

Study of spin-dependent structure functions of ^3He and ^3H at NNLO approximation and corresponding nuclear corrections

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(Dated: November 1, 2017)

We determine polarized parton distribution functions (PPDFs) and structure functions from recent experimental data of polarized deep inelastic scattering (DIS) on nucleons at next-to-next-to-leading order (NNLO) approximation in perturbative quantum chromodynamic (pQCD). The nucleon polarized structure functions are computed using the Jacobi polynomial approach while target mass corrections (TMCs) are included in our fitting procedure. Having extracted the polarized spin structure functions, we extend our study to describe ^3He and ^3H polarized structure functions, as well as the Bjorken sum rule. We also explore the importance of the nuclear corrections on the polarized nuclear structure functions at small and large values of x . Our results are compared with the recent available and high precision polarized ^3He and ^3H experimental data.

PACS numbers: 13.60.Hb, 24.85.+p, 12.39.-x, 14.65.Bt

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I. INTRODUCTION

Precise understanding of parton distribution functions (PDFs) will be a key ingredient in searches for new physics at the LHC through, for example, top-quark and Higgs-boson coupling measurements. Consequently, reliable extraction of information on the unpolarized parton distribution functions (PDFs) [1–5], polarized (PPDFs) [6–10] and nuclear PDFs [11–14] from global QCD analyses of deep inelastic scattering (DIS) data, as well as many related studies [15–41], provides deep understanding of the hadrons structure in term of their quarks and gluon constituents.

The determination of the longitudinal spin structure of the nucleon caused a huge growth of interest in polarized DIS experiments after the surprising EMC [42, 43] result that the quark spin contribution to the nucleon spin of 1/2 might be significantly small. In subsequent measurements by SMC [44], it was confirmed that the quark spin contributes about one third of the spin of the nucleon. Many experiments have been conducted at SLAC, DESY and CERN to extract the nucleon spin-dependent structure functions $g_1(x, Q^2)$ and $g_2(x, Q^2)$. Various analyses of the world data of A_1 or g_1 based on next-to-leading order (NLO) [7–10] and next-to-next-to-leading order (NNLO) [6, 45] calculations in perturbative quantum chromodynamics (pQCD) have been carried out to extract the polarized parton densities along with the estimation of their uncertainties.

Previously in TKAA16 [6], we carried out the first pQCD analysis using Jacobi polynomial approach at NNLO approximation based on only $g_1^{p,n,d}(x, Q^2)$ experimental data. In our latest study KTA-I17 [45], we have extended our NNLO formalism by including target mass corrections (TMCs) and HT terms and enriched it by more data from $g_2^{p,n,d}(x, Q^2)$ observables. In this anal-

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ysis also the Jacobi polynomials were implemented to determine the polarized parton distribution functions. This method has been applied to various QCD calculations [6, 45–49], containing the case of polarized and unpolarized PDF analyses.

In the absence of polarized charged current neutrino experiments, individual light quark sea densities cannot in principle be determined. The inclusive polarized deep inelastic lepton-hadron reactions can only provide information about the $\Delta u + \Delta \bar{u}$, $\Delta d + \Delta \bar{d}$, and $\Delta s + \Delta \bar{s}$, along with the gluon. For many years our group [6, 45, 50], LSS [51, 52], BB10 [53], and other people made simplifying assumptions about the sea quark densities and consequently were able to present results for the valence u_v and d_v . In the present study, we applied all polarized g_1 data including very recent COMPASS16 g_1^p and g_1^d data [54, 55] to determine the sum of quark and anti-quark polarized PDFs $\Delta q(x) + \Delta \bar{q}(x)$. This method provides no information about the individual polarized quark and anti-quark distributions. We focused only on g_1 experimental data due to their smaller uncertainties compared with the g_2 measurements, indicating the lack of knowledge in the g_2 structure function. Before one can precisely extract PPDF, it is important to take into account the target mass corrections arising from purely kinematic effects.

In addition to the scattering of polarized lepton beams from polarized nucleon, the polarized light nuclear targets provided the opportunity to probe the spin structure of the nucleon. Among these are the SMC experiments at CERN [44] and the E143 [56] and E155 [57] experiments at SLAC that used polarized deuterium. Meanwhile, the HERMES Collaboration at DESY [58] and the E154 experiments at SLAC [59, 60] utilized polarized ^3He . Recently both polarized ^3He and ^3H targets were used at E06-014 experiments at Jefferson Lab (JLAB) in Hall A which are the latest and most up-to-date data for the spin-dependent g_1 and g_2 structure functions of ^3He [61]. Hence, we step further from the “structure of nucleon” to “nuclei” in terms of their parton constituents. In order to study the polarized nuclear structure function, $g_{1,2}^A(x, Q^2)$, one needs to consider nuclear corrections. In this paper, we study the nuclear effects in inclusive scattering of polarized leptons from polarized ^3He and ^3H nuclei in the DIS region. We focus in particular on the kinematics at intermediate and large values of Bjorken x where the major contributions come from the incoherent scattering.

This paper is organized as follows. In Sec. II, we provide an overview of data sets. In Sec. III, we review our theoretical framework and summarize the basic formulas relevant for the analysis. Section IV contains the formalism used for computing the χ^2 minimization and PPDF uncertainties. We introduce the nuclear structure functions and corresponding nuclear corrections in Sec. V. In Sec. VI we discuss how well the predictions for the NNLO polarized and the inclusion of TMCs effects into NNLO polarized structure function analysis improves the preci-

sion of the extracted PPDFs as well as nuclear structure functions. Finally, in Sec. VII we summarize our findings.

II. OVERVIEW OF DATA SETS

The combined set of data was included in our NNLO QCD fit to the available g_1^p , g_1^n and g_1^d world data. For the proton data we use E143 [56], HERMES98 [62], HERMES06 [63], SMC [44], EMC [42, 43], E155 [64], COMPASS10 [65] and COMPASS16 [55]; for the neutron data we use E142 [66], E154 [67], HERMES98 [62], HERMES06 [58], JLAB03 [68], JLAB04 [69] and JLAB05 [70]; and for the deuteron data we use E143 [56], E155 [57], SMC [44], HERMES06 [63], COMPASS05 [71], COMPASS06 [72], and COMPASS16 [54]. These data sets are summarized in Table I.

The x and Q^2 nominal coverage of the data considered in our QCD fit is illustrated in Fig. 1. This plot nicely represents the kinematic coverage of x , and Q^2 of the proton, neutron and deuteron polarized structure functions. Recent COMPASS results for protons and deuterons [54, 55] at low x ($x < 0.03$) increase considerably the accuracy compared to the only available result in this region, SMC [44]. Despite recent outstanding experimental efforts, the kinematical coverage of present polarized DIS observables is still rather limited. Our data (511 experimental data points) cover the kinematic range $0.0035 < x < 0.75$, $1 \text{ (GeV}^2) < Q^2 < 100 \text{ (GeV}^2)$ and $W > 4 \text{ GeV}$.

This coverage leads to wider uncertainty for extracted polarized PDFs at small x . Consequently, the polarized gluon distribution $\Delta G(x, Q^2)$ and strange distribution $(\Delta s + \Delta \bar{s})(x, Q^2)$ are still weakly constrained, especially for the case of the polarized gluon distribution. Any conclusion on gluon treatment at $x < 0.01$ relies on the behavior of low x polarized DIS data, which is not accurately known up to now. The statistical and systematic uncertainties are both taken into account. The systematic uncertainties are added quadratically.

III. THEORETICAL ANALYSIS

A. Leading-twist polarized DIS structure functions

The calculations applied in this analysis are all performed at NNLO approximation. Correspondingly, the polarized PDFs evolve from the input scale Q_0^2 using NNLO splitting functions [73] and the NNLO hard scattering cross section expressions are considered. To have a full leading-twist (LT) analysis we shall write

$$g_{1,2}(x, Q^2)_{\text{LT}} = g_{1,2}(x, Q^2)_{\text{pQCD}} + h_{1,2}^{\text{TMCs}}(x, Q^2)/Q^2, \quad (1)$$

where $h_{1,2}^{\text{TMCs}}(x, Q^2)$ is explained in Sec. IIIB. The NNLO spin-dependent proton structure functions,

Table I: Summary of published polarized DIS experimental data points with measured x and Q^2 ranges and the number of data points, the χ^2 for each given data set, and the fitted normalization shifts \mathcal{N}_i .

