

Appendix — Minimal Condition for a Local Repulsive Term Canceling Weight in the HGR Framework

Laurent Besson (original idea)

Appendix writing: ChatGPT

Working version — April 20, 2026

Purpose of this appendix

The purpose of this appendix is to examine, within the spirit of Hypercomplex General Relativity (HGR), the minimal conditions required for the appearance of a *local repulsive term* capable of canceling the weight of a test mass m placed above a source mass M .

The goal here is not to propose a realistic experimental device, but rather to fix a minimal theoretical architecture for *local antigravity*, understood as an activated, anisotropic, and spatially confined geometric correction.

1 Formulation of the local problem

Consider a source mass M and a test mass m placed at a fixed altitude r_0 above M . In the non-relativistic limit, the standard gravitational acceleration is

$$\mathbf{a}_M(r) = -\frac{GM}{r^2} \hat{\mathbf{r}}. \quad (1)$$

To cancel locally the weight of m at the point r_0 , one must introduce a local repulsive contribution such that

$$\mathbf{a}_{\text{eff}}(r_0) = \mathbf{a}_M(r_0) + \mathbf{a}_{\text{rep}}(r_0) = 0. \quad (2)$$

The minimal condition is therefore

$$\boxed{\mathbf{a}_{\text{rep}}(r_0) = +\frac{GM}{r_0^2} \hat{\mathbf{r}}.} \quad (3)$$

2 Geometric translation in terms of an effective connection

In the weak-field approximation and for a slow particle, the spatial acceleration is related to the effective connection by

$$a^i \simeq -c^2 \Gamma_{00,\text{eff}}^i. \quad (4)$$

We then introduce an effective connection of the form

$$\Gamma_{\mu\nu,\text{eff}}^\rho = \Gamma_{\mu\nu}^\rho(g) + \Delta\Gamma_{\mu\nu}^\rho(H), \quad (5)$$

where $\Delta\Gamma_{\mu\nu}^\rho(H)$ represents the correction induced by the hypercomplex sector.

The local weight-cancellation condition then becomes

$$\boxed{-c^2 \Delta\Gamma_{00}^r(r_0) = +\frac{GM}{r_0^2}.} \quad (6)$$

In other words, the hypercomplex sector must locally correct transport in such a way as to produce a radial acceleration opposite to that generated by the source mass.

3 Minimal affine hypercomplex ansatz

In the spirit of the strong affine extension of HGR, one may take as a minimal model

$$\Delta\Gamma_{\mu\nu}^\rho(H) = \frac{\lambda}{2} \left(A_\mu^I \Sigma_{\nu I}^\rho + A_\nu^I \Sigma_{\mu I}^\rho \right), \quad (7)$$

where:

- A_μ^I is an internal hypercomplex connection;
- $\Sigma_{\nu I}^\rho$ is a base-fiber projection tensor;
- λ is an effective geometric coupling.

In this case, the local compensation condition reads

$$\boxed{-c^2 \frac{\lambda}{2} \left(A_0^I \Sigma_{0I}^r + A_0^I \Sigma_{0I}^r \right)_{r=r_0} = +\frac{GM}{r_0^2}.} \quad (8)$$

Equivalently,

$$\boxed{-c^2 \lambda \left(A_0^I \Sigma_{0I}^r \right)_{r=r_0} = +\frac{GM}{r_0^2}.} \quad (9)$$

4 Need for local activation and a spatial domain

A repulsive term useful for local antigravity cannot be merely cosmological. It must be *locally activated* within a spatial region D .

We therefore introduce a domain function

$$\mathcal{A}_D(x) \in [0, 1], \quad (10)$$

which is approximately 1 inside the weight-cancellation region and 0 outside.

A possible smooth form is

$$\mathcal{A}_D(x) = \exp\left(-\frac{d(x, D)^2}{L^2}\right), \quad (11)$$

where L sets the characteristic thickness of the domain.

