

Description of the Module “probabilities”

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The normalized Maxwell-Boltzmann distribution is [e.g. see <http://mathworld.wolfram.com/MaxwellDistribution.html>]:

$$F(v)dv = 4\pi \left(\frac{m}{2\pi k_B T}\right)^{3/2} v^2 \exp\left(-\frac{mv^2}{2k_B T}\right) dv = \sqrt{\frac{2}{\pi}} \left(\frac{m}{k_B T}\right)^{3/2} v^2 \exp\left(-\frac{mv^2}{2k_B T}\right) dv. \quad (1)$$

Using the auxiliary variable:

$$a = \sqrt{\frac{k_B T}{m}}. \quad (2)$$

one can simplify this result to:

$$F(v)dv = \sqrt{\frac{2}{\pi}} a^{-3} v^2 \exp\left(-\frac{v^2}{2a^2}\right) dv. \quad (3)$$

The cumulative probability function is given by:

$$D(v) = \operatorname{erf}\left(\frac{v}{a\sqrt{2}}\right) - \sqrt{\frac{2}{\pi}} \frac{v}{a} \exp\left(-\frac{v^2}{2a^2}\right). \quad (4)$$

$$m = \frac{k_B T}{a^2}. \quad (5)$$

$$E = \frac{mv^2}{2} = \frac{k_B T}{a^2} \frac{v^2}{2} = \frac{k_B T}{2} \left(\frac{v}{a}\right)^2 \rightarrow \frac{v}{a} = \left(\frac{2E}{k_B T}\right)^{1/2}. \quad (6)$$

So:

$$D(E) = \operatorname{erf}\left(\left(\frac{E}{k_B T}\right)^{1/2}\right) - \sqrt{\frac{4}{\pi}} \left(\frac{E}{k_B T}\right)^{1/2} \exp\left(-\frac{E}{k_B T}\right). \quad (7)$$

Lets compute the energy probability distribution function:

$$\begin{aligned} P(E) = \frac{dD}{dE} &= \frac{2}{\sqrt{\pi}} \frac{1}{k_B T} \left(\frac{E}{k_B T}\right)^{-1/2} \exp\left(-\frac{E}{k_B T}\right) - \sqrt{\frac{4}{\pi}} \frac{1}{k_B T} \left(\frac{E}{k_B T}\right)^{-\frac{1}{2}} \exp\left(-\frac{E}{k_B T}\right) \\ &+ \frac{1}{k_B T} \sqrt{\frac{4}{\pi}} \left(\frac{E}{k_B T}\right)^{\frac{1}{2}} \exp\left(-\frac{E}{k_B T}\right) = \frac{1}{k_B T} \sqrt{\frac{4}{\pi}} \left(\frac{E}{k_B T}\right)^{\frac{1}{2}} \exp\left(-\frac{E}{k_B T}\right) \end{aligned}$$

which agrees with the result here:

https://en.wikipedia.org/wiki/Maxwell%E2%80%93Boltzmann_distribution

Consider a harmonic oscillator with a frequency of mode i , ω_i . The HO can be in any vibrational state.

The fraction of the first N states $[0, \dots, N-1]$ in the total sum:

$$Z_N = \sum_{n=0}^N \exp\left(-\frac{\hbar\omega_i\left(n+\frac{1}{2}\right)}{k_B T}\right) = \exp\left(-\frac{\hbar\omega_i}{2k_B T}\right) \sum_{n=0}^N r^n = \exp\left(-\frac{\hbar\omega_i}{2k_B T}\right) \frac{1-r^{N+1}}{1-r} = \exp\left(-\frac{\hbar\omega_i}{2k_B T}\right) \frac{1-\exp\left(-\frac{\hbar\omega_i N}{k_B T}\right)}{1-\exp\left(-\frac{\hbar\omega_i}{k_B T}\right)}. \quad (8)$$

Total number of states:

$$Z_\infty = \exp\left(-\frac{\hbar\omega_i}{2k_B T}\right) \frac{1}{1-\exp\left(-\frac{\hbar\omega_i}{k_B T}\right)}. \quad (9)$$

The probability to be in the vibrational state n is given by:

$$P_n = \frac{\exp\left(-\frac{\hbar\omega_i\left(n+\frac{1}{2}\right)}{k_B T}\right)}{Z_\infty} = \frac{\exp\left(-\frac{\hbar\omega_i\left(n+\frac{1}{2}\right)}{k_B T}\right)}{\exp\left(-\frac{\hbar\omega_i}{2k_B T}\right)} \left(1 - \exp\left(-\frac{\hbar\omega_i}{k_B T}\right)\right) = \exp\left(-\frac{\hbar\omega_i}{k_B T} n\right) \left(1 - \exp\left(-\frac{\hbar\omega_i}{k_B T}\right)\right). \quad (10)$$

The probability to be in any of states below state N :

$$P_{<N} = \frac{Z_N}{Z_\infty} = \frac{\exp\left(-\frac{\hbar\omega_i}{2k_B T}\right) \frac{1-\exp\left(-\frac{\hbar\omega_i N}{k_B T}\right)}{1-\exp\left(-\frac{\hbar\omega_i}{k_B T}\right)}}{\exp\left(-\frac{\hbar\omega_i}{2k_B T}\right) \frac{1}{1-\exp\left(-\frac{\hbar\omega_i}{k_B T}\right)}} = 1 - \exp\left(-\frac{\hbar\omega_i}{k_B T} N\right). \quad (11)$$

The probability that any of the states above N (including it) are occupied is given by:

$$P_{\geq N} = 1 - P_{<N} = \exp\left(-\frac{\hbar\omega_i}{k_B T} N\right). \quad (12)$$