Experiment	Ref.	$[x_{\min}, x_{\max}]$	Q^2 range (GeV ²)	Number of data points	χ^2	\mathcal{N}_n
E143(p)	[56]	[0.031–0.749]	1.27–9.52	28	21.424	1.00034601
HERMES(p)	[62]	[0.028–0.66]	1.01–7.36	39	75.369	1.00186515
SMC(p)	[44]	[0.005–0.480]	1.30–58.0	12	10.803	0.99991146
EMC(p)	[42]	[0.015–0.466]	3.50–29.5	10	3.328	1.00220752
E155	[64]	[0.015–0.750]	1.22–34.72	24	35.170	1.024762208
HERMES06(p)	[63]	[0.026–0.731]	1.12–14.29	51	22.672	1.00018245
COMPASS10(p)	[65]	[0.005–0.568]	1.10–62.10	15	26.670	0.99301000
COMPASS16(p)	[55]	[0.0035–0.575]	1.03–96.1	54	53.912	1.00019414
g_1^p				233		
E143(d)	[56]	[0.031–0.749]	1.27–9.52	28	38.159	0.99916419
E155(d)	[57]	[0.015–0.750]	1.22–34.79	24	18.871	0.99991576
SMC(d)	[44]	[0.005–0.479]	1.30–54.80	12	18.375	0.99998812
HERMES06(d)	[63]	[0.026–0.731]	1.12–14.29	51	47.045	1.00001347
COMPASS05(d)	[71]	[0.0051–0.4740]	1.18–47.5	11	8.490	0.99692499
COMPASS06(d)	[72]	[0.0046–0.566]	1.10–55.3	15	12.874	0.99991619
COMPASS16(d)	[54]	[0.0045–0.569]	1.03–74.1	43	37.297	1.00089129
g_1^d				184		
E142(n)	[66]	[0.035–0.466]	1.10–5.50	8	7.466	0.99899932
HERMES(n)	[62]	[0.033–0.464]	1.22–5.25	9	2.697	0.99995848
E154(n)	[67]	[0.017–0.564]	1.20–15.00	17	9.216	0.99961961
HERMES06(n)	[58]	[0.026–0.731]	1.12–14.29	51	17.974	1.00001347
Jlab03(n)	[68]	[0.14–0.22]	1.09–1.46	4	0.0469	0.99981391
Jlab04(n)	[69]	[0.33–0.60]	2.71–4.8	3	3.651	0.90000096
Jlab05(n)	[70]	[0.19–0.20]	1.13–1.34	2	1.674	1.02232189
g_1^n				94		
Total				511	473.195	

$g_1^p(x, Q^2)_{\text{pQCD}}$, can be written as a linear combination of polarized parton distribution functions Δq , $\Delta \bar{q}$ and Δg as,

$$\begin{aligned}
g_1^p(x, Q^2)_{\text{pQCD}} &= \frac{1}{2} \sum_q e_q^2 \Delta q_{\text{NS}}(x, Q^2) \otimes \\
&\left(1 + \frac{\alpha_s(Q^2)}{2\pi} \Delta C_q^{(1)} + \left(\frac{\alpha_s(Q^2)}{2\pi} \right)^2 \Delta C_{\text{NS}}^{(2)} \right) \\
&+ e_q^2 (\Delta q + \Delta \bar{q})(x, Q^2) \otimes \\
&\left(1 + \frac{\alpha_s(Q^2)}{2\pi} \Delta C_q^{(1)} + \left(\frac{\alpha_s(Q^2)}{2\pi} \right)^2 \Delta C_s^{(2)} \right) \\
&+ \frac{2}{9} \left(\frac{\alpha_s(Q^2)}{2\pi} \Delta C_g^{(1)} + \left(\frac{\alpha_s(Q^2)}{2\pi} \right)^2 \Delta C_g^{(2)} \right) \otimes \Delta g(x, Q^2),
\end{aligned} \tag{2}$$

where the ΔC_q and ΔC_g are the spin-dependent quark and gluon hard scattering coefficient functions, calculable at NNLO approximation [74, 75]. The symbol \otimes represents the typical convolution in Bjorken x space. Considering the polarized proton structure function, we can apply isospin symmetry to achieve the neutron one. The deuteron structure function is related to that of the

proton and neutron via

$$\begin{aligned}
g_1^d(x, Q^2)_{\text{LT}} &= \frac{1}{2} \{ g_1^p(x, Q^2)_{\text{LT}} + g_1^n(x, Q^2)_{\text{LT}} \} \\
&\times (1 - 1.5w_D),
\end{aligned} \tag{3}$$

where $w_D = 0.05 \pm 0.01$ is the probability to find the deuteron in a D -state [76–78].

One can use the Wandzura and Wilczek (WW) [61, 79] approximation for the leading-twist g_2 polarized structure function

$$\begin{aligned}
g_2(x, Q^2)_{\text{pQCD}} &= g_2^{\text{WW}}(x, Q^2) = \\
&-g_1(x, Q^2)_{\text{pQCD}} + \int_x^1 \frac{dy}{y} g_1(y, Q^2)_{\text{pQCD}}.
\end{aligned} \tag{4}$$

Target mass corrections do not affect the WW relation if all powers in (M^2/Q^2) are included [79].

B. Target mass corrections

As illustrated in Fig. 1, in polarized DIS most of the small x experimental data points are at low Q^2 , which is one of the features of polarized DIS. In the unpolarized case we can cut the preasymptotic region data, while it is impossible to perform such a procedure for present data on spin-dependent structure functions without losing too

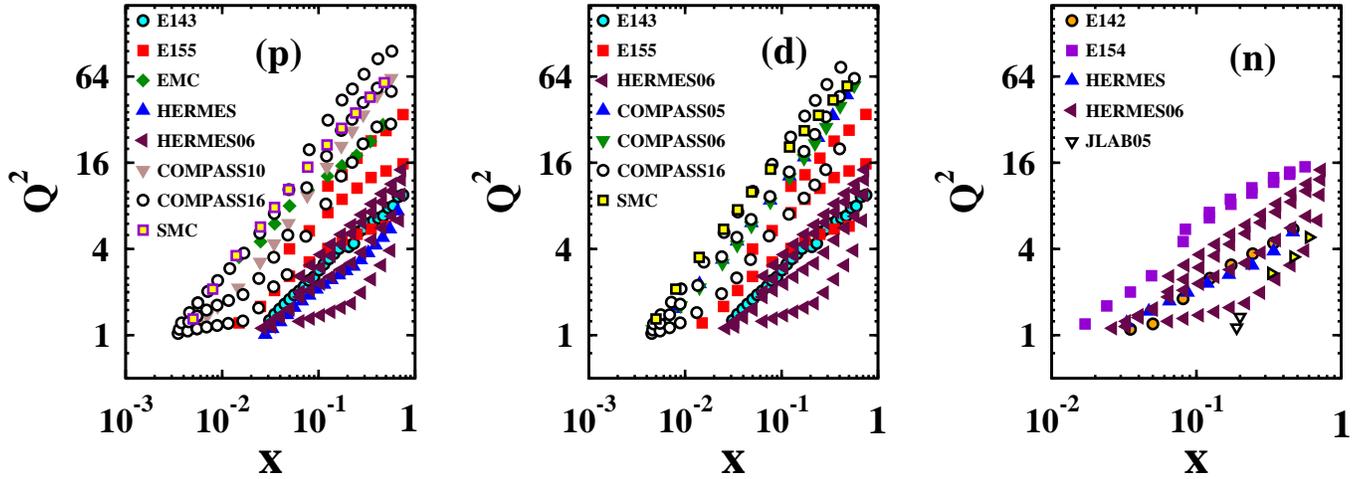


Figure 1: Nominal coverage of the data sets used in the KTA-III17 analysis for proton, neutron, and deuteron observables. DIS data are presented on a logarithmic x and Q^2 scales.

much information. To perform a reliable QCD fit including data at lower Q^2 values, target mass corrections cannot be ignored. The standard approach to calculate TMCs in the case of unpolarized DIS is the one based on the operator product expansion (OPE) in QCD, first formulated by Georgi and Politzer [80]. This closed-form expression is generalized in Ref. [81].

We have performed an analysis to higher terms in the TMC expansion based on the method of Ref. [8] and found that these terms do not change the agreement with g_1 data and the extracted polarized parton densities are insensitive to such a choice. We claim that leading terms in the TMC expansion are reliable in the kinematical range of presently available g_1 data. In our previous study, KTA-I17 [45], we presented the significant effect of considering TMCs and HT contributions while g_2 structure function data are included.

To consider the full LT approach Eq. (1) in our analysis, we applied the method suggested in Refs. [82–84]. This method effectively depends on the LT term (for more details, see our paper [45]). For simplicity of notation, from now to the end of the paper, we will drop the subscript “LT” denoting the leading twist.

