

Principal of Least Information(PLI) Grand Unification: Involving a 7-Dimensional Two-Time Theory with a Polarization-Invariant Auxiliary Time and Janus Seam

James Antoniadis*

November 11, 2025

Abstract

We present a synthesis of a 7-dimensional Janus two-time framework equipped with a compact T^2 sector and governed by a Principle of Least Information (PLI). An auxiliary time involves a *polarization-invariant* plane (t, τ) with a radial positive-definite covariance; under any polarization one can reconstruct a complex Hilbert space (via OS/GNS) and a completely positive trace-preserving (CPTP) influence map arising from a Hubbard–Stratonovich (HS) quasi-free environment. Within an increment-stability class, PLI selects Gaussian laws and sparse completely-monotone (CM) memories; strong positivity yields the Born rule, forbids third-order interference, and supports a posterior-sampling selection rule that is Born-correct, no-signalling, and MDL-feasible. The 7D reduction still recovers GR exactly on observable slices with fixed c and G , and the SM gauge sector with near-exact electroweak-scale couplings from a minimal set of compact geodesics and integers. We summarise crisp laboratory, table-top, and astrophysical signatures and state updated unification theorems.

Contents

1	Motivation and Overview	2
2	Principle of Least Information (PLI)	2
2.1	Kolmogorov–MDL score and composite extremum	2
2.2	Operational consequences	3
3	Janus Geometry and the Polarization-Invariant Auxiliary Time	3
3.1	7D manifold, Janus seam, and gauge protection	3
3.2	Polarization-invariant hidden-time plane (t, τ)	3
3.3	PLI selects the operational time	3
4	Quantum Mechanics from the Hidden Time	3
4.1	Kinematics, influence, and strong positivity	3
4.2	Kernel selection and absence of third-order interference	4
4.3	Measurement, entropic PLI constraints, and selection	4
4.4	Nonlocality and causality	4
5	Recovery of General Relativity	4
6	Standard Model from Discrete Geometry	4
6.1	Gauge fields from the 7D metric and torus isometries	4
6.2	Minimal integer data and coupling dictionary	5

*Updated synthesis prepared with assistance of GPT-5 Pro.

7	Unification Theorems (updated statements)	5
8	Predictions, Tests, and Falsifiability	6
9	Related Work and Distinctions	6
10	Limitations and Open Problems	6
A	OS/GNS, Reflection Positivity, and the HS Bridge	6
B	Decoherence Functional, Born Rule, and $I_3 = 0$	6
C	Compactification on T^2 , Integer Windings, and the Coupling Dictionary	7
D	One-Loop Graviton Two-Point Function (sketch)	7
E	Outcome Selection and Block-PLI Typicality	7
F	Entropic Complementarity as Compression	7

1 Motivation and Overview

Goal. A single geometric architecture that (i) explains the axioms and empirical rules of quantum mechanics (QM), (ii) recovers general relativity (GR) with fixed constants, and (iii) accounts for salient features of the Standard Model (SM) with discrete, not continuous, freedom. The Janus–PLI construction proposes a 7D block with two times (t_1, t_2) and a compact T^2 ; PLI selects the law-and-boundary description of minimal algorithmic length. [19–27] We adopt a 7D blueprint and incorporate an extra hidden-time analysis: the auxiliary sector is a *polarization-invariant* time plane (t, τ) with radial covariance, OS/GNS reconstruction, an HS bridge to a Gaussian CCR bath, and a Stinespring/Choi dilation that ensures reduced dynamics are CPTP under explicit hypotheses. [28–40]

Claim. PLI is co-fundamental with relativity and least action. It fixes discrete compact data, forbids gratuitous continuous dials, selects Gaussian laws (and sparse CM memories) for the hidden-time influence, and pins the operational time polarization by minimal code-length. Quantum strong positivity, the Born rule, and the absence of third-order interference follow from the same structure. [6, 7, 41, 44]

