a vector may move anywise parallel to itself, but a rotor only in its own line.

The adlition of rotors will proceed by the rules which govern the composition of forces and rotations. Here, however, we come upon a very important break in the analogy between rotors and vectors. While
the sum of any number of vectors is always a vector, it will only happen in special cases that the sum of a number of rotors is a rotor. In fact, the composition of two forces whose lines of action do not intersect, or of two rotation-velocities whose axes do not intersect, gives rise to a system of forces on the one hand, and the most general velocity of a rigid body on the other. These still more complex quantities have been studied, and the theory of their addition or composition completely worked out, by Dr. Ball.

A system of forces may be reduced in one way to a single force $P$, and a couple $G$ whose plane is perpendicular to the line of action of the force, or central axis. Dr. Ball speaks of the system of forces as a wrench about a certain screw; the axis of the screw being the central axis, and the pitch being the ratio $\frac{G}{P}$ of the couple to the force. Similarly the velocity of a rigid body may be represented in one way only as a rotation-velocity $\omega$ about a certain axis combined with a translation-velocity $v$ along that axis. Dr. Ball speaks of this velocity as a twist-velocity about a certain screw; the axis of the screw being the axis of rotation, and its pitch the ratio $\frac{v}{\omega}$ of the translation to the rotation. A screw is here a geometrical form resulting from the combination of an axis or straight line given in position with a pitch which is a linear magnitude. A wrench is the association with this geometrical form of a magnitude whose dimensions are those of a force; a twist-velocity the association of a magnitude whose dimensions are those of an angular velocity. The extreme convenience of this nomenclature is well exemplified in the remarkable memoir above referred to.
Just as a vector (translation-velocity or couple) is magnitude associated with direction, and as a rotor (rotation-velocity or force) is magnitude associated with an axis; so this new quantity, which is the sum of two or more rotors (twist-velocity or wrench) is magnitude associated with a screw. Following up the annlogy thus indicated, I propose to call this quantity a motor ; the simplest type of it being the general motion of a rigid body. And we shall say that in general the sum of rotors is a motor, but that in particular cases it may degenerate into a rotor or a vector.

## II.

A quaternion is the ratio of two vectors, or the operation necessary to make one into the other. Let the rectors be AB and AC , as they may both be made to start from any arbitrary point $A$. Then $A B$ is made into $A C$ by turning it round an axis through A perpendicular to
 the plane BAC until its direction coincides with that of AC, and then
magnifying or diminishing it until it is of the same length as AC. The ratio of two vectors then is the combination of an ordinary numerical ratio with a rotation; or, as Hamilton expresses it, a quaternion is the product of a tensor and a versor. Since the point $A$ is perfectly arbitrary, this rotation is not about a definite axis; but is completely specified when its angular magnitude and the direction of its axis are given.
This quaternion* $\frac{A C}{A B}=q$, then, is an operation which, being performed on AB , converts it into AC , so that $q \cdot \mathrm{AB}=\mathrm{AC}$. The axis of the quaternion is perpendicular to the plane BAC; and it is clear that the quaternion operating upon any other vector AD in this plane will convert it into a fourth vector AE in the same plane, the angle DAE being equal to BAC and the lengths of the four lines proportionals. But a quaternion can only operate upon a vector which is perpendicular to its axis. If AF be any vector not in the plane BAC, the expression $q$. AF is absolutely unmeaning. A meaning is indeed subsequently given to an analogous expression in which the signification of AF is different. But it is very important to remark that so long as AF means a vector not perpendicular to the axis of $q$, the expression $q$. AF has no meaning at all.

Let us now consider what is the operation necessary to convert one rotor into another. There is one straight line which meets at right angles the axes of any two rotors, and part of which constitutes the shortest distance between them. Let AC be the shortest distance between the rotors $A B$ and $C D$. Then $A B$ may be converted into CD by a process consisting of three steps. First, turn $A B$ about the axis $A C$ into the position $A B^{\prime}$, parallel
 to CD. Then slide it along this axis into the position CD'. Lastly, magnify or diminish it in the ratio of $\mathrm{CD}^{\prime}$ to CD . The first two operations may be regarded as together forming a twist about a screw whose axis is $A C$ and whose pitch is

$$
\frac{\mathrm{AC}}{\text { circ. meas. of } \mathrm{BAB}^{\prime \prime}}
$$

The ratio of two rotors, then, is the combination of an ordinary numerical ratio with a twist. This twist is associated with a perfectly

* Professor Cayloy, by a very convenient notation, distinguishes $\left.\frac{\mid A C}{A B} \right\rvert\,$ and $\frac{A C \mid}{\mid A B}$; viz., $A B \frac{\mid A C}{A B \mid}=1$, but $\frac{A C \mid}{\mid A B} A B=1$. It should, I think, be a convention that $\frac{X}{Y}$ is always to mean $[\mathbf{X}]$, viz., the operation which converts $\mathbf{Y}$ into $X$, or whieh, coming after the operation $\mathbf{Y}$, is oquivalent to the oneration $\mathbf{X}$.
definite screw, and is only specified when its angular magnitude and the screw (involving direction, position, and pitch) are given. We may say also that just as the rotation (versor) involved in a quaternion is the ratio of two directions, so the twist involved in the ratio of two rotors is really the ratio of their axes.