C. PDF parametrizations and conventions

The method applied to reconstruct the x -dependent quantities from their Mellin moments is the Jacobi polynomial method (the same as our previous QCD analyses [6, 45]). The main difference, as indicated in Sec. I, is that we consider a new input parametrization at the initial scale $Q_0^2 = 1 \text{ GeV}^2$ for the sum of quark and

antiquark polarized PDF instead of the valence and sea distributions, which are more general. We consider the general form of

$$\begin{aligned}
 x(\Delta u + \Delta \bar{u})(x, Q_0^2) &= \eta_{u+} A_{u+} x^{\alpha_{u+}} (1-x)^{\beta_{u+}} \\
 &\quad \times (1 + \epsilon_{u+} \sqrt{x} + \gamma_{u+} x), \\
 x(\Delta d + \Delta \bar{d})(x, Q_0^2) &= \eta_{d+} A_{d+} x^{\alpha_{d+}} (1-x)^{\beta_{d+}} (1 + \gamma_{d+} x), \\
 x(\Delta s + \Delta \bar{s})(x, Q_0^2) &= \eta_{s+} A_{s+} x^{\alpha_{s+}} (1-x)^{\beta_{s+}}, \\
 x\Delta G(x, Q_0^2) &= \eta_G A_G x^{\alpha_G} (1-x)^{\beta_G} \\
 &\quad \times (1 + \epsilon_G \sqrt{x} + \gamma_G x). \tag{5}
 \end{aligned}$$

Here, the notation $q_+ = q + \bar{q}$ is applied for light quarks. The normalization factors, A_i , are fixed such that η_i represent the first moments of the polarized distributions. As usual, the set of free parameters in Eq. 5 is constrained by the well-known sum rules

$$\begin{aligned}
 a_3 &= g_A = F + D = 1.269 \pm 0.003, \\
 a_8 &= 3F - D = 0.585 \pm 0.025. \tag{6}
 \end{aligned}$$

Here, a_3 and a_8 are nonsinglet combinations of the first moments of the polarized parton distributions corresponding to the axial charges for octet baryons [85, 86]. These parameters, F and D, are measured in hyperon and neutron β decay and finally lead to the constraints

$$\begin{aligned}
 a_3 &= (\Delta u + \Delta \bar{u})(Q^2) - (\Delta d + \Delta \bar{d})(Q^2), \\
 a_8 &= (\Delta u + \Delta \bar{u})(Q^2) + (\Delta d + \Delta \bar{d})(Q^2) \\
 &\quad - 2(\Delta s + \Delta \bar{s})(Q^2). \tag{7}
 \end{aligned}$$

So considering Eqs. (6) and (7), the parameters η_{u_+} and η_{d_+} can be extracted versus $\Delta u + \Delta \bar{u}$, $\Delta d + \Delta \bar{d}$, and $\Delta s + \Delta \bar{s}$. Here, we do not make any simplifying assumptions on the equality of the light sea quark distributions.

In our previous papers [6, 45], we have considered the Jacobi polynomial method to yield the structure functions from their Mellin moments in N space.

In the polynomial approach, one can easily expand the polarized structure functions in terms of the Jacobi polynomials $\Theta_n^{\alpha,\beta}(x)$ as follows,

$$x g_1(x, Q^2) = x^\beta (1-x)^\alpha \sum_{n=0}^{N_{\max}} a_n(Q^2) \Theta_n^{\alpha,\beta}(x). \quad (8)$$

Here, $\Theta_n^{\alpha,\beta}(x) = \sum_{j=0}^n c_j^{(n)}(\alpha, \beta) x^j$ and N_{\max} is the maximum order of the expansion.

The Jacobi polynomials $\Theta_n^{\alpha,\beta}(x)$ are a class of classical orthogonal polynomials. They are orthogonal with respect to the weight $x^\beta (1-x)^\alpha$ on the interval $[0, 1]$. In the polynomial approach, the Q^2 dependence of the $x g_1(x, Q^2)$ are codified in the Jacobi polynomial moments, $a_n(Q^2)$. Using the orthogonality relation, one can obtain this moment as

$$\begin{aligned} a_n(Q^2) &= \int_0^1 dx x g_1(x, Q^2) \Theta_n^{\alpha,\beta}(x) \\ &= \sum_{j=0}^n c_j^{(n)}(\alpha, \beta) \mathcal{M}[x g_1, j+2](Q^2). \end{aligned} \quad (9)$$

The Mellin transform $\mathcal{M}[x g_1^{\tau^2}, N]$ is defined via,

$$\mathcal{M}[x g_1, N](Q^2) \equiv \int_0^1 dx x^{N-2} x g_1(x, Q^2). \quad (10)$$

Applying the Jacobi polynomial expansion method, the polarized structure function $x g_1(x, Q^2)$ can be constructed as

$$\begin{aligned} x g_1(x, Q^2) &= x^\beta (1-x)^\alpha \sum_{n=0}^{N_{\max}} \Theta_n^{\alpha,\beta}(x) \\ &\times \sum_{j=0}^n c_j^{(n)}(\alpha, \beta) \mathcal{M}[x g_1, j+2](Q^2). \end{aligned} \quad (11)$$

We also scrutinized the sensitivity of the Jacobi polynomials to its free parameters. These results are discussed in detail in Ref. [45]. In our current analysis by setting the $\{N_{\max}=7, \alpha=3, \beta=0.5\}$, the optimal convergence of this expansion throughout the whole kinematic region covered by the polarized DIS data is achievable.

IV. χ^2 MINIMIZATION AND UNCERTAINTIES OF PHYSICAL PREDICTIONS

The goodness of fit is traditionally determined by the effective global χ^2 minimization algorithm that measures

the quality of fit between theory and experiment. χ_{global}^2 is defined by

$$\chi_{\text{global}}^2 = \sum_{i=1}^{N_{\text{data}}} w_n \times \left[\left(\frac{1 - \mathcal{N}_n}{\Delta \mathcal{N}_n} \right)^2 + \sum_{i=1}^{N_{\text{data}}} \left(\frac{\mathcal{N}_n g_{1,i}^{\text{Exp}} - g_{1,i}^{\text{Theory}}}{\mathcal{N}_n \Delta g_{1,i}^{\text{Exp}}} \right)^2 \right], \quad (12)$$

where n labels the number of different experiments and w_n is a weight factor for the n th experiment. g^{Exp} , Δg^{Exp} , and g^{Theory} indicate the data value, the uncertainty, and the theory value for the data point i of data set n , respectively. The $\Delta \mathcal{N}_n$ are the experimental normalization uncertainty quoted by the experiments. The relative normalization factors, \mathcal{N}_n , appear as free parameters in the fit. These 22 normalization shifts are determined at the prefitting procedure along with the PDF free parameters and strong coupling constant using the CERN program library MINUIT [87]. Afterwards, they are fixed at their best fitted values to further reduce the free parameters. Finally, we minimize the above χ^2 value with 16 free parameters including the strong coupling constant.

To visually evaluate the fit quality, in Fig. 2 we plot the χ^2/dof for individual experiments per nucleon target. This allows us to check that the majority of experiments have a $(\chi^2/\text{dof}) \simeq 1$. It means that most of the experiments satisfy this goodness-of-fit scale parameter. The largest contribution to χ^2 arises from the HERMES and COMPASS10 data for protons, and the SMC data for deuteron structure functions. The smallest contribution comes from the JLAB03 data for $x g_1^p$. From Fig. 2, one can conclude that the HERMES98 and COMPASS10 data for $x g_1^p$ are difficult to be describe based on our fitting scenario. The motivation of considering them in our analysis originates mainly from our goal to have the most complete and up-to-date sets of data for the polarized structure functions.

Different methods to estimate the uncertainties of PDFs obtained from global χ^2 optimization together with technical details were described in Refs. [4, 11, 88–91]. The most common and effective approach is the ‘‘Hessian method’’. In this section the outline of this method is explained because it is used in our analysis. Up to the leading quadratic terms, the increase χ^2 can be written in terms of the Hessian matrix

$$H_{ij} \equiv \frac{1}{2} \frac{\partial^2 \chi^2}{\partial a_i \partial a_j} \Bigg|_{\min} \quad (13)$$

as

$$\Delta \chi_{\text{global}}^2 = \chi_{\text{global}}^2 - \chi_{\min}^2 = \sum_{ij}^{N_{\text{par}}} H_{ij} (a_i - a_i^{(0)}) (a_j - a_j^{(0)}), \quad (14)$$

where a_i ($i=1, 2, \dots, N$) refers to the free parameters for each PDF presented in Eq. (5) and N is the number of

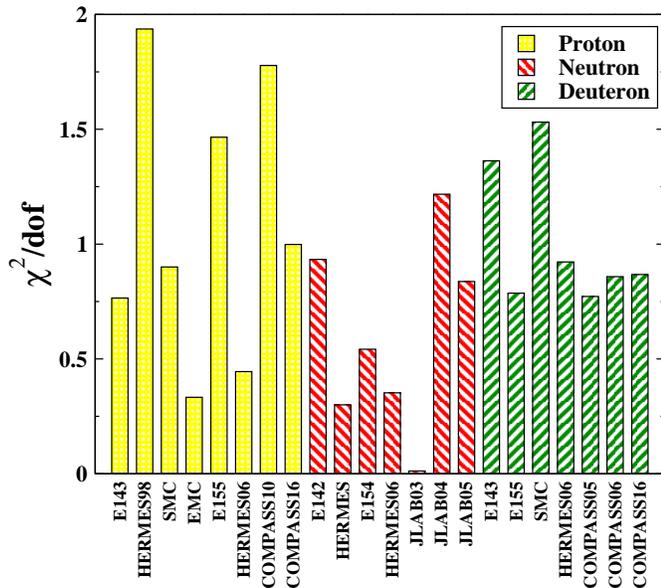


Figure 2: Value of χ^2/dof for individual experiments per nucleon target used in the KTA-III17

parameters. $\Delta\chi_{\text{global}}^2$ illustrates the allowed variation in χ^2 . The standard formula for linear error propagation is

$$(\Delta q)^2 = \Delta\chi_{\text{global}}^2 \sum_{i,j} \frac{\partial \Delta q}{\partial a_i} (H_{ij})^{-1} \frac{\partial \Delta q}{\partial a_j}. \quad (15)$$

The parameter value of polarized PDF, i.e. a_1^0, \dots, a_n^0 , extracted from the NNLO QCD fit, will be presented in Sec. VI. Equation (15) is not convenient to use since the derivative of Δq with respect to each parameter a_i is required. An improved iterative method has been devised in [4, 88, 89] in which the Hessian matrix is diagonalized. We adopt this improved Hessian method in our analysis and work in terms of rescaled eigenvectors and eigenvalues. The PDF uncertainty determination becomes much easier in terms of the appropriately normalized eigenvectors, z_i

$$\Delta\chi_{\text{global}}^2 = \sum_i^{N_{\text{par}}} z_i^2. \quad (16)$$

The uncertainty of an individual PDF at particular values of x and Q^2 can be estimated using

$$\Delta q = \Delta\chi_{\text{global}}^2 \frac{1}{2} \sqrt{\sum_{k=1}^n [\Delta q(S_k^+) - \Delta q(S_k^-)]^2}. \quad (17)$$

Here, S_k^+ and S_k^- are polarized PDF sets displaced along eigenvector directions by the extracted $\Delta\chi_{\text{global}}^2$ value.

