

# Global Analysis of Neutrino Oscillations and Mass Constraints in the Era of Subpercent Precision







# Neutrino oscillation phenomenology: entering the precision era

# Neutrino oscillation phenomenology: entering the precision era

Solar parameters

$$\delta m^2 \sim 7.37 \times 10^{-5} \text{ eV}^2 \quad (2.3\%)$$

$$\sin^2 \theta_{12} \sim 0.303 \quad (4.5\%)$$

Reactor mixing angle

$$\sin^2 \theta_{13} \sim 2.23 \times 10^{-2} \quad (2.4\%)$$



# Neutrino oscillation phenomenology: entering the precision era

Solar parameters

$$\delta m^2 \sim 7.37 \times 10^{-5} \text{ eV}^2 \quad (2.3\%)$$

$$\sin^2 \theta_{12} \sim 0.303 \quad (4.5\%)$$

Reactor mixing angle

$$\sin^2 \theta_{13} \sim 2.23 \times 10^{-2} \quad (2.4\%)$$

Atmospheric parameters

$$\Delta m^2 \sim 2.495 \times 10^{-3} \text{ eV}^2 \quad (0.8\%)$$

$$\Delta m^2 \sim 2.465 \times 10^{-3} \text{ eV}^2$$

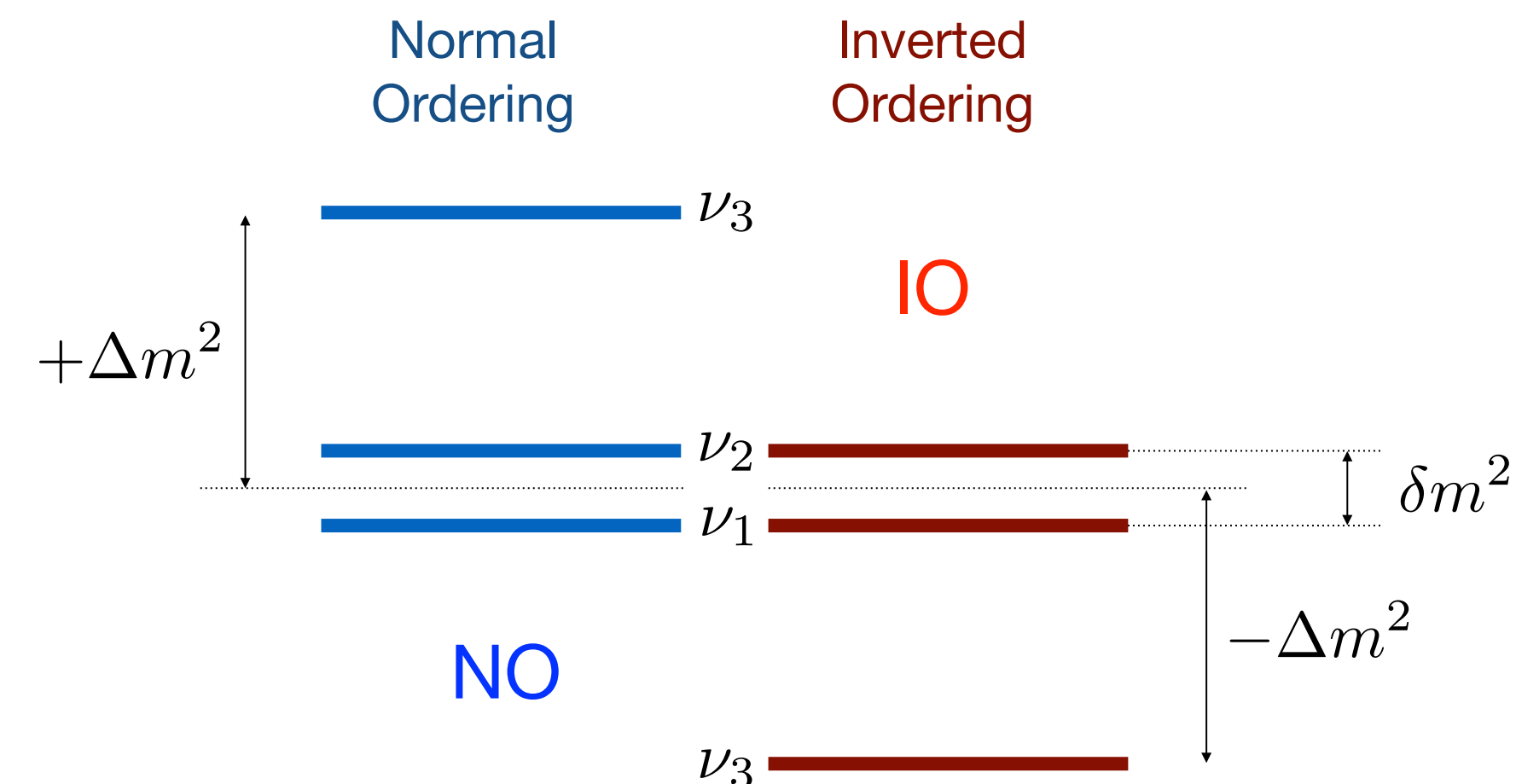
$$\sin^2 \theta_{23} \sim 0.473 \times 10^{-2} \quad (5.1\%)$$

$$\sin^2 \theta_{23} \sim 0.545 \times 10^{-2} \quad (4.3\%)$$



# Neutrino oscillation phenomenology: entering the precision era

Solar parameters		Atmospheric parameters	
$\delta m^2 \sim 7.37 \times 10^{-5} \text{ eV}^2$	(2.3%)	$\Delta m^2 \sim 2.495 \times 10^{-3} \text{ eV}^2$	(0.8%)
$\sin^2 \theta_{12} \sim 0.303$	(4.5%)	$\Delta m^2 \sim 2.465 \times 10^{-3} \text{ eV}^2$	
Reactor mixing angle		$\sin^2 \theta_{23} \sim 0.473 \times 10^{-2}$	(5.1%)
$\sin^2 \theta_{13} \sim 2.23 \times 10^{-2}$	(2.4%)	$\sin^2 \theta_{23} \sim 0.545 \times 10^{-2}$	(4.3%)

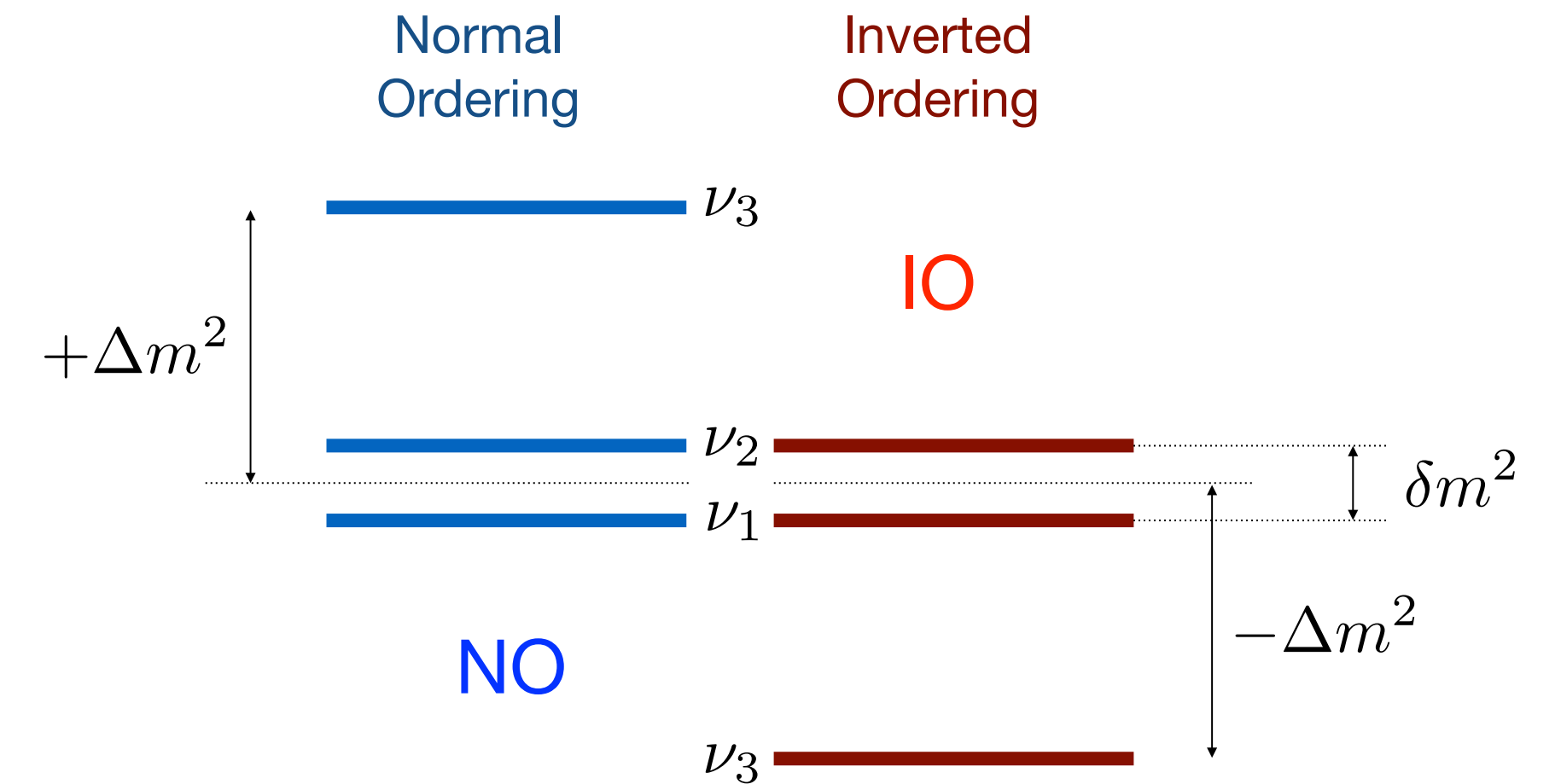


$$\Delta m^2 = \frac{\Delta m_{31}^2 + \Delta m_{32}^2}{2}$$



# Neutrino oscillation phenomenology: entering the precision era

Solar parameters		Atmospheric parameters	
$\delta m^2 \sim 7.37 \times 10^{-5} \text{ eV}^2$	(2.3%)	$\Delta m^2 \sim 2.495 \times 10^{-3} \text{ eV}^2$	(0.8%)
$\sin^2 \theta_{12} \sim 0.303$	(4.5%)	$\Delta m^2 \sim 2.465 \times 10^{-3} \text{ eV}^2$	
Reactor mixing angle		$\sin^2 \theta_{23} \sim 0.473 \times 10^{-2}$	(5.1%)
$\sin^2 \theta_{13} \sim 2.23 \times 10^{-2}$	(2.4%)	$\sin^2 \theta_{23} \sim 0.545 \times 10^{-2}$	(4.3%)



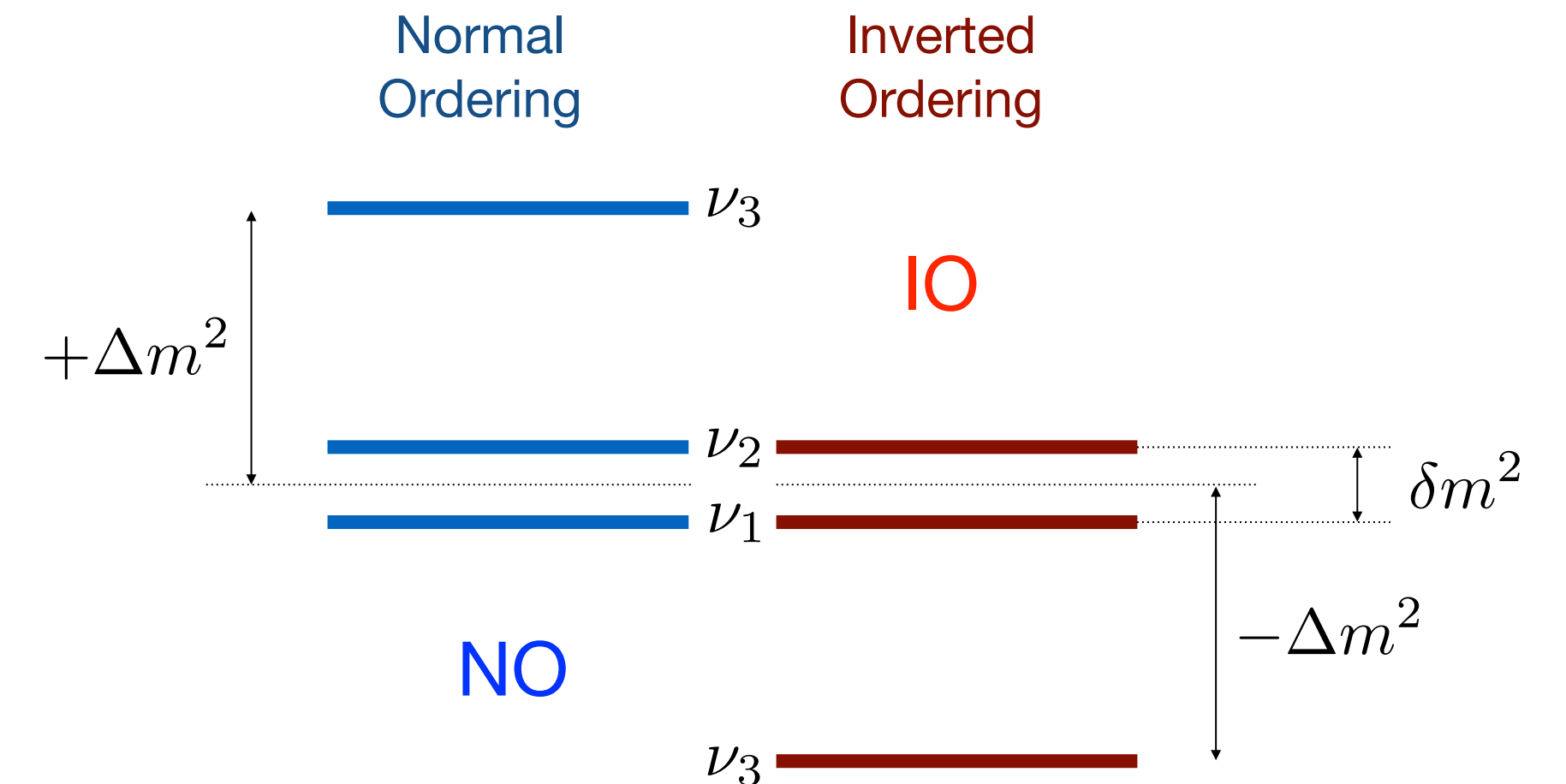
$$\Delta m^2 = \frac{\Delta m_{31}^2 + \Delta m_{32}^2}{2}$$