Here again a remark must be made about the range of this operation. Using the expression tensor-twist to mean the ratio of two rotors (which is in fact a twist multiplied by a tensor), we may say that a tensor-twist can operate upon any rotor which meets its axis at right angles. Let $t$ denote the operation which converts $A B$ into $C D$, so that $t=\frac{\mathrm{CD}}{\mathrm{AB}}$, and $t \cdot \mathrm{AB}=\mathrm{CD}$; then if EF be any other rotor which meets $A C$ at right angles, the expression $t$. EF will have a definite meaning, viz., it will mean a rotor obtained by sliding EF along a distance equal to AC , turning it about AC as axis through an angle equal to $\mathrm{BAB}^{\prime}$, and altering its length in the ratio $\mathrm{AB}: \mathrm{CD}$. But if EF be a rotor not meeting $A C$, or meeting it at any other than a right angle, the expression $t$. EF will have no meaning whatever.

We have now defined the ratio of two rotors, and shown that like a quaternion it has a restricted range of operation. The question naturally arises, what now is the operation which converts one motor into another? We can answer this question very easily in the case in which the two motors have the same pitch; for in this case their ratio is a teusor-twist whose tensor is the ratio of their magnitudes and whose twist is the ratio of their axes. We are led to this by considering each motor as the sum of two rotors which do not intersect. Let $\alpha$ and $\beta$ be two such rotors, $t$ a tensor-twist whose axis meets them both at right angles; then tu is a rotor, say $\gamma$, and $t \beta$ is another rotor, say $\delta$. If therefore we assume the distributive law, we have
or

$$
\begin{gathered}
t(m a+n \beta)=n \gamma+n \delta, \\
t=\frac{m \gamma+n \delta}{m a+n \beta} .
\end{gathered}
$$

It is a mere translation of known theorems to say that the axis of $t$ meets at right angles the axes of the motors $m u+n \beta$ and $m \gamma+n \delta$, and that one of these axes is converted into the other by the same twist that makes $\boldsymbol{a}$ into $\gamma$ or $\beta$ into $\delta$.

The solution of this problem in the general case in which the pitches are different, is not so easy. In the first place, we must remember that every motor consists of a rotor part and a vector part, and that its pitch is determined by the ratio of these two parts. By combining a suitable vector with a motor, therefore, we may make the pitch of it anything we like, without altering the rotor part. Now let it be required to find the operation which will convert a motor A into a motor
B. Let $\mathrm{B}^{\prime}$ be a motor having the same rotor part as B , and the same pitch as $A$; and let $B=B^{\prime}+\beta$, where $\beta$ is a vector parallel to the axis of $B$. Then the ratio $\frac{B}{A}=\frac{B^{\prime}}{A}+\frac{\beta}{A}$; but $\frac{B^{\prime}}{A}$ is a tensor-twist, say $t$, and we may write $\quad \frac{\mathrm{B}}{\mathrm{A}}=t+\frac{\beta}{\mathrm{A}}$,
where it now only remains to find an operation which will convert a motor A into a vector $\beta$.

In order to do this, we must introduce a symbol whose nature and operation will at first sight appear completely arbitrary, bat will bo justified in the sequel. The symbol $\omega$, applied to any motor, changes it into a vector parallel to its axis and proportional to the rotor part of it. That is to say, it changes rotation about any axis into translation parallel to that axis, and a force into a couple in a plane perpendicular to its line of action. But if the rotation is accompanied by translation or the force by a couple, the symbol takes no account whatever of these accompaniments; and if made to operate directly on a vector, reduces it to zero. It follows from this that if it be made to operate twice apon a motor, it reduces it to zero; or $\omega^{2} A=0$ always. The portion of any expression which involves $\omega$ must therefore be treated as an infinitesimal of the first order; all higher orders being uniformly neglected.

Since then $\omega A=a$, a vector, and the ratio $\frac{\beta}{a}$ is a quaternion $q$ so that $q \alpha=\beta$, we may write successively
and then

$$
\begin{aligned}
& \beta=q \alpha=q \omega \mathrm{~A}, \\
& \frac{\beta}{\mathrm{~A}}=q \omega,
\end{aligned}
$$

or the ratio of tuo motors may be expressed as the sum of tue parts, one of which is a tensor-twist, and the other is $\omega$ multiplied by a quaternion.