With the standard “parameter-fitting” criterion, we would expect the uncertainties to be given by the choice

of tolerance parameter $T^2 = \Delta\chi_{\text{global}}^2 = 1$ for the 68% (one-sigma-error) confidence level (C.L.) limit [4, 89]. For the general case with N degrees of freedom, the $\Delta\chi_{\text{global}}^2$ value needs to be calculated to determine the size of the uncertainties. Assuming that $\Delta\chi_{\text{global}}^2$ follows the χ^2 distribution with N degrees of freedom, we have the C.L. P as [4, 89, 92]

$$P = \int_0^{\Delta\chi_{\text{global}}^2} \frac{1}{2\Gamma(N/2)(\frac{S}{2})^{(N/2)-1}} e^{-\frac{S}{2}} dS. \quad (18)$$

In the case of the one-free-parameter fit, one obviously has $\Delta\chi_{\text{global}}^2 = 1$. Since the polarized parton distributions in common QCD fits are considered with several free parameters, $N > 1$, the value of $\Delta\chi_{\text{global}}^2$ should be reevaluated from Eq. (18). Here we calculate the uncertainties of polarized PDFs with $\Delta\chi_{\text{global}}^2 = 1$. For other values of $\Delta\chi_{\text{global}}^2$, one can simply scale our error bands by $(\Delta\chi_{\text{global}}^2)^{1/2}$.

V. POLARIZED STRUCTURE FUNCTIONS FOR NUCLEI

The experimental data on deep inelastic lepton-nucleus ($\ell^\pm + A$) scattering can reveal more information on the behavior of spin structure functions of nucleon and nucleon correlations in nuclei at small and large values of Bjorken x . It will provides more reliable pictures of the nuclear phenomena. Moreover, the absence of free neutron targets means that polarized light nuclei such as deuterium and ^3He must be used as effective polarized neutron targets. The scattering of polarized lepton beams ℓ^\pm from polarized nuclear targets paves the way to detail study of spin structure of the nucleon encoded in the spin structure functions of $g_1(x, Q^2)$ and $g_2(x, Q^2)$. In order to extract these spin structure functions from the spin-dependent DIS data on nuclear targets, one needs to account for the nontrivial nuclear corrections. The nuclear effects that play an important role in the polarized as well as unpolarized DIS on nuclei can be divided into “coherent” and “incoherent” contributions [93, 94]. Incoherent nuclear effects, which are present at all values of Bjorken x , arise from the scattering of the incoming leptons on each individual nucleons. Some examples of incoherent nuclear effects are the well-known Fermi motions, spin depolarizations, binding and the presence of a non-nucleonic degree of freedom.

The coherent nuclear effects are typically important at small values of the momentum fraction x . They result from the interaction of the incoming leptons with two or more nucleons inside the targets. Nuclear shadowing which is very important for the region $10^{-4} \leq x \leq 0.03 - 0.07$, and nuclear anti-shadowing at $0.03 - 0.07 \leq x \leq 0.2$ are examples of the coherent nuclear effects. Nuclear shadowing corrections in polarized ^3He have been argued in detail in Refs. [95, 96]. These effects arise from multiple scattering of the leptons from two or more nucleons

in the ${}^3\text{He}$ nucleus. In addition to the point mentioned above, contributions to the polarized ${}^3\text{He}$ structure function from non-nucleonic degrees of freedom in the nucleus have been discussed in detail in Refs. [93, 94].

For the case of the polarized DIS in which we present in the following analysis, the major contributions come from the incoherent scattering on the nucleons of the targets. The aim of this work is to study these nuclear effects for the case of the polarized DIS at NNLO approximation and assess their importance on the spin structure functions of $g_{1,2}^{3\text{He}}$ and $g_1^{3\text{H}}$.

In following, we describe the nuclear effects in inclusive scattering of polarized leptons from polarized helium ${}^3\text{He}$ and tritium ${}^3\text{H}$ nuclei focusing in particular on kinematics at whole values of momentum fraction x .

A. Weak binding approximation (WBA)

To begin our discussion on the nuclear corrections, we should note that in the standard nuclear structure function analyses, the contributions from the spin depolarization, Fermi motions and binding are described within the framework of the so-called ‘‘convolution approach’’. The polarized nuclear structure function of ${}^3\text{He}$ and ${}^3\text{H}$ can be written as a convolution of the nucleons light-cone momentum distributions with the off-shell polarized nucleon structure functions of the protons $g_1^p(x, Q^2)$ and neutrons $g_1^n(x, Q^2)$. In what follows, we apply the method presented in Refs. [93, 97] and adopt the approach in which the nucleons light-cone momentum distributions can be extracted from the nuclear spectral functions [98, 99].

The well-known coherent effects associated with the multiple scattering from two or more nucleons inside the nucleus give rise to corrections at small values of Bjorken x . In order to study the nuclear effects in the incoherent scattering, we restrict ourselves to the intermediate and large regions of momentum fraction x in which the incoherent scattering from a single nucleon is assumed to dominate [18, 97, 100, 101]. In this framework, in which the nucleus is treated as a non-relativistic system of weakly bound nucleons, the spin-dependent structure functions of ${}^3\text{He}$ can be obtained as [102–107]

$$g_1^{3\text{He}}(x, Q^2) = \int \frac{dy}{y} \left[2\Delta f^p(y, \gamma) g_1^p\left(\frac{x}{y}, Q^2\right) + \Delta f^n(y, \gamma) g_1^n\left(\frac{x}{y}, Q^2\right) \right], \quad (19)$$

where $y = \frac{p \cdot q}{M_N}$ is the nuclear light-cone momentum fraction carried by the interacting nucleons inside the nucleus. The functions $f^{N(=p,n)}(y, \gamma)$ are the nucleon light-cone momentum distributions in the ${}^3\text{He}$ nucleus computed in terms of the nuclear spectral functions [93, 97]

$$\Delta f^N(y, \gamma) = \int \frac{d^4p}{(2\pi)^4} D^N(\varepsilon, p, \gamma) \delta(y - 1 - \frac{\varepsilon + \gamma p_z}{M}), \quad (20)$$

with $N = \text{proton or neutron}$, and D^N is the energy-momentum distribution functions. In the Bjorken limit in which ($\gamma \rightarrow 1$), the nucleon light-cone momentum distributions $f^N(y, \gamma)$ depend only to the y variables, which covers the range between x and $\frac{M_{3\text{He}}}{M} \approx 3$. Finally, the polarized structure functions of $g_1^{3\text{He}}(x, Q^2)$ and $g_1^{3\text{H}}(x, Q^2)$ can be written as

$$g_1^{3\text{He}}(x, Q^2) = \int_x^3 \frac{dy}{y} \Delta f_{3\text{He}}^n(y) g_1^n(x/y, Q^2) + 2 \int_x^3 \frac{dy}{y} \Delta f_{3\text{He}}^p(y) g_1^p(x/y, Q^2), \quad (21)$$

$$g_1^{3\text{H}}(x, Q^2) = 2 \int_x^3 \frac{dy}{y} \Delta f_{3\text{H}}^n(y) g_1^n(x/y, Q^2) + \int_x^3 \frac{dy}{y} \Delta f_{3\text{H}}^p(y) g_1^p(x/y, Q^2). \quad (22)$$

The Fermi motions and their bindings are parametrized through the distributions of $\Delta f_{3\text{He}}^{N(=p,n)}$ and $\Delta f_{3\text{H}}^{N(=p,n)}$ which can be calculated using the ground-state wave functions of ${}^3\text{He}$ and ${}^3\text{H}$ nuclei, respectively. It is worth mentioning here that, because of isospin symmetry, the light-cone momentum distribution $\Delta f_{3\text{He}}^p(y)$ is equal to the $\Delta f_{3\text{H}}^n(y)$.

In order to obtain the polarized light-cone momentum distributions, we used the numerical results presented in Ref. [98]. The polarized light-cone distribution functions for the neutron $\Delta f_{3\text{He}}^n(y)$ as well as for the proton $\Delta f_{3\text{He}}^p(y)$ are shown as a function of y in Figs. 3 and 4, respectively. It has been shown in various studies that the polarized light-cone distribution functions $\Delta f^N(y, \gamma = 1)$ are sharply peaked at $y \approx 1$ [98, 99, 105].

