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Toward new energy-rich molecular systems: from N_{10} to N_{60}

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Abstract

The electronic structure calculations locate the local minima for bicyclic N_{10} and the fullerene analog N_{60} , both as high-energy density species. The bridging N–N bond in N_{10} is quite remarkably strong ($\Delta E [N_{10} \rightarrow 2N_5] = 93$ kcal/ mol), yet flexible to allow a facile rotation of one ring with respect to the other. This property could permit N_{10} to serve as a building block for specific clustering into the nitrogen buckminsterfullerene structure. © 2000 Elsevier Science B.V. All rights reserved.

1. Introduction

In contrast to the same group elements (phosphorus and arsenic), nitrogen has thus far resisted transformation from its basic molecular diatomic phase. The possible existence of polymeric forms of nitrogen, in which the strong triple bond of the diatom is transformed and replaced by single or single and double bonds, is of particular interest due to the energetic properties associated with these materials. As potential high-energy density materials from an abundant natural source, they become leading candidates for alternative energy storage media.

The prospect of destabilizing the strong $N \equiv N$ triple bond under high pressure invited several experimental [1,2] and theoretical [3–5] attempts to locate and characterize bulk polymeric phases. These studies, however, appeared to be contradictory. Calculations suggested polymeric regimes to exist below 100 GPa, at 70 and 50 GPa, while

Quantum chemical-based computational studies at various levels of theory [8–12] have shown that N₄ and N₈ are metastable species, having (several) local minima on their respective potential energy surfaces. Coupled cluster calculations, CCSD(T), by Lee and Rice [10] showed the energy difference ΔE (N₄—2N₂) to be 186 kcal/mol, and an energy barrier to decomposition of 61 kcal/mol. Later work examined reaction pathways to other structures [11], and determined the lowest-energy conformation of N₄ to be an open-chain triplet [13] with a dissociation barrier of 63 kcal/mol [14]. The results suggested tetrahedral N₄ to have a reasonably long lifetime, and could be trapped and studied by matrix isolation techniques.

Computational studies on polyazacubane [15,16] indicated that the high-energy content of

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the diamond-anvil-cell experiments failed to record a transition from the diatomic form for pressures up to 180 GPa. A non-static, shock compressed nitrogen experiment [6] recorded an anomaly at pressure above 30 GPa that was later interpreted as a dissociation to dense atomic phase [7].

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these species, $\Delta E (N_8 - 4N_2) = 423$ kcal/mol for cubic N_8 using DZP CCSD(T) [12], is not due to strain energy but rather to a weak N-N single bond. Five isomers of N₈ have been determined to be metastable [12,17] using highly correlated methods. In order to guide laboratory synthesis, protonated structures of N₄ and N₈ have been optimized and harmonic vibrational frequencies were also calculated [18]. Recent studies reported 10 Hartree-Fock and MP2 isomers of N₈ corresponding to analogous CH structures [19], and the isomerization reaction of N₈ cubane to planar bicyclic (pentalene-like) structure with multiconfigurational self-consistent field and second-order perturbation theory (CASPT2) treatment [20]. Nine single-bonded structures of N_x (x = 8-24) were also considered [21] at the DFT(B3LYP)/ 6-31G** level of computation. Similar computational studies should eventually lead to a viable pathway for the experimental characterization and production of some form of polymeric nitrogen.

We are tempted to address the following: (a) Could a nitrogen fullerene analog, N_{60} be a possible candidate, and (b) if so, what is the reasonable building block for such a cluster. A more general theme for our subsequent work is that computational methods are currently mature enough to help guide experimental efforts in establishing new polymorphs of nitrogen, through a combination of electronic structure and molecular dynamics calculations.