What is still  
Unknown



# Neutrino oscillation phenomenology: entering the precision era

Solar parameters		Atmospheric parameters	
$\delta m^2 \sim 7.37 \times 10^{-5} \text{ eV}^2$	(2.3%)	$\Delta m^2 \sim 2.495 \times 10^{-3} \text{ eV}^2$	(0.8%)
$\sin^2 \theta_{12} \sim 0.303$	(4.5%)	$\Delta m^2 \sim 2.465 \times 10^{-3} \text{ eV}^2$	
Reactor mixing angle		$\sin^2 \theta_{23} \sim 0.473 \times 10^{-2}$	(5.1%)
$\sin^2 \theta_{13} \sim 2.23 \times 10^{-2}$	(2.4%)	$\sin^2 \theta_{23} \sim 0.545 \times 10^{-2}$	(4.3%)



$$\Delta m^2 = \frac{\Delta m_{31}^2 + \Delta m_{32}^2}{2}$$

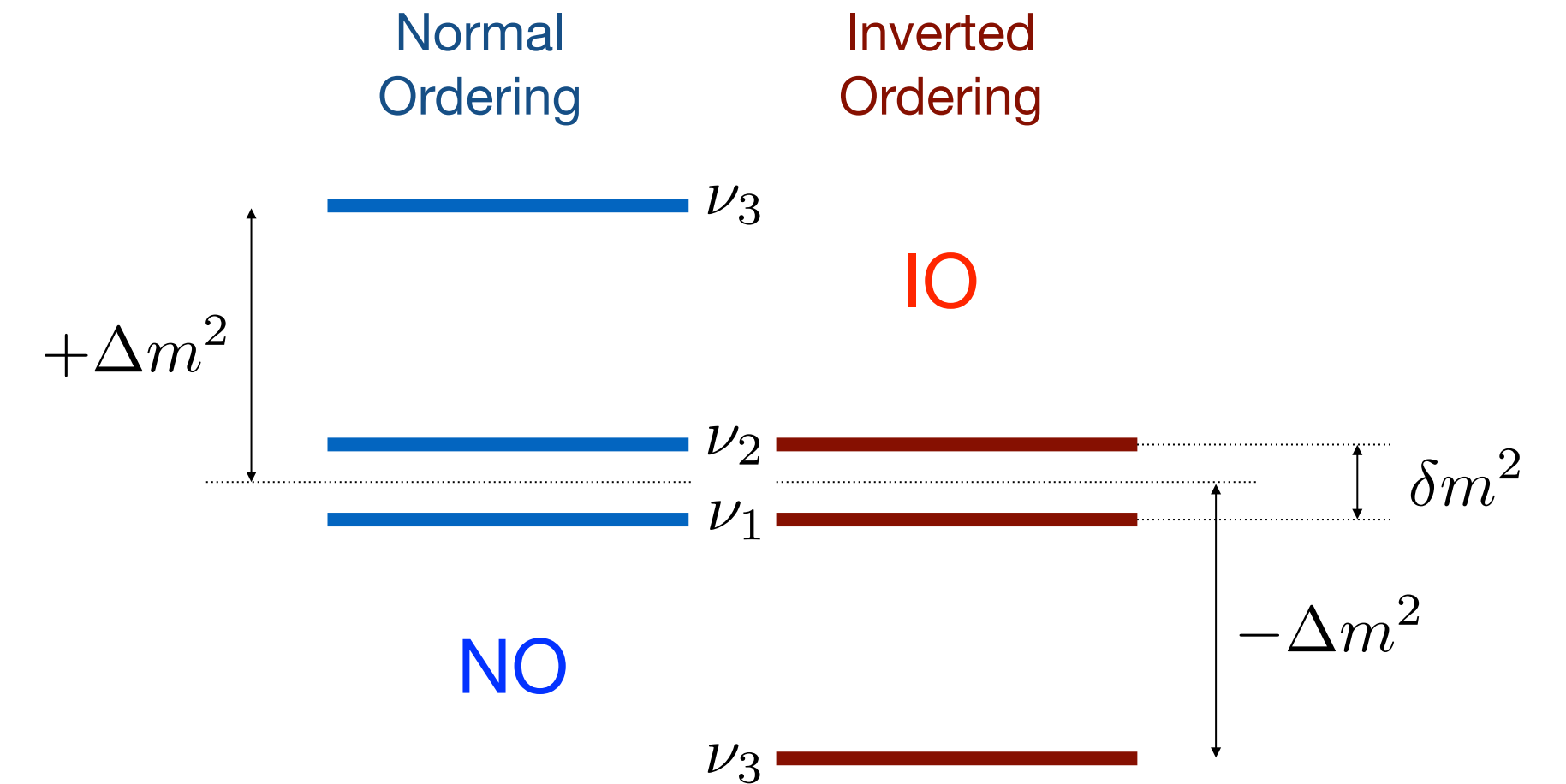
What is still  
Unknown

Mass Ordering  
CP-violating phase  $\delta_{CP}$   
Octant of  $\theta_{23}$



# Neutrino oscillation phenomenology: entering the precision era

Solar parameters		Atmospheric parameters	
$\delta m^2 \sim 7.37 \times 10^{-5} \text{ eV}^2$	(2.3%)	$\Delta m^2 \sim 2.495 \times 10^{-3} \text{ eV}^2$	(0.8%)
$\sin^2 \theta_{12} \sim 0.303$	(4.5%)	$\Delta m^2 \sim 2.465 \times 10^{-3} \text{ eV}^2$	
Reactor mixing angle		$\sin^2 \theta_{23} \sim 0.473 \times 10^{-2}$	(5.1%)
$\sin^2 \theta_{13} \sim 2.23 \times 10^{-2}$	(2.4%)	$\sin^2 \theta_{23} \sim 0.545 \times 10^{-2}$	(4.3%)



$$\Delta m^2 = \frac{\Delta m_{31}^2 + \Delta m_{32}^2}{2}$$

What is still  
Unknown

Mass Ordering  
CP-violating phase  $\delta_{CP}$   
Octant of  $\theta_{23}$

Absolute mass scale  
Nature of  $\nu$  (Dirac/Majorana)





# Methodological sequence for global oscillation analysis

Methodological sequence for global oscillation analysis



# Methodological sequence for global oscillation analysis

## Long Baseline Accelerator + Solar + KamLAND

minimal set sensitive to all oscillation parameters ( $\delta m^2$ ,  $\Delta m^2$ ,  $\theta_{12}$ ,  $\theta_{23}$ ,  $\theta_{13}$ ,  $\delta_{CP}$ ) and to mass ordering

# Methodological sequence for global oscillation analysis

## Long Baseline Accelerator + Solar + KamLAND

minimal set sensitive to all oscillation parameters ( $\delta m^2$ ,  $\Delta m^2$ ,  $\theta_{12}$ ,  $\theta_{23}$ ,  $\theta_{13}$ ,  $\delta_{CP}$ ) and to mass ordering

## Long Baseline Accelerator + Solar + KamLAND + Short Baseline Reactor

will add sensitivity to  $\Delta m^2$ ,  $\theta_{13}$  + correlations



# Methodological sequence for global oscillation analysis

## Long Baseline Accelerator + Solar + KamLAND

minimal set sensitive to all oscillation parameters ( $\delta m^2$ ,  $\Delta m^2$ ,  $\theta_{12}$ ,  $\theta_{23}$ ,  $\theta_{13}$ ,  $\delta_{CP}$ ) and to mass ordering

## Long Baseline Accelerator + Solar + KamLAND + Short Baseline Reactor

will add sensitivity to  $\Delta m^2$ ,  $\theta_{13}$  + correlations

## Long Baseline Acc. + Solar + KamLAND + Short Baseline Reactor + Atmospheric

will add sensitivity to  $\Delta m^2$ ,  $\theta_{23}$ ,  $\delta_{CP}$  and mass ordering

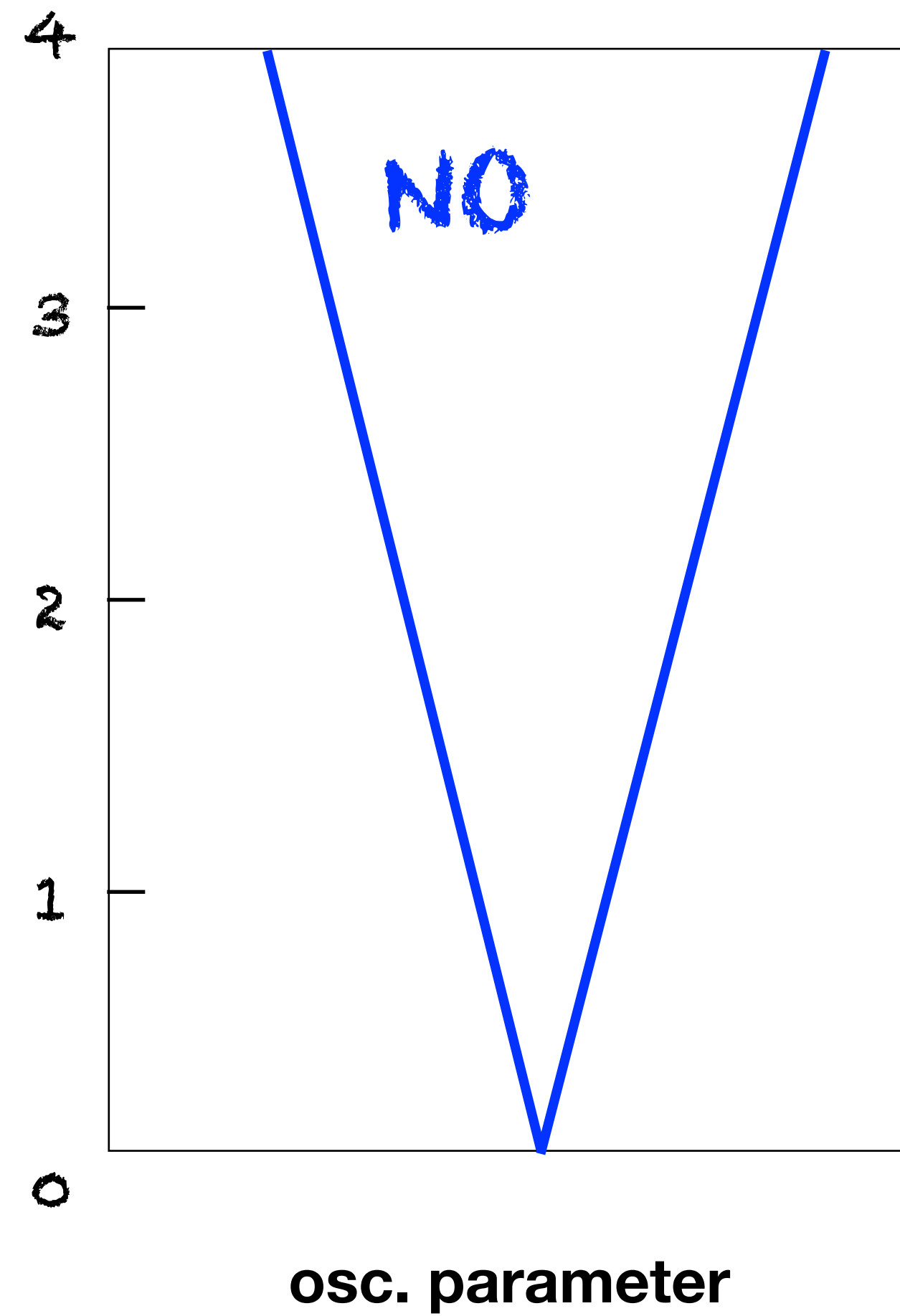




Bounds on single parameters, after marginalisation over all other parameters, shown in terms of  $N\sigma = \sqrt{\Delta\chi^2}$

Bounds on single parameters, after marginalisation over all other parameters, shown in terms of  $N\sigma = \sqrt{\Delta\chi^2}$

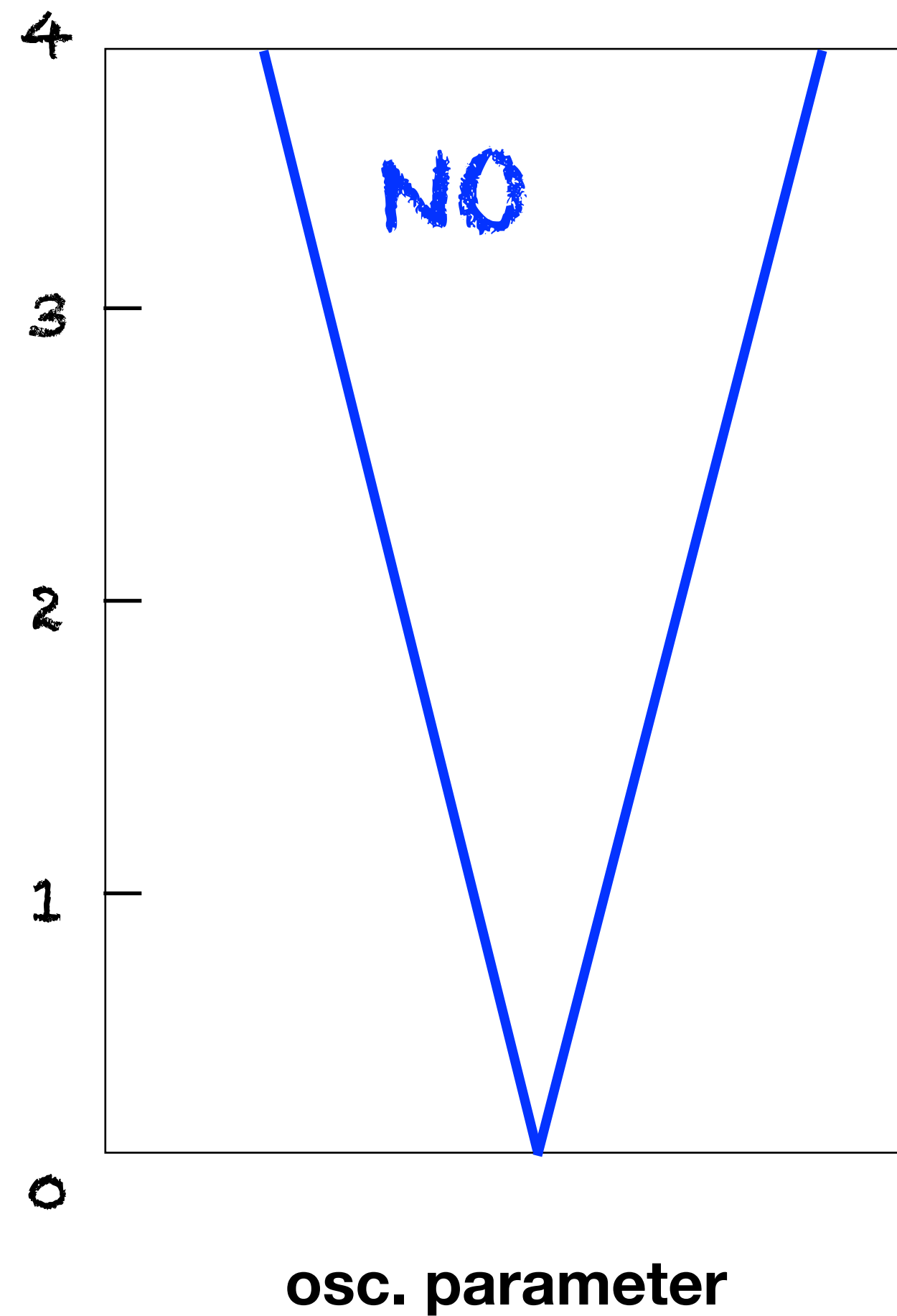
Bounds linear and symmetric for  
gaussian errors



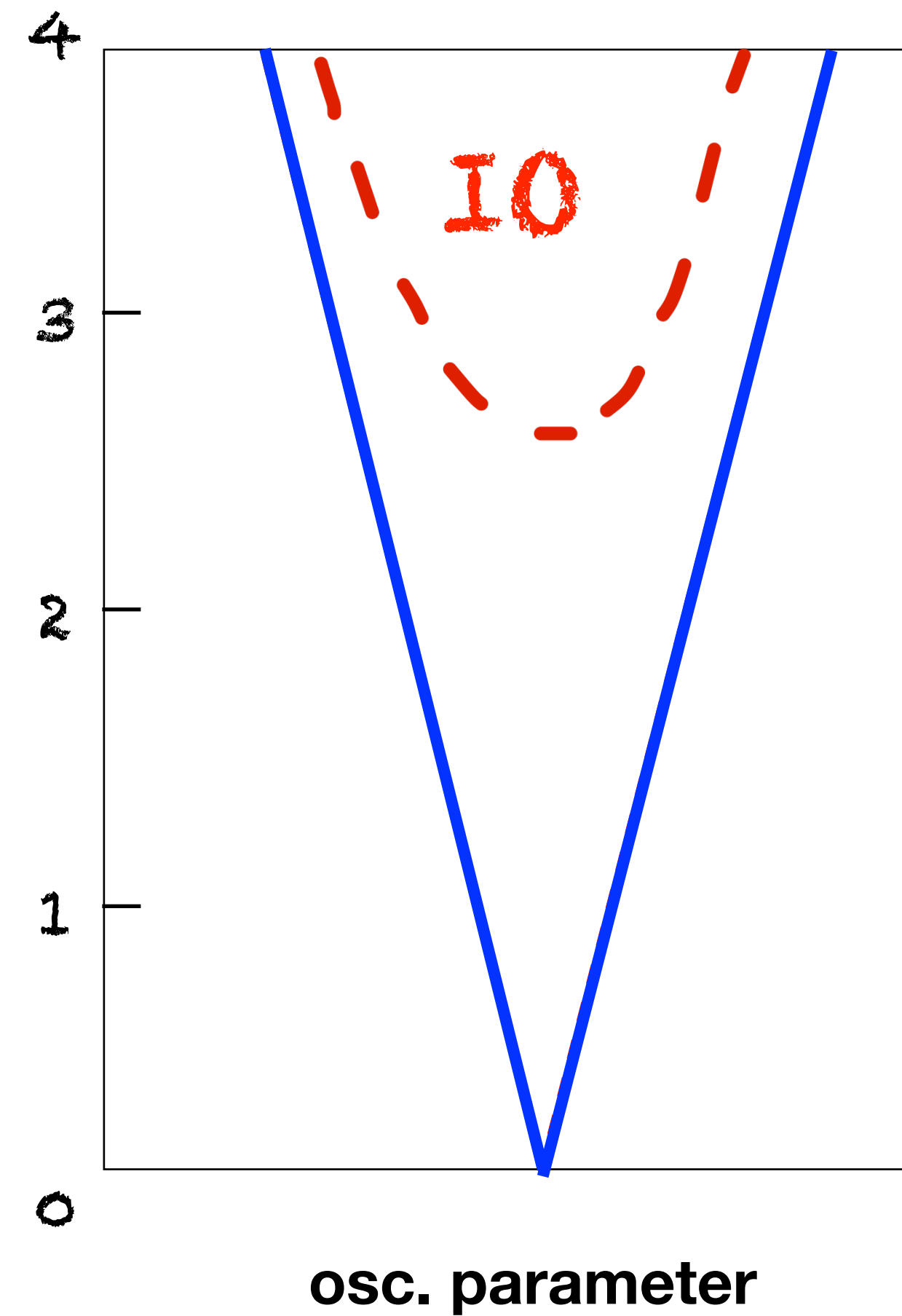


Bounds on single parameters, after marginalisation over all other parameters, shown in terms of  $N\sigma = \sqrt{\Delta\chi^2}$