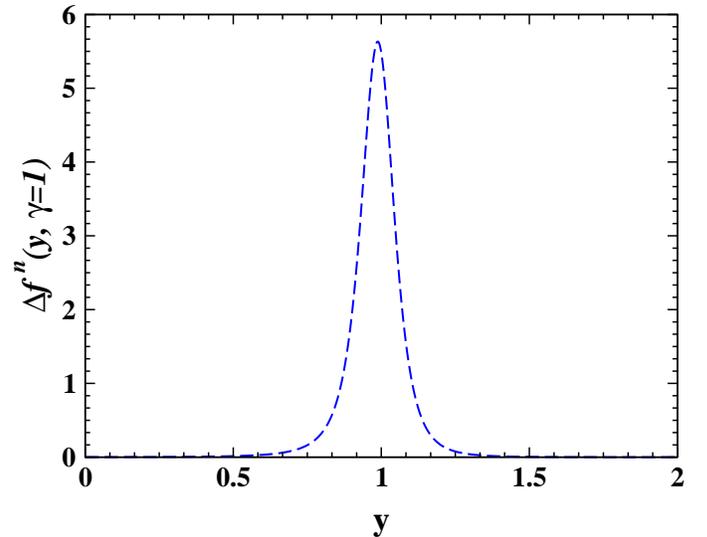


Figure 3: The polarized neutron light-cone momentum distribution functions at the Bjorken limit ($\gamma \rightarrow 1$) in ${}^3\text{He}$, based on the results presented in Ref. [98].

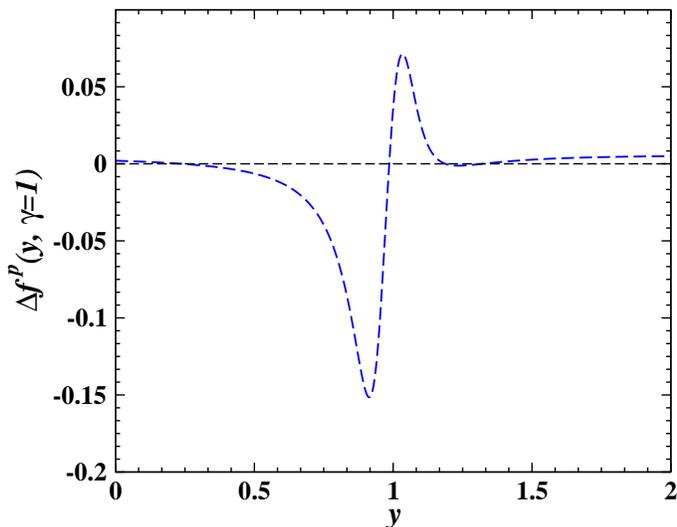


Figure 4: The polarized proton light-cone momentum distribution functions at the Bjorken limit ($\gamma \rightarrow 1$) in ${}^3\text{He}$, based on the results presented in Ref. [98].

B. Effective polarizations of the nucleons

In the limit of zero nuclear binding, and in the Bjorken limits in which ($\gamma \rightarrow 1$), the polarized proton light-cone momentum distribution functions $\Delta f^N(y)$ become infinitesimally narrower and are sharply peaked around the points $y \approx 1$ $\{f^N \sim \delta(1-y)\}$, due to the small average separation energy per nucleons inside the nuclei.

Thus, in this approximation, one can express the polarized nuclear structure functions $g_1^A(x, Q^2)$ as linear combinations of the polarized proton $g_1^p(x, Q^2)$ and neutron $g_1^n(x, Q^2)$ structure functions weighted by effective polarizations. Consequently, Eq. (19) is often approximated by

$$g_1^{3\text{He}}(x, Q^2) = 2P^p g_1^p(x, Q^2) + P^n g_1^n(x, Q^2). \quad (23)$$

Here P^p and P^n are the proton and neutron effective polarizations inside the polarized ${}^3\text{He}$ nucleus. The proton effective polarizations P^p are assumed to be the average polarizations of the two protons inside the ${}^3\text{He}$ nuclei. The effective polarizations can be described in terms of integrals of the diagonal light-cone momentum distribution functions at the Bjorken limit in which ($\gamma \rightarrow 1$),

$$\begin{aligned} P^p &= \int_0^3 dy \Delta f_{3\text{He}}^p(y, \gamma = 1), \\ P^n &= \int_0^3 dy \Delta f_{3\text{He}}^n(y, \gamma = 1). \end{aligned} \quad (24)$$

The effective polarizations presented in our analysis can be computed numerically from models of the ${}^3\text{He}$ nucleus wave functions. The calculations of Refs. [98, 108] have shown that $P^p = -0.028 \pm 0.004$ and $P^n = 0.86 \pm 0.02$.

C. Non-nucleonic contributions

As we have discussed earlier, free nucleons behave differently from nucleons in which bound in the nuclei. This is due to the some important nuclear corrections such as the nuclear bindings, Fermi motions, nuclear shadowing, and anti-shadowing as well as the non-nucleonic degrees of freedom. Consequently, the descriptions of the nucleon as a mere collections of the protons and neutrons (descriptions of nuclear properties in terms of the nucleon degrees of freedom alone) may not be complete, and, hence the nuclear corrections due to the non-nucleonic degrees of freedom have to be considered. For the spin-dependent observables $g_1^A(x, Q^2)$, small admixtures of the $\Delta(1232)$ isobar in three-body wave functions were found to be necessary for a better descriptions of the polarized nuclear structure functions [109].

In order to consider the nuclear effects in which originate from the non-nucleonic degrees of freedom, we utilized the works presented in Ref. [93] which provides a description of the $g_1(x, Q^2)$ spin structure functions of helium ${}^3\text{He}$ as well as tritium ${}^3\text{H}$ nuclei over the range of $10^{-4} \leq x \leq 0.8$. It models the ${}^3\text{He}$ wave function including the S , S' , and D states, and the non-nucleonic degrees of freedoms from the effect of $\Delta(1232)$ isobar. The corresponding results for the ${}^3\text{He}$ and ${}^3\text{H}$ can be written,

$$\begin{aligned} g_1^{3\text{He}} &= \int_x^3 \frac{dy}{y} \Delta f_{3\text{He}}^n(y) g_1^n(x/y) + 2 \int_x^3 \frac{dy}{y} \Delta f_{3\text{He}}^p(y) g_1^p(x/y) \\ &- 0.014 (g_1^p(x) - 4g_1^n(x)), \end{aligned} \quad (25)$$

and

$$\begin{aligned} g_1^{3\text{H}} &= 2 \int_x^3 \frac{dy}{y} \Delta f_{3\text{H}}^n(y) g_1^n(x/y) + \int_x^3 \frac{dy}{y} \Delta f_{3\text{H}}^p(y) g_1^p(x/y) \\ &+ 0.014 (g_1^p(x) - 4g_1^n(x)). \end{aligned} \quad (26)$$

The last terms in Eqs. (25) and (26) arise from the $\Delta(1232)$ component in the ${}^3\text{He}$ and ${}^3\text{H}$ wave functions. They will have sizable contributions at large values of Bjorken x , $0.2 \leq x \leq 0.8$.

The same formula can be applied for the g_2^A nuclear structure functions as

$$\begin{aligned} g_2^{3\text{He}} &= \int_x^3 \frac{dy}{y} \Delta f_{3\text{He}}^n(y) g_2^n(x/y) + 2 \int_x^3 \frac{dy}{y} \Delta f_{3\text{He}}^p(y) g_2^p(x/y) \\ &- 0.014 (g_2^p(x) - 4g_2^n(x)), \end{aligned} \quad (27)$$

and

$$\begin{aligned} g_2^{3\text{H}} &= 2 \int_x^3 \frac{dy}{y} \Delta f_{3\text{H}}^n(y) g_2^n(x/y) + \int_x^3 \frac{dy}{y} \Delta f_{3\text{H}}^p(y) g_2^p(x/y) \\ &+ 0.014 (g_2^p(x) - 4g_2^n(x)). \end{aligned} \quad (28)$$

D. Shadowing and anti-shadowing corrections

At high energy or small values of momentum fraction x , the virtual photons can interact “coherently” with several nucleons in the nuclear targets. These behaviors are manifested in the nuclear shadowing and anti-shadowing effects, and breaks down the well-known convolution approximations [93]. Considering these nuclear corrections, one can write the polarized nuclear structure functions as

$$g_1^{3\text{He}} = \int_x^3 \frac{dy}{y} \Delta f_{3\text{He}}^n(y) g_1^n(x/y) + 2 \int_x^3 \frac{dy}{y} \Delta f_{3\text{He}}^p(y) g_1^p(x/y) - 0.014 \left(g_1^p(x) - 4g_1^n(x) \right) + a(x)g_1^n(x) + b(x)g_1^p(x), \quad (29)$$

and

$$g_1^{3\text{H}} = 2 \int_x^3 \frac{dy}{y} \Delta f_{3\text{H}}^n(y) g_1^n(x/y) + \int_x^3 \frac{dy}{y} \Delta f_{3\text{H}}^p(y) g_1^p(x/y) + 0.014 \left(g_1^p(x) - 4g_1^n(x) \right) + a(x)g_1^n(x) + b(x)g_1^p(x), \quad (30)$$

The functions $a(x)$ and $b(x)$ describe both the corrections of nuclear shadowing and anti-shadowing, and are functions of x and Q^2 .