In this Letter, several computational methods are employed to locate the extrema of bicyclic N_{10} . Hartree-Fock self-consistent field (SCF), secondorder Møller-Plesset (MP2), quadratic configuration interaction with single and double substitution (QCISD), and density functional (B3LYP) are used as implemented in GAUSSIAN 98 [22] and the parallel NWChem [23] packages of codes. Two correlation consistent basis sets of Dunning [24], cc-PVDZ and cc-PVTZ were used in the expansion of the electronic wave functions. We also report single-point QCISD(T) calculations in which approximate triples' contributions to energy are included. Geometry search, not constrained by symmetry, located a planar structure (D₂h symmetry) as a saddle point and an orthogonal

structure (D₂d symmetry) as a local minimum. Only 3–6 kcal/mol, depending on the level of computation, separate the two structures. This molecule is postulated herein to serve as a building block for yet a super-high-energy content N₆₀ cluster, which is shown to be a local minimum at the AM1 [25] and SCF levels of theory. Harmonic vibrational frequencies of N₁₀ and infrared active modes of N₆₀ are reported.

2. Results and discussion

$2.1. N_{10}$

Table 1 lists the optimized geometry of the minimum and transition state structures of N₁₀ cluster at various levels of computational treatment, as depicted in Fig. 1. The bond lengths in the two structures are very similar except in the bridging bond, \mathbf{R}_2 , where it is stretched by about 0.02 A in the transition state structure from that of the equilibrium. SCF treatment invariably provided shorter bond lengths by as much as 0.04 A than the higher level methods. Bond angles, however, are consistent in all levels of treatment. Remarkable agreement should be noted between the DFT(B3LYP) and QCISD results. The cc-PVDZ basis seems to provide adequate values when compared with the cc-PVTZ basis. Furthermore, the geometry and frequencies for the minimum structure are in good agreement when compared with the unpublished B3LYP results of Bartlett and Fau [26] using the augmented cc-PVDZ basis.

One expects single and double N—N bonds to exist in the N₁₀ molecule. The bond distances in Table 1 show that \mathbf{R}_1 , \mathbf{R}_2 and \mathbf{R}_4 have similar magnitudes, while the \mathbf{R}_3 bond is much shorter. Typical experimental bond lengths for N—N single is 1.449 Å and for N=N double 1.252 Å [27]. The calculated bonds \mathbf{R}_1 , \mathbf{R}_2 and \mathbf{R}_4 are very close to the *averaged* experimental single and double bonds, indicative of substantial charge delocalization in this system. The \mathbf{R}_3 bond, however, is distinctly close to the N=N bond value. The absence of purely single-bonded N—N moiety, and the existence of higher bond order in this system are encouraging for the stabilization effect, since N—N

| | SCF | B3LYP | MP2 | QCISD | |
|----------------|---------------|---------------|---------------|-------|--|
| R ₁ | 1.309 (1.305) | 1.346 (1.338) | 1.342 (1.322) | 1.345 | |
| | 1.298 (1.293) | 1.335 (1.328) | 1.339 (1.319) | 1.333 | |
| \mathbf{R}_2 | 1.332 (1.332) | 1.353 (1.350) | 1.357 (1.344) | 1.351 | |
| | 1.354 (1.353) | 1.371 (1.369) | 1.370 (1.356) | 1.374 | |
| R ₃ | 1.246 (1.242) | 1.281 (1.275) | 1.316 (1.304) | 1.282 | |
| | 1.259 (1.256) | 1.294 (1.288) | 1.323 (1.309) | 1.296 | |
| \mathbf{R}_4 | 1.346 (1.345) | 1.378 (1.373) | 1.362 (1.344) | 1.388 | |
| | 1.325 (1.322) | 1.359 (1.355) | 1.354 (1.337) | 1.362 | |
| α1 | 111.9 (111.8) | 112.5 (112.3) | 114.4 (114.5) | 112.6 | |
| | 112.9 (112.9) | 113.5 (113.3) | 114.9 (114.9) | 113.8 | |
| α2 | 109.3 (109.2) | 109.6 (109.5) | 109.8 (109.7) | 109.4 | |
| | 109.4 (109.4) | 109.8 (109.6) | 110.0 (109.8) | 109.6 | |
| α3 | 124.0 (124.1) | 123.7 (123.8) | 122.8 (122.8) | 123.7 | |
| | 123.5 (123.5) | 123.2 (123.3) | 122.5 (122.5) | 123.1 | |

Optimized molecular geometry of N₁₀ minimum and transition state structures as in Fig. 1 at various levels of theory^a

^a Distances are in Å, and angles in degrees. Calculated parameters for the transition state structure are below that of the minimum. Results using cc-PVTZ basis set are in parenthesis.