The same ratio may be expressed in another form. Let an arbitrary point $O$ be assumed as the origin; then every motor may be expressed in one way as the sum of a rotor passing through 0 and a vector. Now the theory of rotors passing through a fixed point is exactly the same as that of vectors in general, and the ratio of any two of them is a tensor-twist whose pitch is zero, or what is the same thing, a quaternion whose axis is constrained to pass throngh the fixed point. If we use cursive Greek letters (as $a, \dot{\beta}^{\prime}$ ) in general to represent rotors through the origin, we may distinguish vectors from them by pretixing the symbol $\omega$; thas $\omega$ a denotes a vector parallel and proportional to the rotor $a$. The ratio $\frac{\beta}{\alpha}$ will then be a quaternion $q$, which is also tho

$$
\text { vol. Iv. }- \text { No. } 65 .
$$ now be requircd to find the ratio of two motors $a+\omega \beta, \gamma+\omega \delta$; or the value of the expression $\quad \frac{\gamma+\omega \delta}{a+\omega \bar{\beta}}$.

$$
\text { First, let } \frac{\gamma}{a}=q ; \text { then } q(\alpha+\omega \beta)=\gamma+q \omega \beta=\gamma+\omega q \beta \text {. }
$$

The symbol $q \beta$ has at present no geometrical meaning; for in general the rotors $a, \beta, \gamma$ will not be coplanar, and cannot therefore be operated on by the same quaternion $q$. If however (as in the Calculus of Quaternions) we consider all these quantities as expressed in terms of three rectangular unit rotors through the origin, $\frac{\delta-q l}{a}$ will be a perfectly definite quaternion $r$. The equation

$$
r a=\delta-q \beta
$$

is, like the equation $\quad q(\alpha+\omega \beta)=\gamma+\omega q \beta$,
at present purely literal and devoid of meaning. Yet if (remembering the properties of the symbol $\omega$ ) we add $\omega$ times the first equation to the second and assume the distributive law, we obtain

$$
(q+\omega r)(a+\omega \beta)=\gamma+\omega \delta
$$

In this way the ratio $\frac{\gamma+\omega \delta}{a+\omega \beta}$ is expressed in the form $q+\omega r$, which expression may conveniently be called a biquaternion. $\dagger$ The final equation, however, is not susceptible of interpretation in the same sense as the equation $q a=\gamma$. The expression $q+\omega r$ does not denote the sum of geometrical operations which can be applied to the motor $a+\omega \beta$ as a whole; and the ratio of two motors is only expressed by a symbol as the sum of two parts, each of which separately has a definite meaning in certain wther cases, but nut in the case in point. In following sections this difficulty will be partly overcome by showing that the system here sketehed is the limit of another in which it does not occur.

The preceding remarks may however explain, and be illustrated by, the following table:-

| (feomitheal Fomas | Quantity | Example | Ratio |
| :---: | :---: | :---: | :---: |
| Stuse con st. line | Vector on st. line | Addition or Subtraction | Signed Ratio |
| Dirction in plane | Vector in plane | Complex quautity | Complex Ratio |
| Dircetion in space | Vector in splace | Tranelation, Couplo | Quaternion |
| Axis | Rotor | 1Rutation-Velocity, Force | Twist |
| Screw | Mlotor | Twist-Velocity, System of Forces | Biquaternion |

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## III.

That geometry of three-dimensional space which assumes the Euclidian postulates has been called by Dr. Klein the parabolic geometry of space, to distinguish it from two other varieties, which assume uniform positive and negative curvature respectively, and which he calls the elliptic and hyperbolic geometry of space. The investigations which follow involve the postulates of elliptic geometry. As, however, the postulate of uniform positive curvature is not sufficient to define this, it may be worth while to devote a short space to an explanation of its nature.
Space of three dimensions is that the points of which may be associated with systems of values of three variables $x, y, z$. It is not in general possible, however, so to make this association that to every system of values there shall correspond in general one point, and to every point in general one system of values. When this is the case, the space is called unicursal. An algebraic space is one in which the position of a point may be uniquely defined by a set of values of periodic algebraic integrals, without exceptions which form a part of the space. Thus, unicursal spaces are a particular case of algebraic. Attending now to unicursal spaces only, we must observe that there are in general exceptions to the unique correspondence of points and valuesystems; namely, there are certain points to each of which correspond an infinite number of values. of the coordinates satisfying a certain equation or equations; and there are certain value-systems to which correspond, not points, but loci in the space. The assignment of these point-equations and loci-values and of their relations with one auother serves to determine the projective-connection of the space; and when once these are known, the whole of its projective geometry may be worked out. The point-equations and loci-values may or may not involve imaginary values of the variables or their coefficients; but in all cases they must be taken into account. The points which correspond to real systems of values are called real points; those which correspond to imaginary systems, imaginary poiuts : the study of these latter, which does not strictly belong to that of three-dimensional space, is undertaken only for the sake of the former.
Loci which correspond to linear equations between the coordinates may at prosent be called planes, and their intersections lincs; this is a purely projective defiuition, and these loci are not necessarily flat plancs and straight lines in the metrical sense. Points, lines, and planes are iucluded in the name elements.
The metric geometry of spaco* is the theory of the projective relations of certain fixed geometrical forms with all other geometrical forms, or of the invariant relations of certain fixed algebraic forms with all other