Since in the most nuclear polarized DIS experiments the x coverage does not drop below the limits $x \sim 0.2$, the shadowing and anti-shadowing corrections can be ignored [61]. However the calculations of Ref. [93] have shown that these two effects are quite significant and can affect the extraction of the nucleon spin functions at small values of Bjorken x . As we mentioned earlier, current experimental data do not reach to very small values of x . Consequently, the corrections from shadowing ($10^{-4} \leq x \leq 0.03 - 0.07$) and anti-shadowing ($0.07 - 0.03 \leq x \leq 0.2$) can be completely ignored in the analysis of the present DIS data on polarized nuclei.

VI. DISCUSSION OF FIT RESULTS

The best values for parton distribution functions are demonstrated in Table II. Parameters marked with (*) are fixed after an initial minimization step to their best values. Accordingly, there are 16 unknown parameters, including the strong coupling constant, that provide enough flexibility to have a reliable fit. We achieve $\chi^2/dof = 473.195/495 = 0.955$, that provides an acceptable fit to data. We extract the strong coupling constant simultaneously with polarized PDF parameters to study its correlation with the others. We obtain the value of $\alpha_s(Q_0^2) = 0.30998 \pm 0.0113$ at a 0.68% confidence level. Rescaling the coupling constant to the mass of the Z boson, we achieve $\alpha_s(M_Z^2) = 0.1106 \pm 0.0010$. The present world average value is $\alpha_s(M_Z^2) = 0.1185 \pm 0.0006$ [110].

A. Polarized parton distribution functions

Our NNLO polarized PDF along with their corresponding 0.68% C.L. uncertainties are presented in Fig. 5. Various parametrizations from the literature [9, 10, 53] and [45] obtained from NLO and NNLO QCD analyses of the inclusive data are presented for comparison. The $x(\Delta u + \Delta \bar{u})$ and $x\Delta G$ distributions are positive, while the $x(\Delta d + \Delta \bar{d})$ and $x(\Delta s + \Delta \bar{s})$ distributions are negative.

For the $x(\Delta u + \Delta \bar{u})$ distributions, all of the curves are comparable. Examining the $x(\Delta d + \Delta \bar{d})$ distributions we see that most of the fits are in agreement, with the possible exception of the LSS15 and NNPDF1.0.

All analyses of the polarized inclusive DIS data have extracted significantly negative results for the polarized strange quark distribution functions, $x(\Delta s + \Delta \bar{s})$, for all x values, even though the parametrization allowed a sign change as a function of x .

Results for polarized gluon distribution from the various fits on the present polarized inclusive DIS data are quite spread. The difficulties in constraining $x\Delta G$ cannot be ruled out with present data. Our gluon distribution tends to zero more quickly than the others.

B. g_1 structure functions

Figure 6 represents results for the polarized structure function xg_1^p and xg_1^d as a function of x at $Q^2 = 15 \text{ GeV}^2$ and $Q^2 = 14 \text{ GeV}^2$, respectively. For comparison, we illustrate the results extracted in BB10 [53], DNS05 [111] at the NLO approximation, and KTA-I17 [45] at the NNLO approximation. Our curves stay compatible with the recent COMPASS16 published data [54, 55] within statistical uncertainties. Note that the experimental observables belong to the scale region of $1.03 < Q^2 < 96.1 \text{ GeV}^2$ and $1.03 < Q^2 < 74.1 \text{ GeV}^2$ for proton and deuteron polarized structure functions.

In Fig. 7 we present our result for the polarized structure function of neutron xg_1^n as a function of x at $Q_0^2 = 4 \text{ GeV}^2$. We observe that our result coincides with those of BB10 [53], DNS05 [111] at the NLO approximation, and KTA-I17 [45] at the NNLO approximation. The HERMES06 experimental data [58] are well described within errors by all the curves. The experimental observables belong to the scale region of $1.12 < Q^2 < 14.29 \text{ GeV}^2$. Generally the xg_1^n data have larger uncertainties compared with the xg_1^p and xg_1^d data. More accurate experimental measurements on light nuclear targets are required to allow us to scrutinize neutron structure functions.

C. Nuclear polarized structure functions

We are in a position to apply the formalism developed in Sec. V to compute the $g_{1,2}^{3\text{He}}$ and $g_1^{3\text{H}}$ structure func-

Table II: Obtained parameter values and their statistical errors at the input scale $Q_0^2 = 1 \text{ GeV}^2$ determined from leading-twist analyses in the NNLO approximation. Those marked with (*) are fixed. Note that the TMCs are included in our QCD analysis.

Flavor	η	α	β	ϵ	γ
$u + \bar{u}$	0.807*	0.259 ± 0.007	2.857 ± 0.049	-4.95 ± 0.324	38.12 ± 1.85
$d + \bar{d}$	-0.461*	0.332 ± 0.72	3.139 ± 0.75	0	4.53 ± 1.85
$s + \bar{s}$	-0.119 ± 0.008	0.249 ± 0.048	15.68 ± 4.21	0	0
G	0.133 ± 0.027	23.33 ± 4.32	86.57 ± 6.35	1.434 ± 0.23	-4.992 ± 0.574

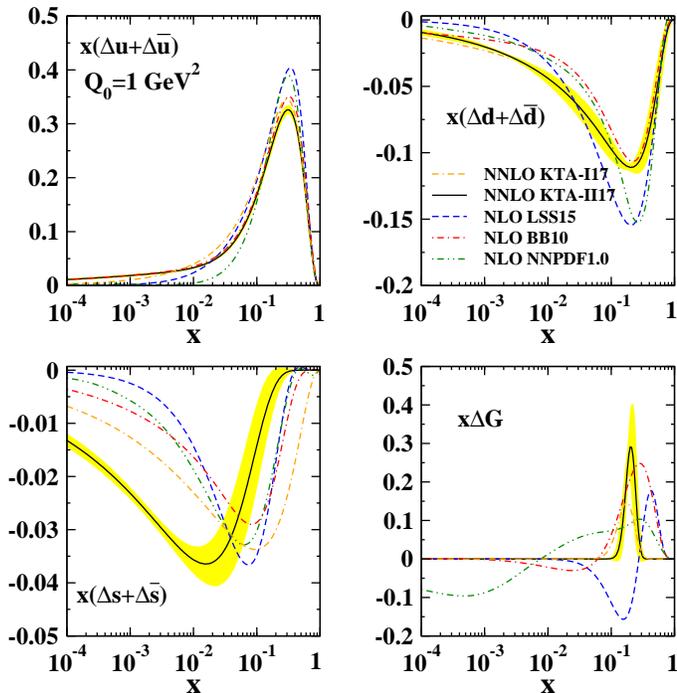


Figure 5: Our results KTA-III17 for the polarized PDFs at $Q_0^2 = 1 \text{ GeV}^2$ as a function of x in the NNLO approximation plotted as a solid curve along with their 68% C.L. uncertainties, as described in the text. We also show the results obtained in earlier global analyses of KTA-II17 (dashed-dashed-dotted) [45], LSS15 (dashed) [9], BB10 (dashed-dotted) [53], NNPDF1.0 (dashed-dotted-dotted) [10].

tions and corresponding nuclear corrections based on the extracted PPDFs. In particular, we study the impact of nuclear effects originating from the “non-nucleonic degrees of freedom” on the extraction of the spin structure of the ^3He and ^3H .

In Figs. (8) and (9), we show our results for the $g_1^{^3\text{He}}$ and $g_1^{^3\text{H}}$ polarized structure functions at NNLO approximation based on Eqs. (25) and (26) respectively, and compare with the curves based on KTA-II17 [45]. The experimental data from E142 [112] and JLAB04 [113]

are well described by the fit in Fig. 8. One can conclude that our results for the $g_1^{^3\text{He}}$ based on both of our spin-dependent PDFs (KTA-II17 and KTA-III17) reproduce the trend of the existing data.

We would like to stress again that the small x ($10^{-4} \leq x \leq 0.2$) effects from nuclear shadowing and antishadowing were not taken into account in the results presented in Figs. 8 and 9. At large x values, $x \geq 0.8$, all the results coincide. The nuclear corrections lead to a sizable difference in the small x values. For intermediate and low x values, including nuclear corrections underestimates the results.

In Figs. 10 and 11, our theory predictions for the polarized $g_1^{^3\text{He}}(x, Q^2)$ and $g_2^{^3\text{He}}(x, Q^2)$ structure functions are displayed as a function of x at NNLO approximation and compared with the recent data from the JLAB16 Collaboration [61]. The left plots correspond to $Q^2 = 4.74 \text{ GeV}^2$ and the right ones correspond to $Q^2 = 5.89 \text{ GeV}^2$. From Fig. 10, we can conclude that applying nuclear corrections to the spin-dependent ^3He structure functions decreases $g_1^{^3\text{He}}(x, Q^2)$ at small values of x . We observe similar trends for $g_2^{^3\text{He}}(x, Q^2)$ in Fig. 11. Figure 11 also indicates that the results are strongly model dependent in the whole Bjorken x region. The $g_2^{^3\text{He}}(x, Q^2)$ structure functions based on KTA-III17 polarized PDF describe the JLAB16 data better than KTA-II17 PPDFs.