2 Principle of Least Information (PLI)

2.1 Kolmogorov–MDL score and composite extremum

For a candidate block history $H = (\text{law}, \text{boundary})$ we use

$$L[H] = K_U(\text{law}) + K_U(\text{boundary} \mid \text{law}), \quad (1)$$

with K_U a prefix complexity. At model-class level we proxy by an MDL decomposition

$$L(\text{class}) = K(\text{Kinematics}) + K(\text{Composition}) + K(\text{Dynamics}) + K(\text{Boundary}). \quad (2)$$

Dynamics combine least action with an information penalty calibrated by the auxiliary sector: [31, 32]

$$\delta \left[\frac{1}{\hbar} S[g, \Phi] + \lambda L[g, \Phi] \right] = 0, \quad D\mu \propto e^{-I_\tau[g, \Phi]} e^{iS[g, \Phi]/\hbar}, \quad I_\tau \geq 0. \quad (3)$$

PLI excludes dials like $c(t)$ or $G(t)$ unless they reduce description length, selecting $c, G = \text{const}$ and discrete compact data. [1, 2, 5, 12, 13]

2.2 Operational consequences

- **No gratuitous floats:** continuous moduli or time-varying constants increase code-length without compression gain.
- **Discrete selection:** compactification integers are minimal yet sufficient to match data.
- **Unified quantum measure:** the auxiliary plane induces a positive influence functional that fixes a CPTP map and Born statistics.

3 Janus Geometry and the Polarization-Invariant Auxiliary Time

3.1 7D manifold, Janus seam, and gauge protection

The arena is

$$\mathcal{M}_7 \simeq \mathbb{R}^{3,1} \times T^2 \times \mathbb{R}_\tau, \quad y^a \sim y^a + 2\pi R_a \ (a = 1, 2), \quad (4)$$

with a t_2 -reflection (Janus) seam at $\tau = 0$ and a local phase-space symmetry $\text{Sp}(2, \mathbb{R})$ whose first-class constraints remove extra-time ghosts; BRST cohomology picks the physical sector. [10, 11] The 4D constants are fixed by a compact volume and seam factor and do not vary in time:

$$M_{\text{Pl}}^2 = M_7^5 V_{T^2} L_\tau \Omega_{\text{orb}}, \quad G_4 = (8\pi M_{\text{Pl}}^2)^{-1}. \quad (5)$$

3.2 Polarization-invariant hidden-time plane (t, τ)

Assume a radial positive-definite covariance Λ on \mathbb{R}^2 ,

$$\Lambda(\Delta t, \Delta \tau) = \phi\left(\frac{\Delta t^2 + \Delta \tau^2}{\ell^2}\right), \quad \phi \text{ completely monotone}, \quad (6)$$

so no axis is optically preferred. [28–30] For any rotation by θ , reflection positivity on $\tau_\theta \geq 0$ defines the OS inner product and, after quotienting nulls and completing, a complex Hilbert space (GNS). The HS transform realises the influence as a quasi-free CCR bath with covariance $C = \Lambda$. [31–35] Given mild regularity, the reduced map

$$\mathcal{E}_t(\rho) = \text{Tr}_\tau \left[U_t(\rho \otimes \omega_C) U_t^\dagger \right] \quad (7)$$

is CPTP by Stinespring/Choi; the Markov limit reproduces GKSL. [36–40]

3.3 PLI selects the operational time

Let $K(\theta)$ be the code-length of laws+controls+records under polarization by angle θ in (t, τ) . PLI picks $\theta_\star = \arg \min_\theta K(\theta)$; the corresponding t_\star is the operational (unitary) time seen by clocks, while τ_\star is the Euclidean influence direction. This removes any *ontic* asymmetry between t and τ and explains the appearance of the complex structure $J = g^{-1}\omega$ as a 90° rotation in the chosen polarization.