Bounds linear and symmetric for gaussian errors

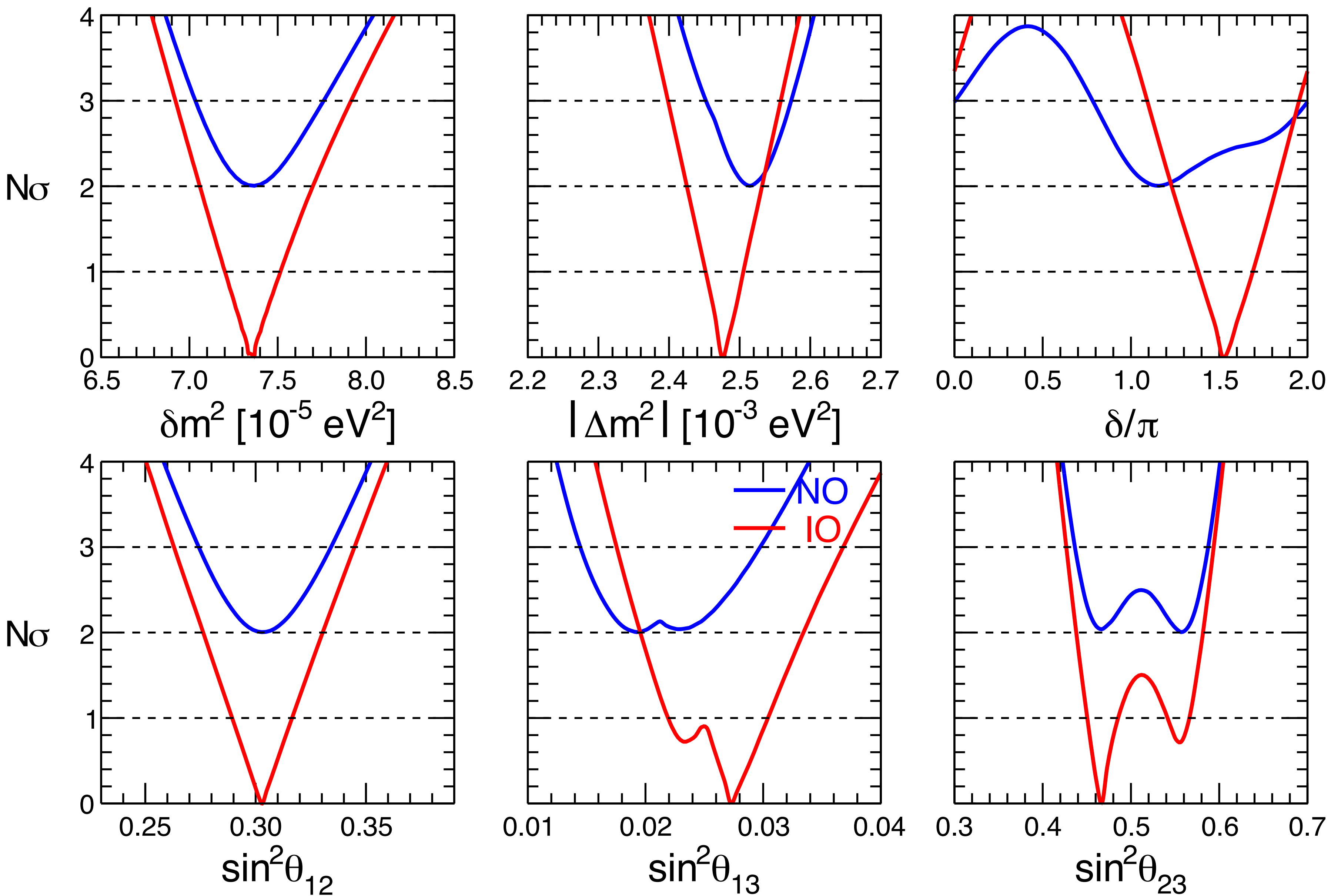


Bounds for IO move upwards (currently best fit in NO)



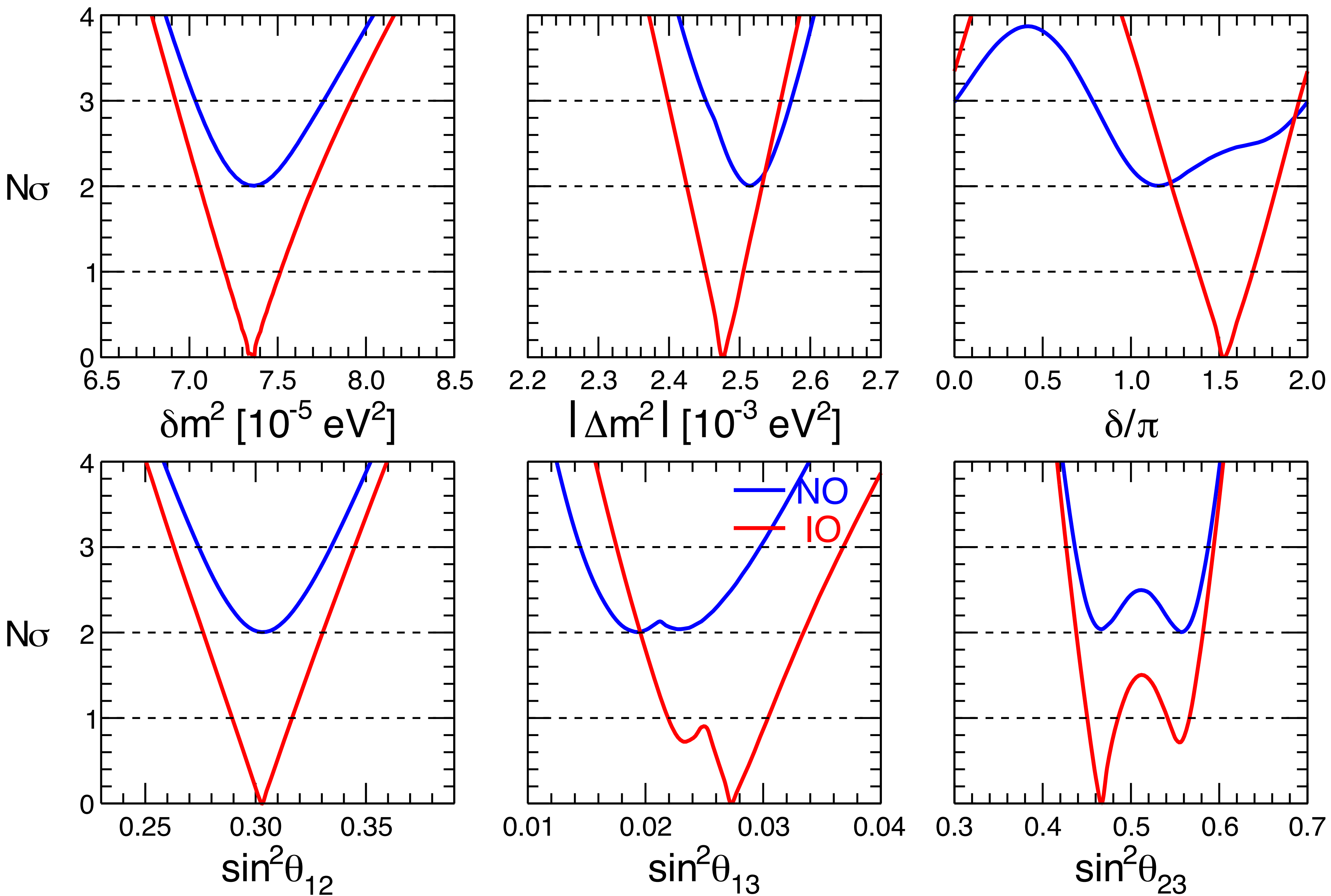


# LBL Acc + Solar + KamLAND



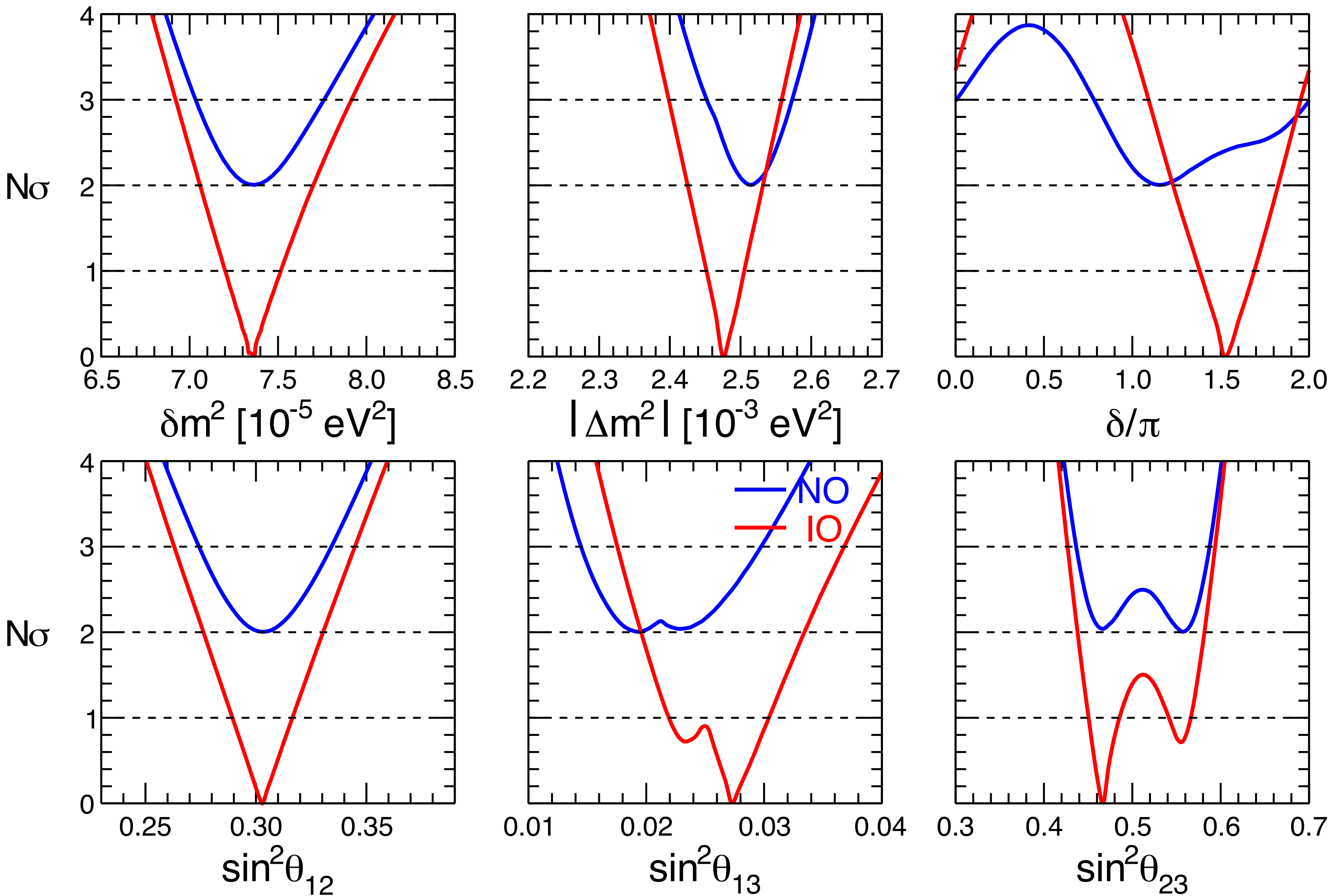


# LBL Acc + Solar + KamLAND



Gaussian errors for  
 $(\delta m^2, |\Delta m^2|, \sin^2 \theta_{12})$

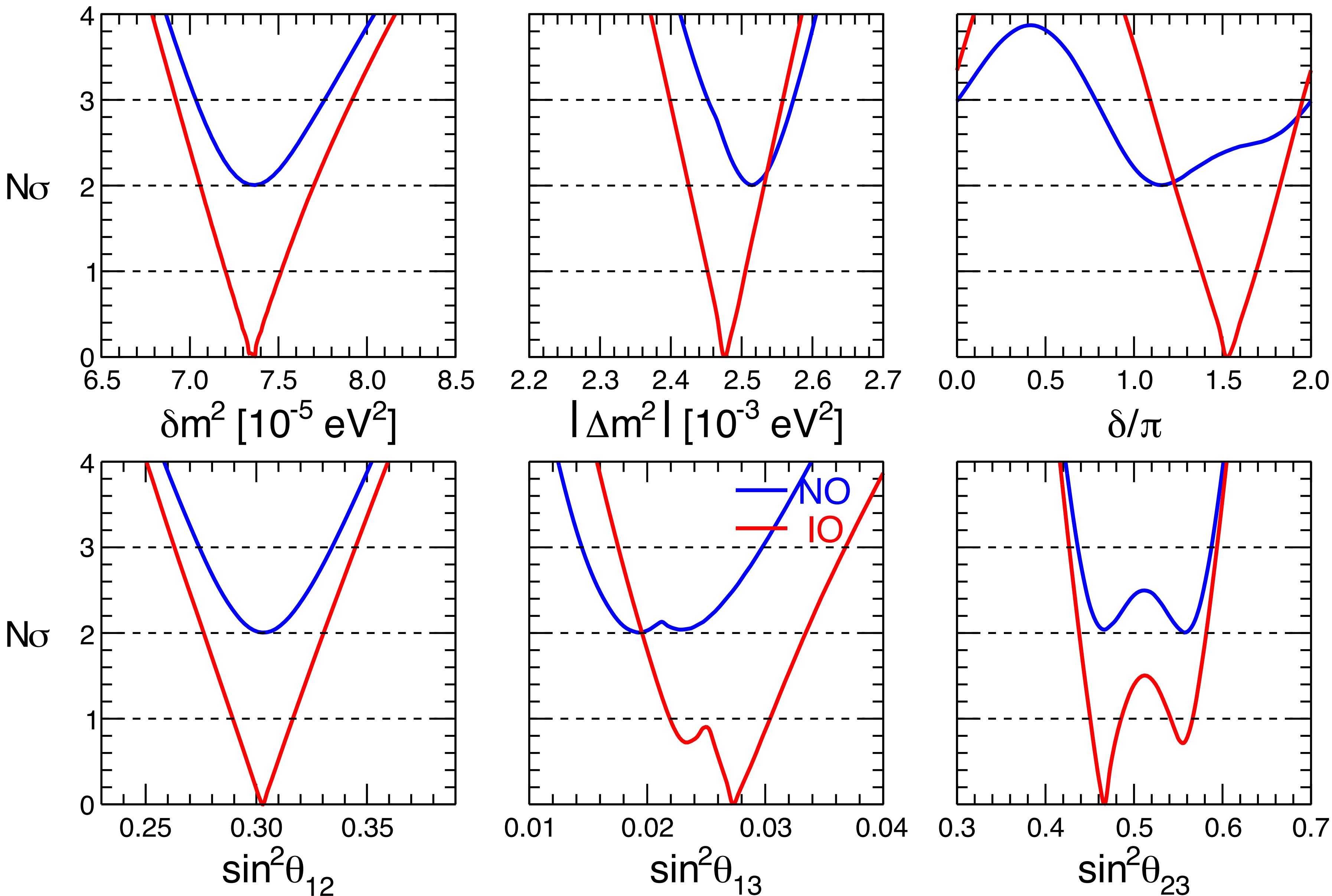
# LBL Acc + Solar + KamLAND



Gaussian errors for  
( $\delta m^2, |\Delta m^2|, \sin^2 \theta_{12}$ )