[^1]algebraic forms. The word power will be explained as much as is wanted in the sequel; meanwhile it may be said that these fixed forms (called all together the absolute) are given when we know the points, the lines, and the planes of the absolute, or say the elements of the absolute; and that the power of an element of the absolute in regard to any arbitrary element is infinite. In other words, we require in general equations of the absolute in point-, line-, and plane-coordinates respectively.

A unicursal space the points of which may be represented uniquely by value-systems of the coordinates $x, y, z$, without the exception of any point-equations or loci-values, is called a linear space. This is merely a projective definition, and leaves the absolute, therefore the whole of metric geometry, undetermined.
There is a particular determination of the absolute in a linear space which is of the utmost importance. It is that in which the points of the absolute are those of a certain quadric surface, while the lines and planes of the absolute are those which touch this surface; or in which the three equations of the absolute are of the second degree. There are three cases* to be considered, as being the only ones of which observed space can form a part:-
(1) Elliptic geometry; all the elements of the absolute are imaginary.
(2) Hyperbolic geometry; the absolute contains no real straight lines, and surrounds us. In this case, real points situate on the other side of the surface are called ideal.
(3) P'urabolic geometry; the surface degenerates into an imaginary conic in a real plane. The points of the nbsolute are points in the (real) plane of this conic; the lines and planes are the imaginary lines and planes which meet and touch the conic respectively.
The first of these suppositions will be made in what follows. It may be well here to set down in what it consists.
(1) The space to be considered is such that there is one point of it for every set of values of the coordinates $x, y, z$, and one set of values for evcry point, without any exception whatever.
(2) There is a certain quadric surface, called the absolute, all whose points and tangent planes are imaginary. If the line joining two poiuts $a, b$ meet the absolute in $i, j$, the quantity

$$
\frac{a b \cdot i j}{\sqrt{ }(a i \cdot a j \cdot l i . l j)} \equiv \overline{a b},
$$

(which is a function of anharmonic ratios, and therefore an invariant,) is called the power of the points $a, b$ in regard to one another, or of