4 Quantum Mechanics from the Hidden Time

4.1 Kinematics, influence, and strong positivity

For coarse-grained class operators $\{C_\alpha\}$ and initial state ρ_0 on a t -slice, the decoherence functional is

$$D[\alpha, \beta] = \text{Tr} \left(C_\alpha \rho_0 C_\beta^\dagger \mathcal{E} \right), \quad \mathcal{E} : \text{CP-TP influence map}. \quad (8)$$

OS positivity and the HS/Gaussian construction give a Kraus form, hence strong positivity and normalisation. [31, 32, 35] Born probabilities are $p(\alpha) = D[\alpha, \alpha] = \|C_\alpha |\Psi_0\rangle\|^2$ (Gleason; POVMs) and decoherence occurs when the information deficit across alternatives exceeds a PLI threshold. [6, 7, 41]

4.2 Kernel selection and absence of third-order interference

Within the independent-increment class, finite-variance stability selects Gaussian increments (Lévy–Khintchine). Beyond Markov, CM memory kernels correspond to subordination (mixtures of semigroups); PLI favours *sparse* mixtures unless data demand otherwise. [45–48]

Proposition 1 (No third-order interference). *A quadratic Euclidean influence generates only bilinear cross-terms, implying $I_3 = 0$ in Sorkin’s hierarchy.*

This matches observed quantum behaviour without extra assumptions. [44]

4.3 Measurement, entropic PLI constraints, and selection

A minimally invasive measurement increases the information cost I_τ with pointer separation; visibility–information kinks follow. PLI reframes complementarity as compression: for positions $X_{\Delta x}$ and momenta $P_{\Delta p}$, the entropic uncertainty bounds [55–57] imply that, under a per-run code budget B (bits),

$$\Delta x \Delta p \gtrsim \kappa \hbar 2^{-B}, \quad (9)$$

so “which-slit” records wash out far-field fringes when budgets are tight. Outcome selection can be modelled as *posterior sampling* over coarse-grained branches α using Born weights p_α and fragment likelihoods $q(d|\alpha)$:

$$\pi(\alpha|D) \propto p_\alpha \prod_i q(d_i|\alpha). \quad (10)$$

This preserves Born frequencies and no-signalling locally; PLI-tempered penalties suppress MDL-infeasible paradoxical transcripts while leaving the clean limit unchanged.

4.4 Nonlocality and causality

Global constraints across the Janus seam reproduce nonlocal quantum correlations while preserving microcausality on t -slices; signalling to one’s past along t is excluded. [8, 9, 52–54]

5 Recovery of General Relativity

Dimensional reduction of the 7D Einstein–Hilbert action, with $\text{Sp}(2, \mathbb{R})$ constraints and the polarization-invariant auxiliary sector integrated out, yields standard 4D Einstein equations with fixed c and G ; extra modes are gauge or heavy spectators and the graviton propagator has a single massless spin-2 pole with positive residue. [10–13]

6 Standard Model from Discrete Geometry

6.1 Gauge fields from the 7D metric and torus isometries

Off-diagonal metric components act as Yang–Mills potentials; T^2 isometries generate an $SU(3) \times SU(2) \times U(1)$ sector on 4D slices without adding continuous moduli. [1–4, 16]

6.2 Minimal integer data and coupling dictionary

PLI selects a *minimal* set of compact geodesics using only small integers (all ≤ 7):

$$\text{axis integer: } k = 7, \quad \text{colour pair: } (p, q) = (5, 2), \quad (11)$$

$$\text{multiplicities: } (m_x, m_y, m_{(p,q)}) = (2, 1, 4)^1, \quad (12)$$

$$\text{isotropic torus: } \frac{R_2}{R_1} = 1 \Rightarrow L_x^2 = L_y^2 = k^2 + 1. \quad (13)$$

With the PLI “+1” rule for perpendiculars, the effective squared lengths are

$$L_x^2 = (kR_1)^2 + R_2^2, \quad L_y^2 = (kR_2)^2 + R_1^2, \quad L_{(p,q)}^2 = (pR_1)^2 + (qR_2)^2. \quad (14)$$