Two minima for  $\theta_{13}$  from  
some residual degeneracy

# LBL Acc + Solar + KamLAND



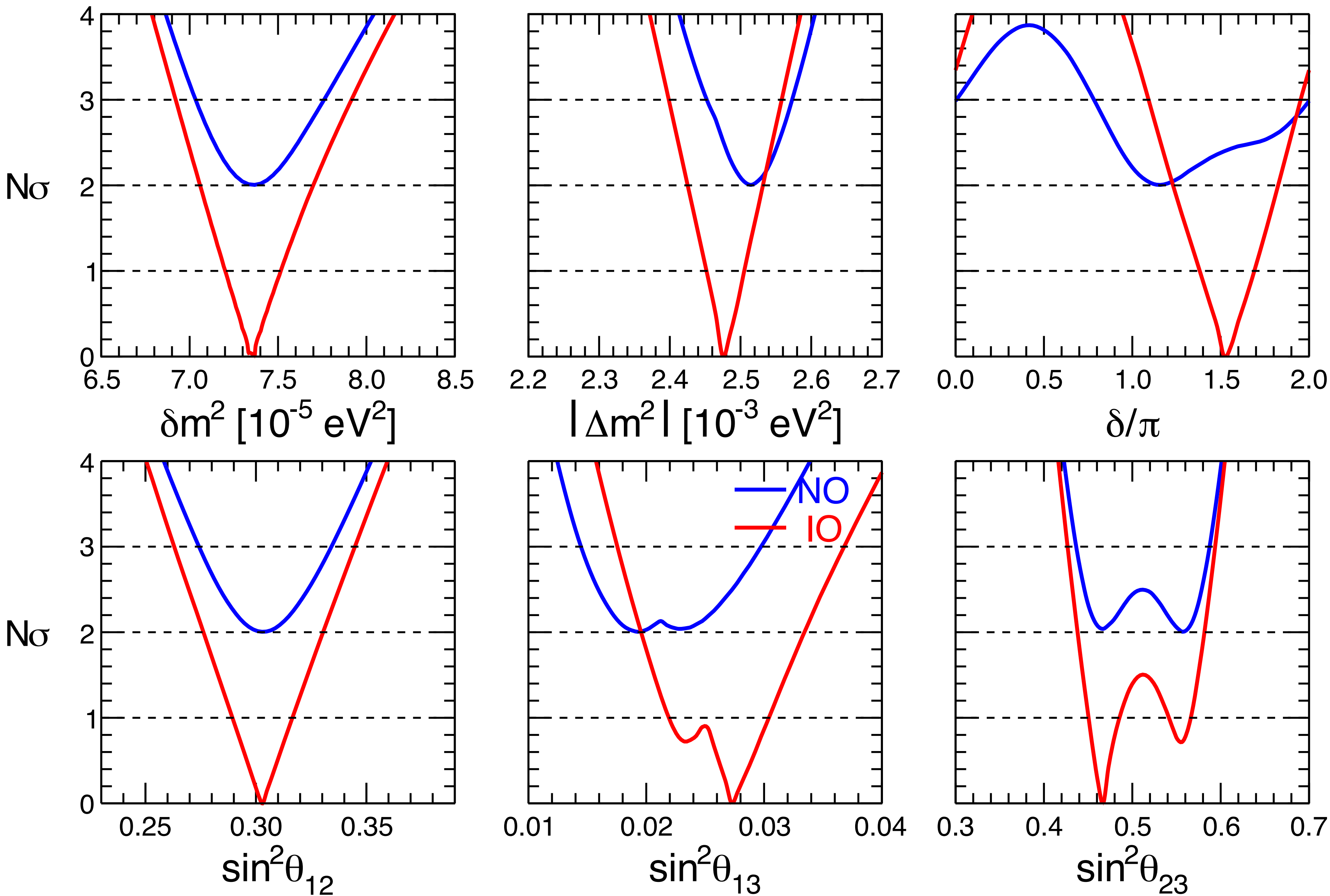
Gaussian errors for  
( $\delta m^2, |\Delta m^2|, \sin^2 \theta_{12}$ )

Two minima for  $\theta_{13}$  from  
some residual degeneracy

While T2K and NOvA  
individually point to NO,  
the combined dataset  
favors IO at  $\sim 2\sigma$ .



# LBL Acc + Solar + KamLAND



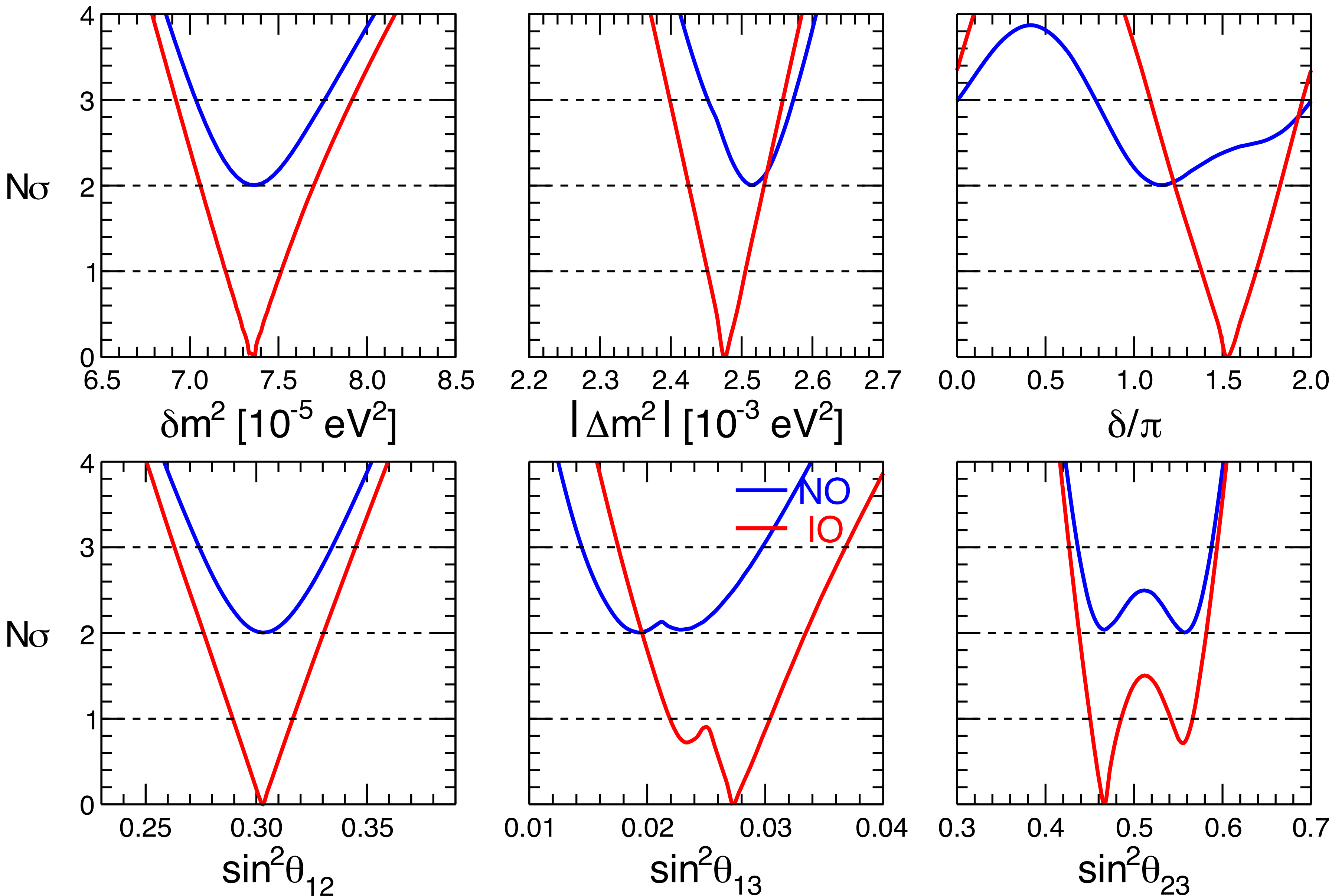
Gaussian errors for  
( $\delta m^2, |\Delta m^2|, \sin^2 \theta_{12}$ )

Two minima for  $\theta_{13}$  from  
some residual degeneracy

While T2K and NOvA  
individually point to NO,  
the combined dataset  
favors IO at  $\sim 2\sigma$ .

Octant ambiguity for  $\theta_{23}$

# LBL Acc + Solar + KamLAND



Gaussian errors for  
( $\delta m^2, |\Delta m^2|, \sin^2 \theta_{12}$ )

Two minima for  $\theta_{13}$  from  
some residual degeneracy

While T2K and NOvA  
individually point to NO,  
the combined dataset  
favors IO at  $\sim 2\sigma$ .

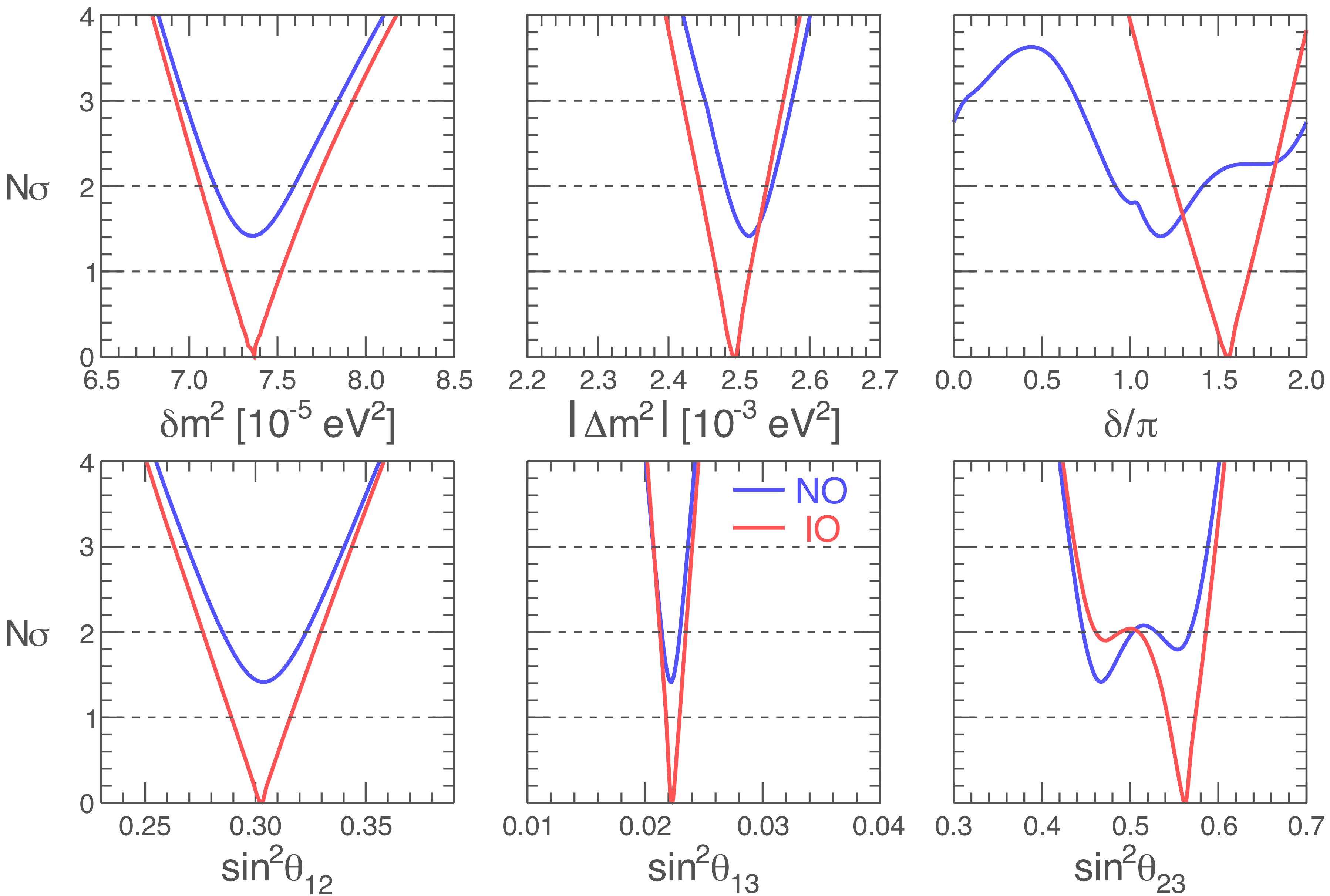
Octant ambiguity for  $\theta_{23}$

Due to some tension,  
in IO, indications for  
CP violation  $> 3\sigma$

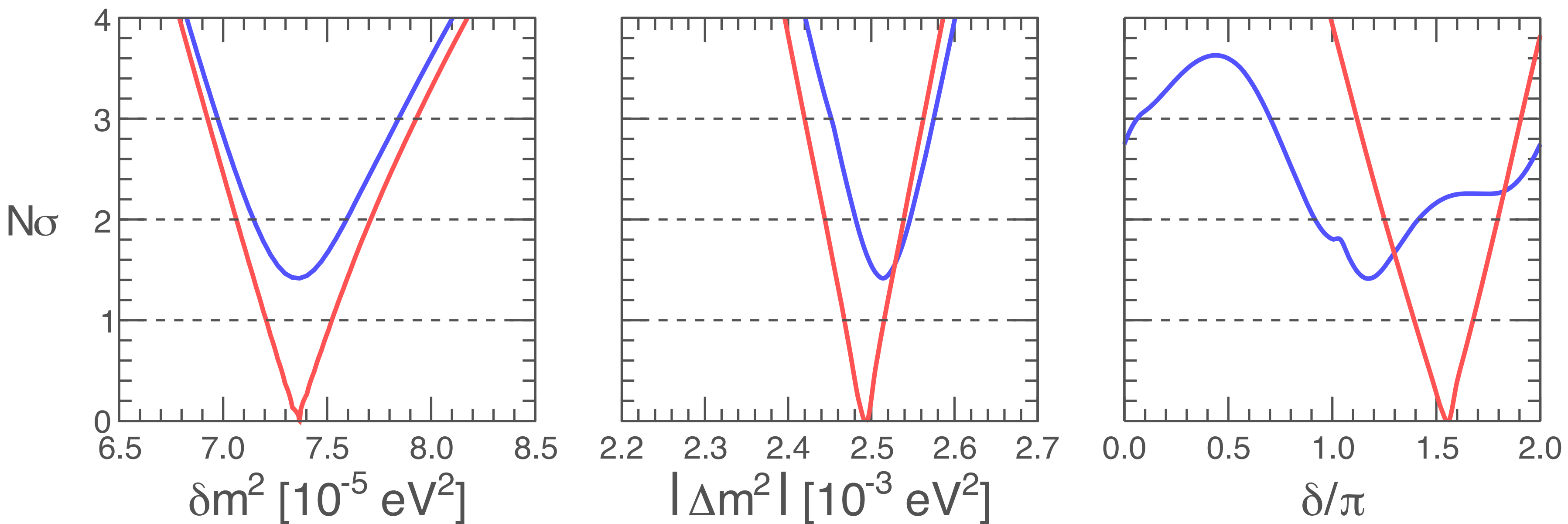




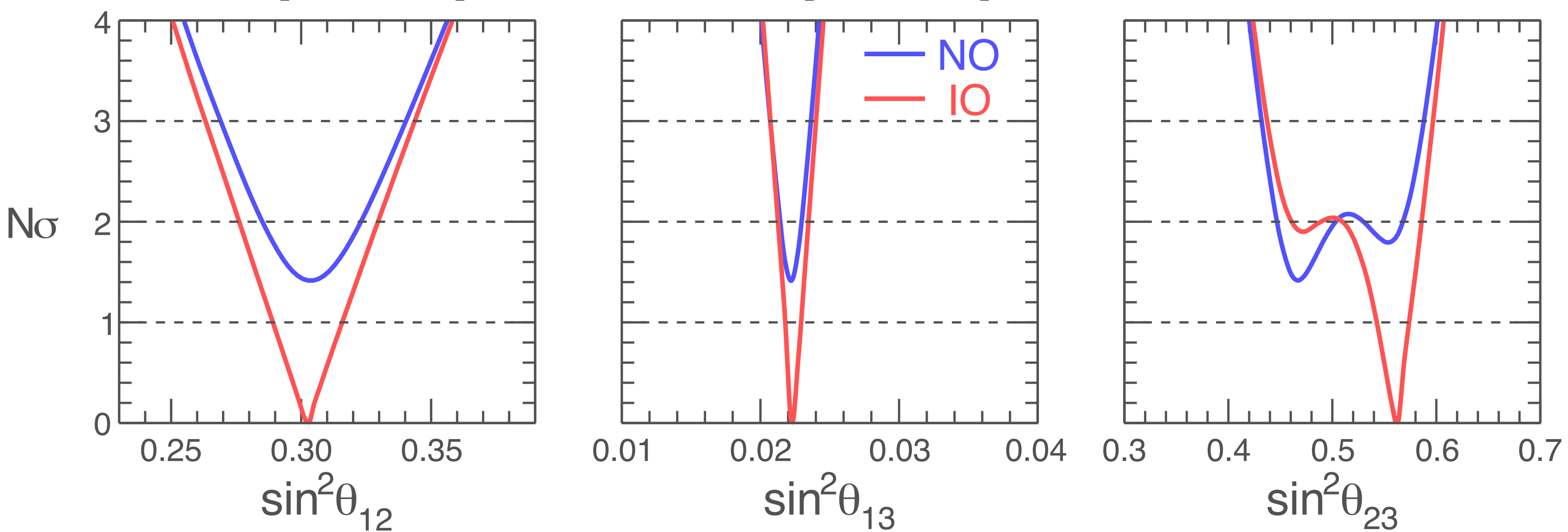
LBL Acc + Solar + KamLAND + SBL Reactors