The DIS data reported by E06-014 experiments at Jefferson Lab (JLAB) in Hall A are the latest and most up-to-date data for the spin-dependent g_1 and g_2 structure functions of ^3He [61]. These data sets were obtained from the scattering of a longitudinally polarized electron beam from a transversely and longitudinally polarized ^3He target. This measurement covers the kinematic regions of $0.25 \leq x \leq 0.9$ and $2 \text{ GeV}^2 \leq Q^2 \leq 6 \text{ GeV}^2$.

In Fig. 12, the spin-dependent $x^2 g_1^{^3\text{He}}(x, Q^2)$ structure function is plotted and compared to the world DIS data from E142 [112], JLAB04 [113], JLAB03 [115] and recent data from JLAB16 collaboration [61]. KTA-III17 predictions follow the trend of existing data.

D. Bjorken sum rule

Having at hand the spin-dependent structure of the proton and neutron, one may also examine in more detail

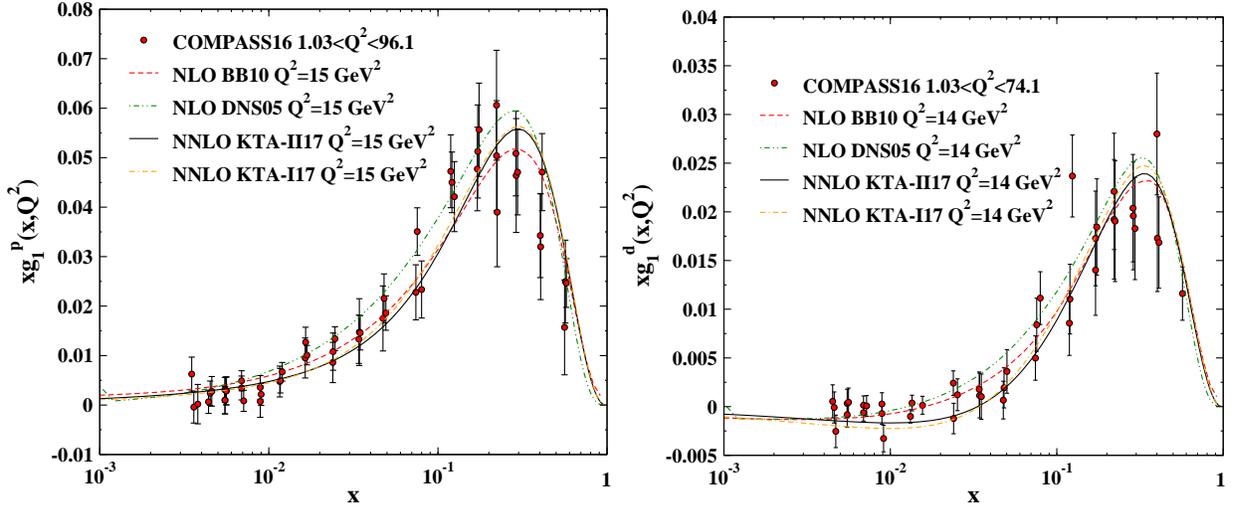


Figure 6: KTA-III7 prediction (solid curve) for the polarized structure function of the proton (left) and deuteron (right) as a function of x for the mean value of $Q^2 = 15 \text{ GeV}^2$ and $Q^2 = 14 \text{ GeV}^2$, respectively. Also shown are BB10 (dashed) [53], DNS05 (dashed-dotted-dotted) [111] extracted at the NLO approximation, and KTA-I17 (dashed-dashed-dotted) [45] obtained at the NNLO approximation together with the recent experimental data from the COMPASS16 collaborations [54, 55].

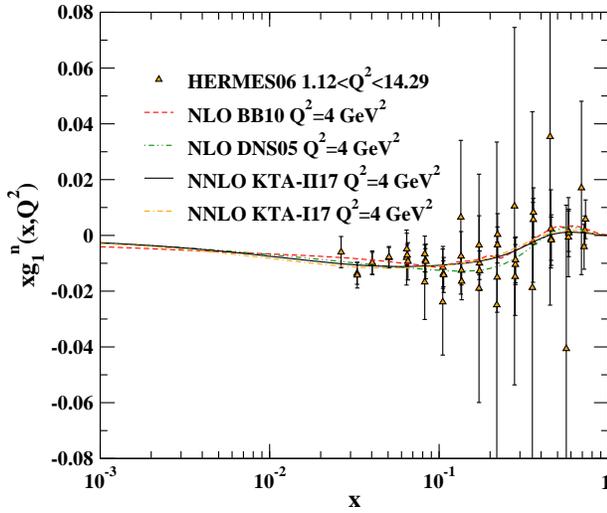


Figure 7: KTA-III7 prediction (solid curve) for the polarized structure function of neutron for the mean value of $Q^2 = 4 \text{ GeV}^2$ as a function of x in the NNLO approximation. Also shown are BB10 (dashed) [53], DNS05 (dashed-dotted-dotted) [111] extracted at the NLO approximation, and KTA-I17 (dashed-dashed-dotted) [45] obtained at the NNLO approximation together with the experimental data from the HERMES06 Collaboration [58].

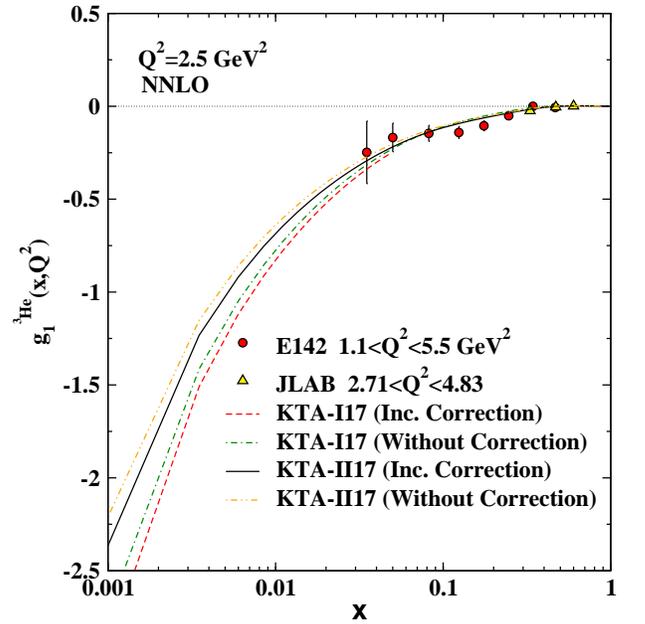


Figure 8: Analytical results with and without nuclear corrections for the polarized structure function of $g_1^{3\text{He}}(x, Q^2)$ as a function of x at NNLO approximation. The current fit is the solid curve. Also shown are data from E142 [112] and JLAB04 [113].

the Bjorken sum rule [116]

$$\int_0^1 [g_1^p(x, Q^2) - g_1^n(x, Q^2)] dx = \frac{1}{6} g_A [1 + \mathcal{O}(\frac{\alpha_s}{\pi})], \quad (31)$$

which relates the difference of the first moments of the proton $\int_0^1 g_1^p(x, Q^2) dx$ and neutron $\int_0^1 g_1^n(x, Q^2) dx$ spin structure functions to the axial vector coupling constant

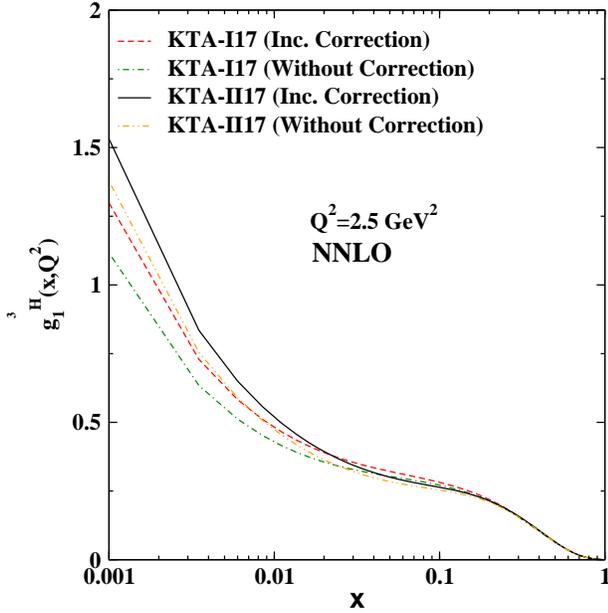


Figure 9: Analytical results with and without nuclear corrections for the polarized structure function of $g_1^{3\text{H}}(x, Q^2)$ as a function of x at NNLO approximation. The current fit is the solid curve. Also shown are the QCD NNLO curves obtained by KTA-II7 [45] for comparison.

measured in the β -decay of neutrons, $g_A = 1.2670 \pm 0.0035$ [110]. This sum rule can be straightforwardly generalized for the difference of the spin structure functions of ${}^3\text{He}$ and ${}^3\text{H}$ as follows [93, 117]

$$\int_0^3 [g_1^{3\text{H}}(x, Q^2) - g_1^{3\text{He}}(x, Q^2)] dx = \frac{1}{6} g_A|_{\text{triton}} [1 + \mathcal{O}(\frac{\alpha_s}{\pi})], \quad (32)$$

where $g_A|_{\text{triton}}$ is the axial vector coupling constant measured in the β decay of the triton, with $g_A|_{\text{triton}} = 1.211 \pm 0.002$ [?]. Finally, taking the ratio of Eqs. (31) and (32), one can find [93, 117]

$$\eta \equiv \frac{g_A|_{\text{triton}}}{g_A} = \frac{\int_0^3 [g_1^{3\text{H}}(x, Q^2) - g_1^{3\text{He}}(x, Q^2)] dx}{\int_0^1 [g_1^p(x, Q^2) - g_1^n(x, Q^2)] dx} = 0.9937 \pm 0.004. \quad (33)$$

We compute the above ratio to be $\eta = 0.923$ based on Eqs. (21) and (22). Including nuclear corrections Eqs. (25) and (26) modifies this value to $\eta = 0.970$.