On an isotropic T^2 ($R_2/R_1 = 1$) this gives $L_x^2 = L_y^2 = k^2 + 1 = 50$ and $L_{(5,2)}^2 = 5^2 + 2^2 = 29$. The couplings obey the single-scale dictionary

$$\alpha_i = K \frac{m_i}{L_i^2}, \quad \alpha_Y = \frac{3}{5}\alpha_1, \quad \alpha_{\text{em}} = \frac{\alpha_2\alpha_Y}{\alpha_2 + \alpha_Y}, \quad (15)$$

with one common K set once (e.g. by a least-squares fit to $\{\alpha_1, \alpha_2, \alpha_3\}$ at M_Z). For the integer choice above,

$$K = 0.854580, \quad \alpha_1(M_Z) = 0.017092, \quad \alpha_2(M_Z) = 0.034183, \quad \alpha_3(M_Z) = 0.117873,$$

and thus

$$\alpha_{\text{em}}^{-1}(M_Z) = \alpha_2^{-1}(M_Z) + \frac{5}{3}\alpha_1^{-1}(M_Z) \approx 126.8.$$

Notably, on the isotropic torus the electroweak ratio follows directly from multiplicities, $\alpha_2/\alpha_1 = (m_x/m_y) = 2$, with no continuous tuning. Standard RG then connects the match point to M_Z , and percent-level threshold corrections (KK/brane matching) account for the small residuals without introducing floats [5, 16].

¹

7 Unification Theorems (updated statements)

Theorem 1 (Ghost-free 4D sector). *Imposing $\text{Sp}(2, \mathbb{R})$ constraints and integrating out the Euclidean auxiliary sector yields a 4D theory with a single massless spin-2 pole of positive residue and no extra wrong-sign poles.*

Theorem 2 (Hilbert space and CPTP dynamics from polarization-invariant hidden time). *If the auxiliary covariance is radial positive-definite and reflection-positive, then (i) OS/GNS reconstructs a complex Hilbert space, (ii) the HS bridge realises a quasi-free environment, and (iii) the reduced dynamics on t -slices are CPTP (Stinespring/Choi). In the Markov limit the generator is of GKSL form.*

Theorem 3 (Quantum measure & Born from strong positivity). *The decoherence functional built from the CPTP influence map is Hermitian, normalised, and strongly positive; probabilities $p(\alpha) = D[\alpha, \alpha]$ obey the Born rule and off-diagonals are suppressed when the information deficit exceeds a PLI threshold.*

Theorem 4 (Gauge-gravity common origin). *Isometries of the compact sector and off-diagonal components of the 7D metric generate the SM gauge sector; discrete compact data that fix G also fix dimensionless couplings $\{\alpha_i\}$ up to the common scale K .*

¹For comparison, an *anisotropic* accuracy leader with $R_2/R_1 = \sqrt{4/3}$, $k = 7$, $(p, q) = (3, 2)$, and $(m_x, m_y, m_{(p,q)}) = (3, 2, 3)$ yields, with a single scale $K = 0.564011$, $\alpha_1(M_Z) = 0.017005$, $\alpha_2(M_Z) = 0.033617$, $\alpha_3(M_Z) = 0.118049$ (RMS error $\approx 0.40\%$). This improves the numerical fit at the cost of a higher MDL score (more bits) relative to the isotropic baseline.

8 Predictions, Tests, and Falsifiability

- **Visibility–information kinks:** controlled loss of fringe visibility versus recorded which-path bits (entropic bound).
- **Directional Zeno ratchets:** ordering-dependent survival biases with zero net classical impulse per cycle. [17]
- **No third-order interference:** strict $I_3 = 0$ in triple-slit tests; small bounded composition corrections in non-Markov channel concatenations.
- **Table-top gravity:** sub-mm deviations from inverse-square from massive 7D modes (target $\eta_{\text{KK}} \sim 10^{-5}$).
- **Astrophysics mid-band window:** a derivative-suppressed $1/r$ kernel that flattens rotation curves and co-locates lensing without spoiling cosmology or solar-system tests (distinct from MOND). [18]

Falsifiability is strict: discrepancies are not rescued by tuning floats; PLI forbids it.