Fig. 1. (a) The minimum and (b) the transition state structures with molecular parameters of N_{10} .

single bonds are in general weak, with greater likelihood to entice decomposition.

The harmonic vibrational frequencies for the minimum structure are reported in Table 2. SCF modes exhibit tendencies toward producing higher frequencies, consistent with the calculated shorter bond distances at this level of treatment. The B3LYP and MP2 calculations are in more general agreement with one another. The strongest IR bands are exhibited in the B₂ mode at 856 (969) and 1082 (1070) wavenumbers at the B3LYP (MP2) computational levels. Again, use of the larger cc-PVTZ basis set did not seem to substantially alter the results as determined by the SCF and B3LYP methods.

The planar, D_2h symmetry, structure (Fig. 1b) was confirmed to be a saddle point with the imaginary (negative) frequency being 93, 76, and 75 cm⁻¹ at the cc-PVDZ SCF, B3LYP and MP2 levels, respectively. This mode is associated with a rocking motion of one ring with respect to the other, and intrinsic reaction coordinate (IRC) calculations at the B3LYP level proved this internal motion. Bonding in this structure does not differ, as expected, from the minimum D_2d structure except in about 0.02 Å stretch for the ring-

Table 1

connecting bond, \mathbf{R}_2 . Energetically, the saddle point is placed 4.5, 3.3, 6.1 and 5.4 kcal/mol higher than the minimum structure using the cc-PVDZ B3LYP, MP2, QCISD, and QCISD(T) methods, respectively. The triple-zeta basis, cc-PVTZ does not change these values appreciably, giving 4.9 and 3.2 kcal/mol for B3LYP and MP2 treatment, respectively. Further calculations at the MP2 and B3LYP levels showed that there is no other barrier separating these two structures, indicative of the flexibility of this bond with respect to rotation.

We considered the decomposition of N_{10} into two rings, $N_{10} \rightarrow 2N_5$, and optimized the N₅ ring radical at the cc-PVDZ QCISD configuration interaction with single and double excitations (CISD), MP2 and B3LYP. The C_2v symmetryoptimized structure for the single ring has similar bonding in \mathbf{R}_1 and \mathbf{R}_4 as in N_{10} . The significant difference is that the \mathbf{R}_3 bond is now more reminiscent of a purely single bond, $\mathbf{R}_3 = 1.607$, 1.568, and 1.536 A at the QCISD, CISD, and B3LYP levels, respectively. MP2 optimizations tended toward opening of the ring, and thus, a highly stretched (non-bonded) \mathbf{R}_3 distance. Furthermore, energy second-derivative calculations at all the above levels confirmed that the optimized structure is a first-order saddle point, so that it is unlikely to be stable and further decomposition is expected. The reaction energy, $\Delta E (N_{10} \rightarrow 2N_5)$ was calculated to be 93 and 84 kcal/mol at the QCISD and B3LYP levels, indicative of a remarkably strong connecting (\mathbf{R}_2) bond. For the purpose of comparison, we note that the C-N and N–N bond energies in model energetic materials such as nitromethane (H_3C-NO_2) and nitramine (H_2N-NO_2) are 60 and 45 kcal/mol, respectively. Finally, we determined the reaction energy $\Delta E (N_{10} \rightarrow 5N_2)$ to be 286, 270, and 224 kcal/mol at the QCISD, MP2, and B3LYP levels, respectively. These values are reflective of the interest in polynitrogen as high-energy density materials.