One can conclude that the corrections associated with the presence of the Δ resonance change the value of the Bjorken sum rule in the $A = 3$ nuclei. It has been shown in Refs. [93, 117] that the contributions to the spin-dependent structure functions of ${}^3\text{He}$ from non-nucleonic degrees of freedom in the nucleus, such as the $\Delta(1232)$

isobar, lead to the $\approx 4\%$ difference between the g_A in the free nucleon and $g_A|_{\text{triton}}$ in the $A = 3$ nuclei.

VII. SUMMARY AND CONCLUSIONS

Our new NNLO analysis of the inclusive data included for the first time the extremely accurate COMPASS16 data on protons and deuterons [54, 55]. During the analysis, we considered TMCs to extract the polarized PDFs inside the nucleon. We adopt more general input parametrizations for the sum of quark and antiquark polarized PDFs instead of the valence and sea quark distributions. Similar to our previous papers, we use the Jacobi polynomial method to yield the structure functions $g_1^{N=p,n,d}(x, Q^2)$ from its moments in the whole x region. Having extracted the polarized structure functions, we estimated the nuclear structure functions of $g_{1,2}^{3\text{He}}$ and $g_1^{3\text{H}}$.

Due to increasing levels of precision attained in new generations of polarized DIS experiments, spin-dependent ${}^3\text{He}$ and ${}^3\text{H}$ targets become essential tools for studying the spin structure of the nucleons. They are also providing the most direct means of probing the polarized quarks and gluon distributions in the free neutron.

We also have performed a detailed analysis of nuclear corrections to the spin-dependent $g_{1,2}^{3\text{He}}$ and $g_1^{3\text{H}}$ structure functions. In addition to the nuclear effects arising from the “incoherent” scattering on nuclear targets, we have also examined the contributions from “non-nucleonic degrees of freedom” and have related the strength of these corrections to the Bjorken sum rule in the $A=3$ nuclei.

In this paper, we carry out an approximate way of assessing the importance of the nuclear effects. Neutron data are coming from deuterium and helium-3, and their corresponding nuclear corrections are ignored. We believe that the neutron data are only reliable, if, we assume nuclear corrections are negligible whereas they are not. We suggest re-extracting the $g_1^n(x, Q^2)$ from $g_1^{3\text{He}}(x, Q^2)$ data as previously argued in Ref. [93].

ACKNOWLEDGMENTS

The authors would like to thank Elliott Leader and Emanuele Nocera for their careful reading of the manuscript and helpful discussions. The authors are also grateful to the School of Particles and Accelerators, Institute for Research in Fundamental Sciences (IPM) for financial support of this project. Hamzeh Khanpour is grateful to the University of Science and Technology of Mazandaran for financial support provided for this research.

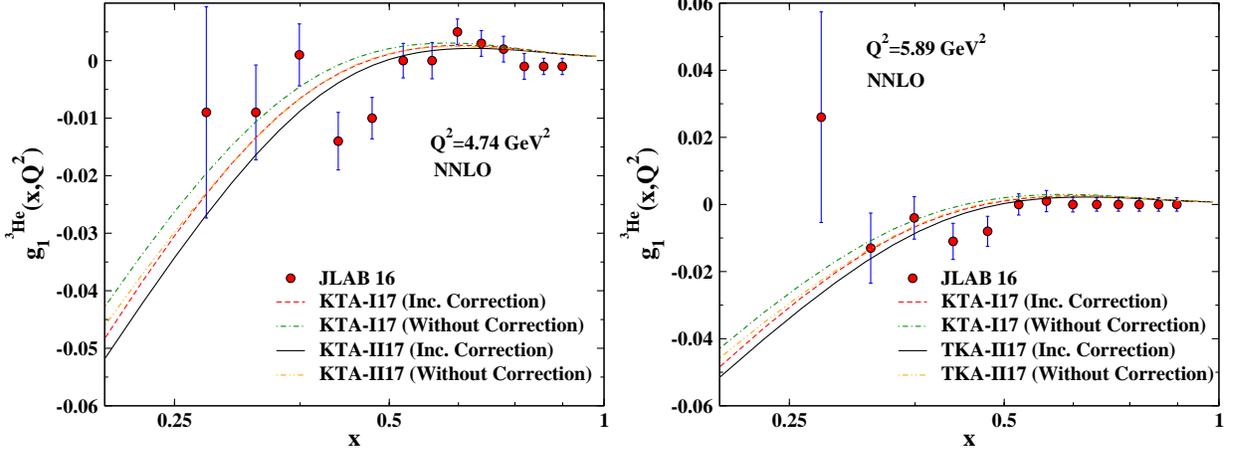


Figure 10: Analytical result for the polarized $g_1^{3\text{He}}(x, Q^2)$ structure function as a function of x at NNLO approximation that has been compared with the recent and up-to-date experimental data from JLAB16 Collaboration [61]. The left plot corresponds to $Q^2 = 4.74 \text{ GeV}^2$ and the right one corresponds to $Q^2 = 5.89 \text{ GeV}^2$.

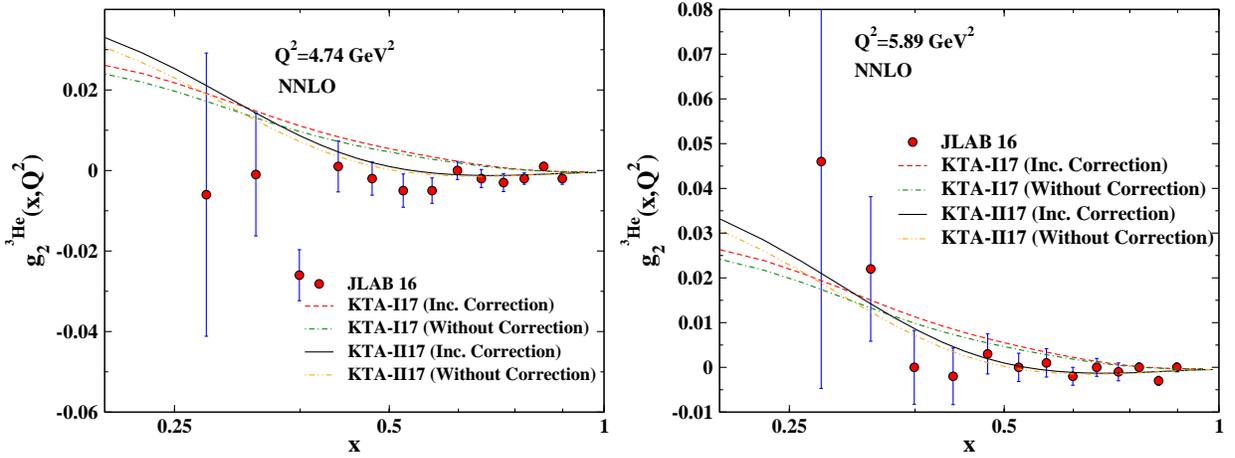


Figure 11: Analytical result for the polarized $g_2^{3\text{He}}(x, Q^2)$ structure function as a function of x at the NNLO approximation, which has been compared with the recent and up-to-date experimental data from the JLAB16 Collaboration [61]. The left plot corresponds to $Q^2 = 4.74 \text{ GeV}^2$ and the right one corresponds to $Q^2 = 5.89 \text{ GeV}^2$.

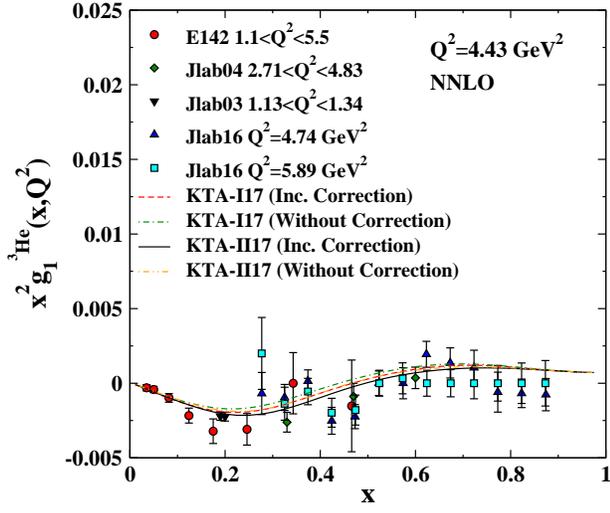


Figure 12: KTA-III7 results for the $x^2 g_1^{3\text{He}}(x, Q^2)$ structure function as a function of x at the NNLO approximation. Our curves are compared to the world data from E142 [112], JLAB04 [113], JLAB03 [115] and recent data from the JLAB16 Collaboration [61].

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