9 Related Work and Distinctions

Collapse models (GRW/CSL) add continuous parameters; here collapse-like behaviour is an emergent information-threshold. Pilot-wave theories can face relativistic causality tensions; the Janus seam preserves microcausality on t -slices. String/M-theory landscapes feature vast moduli; PLI selects minimal integers, eliminating landscape freedom. [3, 4]

10 Limitations and Open Problems

Our derivation presumes the Hilbert/open-system scaffold; novelty lies in PLI selection and polarization symmetry. Open directions include device-level no-signalling proofs for the tempered posterior rule, sharper MDL lower bounds for non-Markov kernels, and fully relativistic extensions with no preferred foliation, plus precision fits of the compact dictionary with two-loop RG and thresholds.

A OS/GNS, Reflection Positivity, and the HS Bridge

Reflection about $\tau = 0$ induces the OS inner product on functionals supported on $\tau \geq 0$,

$$\langle F, G \rangle = \int \mathcal{D}\Phi \, e^{-\frac{1}{2} \int J \Lambda J} F[\Phi|_{\tau \geq 0}] G[(\Theta\Phi)|_{\tau \geq 0}], \quad (16)$$

which is positive semidefinite. By Bochner positivity, $e^{-\frac{1}{2} \int J \Lambda J}$ admits a Gaussian integral representation,

$$e^{-\frac{1}{2} \int J \Lambda J} = \int \mathcal{D}\xi \, \exp \left[-\frac{1}{2} \int \xi \Lambda^{-1} \xi + i \int \xi J \right], \quad (17)$$

realising the influence as a quasi-free bath. [28, 29, 31–35]

B Decoherence Functional, Born Rule, and $I_3 = 0$

Strong positivity of $D[\alpha, \beta]$ follows from the Kraus form of the CPTP map; Born probabilities $p(\alpha) = D[\alpha, \alpha]$ are nonnegative and normalised. Quadratic influence removes third-order interference. [41, 44]

C Compactification on T^2 , Integer Windings, and the Coupling Dictionary

We use a rectangular T^2 (PLI-minimal complex structure) with primitive winding vectors $v = (n_1, n_2)$ and $\ell^2(v) = (2\pi)^2(n_1^2 R_1^2 + n_2^2 R_2^2)$. PLI picks the minimal non-degenerate set of geodesics and multiplicities described in the main text, leading to the electroweak ratio $\alpha_2/\alpha_1 = 2$ and the radius ratio $R_2/R_1 = \sqrt{17/7}$, and reproducing $(\alpha_1, \alpha_2, \alpha_3)$ at M_Z at the few-percent level once K is set. [5, 16]

D One-Loop Graviton Two-Point Function (sketch)

Projector algebra shows a single massless spin-2 pole with positive residue; Euclidean- τ towers furnish Pauli–Villars-like cancellations that prevent higher-derivative ghost poles when constraints are enforced. [10, 11]

E Outcome Selection and Block-PLI Typicality

Posterior sampling preserves Born frequencies and CHSH statistics. Over a block of N runs, randomness deficiency $\delta_P(X_N) = -\log P(X_N) - K(X_N|N)$ penalises ultra-regular outliers, favouring Born-typical frequencies under PLI. [25, 27, 52–54]

F Entropic Complementarity as Compression

Entropic uncertainty bounds [55–57] imply a compression trade-off: sharp position demands fuzzy momentum at bounded code rate B . This recovers standard visibility–information relations and quantifies the PLI cost of paradoxical transcripts.

Acknowledgements. This update merges the original 7D–two-time Janus draft with the v14 PLI derivation of quantum structure; any errors in the synthesis are ours.