[^2]either in regard to the other. The distance of these two points is an angle $\theta$ such that $\sin \theta=\overline{a b}$.
Similarly, if through the line of intersection of the planes A, B there be drawn the tangent planes. I, J to the absolute, the power of the planes $A, B$ in regard to one another is the quantity
$$
\frac{\mathrm{AB} \cdot \mathrm{IJ}}{\sqrt{(\mathrm{AI} \cdot \mathrm{AJ} \cdot \mathrm{BI} \cdot \mathrm{BJ})}}=\overline{\mathrm{AB}},
$$
and the angle between them is an angle $\phi$ such that
$$
\sin \phi=\overline{\mathrm{AB}} .
$$
(3) If two points are conjugate in regard to the absolute, they are distant a quadrant from one another; if two lines or planes are conjugate in regard to the absolute, they are at right angles. Thus all thie points at a quadrant distance from a given point are situate on its polar plane in regard to the absolute, and every plane through it cuts this polar plane at right angles. Every line has a polar line in regard to the absolute, such that every point on the polar line is distant a quadrant from every point on the line; aud every line which is at right angles to either meets the other. Through an arhitrary point can in general be drawn one line perpendicular to a given plane; namely, the line joining the point to the pole of the plane. If, however, the point is the pole of the plane, every line through it is perpendicular to the plane. Similarly, from a point not on the polar of a given line can be drawn one and ouly one perpendicular to the line; namely, the line through the point which neets the given line and its polar.
(4) In general, two lines can be drawn so that each meets two given lines at right angles, and these are polars of one another. One line may therefore be converted into another by rotation about two polar axcs. These axes are determincd as the lines which meet the troo given lines and their polars. If we travel continuously along one of these lines and draw perpendiculars on the other, one of these axes determines the shortest distance between the lines, and the other the longest. If thin these two are equal, the lines are equidistant along their whole length. Thus there is a case of exception in which two lines and their polars belong to the same set of generators of a hyperboloid; the lines are then equidistant along their whole length, and meet the same two generators of one system of the absolute. I shall use the word parallel to denote two lines so situated; and they shall be called right parallel or left parallel according as one is converted into the other by a right-landed or left-handed twist. Through an arbitrary point can be drawn ono right parallel and one left parallel to $\Omega$ given line; the angle between them is twice the distance of the point from the line. There are many points of analogy between the parallels here defined and those of parabolic geometry. Thus, if a line meet two parallel lines, it makes equal
angles with them; and a series of parallel lines meeting a given line constitute a raled surface of zero curvature. The geometry of this surface is the same as that of a finite parallelogram whose opposite sides are regarded as identical.
(5) A twist-velocity of a rigid body must be regarded as having two axes. For a motion of translation along any axis is the same thing as a rotation about the polar axis, and vice verst. Hence a twist-velocity is compounded of rotation-velocities about two polar axes; say these are $\theta, \phi$. Then the motion may be regarded either as a twist-velocity about a screw whose pitch is $\frac{\phi}{\theta}$ and whose axis is the first axis, or about a screw whose pitch is $\frac{\theta}{\phi}$ and whose axis is the polar axis. In general, then, a motor has two axes, and is expressible in one way only as the sum of two polar rotors. There is, however, one case of exception in which the axes of a motor are indeterminate; that, namely, in which the magnitudes of the two polar rotors are equal.* If a rigid body receive at the same time a rotation about an axis and an equal translation along it, all the points of the body will describe parallel straight lines; and the motion of the body is at the same time a rotation about auy one of these lines combined with an equal translation along it. Such a motion may be adequately represented by a line of given length drawn through any point whatever parallel to a given line. A motor of pitch unity, or which is its own polar, may therefore be regarded as having the nature of a vector, and shall in future be denoted by that name. For we may define a vector as a motor whose axes are indeterminate; and the case we are now considering is the only case of such indetermination which occurs in elliptic geometry. Vectors will be called right or left according as the twist of them is right- or left-handed.

Prop.: Every motor is the sum of a right and a left vector. For let A be a motor, and $A^{\prime}$ the polar motor; then we have $A=\frac{1}{2}\left(A+A^{\prime}\right)$ $+\frac{1}{2}\left(A-A^{\prime}\right)$. Now $A+A^{\prime}$ and $A-A^{\prime}$ are both motors of pitch unity, but one right-handed and the other left-handed.

## IV.

A fixed point being chosen as origin, let three lines perpendicular to one another be drawn through it, and let three onit-rotors haring these lines as axes be denoted by the symbols $i, j, l$. Then every rohur through the origin will be denoted by an expression of the form $i x+j y+k z$, where $x, y, z$ are scalar quautities, or the ratios of magnitudes. The symbols $i, j, \not, j$ shall have also another meaning ; viz., each

[^3]shall siguify the rotation through a right angle about its axis of any rotor which meets that axis at right augles. When they are performed on rotors passing through the origin, these operations satisfy the equations $i^{3}=j^{2}=k^{3}=i j k=-1$, by the ordinary rules of quaternions ; and it is easy to see that the same equations huld good when the operations are performed on rotors not passing through the origin. The compound symbol ix+jy+liz is also to have an aualogous secondary meaning; viz., a rectangular rotation about the axis of the rotor which it previously denoted, combined with a tensor $\sqrt{ }\left(x^{2}+y^{2}+z^{2}\right)$. It can operate only on a rotor which meets its axis at right angles. This being so, the ratio of any two rotors throngh the origin is a quaternion of the form $q \equiv w+i x+j y+l c z \equiv w+\rho$, say. The axis $\rho$ of this quaternion is perpendicular to the plane of the two rotors. If $\boldsymbol{a}$ be a rotor through the origin and $q$ a quaternion, the product $q$ a can be formed according to the Hamiltonian rules of multiplication, and is in general a quaternion $r$. In this general case the equation $q a=r$ can only be interpreted by giving to $a$ its secondary meaning; and the translation of this statement into words is as follows :-If a rotor be capable of being successively operated upon by the rectangular versor $\boldsymbol{a}$ and the quaternion $q$, the final result will be the same as if it had been originally operated upon by the quaternion $r$. If, however, the axes of $q$ and $a$ are at right angles, the scalar part of $r$ will be wanting, and we may write the equation $q a=\rho$. This equation is now susceptible of a primary interpretation; viz., the quaternion $q$ operating on the rotor a produces the rotor $\rho$; although the secondary interpretation does not cease to be true.