The local effective connection then becomes

$$\Gamma_{\mu\nu, \text{eff}}^\rho = \Gamma_{\mu\nu}^\rho(g) + \mathcal{A}_D(x) \Delta\Gamma_{\mu\nu}^\rho(H). \quad (12)$$

The local condition at the center x_0 of the domain is

$$\boxed{-c^2 \mathcal{A}_D(x_0) \Delta\Gamma_{00}^r(x_0) = +\frac{GM}{r_0^2}.} \quad (13)$$

5 Need for directional anisotropy

A simple isotropic scalar correction is insufficient to cancel weight in a given direction. To compensate the gravitational field of M above the source, one needs a structure capable of selecting the radial direction $\hat{\mathbf{r}}$.

The local repulsive term must therefore be at least:

- either vectorial,
- or affinely anisotropic,
- or a mixed base–fiber tensorial structure.

In the ansatz (7), this anisotropy is carried by the projection $\Sigma_{\nu I}^\rho$ and by the internal orientation of the sector A_μ^I .

6 Local stability condition

Canceling the weight at a single point is not sufficient to obtain a useful static levitation. One must also require stability under small radial perturbations.

If $a_{\text{eff}}(r)$ denotes the effective radial acceleration, the minimal conditions are

$$a_{\text{eff}}(r_0) = 0, \quad (14)$$

and

$$\left. \frac{da_{\text{eff}}}{dr} \right|_{r=r_0} < 0. \quad (15)$$

The second condition ensures that a small deviation near r_0 does not turn the equilibrium into an explosive instability.

7 Minimal radial toy model

A minimal effective form for the acceleration is

$$\mathbf{a}_{\text{eff}}(r) = -\frac{GM}{r^2} \hat{\mathbf{r}} + \mathcal{A}_D(r) a_H(r) \hat{\mathbf{r}}. \quad (16)$$

The compensation condition at the target altitude r_0 is

$$\boxed{a_H(r_0) = \frac{GM}{r_0^2}}. \quad (17)$$

A simple form for the local repulsive profile is

$$a_H(r) = a_0 f(r), \quad (18)$$

where $f(r)$ is centered on r_0 and decays rapidly outside the domain.

For instance,

$$f(r) = \exp\left(-\frac{(r-r_0)^2}{\ell^2}\right), \quad (19)$$

with the adjustment condition

$$a_0 = \frac{GM}{r_0^2} \quad (20)$$

if $f(r_0) = 1$.

8 Effective tensorial writing

At the level of the field equation, the idea may be summarized by

$$G_{\mu\nu} = 8\pi G T_{\mu\nu}^{(m)} + \mathcal{A}_D(x) \Theta_{\mu\nu}^{\text{rep}}(x). \quad (21)$$

But if one wants an oriented local compensation of weight, the tensor $\Theta_{\mu\nu}^{\text{rep}}$ should not be assimilated to an ordinary isotropic fluid. It must contain anisotropic components capable of locally correcting the component Γ_{00}^r of the effective connection.

In other words, the correct physical language is no longer merely that of a “repulsive density”, but that of an *activated anisotropic geometric domain*.

Conclusion

Within the HGR framework, the appearance of a static local antigravity effect above a source mass would require at minimum:

- a local affine correction of the type $\Delta\Gamma_{\mu\nu}^\rho(H)$;
- a spatial activation confined within a domain D ;
- a directional anisotropy capable of selecting the radial axis;
- an exact compensation condition,

$$-c^2 \Delta\Gamma_{00}^r(r_0) = +\frac{GM}{r_0^2},$$

or, with an activated domain,

$$-c^2 \mathcal{A}_D(x_0) \Delta\Gamma_{00}^r(x_0) = +\frac{GM}{r_0^2};$$

- a local stability condition on the radial profile of the repulsive term.

The minimal formulation therefore suggests that a *local cancellation of weight* would not result from a simple homogeneous cosmological term, but from a genuine *local, activated, and oriented geometric bubble* producing an affine correction opposite to the gravitational field of the source mass.