The relative stability of the N–N (\mathbf{R}_2) bond and the internal rotation flexibility in N₁₀ as shown above afford a plausible avenue for the synthesis of such materials. The only stable molecular forms of nitrogen known to date are the diatomic N_2 and the azide ion N_3^- , although the synthesis of the ion N_5^+ (C_{2v} symmetry) in a complex with As_xF_v anion seems to have been successful [28]. A stable ring structure of N_5^- ion (D₅h symmetry) with bond lengths of 1.329 A has been determined with coupled cluster (CCSD(T)) theory [26]. Combination of these two ions might prove fruitful towards the synthesis of bicyclic N_{10} .

2.2. N₆₀

The surprising discovery of icosahedral C_{60} [29] as a new allotropic form of carbon not only opened the door to new technological advances in the field, but also directed the search for new exotic molecular structures. From the electronic structure perspective, there is no reason in principle why the three-dimensional π -bonding system in C₆₀ cannot be extended to other elements that exhibit similar bonding such as nitrogen and boron. The electron delocalization in C_{60} , albeit, would give way to (weaker) singly bonded N-N or B–B structures. The metastability of N_{60} would establish this hypothetical molecule as a superhigh-energy containing material.

Table 3 lists the equilibrium bond distances and IR active vibrational frequencies of N₆₀ as determined at the SCF and AM1 levels of theory. The

| various levels of theory from cc-PVDZ basis | | | | | | |
|---|--------------|-------------|-------------|--|--|--|
| Mode | SCF | B3LYP | MP2 | | | |
| B ₁ | 82 (0.0) | 60 (0.0) | 37 (0.0) | | | |
| E | 116 (3.8) | 105 (3.2) | 109 (3.6) | | | |
| A_1 | 487 (0.0) | 427 (0.0) | 427 (0.0) | | | |
| E | 495 (0.9) | 430 (0.8) | 438 (0.9) | | | |
| E | 747 (0.3) | 669 (0.1) | 670 (0.0) | | | |
| A_2 | 828 (0.0) | 739 (0.0) | 733 (0.0) | | | |
| \mathbf{B}_1 | 850 (0.0) | 765 (0.0) | 755 (0.0) | | | |
| A_1 | 1022 (0.0) | 834 (0.0) | 948 (0.0) | | | |
| B_2 | 1028 (120.0) | 856 (55.0) | 969 (44.5) | | | |
| Е | 1196 (2.5) | 1051 (0.1) | 1047 (1.4) | | | |
| B_2 | 1254 (21.1) | 1082 (66.9) | 1070 (29.1) | | | |
| A_1 | 1288 (0.0) | 1088 (0.0) | 1080 (0.0) | | | |
| E | 1337 (23.8) | 1126 (1.9) | 1140 (3.1) | | | |
| B_2 | 1383 (3.8) | 1132 (0.0) | 1157 (0.0) | | | |
| A_1 | 1507 (0.0) | 1269 (0.0) | 1180 (5.9) | | | |
| \mathbf{B}_2 | 1551 (19.0) | 1319 (13.4) | 1228 (4.9) | | | |
| Е | 1714 (3.6) | 1448 (8.0) | 1312 (4.7) | | | |
| A_1 | 1776 (0.0) | 1512 (0.0) | 1519 (0.0) | | | |

Harmonic vibrational frequencies (in cm⁻¹) and IR intensities

(in km/mol) of N₁₀ minimum structure (D₂d symmetry) at

Table 2

| | R ₁ | \mathbf{R}_2 | Frequencies | | |
|-------------|-----------------------|----------------|------------------------------|--|--|
| SCF/cc-PVDZ | 1.430 | 1.437 | 251 (3), 483 (1), 1298 (0.6) | | |
| SCF/6-31G* | 1.430 | 1.436 | | | |
| AM1 | 1.441 | 1.482 | 608 (56), 701 (94), 1153 (4) | | |

Equilibrium bond distances (in Å) and IR active vibrational frequencies (in cm^{-1}) of N₆₀ from AM1 and SCF methods^a

^a \mathbf{R}_1 refers to pentagonal bonds, and \mathbf{R}_2 to bonds shared in six-member rings. IR intensities are in parenthesis.