With such conventions, the two sides of the equation

$$
(q+r) s=q s+r s
$$

(in which $q, r, s$ are quaternions) have always the same meaning when both are interpretable; which is what is meant by saying that the distributive law holds good for these symbols.

The ratio of two rotors which do not meet is a twist which in general has perfectly definite axes. But when the rotors are polars of one another, the axes of the twist are indeterminate; for any line meeting both meets them at right angles, and will serve for an axis. It is therefore always possible to find a twist which shall simultaneously convert two given rotors into their polars; and any two rectangalar twists with pitch 1 or -1 have a pair of common rotors on which they can operate, and which they convert into one another. Hence we may consider that

All rectangular twists of pitch 1 are equivalent to one another; and all rectangular twists of pitch -1 are equivalent to one another.

The rectangular twist of pitch 1 shall bo denoted by the symbol $\omega$; the exprossion wa will denote tho rotor polar to a and equal to it in magnitude, 8 btained from it by a left-handed twist. During the
operation of this twist, every point of the rotor describes a straight line; if therefore the twist be continued through two right angles, the rotor will be replaced in its original position, not reversed; we have therefore

$$
\omega^{2}=1
$$

Every motor can be expressed as the sum of two rotors, one passing through the origin and the other being polar to a rotor through the origin. The general expression for a motor is therefore

$$
a+\omega \beta .
$$

This will represent a rotor if the two rotor constituents intersect, or if cach is perpendicular to the polar of the other; or if $S a \beta=0$.

Let now

$$
\xi=\frac{1+\omega}{2}, \quad \eta=\frac{1-\omega}{2} ;
$$

then

$$
\begin{gathered}
\xi^{2}=\frac{1+2 \omega+\omega^{2}}{4}=\frac{2+2 \omega}{4}=\xi, \\
\eta^{2}=\frac{1-2 \omega+\omega^{2}}{4}=\frac{2-2 \omega}{4}=\eta, \\
\xi \eta=\frac{1-\omega^{2}}{4}=0 .
\end{gathered}
$$

Any motor $a+\omega \beta$ can also be expressed in the form $\xi \gamma+\eta \delta$. It is clear that $\xi \gamma$ is the right vector part of this motor, and that $\eta \delta$ is the left vector part. If we multiply $\xi \gamma+\eta \delta$ by $\xi$, the result is merely $\xi_{\gamma}$; so the effect of multiplying a motor by $\xi$ is merely to pick out the right vector part of it. The symbols $\xi, \eta$ are thus in a certain sense selective symbols, and are analogons to the S and V of quaternions.

Ratio of two motors.-We can find immediately now the operation which converts a motor $\xi \gamma+\eta \delta$ into a motor $\xi a+\eta \beta$. For if we perform the operation $\quad\left(\xi \frac{a}{\gamma}+\eta \frac{\beta}{\delta}\right)(\xi \gamma+\eta \delta)$,
remembering the laws of multiplication of $\xi, \eta$, we obtain the result $\xi a+\eta \beta$. If then $\frac{a}{\gamma}=q, \frac{\beta}{\delta}=r$, we may write

$$
\frac{\xi \alpha+\eta \beta}{\xi \gamma+\eta \delta}=\xi \frac{\alpha}{\gamma}+\eta \frac{\beta}{\delta}=\xi q+\eta r,
$$

and the latter may be written in the form

$$
\frac{q+r}{2}+\omega \cdot \frac{q-r}{2}=s+\omega t
$$

showing that the ratio of two motors is a liquaternion.
The motor $\xi a+\eta \beta$ will be a rotor if
or if

$$
S(\alpha+\beta)(\alpha-\beta)=0,
$$

and it is easy to see from this that the biquaternion $\xi q+\eta r$ will be a twist, or the ratio of two rotors, if $\mathrm{T} q=\mathrm{T} r$.

## V.

1. Position-Rotor of a Point.-The coordinates of a point in regard to a quadrantal tetrahedron 1234 keing $x_{1}, x_{2}, x_{3}, x_{4}$, the equation to the absolute is $\Sigma x^{2}=0$. The rotor from the origin (the point 4) to the point $x$ is represented by $i_{1} \frac{x_{1}}{x_{4}}+i_{2} \frac{x_{2}}{x_{4}}+i_{3} \frac{x_{3}}{x_{4}}$, or $\Sigma i_{k} \frac{x_{k}}{x_{4}}(k=1,2,3)$, where $i_{1}, i_{2}$, $i_{3}$ are rotors along the edges of the tetrahedron from the origin to the middle points of the edges. The tensor of this rotor is the tangeut of the angular distance from the origin to the point it represents. For if

$$
\rho=i_{1} \frac{x_{1}}{x_{4}}+i_{2} \frac{x_{2}}{x_{4}}+i_{3} \frac{x_{3}}{x_{4}}
$$

$$
[\mathrm{T} p]^{2}=\frac{x_{1}^{2}+x_{2}^{2}+x_{3}^{2}}{x_{4}^{2}}=\tan ^{2} \overparen{O x} \text {, where } o \text { is the origin. }
$$