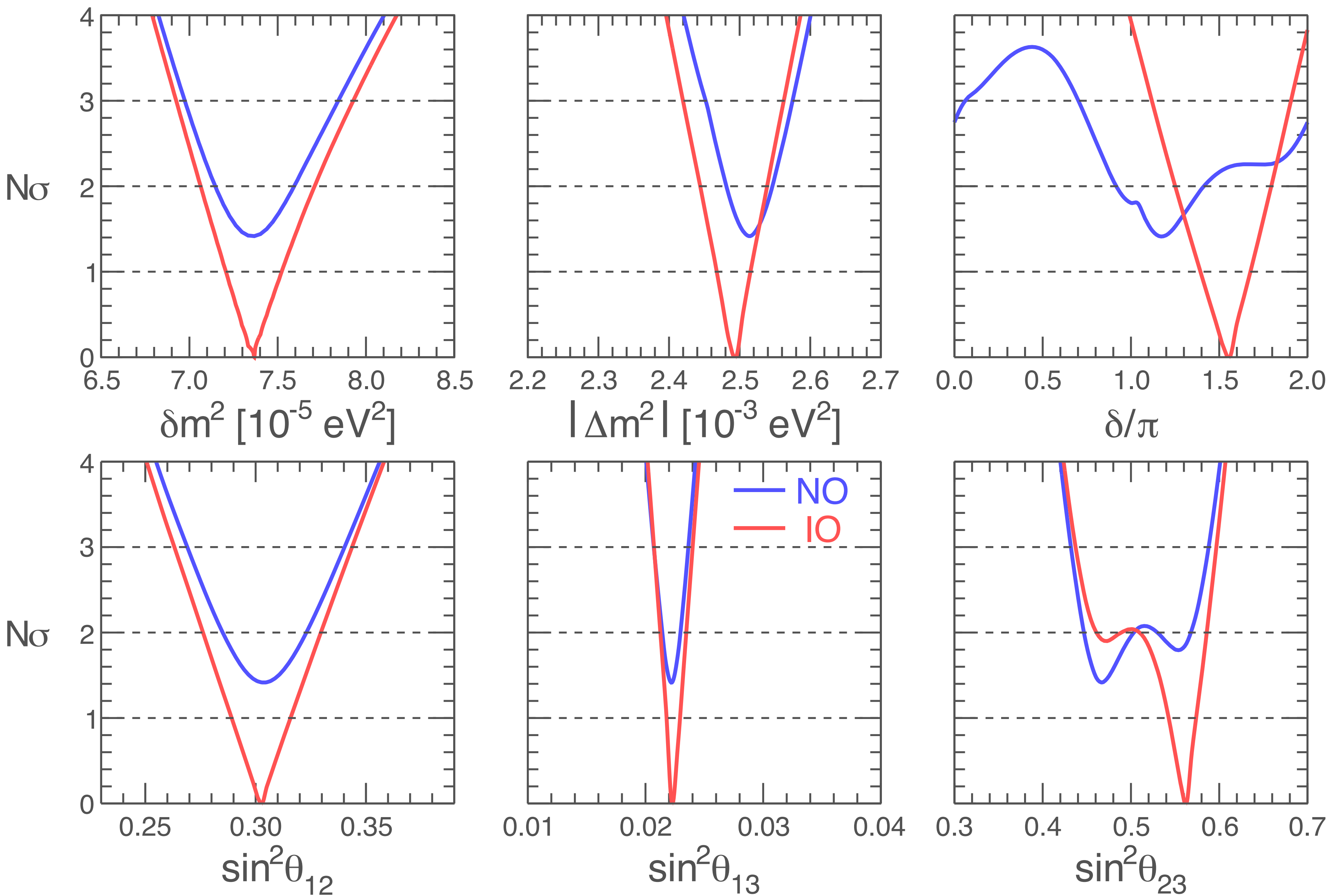
# LBL Acc + Solar + KamLAND + SBL Reactors



Adding SBL reactors:



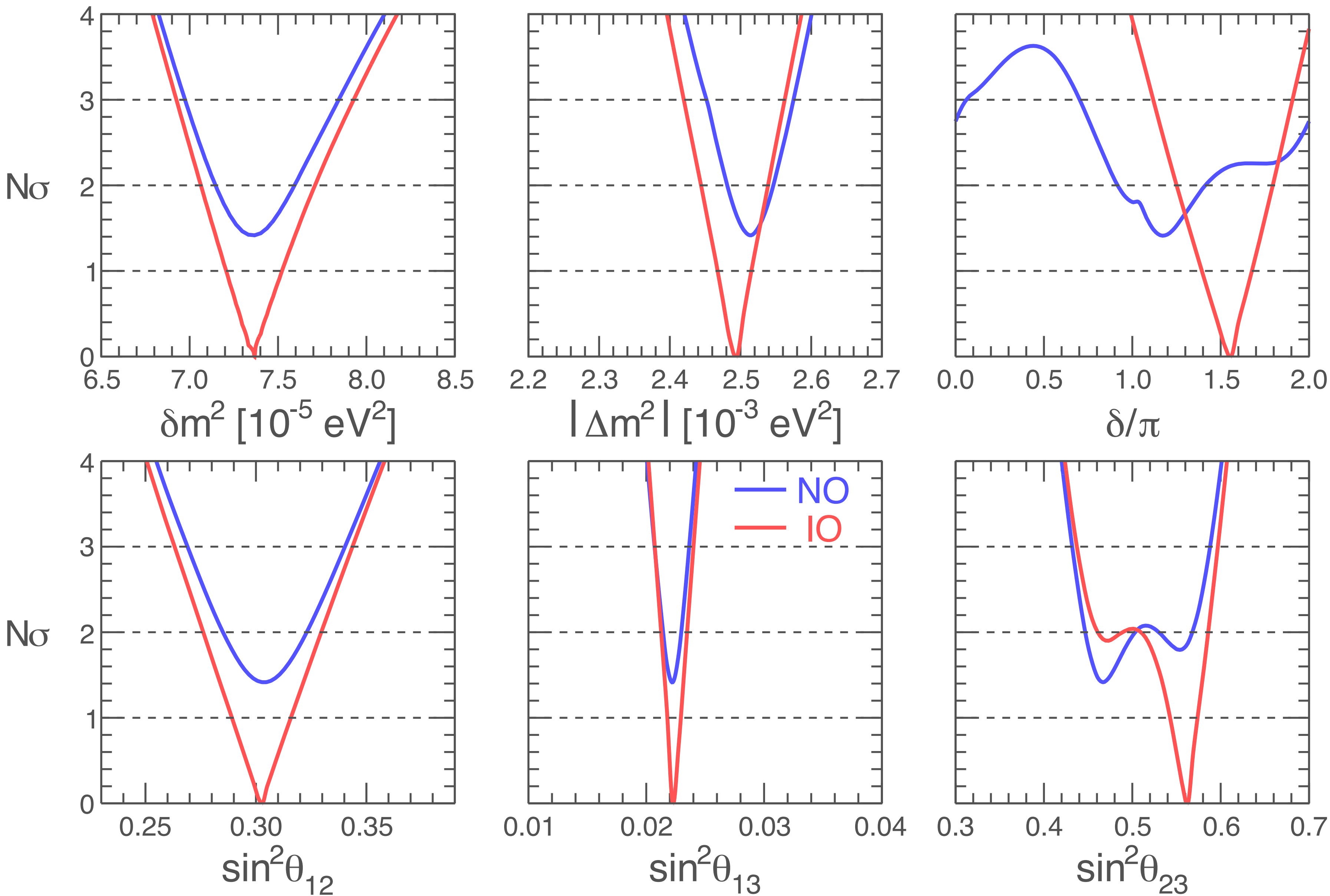
# LBL Acc + Solar + KamLAND + SBL Reactors



Adding SBL reactors:

Lower value of  $\theta_{13}$   
preferred

# LBL Acc + Solar + KamLAND + SBL Reactors



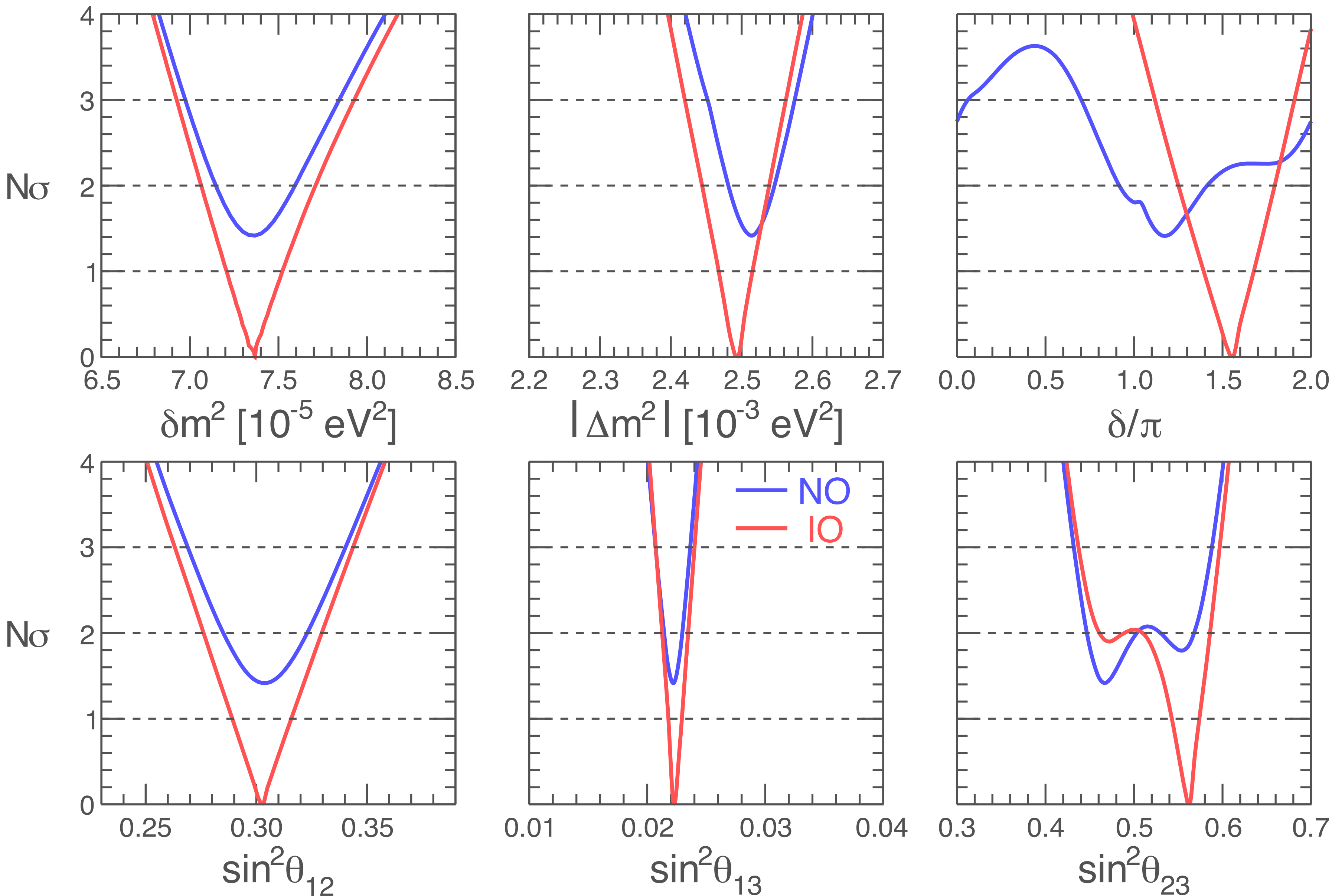
Adding SBL reactors:

Lower value of  $\theta_{13}$   
preferred

Octant ambiguity for  
 $\theta_{23}$  decreased, but  
still at  $\sim 2\sigma$  for IO



# LBL Acc + Solar + KamLAND + SBL Reactors



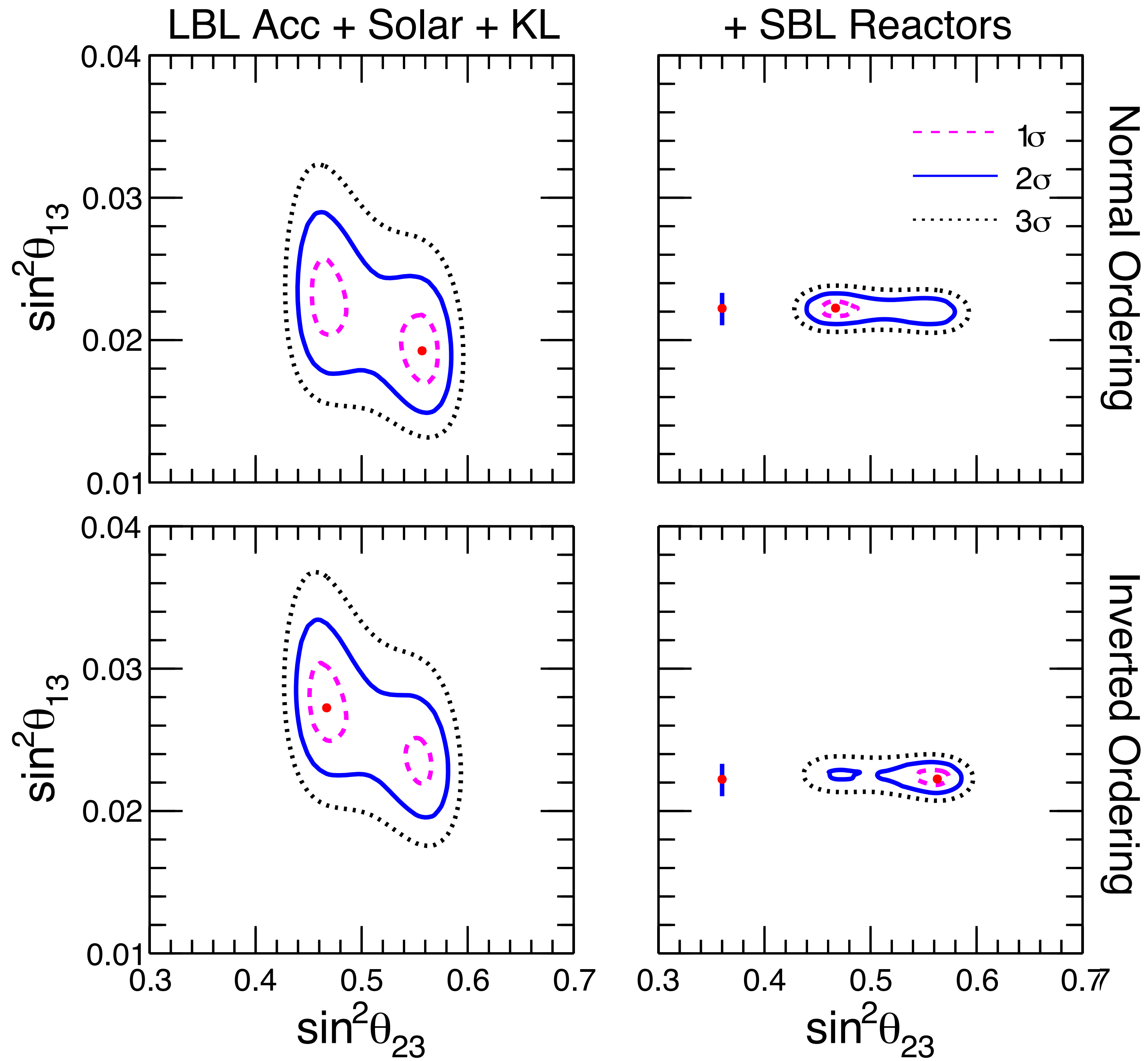
Adding SBL reactors:

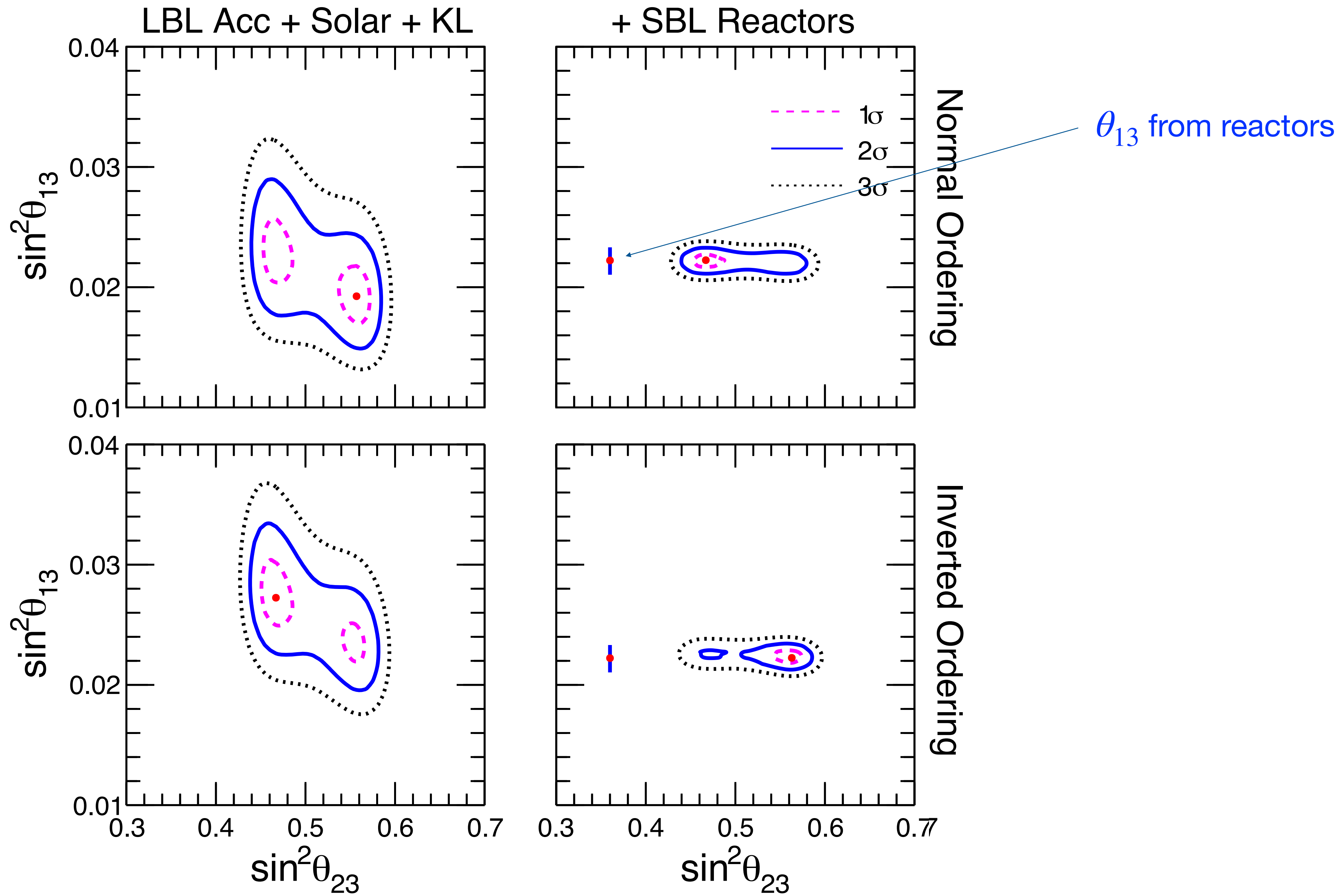
Lower value of  $\theta_{13}$   
preferred

Octant ambiguity for  
 $\theta_{23}$  decreased, but  
still at  $\sim 2\sigma$  for IO

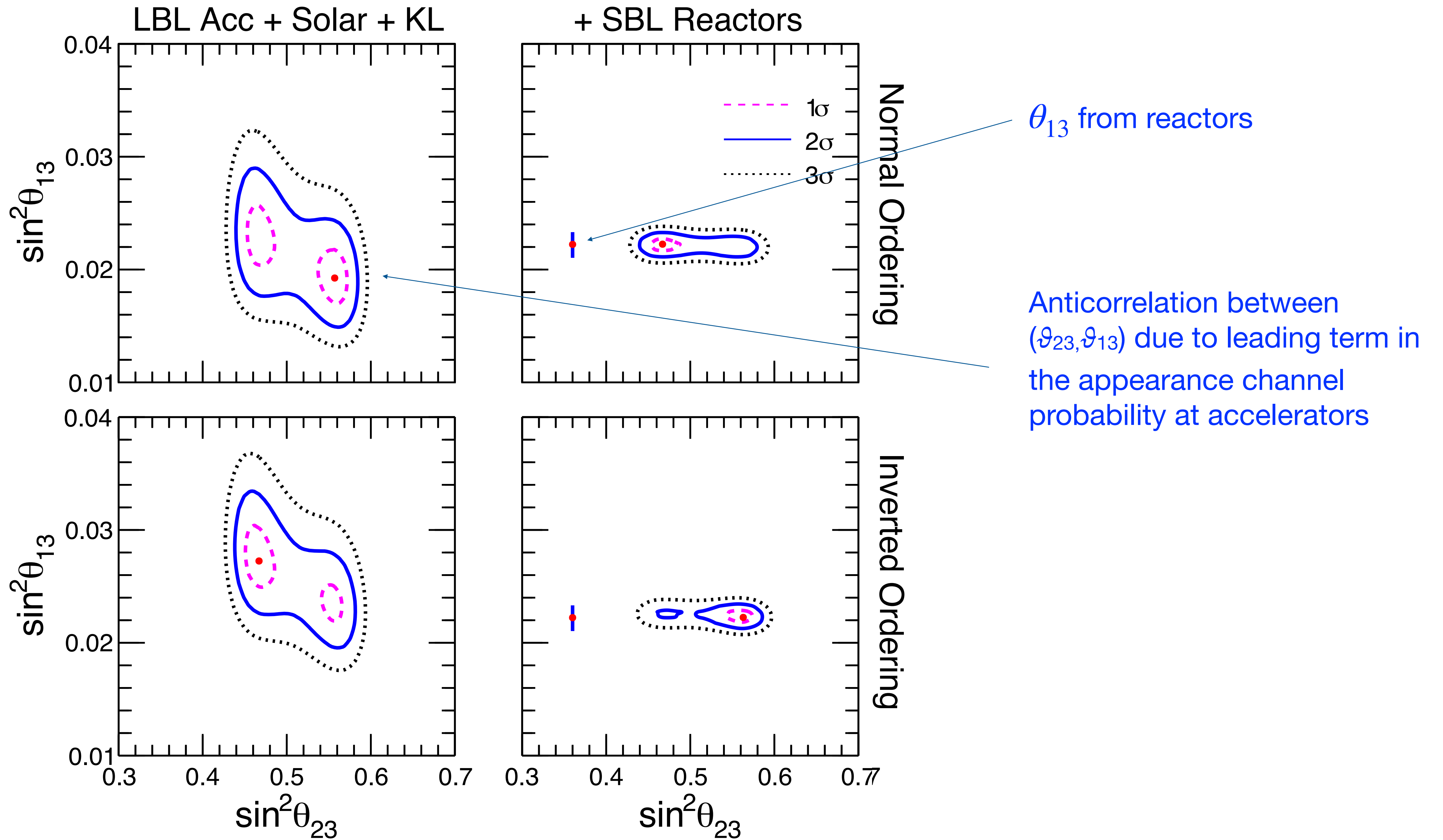
Adding SBL reactors:  
Still preference for IO, but  
at lower CL ( $\sim 1.4\sigma$ )

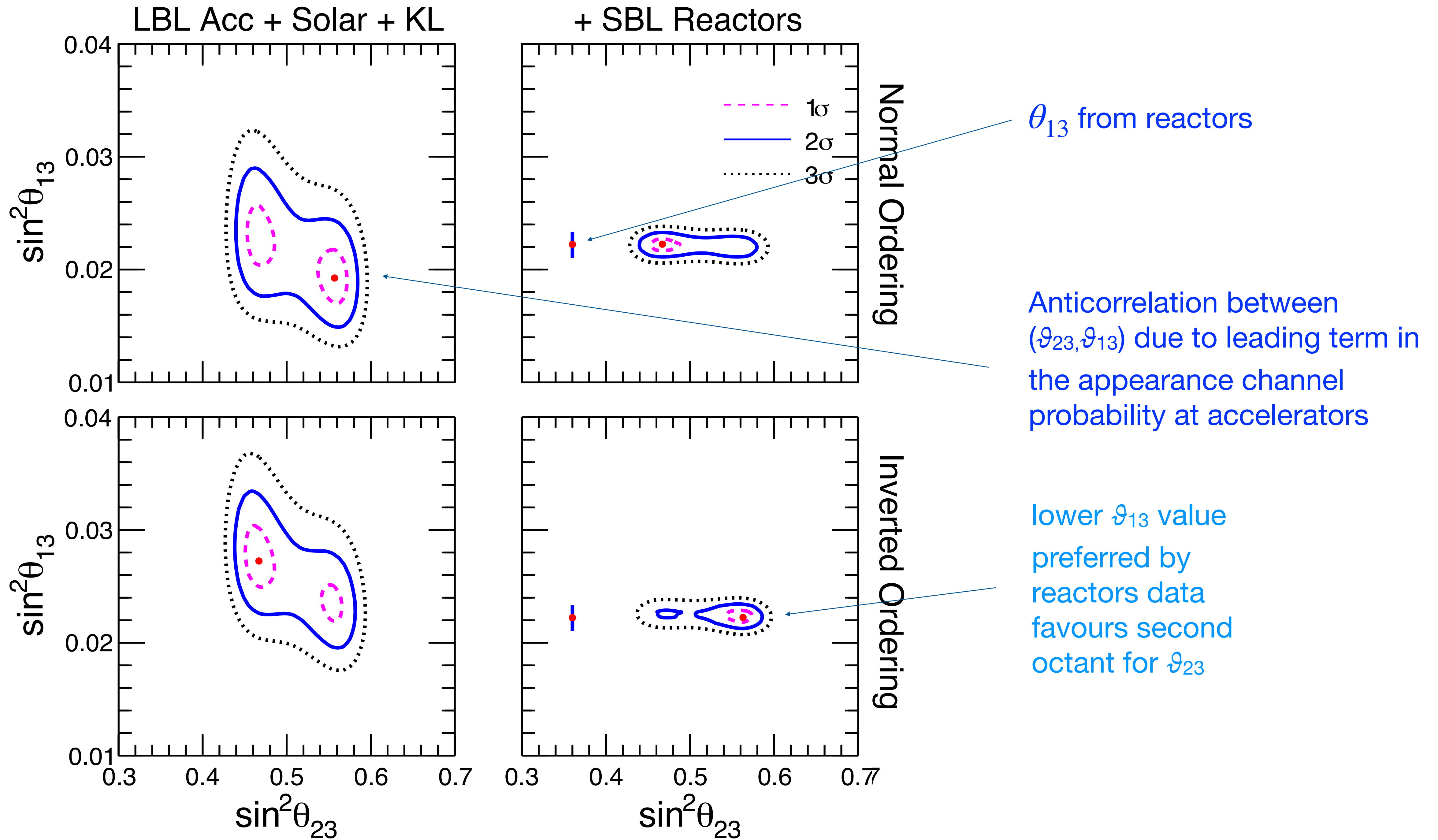








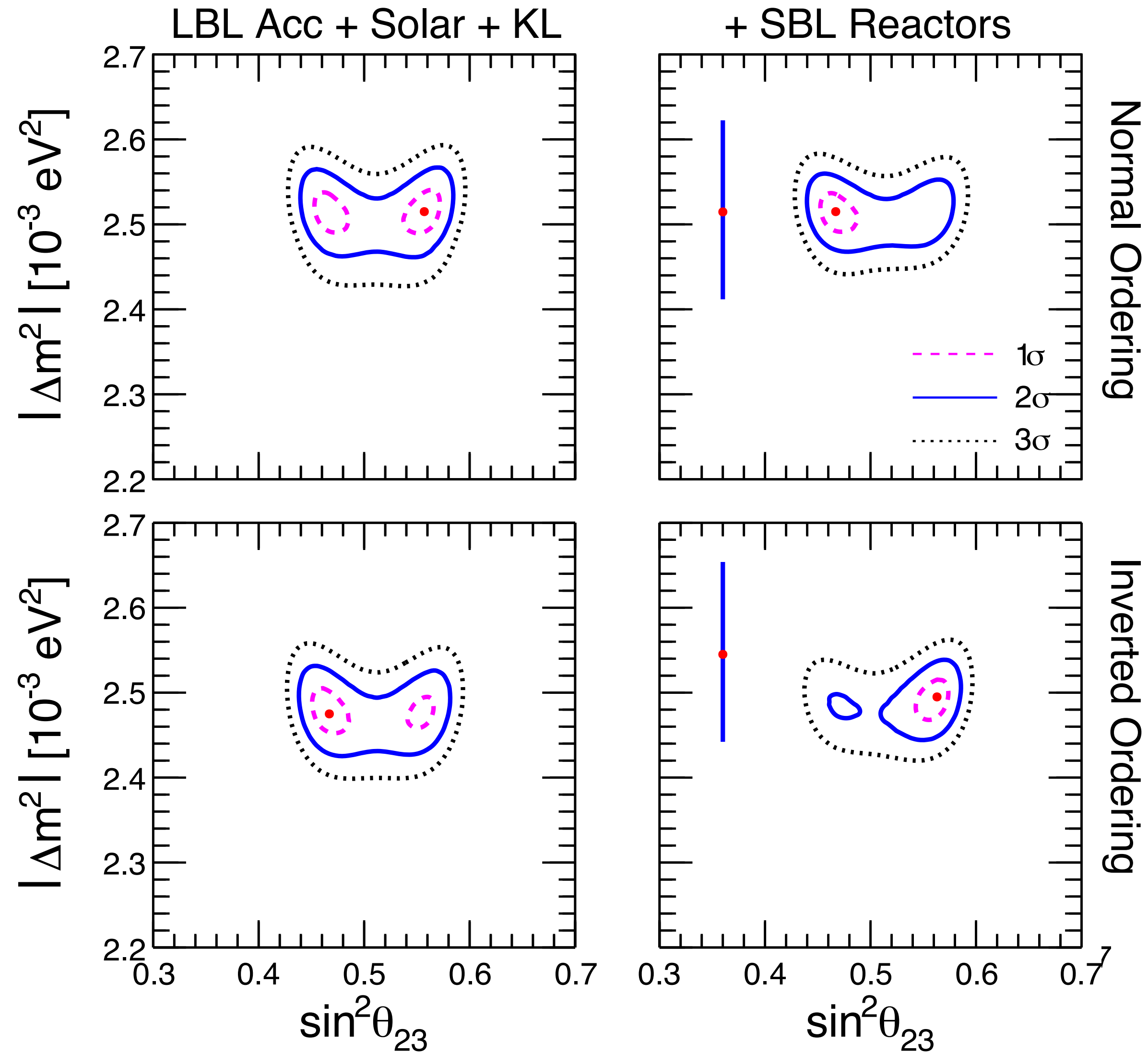






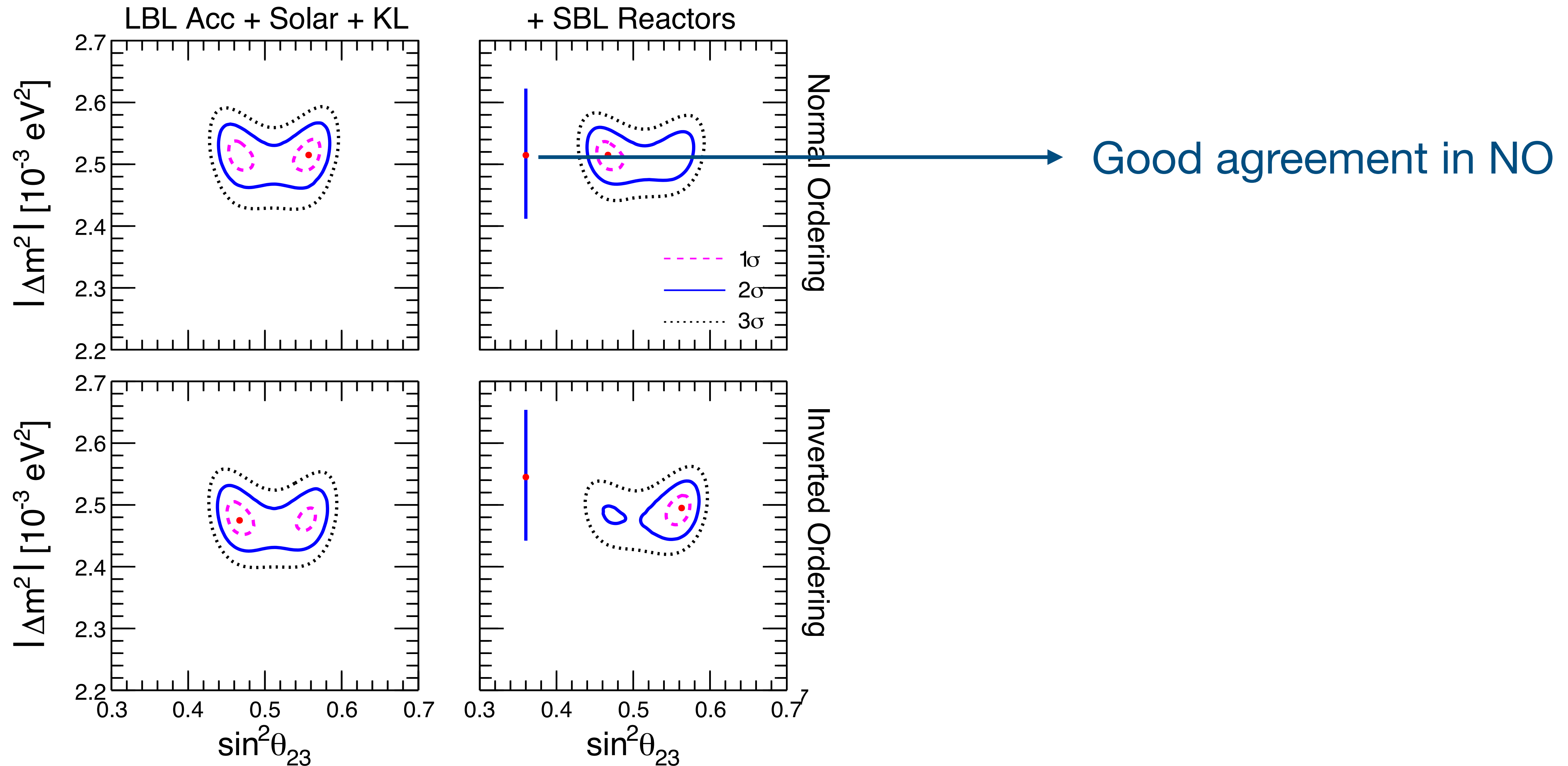
SBL reactor measurement of  $\Delta m^2$  more in agreement with LBL accel. in NO than in IO