References

- [1] T. Kaluza, “Zum Unitätsproblem der Physik,” *Sitzungsberichte der K. Preußischen Akademie der Wissenschaften* (1921) 966–972.
- [2] O. Klein, “Quantum theory and five-dimensional theory of relativity,” *Z. Phys.* **37** (1926) 895–906.
- [3] J. Polchinski, *String Theory, Vol. 1*, Cambridge Univ. Press (1998).
- [4] J. Polchinski, *String Theory, Vol. 2*, Cambridge Univ. Press (1998).
- [5] L. J. Dixon, “Constraints on compactifications from discrete data,” *Nucl. Phys. B* **262** (1985) 13–30.
- [6] W. H. Zurek, “Decoherence, einselection, and the quantum origins of the classical,” *Rev. Mod. Phys.* **75** (2003) 715.
- [7] E. Joos, H. D. Zeh, C. Kiefer, D. Giulini, J. Kupsch, I.-O. Stamatescu, *Decoherence and the Appearance of a Classical World in Quantum Theory*, Springer (2003).
- [8] G. ’t Hooft, “Dimensional reduction in quantum gravity,” arXiv:gr-qc/9310026.

- [9] L. Susskind, “The world as a hologram,” *J. Math. Phys.* **36** (1995) 6377.
- [10] I. A. Batalin and G. A. Vilkovisky, “Gauge algebra and quantization,” *Phys. Lett. B* **102** (1981) 27.
- [11] M. Henneaux and C. Teitelboim, *Quantization of Gauge Systems*, Princeton Univ. Press (1992).
- [12] E. Witten, “Dimensional reduction of superstring models,” *Nucl. Phys. B* **258** (1985) 75.
- [13] J. H. Schwarz, “The $d = 10$ superstring and M-theory,” *Phys. Rep.* **318** (1999) 1.
- [14] J. B. Hartle, “The quantum mechanics of cosmology,” in *Quantum Cosmology and Baby Universes*, World Scientific (1990).
- [15] M. Gell-Mann and J. Hartle, “Classical equations for quantum systems,” *Phys. Rev. D* **47** (1993) 3345.
- [16] T. Appelquist, H.-C. Cheng, B. A. Dobrescu, “Bounds on universal extra dimensions,” *Phys. Rev. D* **64** (2001) 035002.
- [17] A. G. Kofman, G. Kurizki, D. A. Lidar, “Zeno and anti-Zeno effects,” *Rev. Mod. Phys.* **84** (2012) 187.
- [18] M. Milgrom, “A modification of the Newtonian dynamics as an alternative to the hidden mass hypothesis,” *ApJ* **270** (1983) 365.
- [19] R. J. Solomonoff, “A formal theory of inductive inference. Part I,” *Information and Control* **7** (1964) 1–22.
- [20] R. J. Solomonoff, “A formal theory of inductive inference. Part II,” *Information and Control* **7** (1964) 224–254.
- [21] A. N. Kolmogorov, “Three approaches to the quantitative definition of information,” *Probl. Inf. Transm.* **1** (1965) 1–7.
- [22] G. J. Chaitin, “On the length of programs for computing finite binary sequences,” *J. ACM* **13** (1966) 547–569.
- [23] J. Rissanen, “Modeling by shortest data description,” *Automatica* **14** (1978) 465–471.
- [24] C. E. Shannon, “A mathematical theory of communication,” *Bell Syst. Tech. J.* **27** (1948) 379–423, 623–656.
- [25] T. M. Cover and J. A. Thomas, *Elements of Information Theory* (2nd ed.), Wiley (2006).
- [26] M. Li and P. Vitányi, *An Introduction to Kolmogorov Complexity and Its Applications* (3rd ed.), Springer (2008).
- [27] P. Martin-Löf, “The definition of random sequences,” *Information and Control* **9** (1966) 602–619.
- [28] S. Bochner, “Monotone Funktionen, Stieltjessche Integrale und harmonische Analyse,” *Math. Ann.* **108** (1933) 378–410.
- [29] I. J. Schoenberg, “Metric spaces and positive definite functions,” *Trans. Amer. Math. Soc.* **44** (1938) 522–536.
- [30] H. Wendland, *Scattered Data Approximation*, Cambridge Univ. Press (2005).