The angular distance from the origin to a point has an infinite number of values, which differ by multiples of $\pi$. If therefore a rotor be considered to have this angular distance as its length, the rotor of a point can only be defined by such an equation as $\breve{\rho} \equiv \breve{\boldsymbol{\alpha}}$ (mod. $\breve{\pi}_{\text {. }}$ ). To obviate this indetermination, there is required a one-valued anicarsal function having the period $\pi$; the tangent of the angular distance is hereby completely singled out.
2. Equation of a Straight Line.-Let OM be the perpendicular from the origin 0 apon the straight line MP; and let ON be a line perpendicular to OM in the plane MOP. Then from the triangle MOP we have $\frac{\tan \mathrm{OM}}{\tan \mathrm{OP}}=\cos \mathrm{MOP}$;
or if

$$
\begin{gathered}
\mathrm{OM}=a, \quad \mathrm{OP}=\rho, \quad \mathrm{ON}=\beta, \\
\mathrm{T} \alpha=\mathrm{T} \mu \cos \mathrm{MOP} ;
\end{gathered}
$$

so that $a$ is the component of $\rho$ in the direction
 OM, and we have $\rho=a+\beta x$, where $x$ is some scalar.

By varying $x$, then, we get all the points in the line MP. But if $a_{1}$ is any particular value of $\rho$, the equation may just as well be written

$$
\rho=\alpha_{1}+\beta x,
$$

where now $a_{1}$ is not necessarily perpendicular to $\beta$.
This form may be reduced to the preceding as follows:
To find the perpendicular from $O$, put $\delta \mathrm{T} \rho=0$; this gives

$$
S a_{1} \beta+\beta^{2} x=0,
$$

and the equation becomes $\rho=a_{1}-\beta S \frac{a_{1}}{\beta}-\beta x$,
where $a_{1}-\beta \mathrm{S} \frac{a_{1}}{\beta}=a$ of the former equation.

## 3. Rntor along Straight Line whose Equation is given.

Let OR be the rotor through the origin which has right parallelism with MP. Then $\angle \mathrm{NOR}=\mathrm{OM}$. Let OK be perpendicalar to $O N$ and $O M$, aud of such length that $\frac{\tan \mathrm{OK}}{\tan \mathrm{ON}}=\tan$ NOR. Then, if $\gamma=\mathrm{OK}$, $\mathrm{OR}=\beta+\gamma . \quad$ Now $\quad \frac{\mathrm{T} \gamma}{\mathrm{T} \beta}=\mathrm{T} a, \quad$ and $\mathrm{U}_{\gamma}=\mathrm{Ua} \beta$, since $\gamma$ is perpendicular to $a$ and $\beta$. Hence $\gamma=a \beta$; and if R be a
 rotor along MP, $m$ a scalar,
right vector of $R=\xi R=m \xi(\beta+\gamma)=m \xi(\beta+a \beta)$, so $\quad$ left vector of $R=\eta R=m \eta(\beta-\gamma)=m \eta(\beta-a \beta)$;
therefore

$$
\mathrm{R}=m(\beta+\omega a \beta)
$$

Now if R have the same length as $\beta$, we have

$$
\beta^{2}=\mathrm{R}^{2}=m^{2}\left(\beta^{2}+\overline{a \beta^{2}}\right)=m^{2} \beta^{3}\left(1-\alpha^{2}\right) ;
$$

therefore

$$
\mathrm{R}=\frac{\beta+\omega a \beta}{\sqrt{ }\left(1-a^{2}\right)}
$$

Conversely, equation to axis of rotor $\gamma+\omega \delta$ is

$$
\rho=\frac{\delta}{\gamma}+\gamma x
$$

This finds the rotor in the case in which $\rho=\alpha+\beta x$, where $\mathrm{Sa} \beta=0$. But in the general case we have only to write the equation in the form
whence

$$
\begin{aligned}
& \rho=a-\beta \mathrm{S} \frac{a}{\beta}+\beta x, \\
& R=\frac{\beta+\omega\left(a-\beta \mathrm{S} \frac{a}{\beta}\right) \beta}{\sqrt{ }\left(1-a^{2}-\beta^{2} \mathrm{~S}^{2} \frac{a}{\beta}+2 \mathrm{Sa} \beta \mathrm{~S} \frac{a}{\beta}\right)} \\
&= \frac{\beta+\omega \mathrm{Va} \beta}{\sqrt{ }\left(1+\operatorname{Sa} \beta \mathrm{S} \frac{a}{\beta}-a^{2}\right)}
\end{aligned}
$$