# SBL reactor measurement of $\Delta m^2$ more in agreement with LBL accel. in NO than in IO

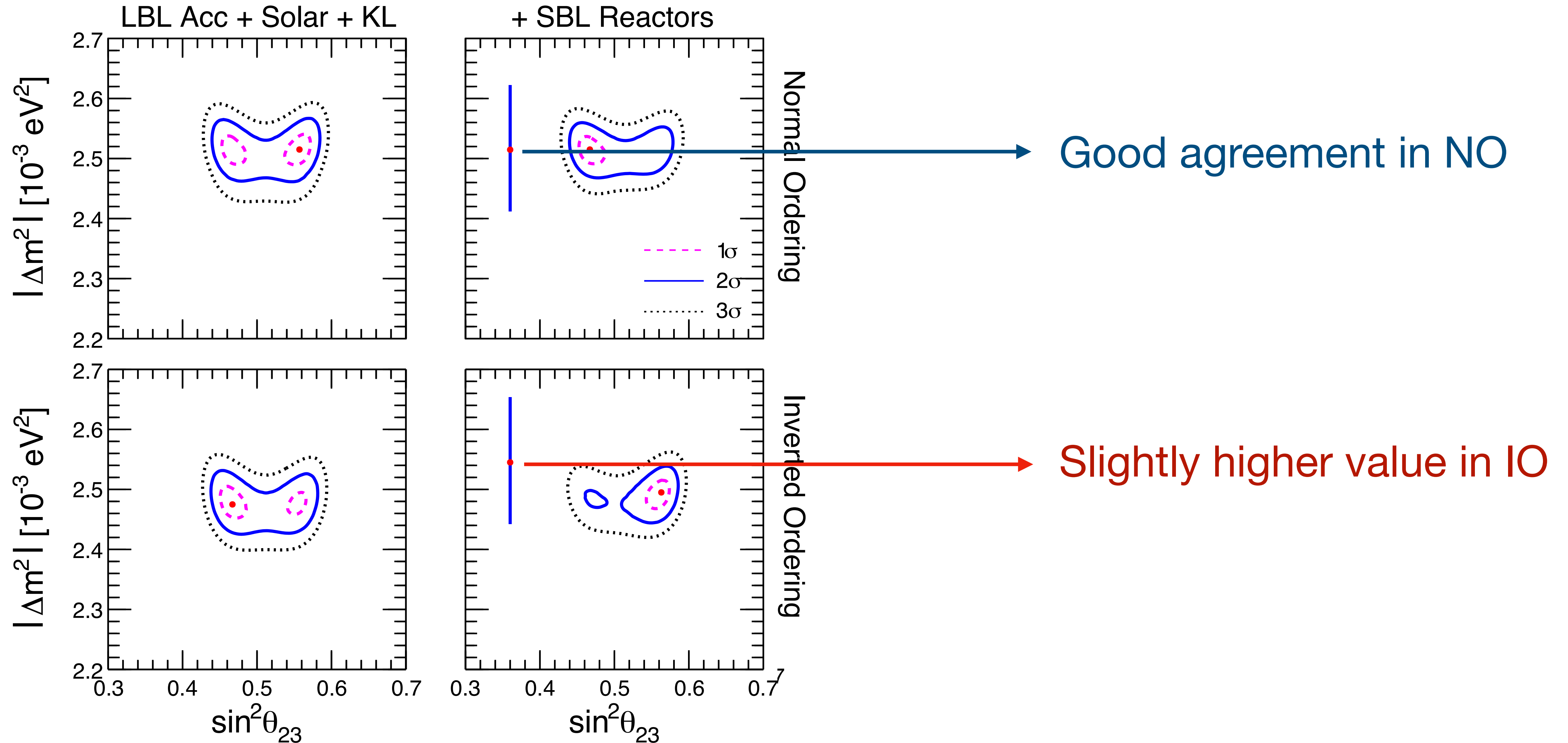




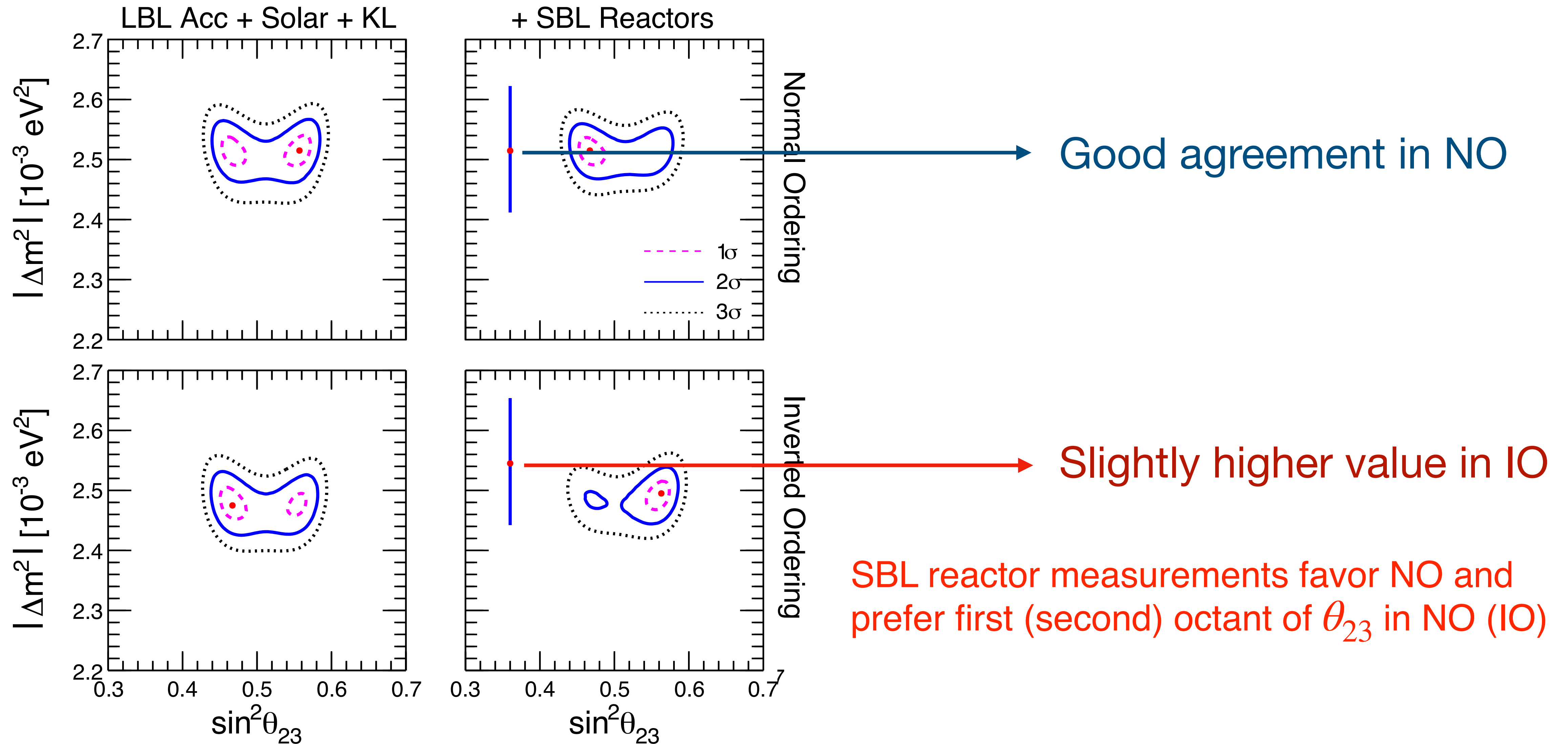
# SBL reactor measurement of $\Delta m^2$ more in agreement with LBL accel. in NO than in IO