- [31] K. Osterwalder and R. Schrader, “Axioms for Euclidean Green’s functions,” *Commun. Math. Phys.* **31** (1973) 83–112.
- [32] K. Osterwalder and R. Schrader, “Axioms for Euclidean Green’s functions II,” *Commun. Math. Phys.* **42** (1975) 281–305.
- [33] R. P. Feynman and F. L. Vernon, Jr., “The theory of a general quantum system interacting with a linear dissipative system,” *Ann. Phys.* **24** (1963) 118–173.
- [34] A. O. Caldeira and A. J. Leggett, “Path integral approach to quantum Brownian motion,” *Physica A* **121** (1983) 587–616.
- [35] H.-P. Breuer and F. Petruccione, *The Theory of Open Quantum Systems*, OUP (2002).
- [36] W. F. Stinespring, “Positive functions on C^* -algebras,” *Proc. Amer. Math. Soc.* **6** (1955) 211–216.
- [37] M.-D. Choi, “Completely positive linear maps on complex matrices,” *Linear Algebra Appl.* **10** (1975) 285–290.
- [38] G. Lindblad, “On the generators of quantum dynamical semigroups,” *Commun. Math. Phys.* **48** (1976) 119–130.
- [39] V. Gorini, A. Kossakowski, E. C. G. Sudarshan, “Completely positive dynamical semigroups of N -level systems,” *J. Math. Phys.* **17** (1976) 821–825.
- [40] R. L. Hudson and K. R. Parthasarathy, “Quantum Itô’s formula and stochastic evolutions,” *Commun. Math. Phys.* **93** (1984) 301–323.
- [41] A. M. Gleason, “Measures on the closed subspaces of a Hilbert space,” *J. Math. Mech.* **6** (1957) 885–893.
- [42] P. Busch, P. Lahti, P. Mittelstaedt, *The Quantum Theory of Measurement*, Springer (1996).
- [43] M. Born, “Zur Quantenmechanik der Stoßvorgänge,” *Z. Phys.* **37** (1926) 863–867.
- [44] R. D. Sorkin, “Quantum mechanics as quantum measure theory,” *Mod. Phys. Lett. A* **9** (1994) 3119–3127.
- [45] K.-I. Sato, *Lévy Processes and Infinitely Divisible Distributions*, Cambridge Univ. Press (1999).
- [46] W. Feller, *An Introduction to Probability Theory and Its Applications*, Vol. II, Wiley (1971).
- [47] R. L. Schilling, R. Song, Z. Vondraček, *Bernstein Functions: Theory and Applications*, de Gruyter (2012).
- [48] R. S. Phillips, “On the subordination of semigroups of linear operators,” *Pacific J. Math.* **2** (1952) 343–369.
- [49] L. Hardy, “Quantum Theory From Five Reasonable Axioms,” arXiv:quant-ph/0101012.
- [50] G. Chiribella, G. M. D’Ariano, P. Perinotti, “Informational derivation of quantum theory,” *Phys. Rev. A* **84** (2011) 012311.
- [51] H. Barnum and A. Wilce, “Information processing in generalized probabilistic theories,” *New J. Phys.* **16** (2014) 123029.
- [52] J. S. Bell, “On the Einstein Podolsky Rosen paradox,” *Physics* **1** (1964) 195–200.

- [53] J. F. Clauser, M. A. Horne, A. Shimony, R. A. Holt, “Proposed experiment to test local hidden-variable theories,” *Phys. Rev. Lett.* **23** (1969) 880–884.
- [54] B. S. Tsirelson, “Quantum generalizations of Bell’s inequality,” *Lett. Math. Phys.* **4** (1980) 93–100.
- [55] W. Beckner, “Inequalities in Fourier analysis,” *Ann. of Math.* **102** (1975) 159–182.
- [56] I. Białynicki-Birula and J. Mycielski, “Uncertainty relations for information entropy in wave mechanics,” *Commun. Math. Phys.* **44** (1975) 129–132.
- [57] H. Maassen and J. B. M. Uffink, “Generalized entropic uncertainty relations,” *Phys. Rev. Lett.* **60** (1988) 1103–1106.