4. Rotor ab joining Points whose Position-Rotors are a, $\beta$.

The equation of this rotor is

$$
\begin{gathered}
\rho=\alpha+(\beta-\alpha) x \\
m \mathrm{R}=\beta-\alpha+\omega \nabla \alpha \beta .
\end{gathered}
$$

whence.
Now if $a_{1}, a_{2}, a_{3}, a_{4} ; b_{1}, b_{2}, b_{3}, b_{4}$ are the coordinates of the points, we have $[T R]^{2}=\tan ^{2} a b=\frac{\Sigma\left(a_{h} b_{k}-a_{k} b_{h}\right)^{2}}{\left(\Sigma a_{h} b_{h}\right)^{2}}=-\frac{(\alpha-\beta)^{2}+(\nabla a \beta)^{2}}{(1-S a \beta)^{2}}$,
therefore

$$
\mathrm{R}=\frac{\beta-a+\omega V a \beta}{1-\operatorname{Sa} \beta}
$$

Cor.-If $\rho$ be the rotor of a variable point on a carve, $d \lambda$ a rotor along the tangent of length equal to the arc of the curve between $\rho$ and $\rho+d \rho$, we have

$$
d \lambda=\frac{d \rho+\omega \nabla \rho d \rho}{1-\rho^{2}}
$$


5. Rotor parallel to $\beta$ through Point whose Position-Rotor is a.

The general equation to a line through the point $\alpha$ is $\rho=a+\lambda x$, where $\lambda$ is any rotor through the origin. A rotor along this line is $\lambda+\omega \mathrm{Va} \lambda$; if this is right parallel to $\beta$, we have
or

$$
\begin{aligned}
\xi(\lambda+\nabla a \lambda) & =\xi \beta, \quad[\xi \omega=\xi] \\
\lambda+\nabla a \lambda & =\beta .
\end{aligned}
$$

Operating by $S a$, we have, since $S . a \nabla a \lambda=0$,

$$
S a \lambda=S a \beta
$$

whence, by addition,

$$
\lambda+a \lambda=\beta+S a \beta,
$$

$$
\text { and } \quad \lambda=(1+a)^{-1}(\beta+S a \beta)=\beta-(1+a)^{-1} \nabla a \beta
$$

The rotor required is $\lambda+\omega \nabla u \lambda$, or $\lambda+\omega(\beta-\lambda)$. This becomes, then,

$$
\beta-(1+a)^{-1} \nabla a \beta+\omega(1+a)^{-1} \nabla a \beta=\beta-2 \eta(1+a)^{-1} \nabla a \beta .
$$

Instead of operating by $S a$ on the equation

$$
\lambda+\nabla a . \lambda=\beta,
$$

we might have operated with $\nabla a$, and got

$$
\begin{gathered}
\nabla a \lambda+a \mathrm{~V} a \lambda=\mathrm{V} a \beta, \text { since } \nabla . a \nabla a \lambda=a \nabla a \lambda, \\
\nabla_{r} \lambda=(1+a)^{-1} \mathrm{~V} a \beta,
\end{gathered}
$$

therefore
and

$$
\lambda=\beta-\nabla a \lambda=\beta-(1+a)^{-1} \nabla a \beta
$$

Similarly, we have for the rotor left parallel to $\beta$,

$$
\lambda=\beta+(1-a)^{-1} \nabla a \beta,
$$

and the rotor is

$$
\begin{aligned}
\lambda+\omega(\lambda-\beta) & =\beta+(1-a)^{-1} \nabla a \beta+\omega(1-a)^{-1} \nabla a \beta \\
& =\beta+2 \xi(1-a)^{-1} \nabla a \beta .
\end{aligned}
$$


[^0]:    * It folluws from this that $\omega \ell=\eta \omega$, or $\omega$ is commutative with quaternions.
    + Hamilton's liquetetrmion is a quatemion with complex coefficients; but it is conveniant (as l'ruf. Pioreo remarks) to suppose from the berbining that all scalars may be complex. As the word is thus no longer wanted in its old meaning, I have made bold to use it in a new one.

[^1]:    * This theory of metric geometry is due to Prof. Cayluy: Sisth Mromoir on Quantics, Phil. 'Truns., 1850.

[^2]:    * On this division sce Dr. Klein, "Uclor diu so-genannte Nicht-Euklidischo Geomethie," Math. Annalen, Bd. 4. The second case is the geumetry of Lobatsehewsky and Bolyai.

[^3]:    * This motion is described in another connection by Drs. Klein and Lio, Math. Annalen, Bd. 4; it is a transformation of the absolute into itself in which two generators remain unaltered.