6-31G^{*} basis set was used in an optimization procedure without symmetry constraints. The two basis sets provide the SCF results that are in very close agreement; with the pentagonal bond (\mathbf{R}_1) being 0.006 Å shorter than the bonds shared in the six-member rings (\mathbf{R}_2). In the case of C_{60} , SCF treatment provided comparable bond lengths $(\mathbf{R_1} = 1.450, \ \mathbf{R_2} = 1.375)$ [30] to those obtained from MP2 calculations ($\mathbf{R_1} = 1.451, \ \mathbf{R_2} = 1.412$) [31] using double-zeta plus polarization (DZP) basis set. Both these results were in close agreement with the experimental determinations $(\mathbf{R_1} = 1.432, \ \mathbf{R_2} = 1.388)$ [32]. We expect the same trend for N₆₀, and will report the results of MP2 and DFT calculations being currently conducted on parallel platforms in future publication. We note that the AM1 bond distances in Table 3 are somewhat larger than the corresponding SCF ones, particularly in the case of the \mathbf{R}_2 bond, which experienced an almost 0.05 Å stretch. More important, however, is that both the SCF and AM1 calculations confirm that the order of the bond lengths in N_{60} has been switched from that of C_{60} : whereas \mathbf{R}_2 is the delocalized (type- π) shorter bond in C_{60} , and it becomes the slightly longer one in N_{60} . This should not be surprising due to the fact that in N_{60} a *p*-type orbital of all nitrogen atoms is the dominant contribution to bonding.

The SCF infrared active low modes listed in Table 3 differ significantly from those obtained from the AM1 theory. All the modes have T_1u symmetry character. The difference could be attributed to the difference in the optimized bond lengths. Due to incorrect description of bond dissociation, SCF calculated frequencies are known to be usually up to 15% as high as the experimental ones. MP2 results with more extensive basis sets are usually more accurate [33] and will be compared to those obtained in this study once available.

The production of new forms of polynitrogen can be achieved through direct synthetic methods or through high pressure-induced transformation. An intriguing possibility of creating N_{60} from six unit blocks of N_{10} is inviting (Fig. 2). A conformation in a cell containing these basic units can be



Fig. 2. (a) Combination of molecular unit blocks of N_{10} to form (b) nitrogen fullerene, N_{60} .

Table 3

envisioned to undergo a transformation to N₆₀ as intermolecular bonds are compressed under high pressure loading, either static (diamond anvil cell) or dynamic (shock compression). Ab initio-based methods in combination with molecular dynamics simulations under these extreme conditions are currently being conducted to investigate the feasibility of these processes and will be reported elsewhere. Such an event might be considered minute and involve considerable energy and entropy constraints. It is, however, the rare events that distinguish chemical research, such as the synthesis of one particular toxin isomer out of 5×10^{21} available configurational choices in the case of palytoxin C₁₂₉H₂₂₃N₃O₅₄ [34,35]. Pending more reliable calculations with correlated methods, the reaction energy of $N_{60} \rightarrow 6 N_{10}$ is estimated at the HF/cc-pVDZ level to be -2430 kcal/ mol. Finally, calculations on chemical reaction mechanism involving the known basic units of nitrogen that could lead to the formation to N_{10} are currently underway.

3. Summary

In this Letter, several computational methods were employed to locate the extrema of bicyclic N_{10} , a high-energy density molecule, and the minimum structure of N60. SCF, DFT, MP2, and QCISD methods located a planar bicyclic structure (D_2h symmetry) as a saddle point for N_{10} , and a corresponding orthogonal structure (D₂d symmetry) as the equilibrium structure. Only 3-6 kcal/ mol separated the two structures, without any further barrier. The structure is remarkably stable with respect to ring connecting bond homolysis. This molecule is postulated herein to serve as a building block for yet a super-high-energy content N_{60} cluster, which is shown to be a local minimum at the AM1 and SCF levels of theory. Harmonic vibrational frequencies of N₁₀ and infrared active modes of N_{60} are reported. This study is the first in a series with the purpose of identifying possible synthetic and physical routes for N₁₀ from the known forms of nitrogen, and for the possibility of producing N₆₀ from a combination of N₁₀ units under extreme conditions.

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