# SBL reactor measurement of $\Delta m^2$ more in agreement with LBL accel. in NO than in IO

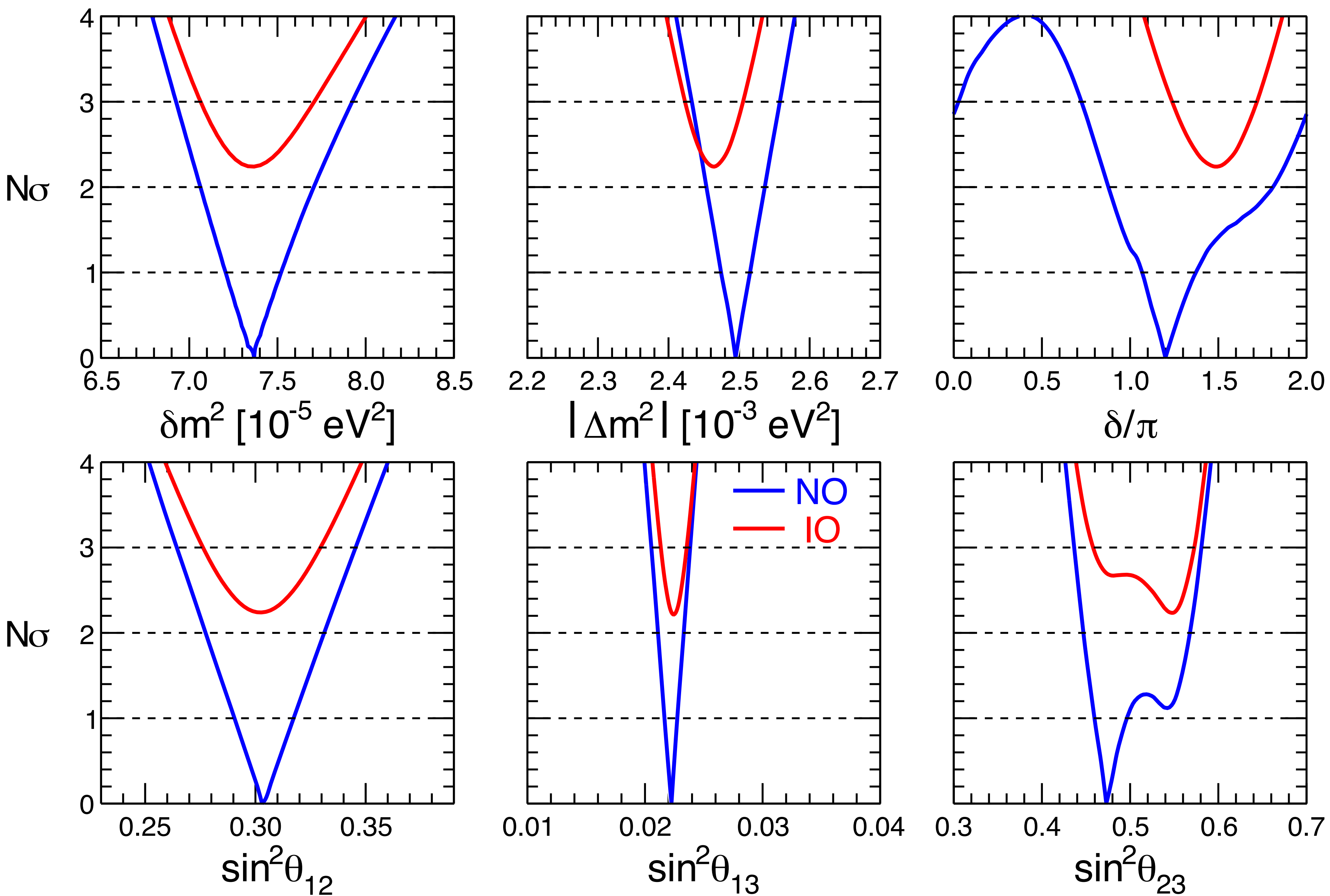


# SBL reactor measurement of $\Delta m^2$ more in agreement with LBL accel. in NO than in IO



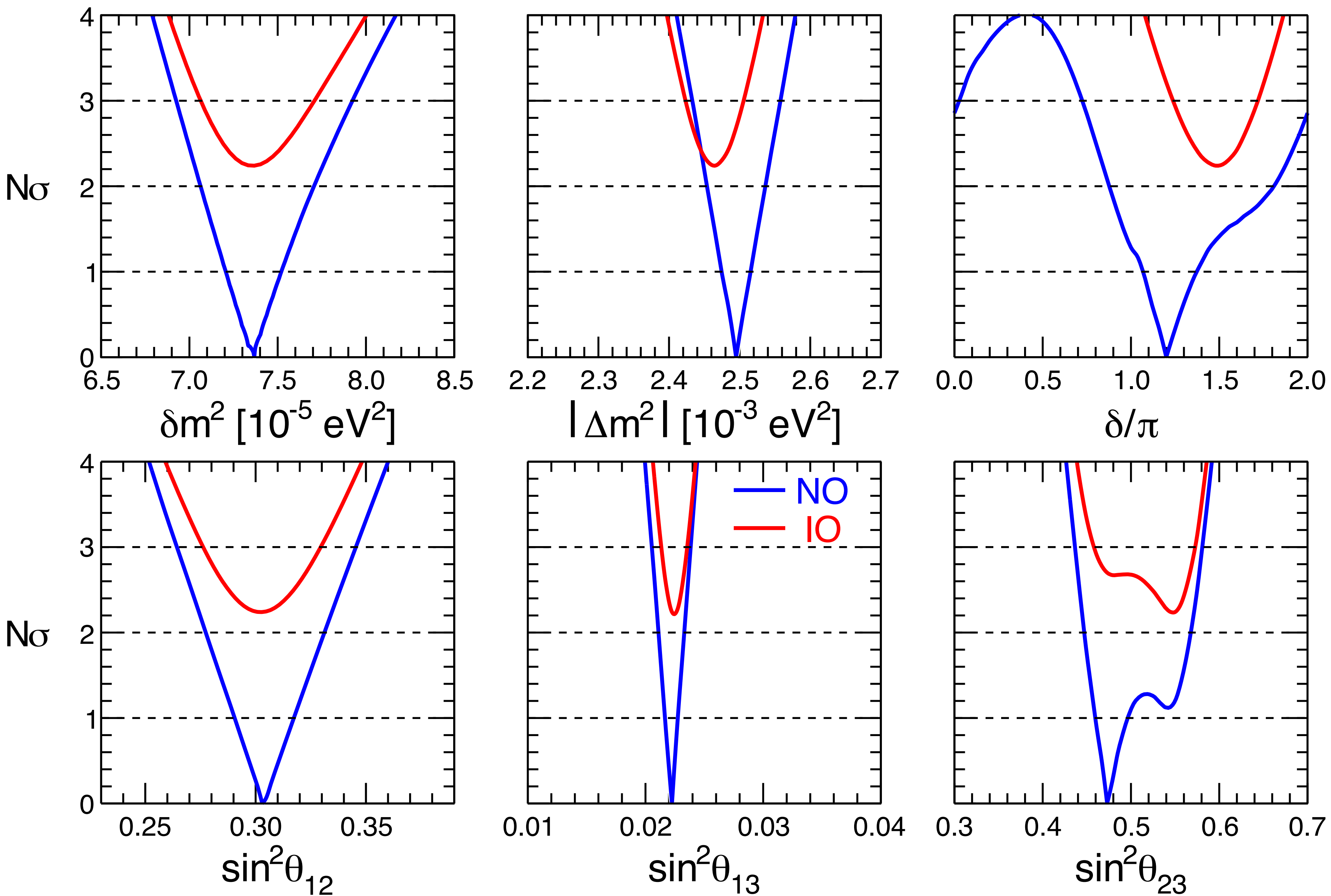


LBL Acc + Solar + KamLAND + SBL Reactors + Atmos



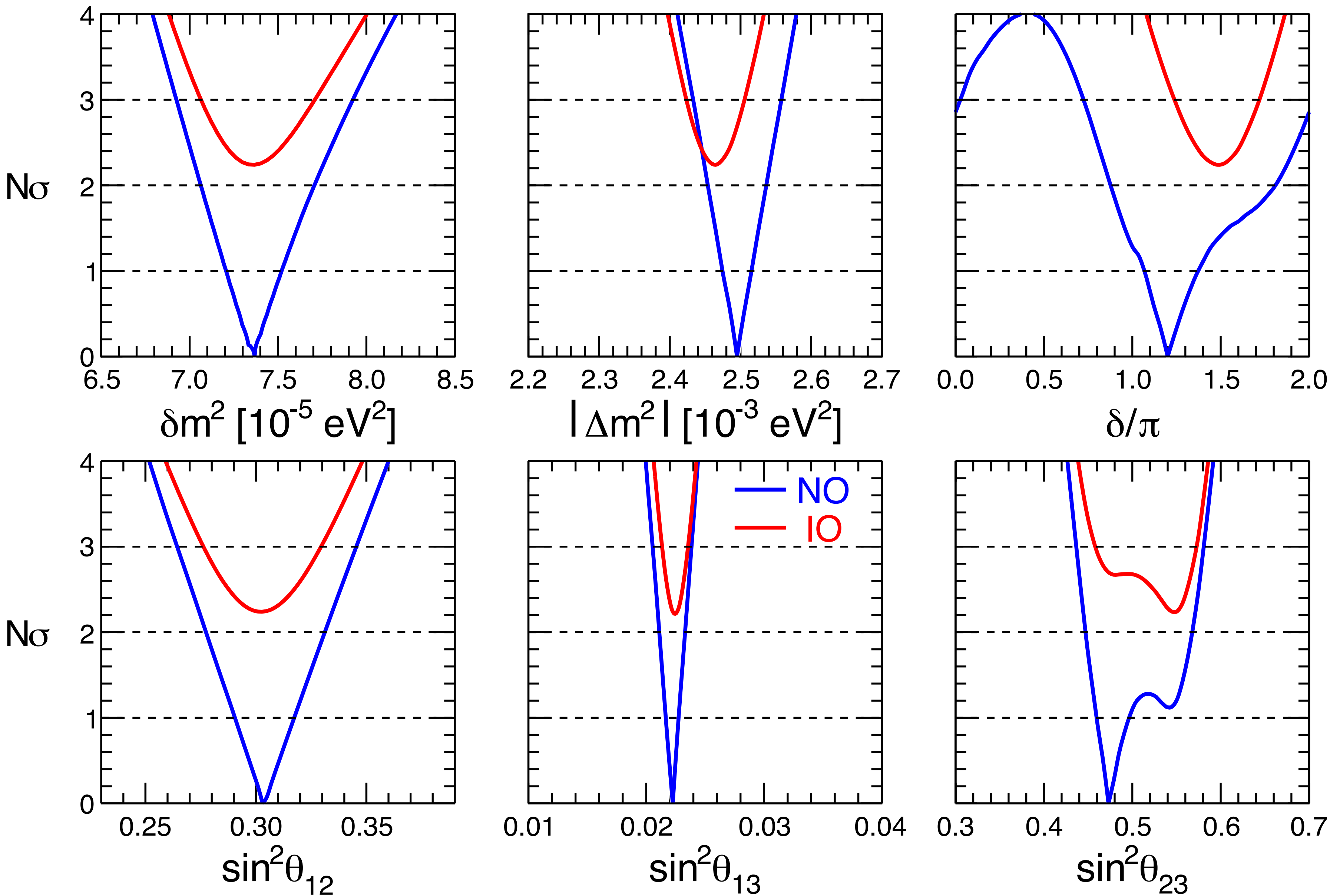


# LBL Acc + Solar + KamLAND + SBL Reactors + Atmos



Including atmospheric data favors NO over IO at  $\sim 2.2\sigma$

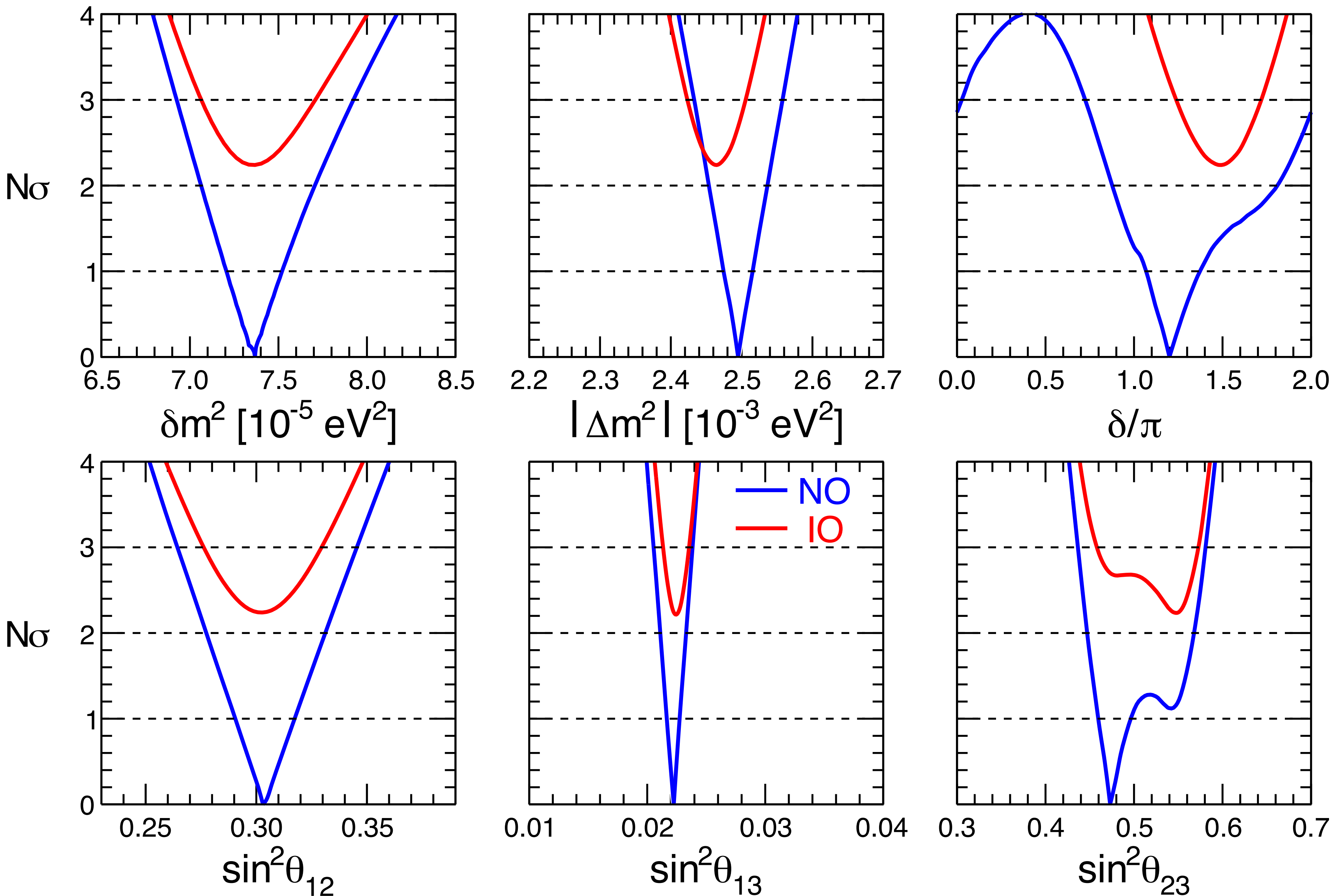
# LBL Acc + Solar + KamLAND + SBL Reactors + Atmos



Including atmospheric data favors NO over IO at  $\sim 2.2\sigma$

Weak hint for CP violation ( $\sim 1.3\sigma$ ) and first octant ( $\sim 1.1\sigma$ )

# LBL Acc + Solar + KamLAND + SBL Reactors + Atmos



Including atmospheric data favors NO over IO at  $\sim 2.2\sigma$

Weak hint for CP violation ( $\sim 1.3\sigma$ ) and first octant ( $\sim 1.1\sigma$ )

Overall status of oscillation unknowns is more uncertain than in older analyses



Absolute  $\nu$  masses  $\Rightarrow (m_\beta, m_{\beta\beta}, \Sigma)$



Absolute  $\nu$  masses  $\Rightarrow (m_\beta, m_{\beta\beta}, \Sigma)$

$\beta$  decay experiments, sensitive to the “effective electron neutrino mass”

$$m_\beta = [c_{13}^2 c_{12}^2 m_1^2 + c_{13}^2 s_{12}^2 m_2^2 + s_{13}^2 m_3^2]^{1/2}$$

Absolute  $\nu$  masses  $\Rightarrow (m_\beta, m_{\beta\beta}, \Sigma)$

$\beta$  decay experiments, sensitive to the “effective electron neutrino mass”

$$m_\beta = [c_{13}^2 c_{12}^2 m_1^2 + c_{13}^2 s_{12}^2 m_2^2 + s_{13}^2 m_3^2]^{1/2}$$

$0\nu\beta\beta$  decay experiments sensitive to the “Effective Majorana mass”:

$$m_{\beta\beta} = |c_{13}^2 c_{12}^2 m_1 + c_{13}^2 s_{12}^2 m_2 e^{i\phi_2} + s_{13}^2 m_3 e^{i\phi_3}|$$

# Absolute $\nu$ masses $\Rightarrow (m_\beta, m_{\beta\beta}, \Sigma)$

$\beta$  decay experiments, sensitive to the “effective electron neutrino mass”

$$m_\beta = [c_{13}^2 c_{12}^2 m_1^2 + c_{13}^2 s_{12}^2 m_2^2 + s_{13}^2 m_3^2]^{1/2}$$

$0\nu\beta\beta$  decay experiments sensitive to the “Effective Majorana mass”:

$$m_{\beta\beta} = |c_{13}^2 c_{12}^2 m_1 + c_{13}^2 s_{12}^2 m_2 e^{i\phi_2} + s_{13}^2 m_3 e^{i\phi_3}|$$

Cosmology and Astrophysics observations, dominantly sensitive to the sum of neutrino masses:

$$\Sigma = m_1 + m_2 + m_3$$

# Absolute $\nu$ masses $\Rightarrow (m_\beta, m_{\beta\beta}, \Sigma)$

$\beta$  decay experiments, sensitive to the “effective electron neutrino mass”

$$m_\beta = [c_{13}^2 c_{12}^2 m_1^2 + c_{13}^2 s_{12}^2 m_2^2 + s_{13}^2 m_3^2]^{1/2}$$

$0\nu\beta\beta$  decay experiments sensitive to the “Effective Majorana mass”:

$$m_{\beta\beta} = |c_{13}^2 c_{12}^2 m_1 + c_{13}^2 s_{12}^2 m_2 e^{i\phi_2} + s_{13}^2 m_3 e^{i\phi_3}|$$

Cosmology and Astrophysics observations, dominantly sensitive to the sum of neutrino masses:

$$\Sigma = m_1 + m_2 + m_3$$

These observables may provide handles to distinguish NO/IO

# Absolute $\nu$ masses $\Rightarrow (m_\beta, m_{\beta\beta}, \Sigma)$

$\beta$  decay experiments, sensitive to the “effective electron neutrino mass”

$$m_\beta = [c_{13}^2 c_{12}^2 m_1^2 + c_{13}^2 s_{12}^2 m_2^2 + s_{13}^2 m_3^2]^{1/2}$$

$0\nu\beta\beta$  decay experiments sensitive to the “Effective Majorana mass”:

$$m_{\beta\beta} = |c_{13}^2 c_{12}^2 m_1 + c_{13}^2 s_{12}^2 m_2 e^{i\phi_2} + s_{13}^2 m_3 e^{i\phi_3}|$$

Cosmology and Astrophysics observations, dominantly sensitive to the sum of neutrino masses:

$$\Sigma = m_1 + m_2 + m_3$$

These observables may provide handles to distinguish NO/IO

Majorana phases give a new source of CP violation



# Absolute $\nu$ masses $\Rightarrow (m_\beta, m_{\beta\beta}, \Sigma)$

$\beta$  decay experiments, sensitive to the “effective electron neutrino mass”

$$m_\beta = [c_{13}^2 c_{12}^2 m_1^2 + c_{13}^2 s_{12}^2 m_2^2 + s_{13}^2 m_3^2]^{1/2}$$

$0\nu\beta\beta$  decay experiments sensitive to the “Effective Majorana mass”:

$$m_{\beta\beta} = |c_{13}^2 c_{12}^2 m_1 + c_{13}^2 s_{12}^2 m_2 e^{i\phi_2} + s_{13}^2 m_3 e^{i\phi_3}|$$

Cosmology and Astrophysics observations, dominantly sensitive to the sum of neutrino masses:

$$\Sigma = m_1 + m_2 + m_3$$

These observables may provide handles to distinguish NO/IO

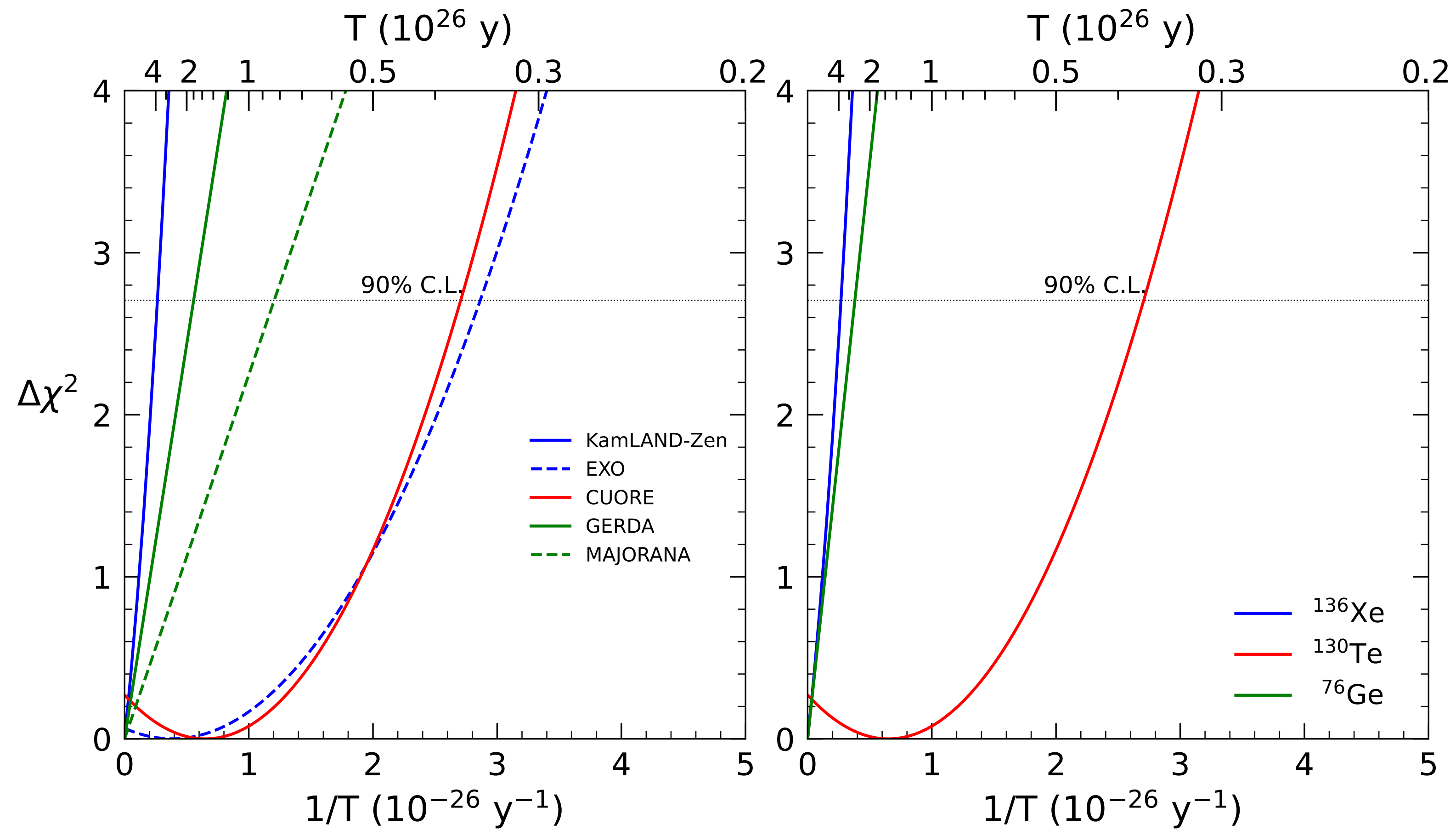
Majorana phases give a new source of CP violation

Note that the three observables are correlated by oscillation data

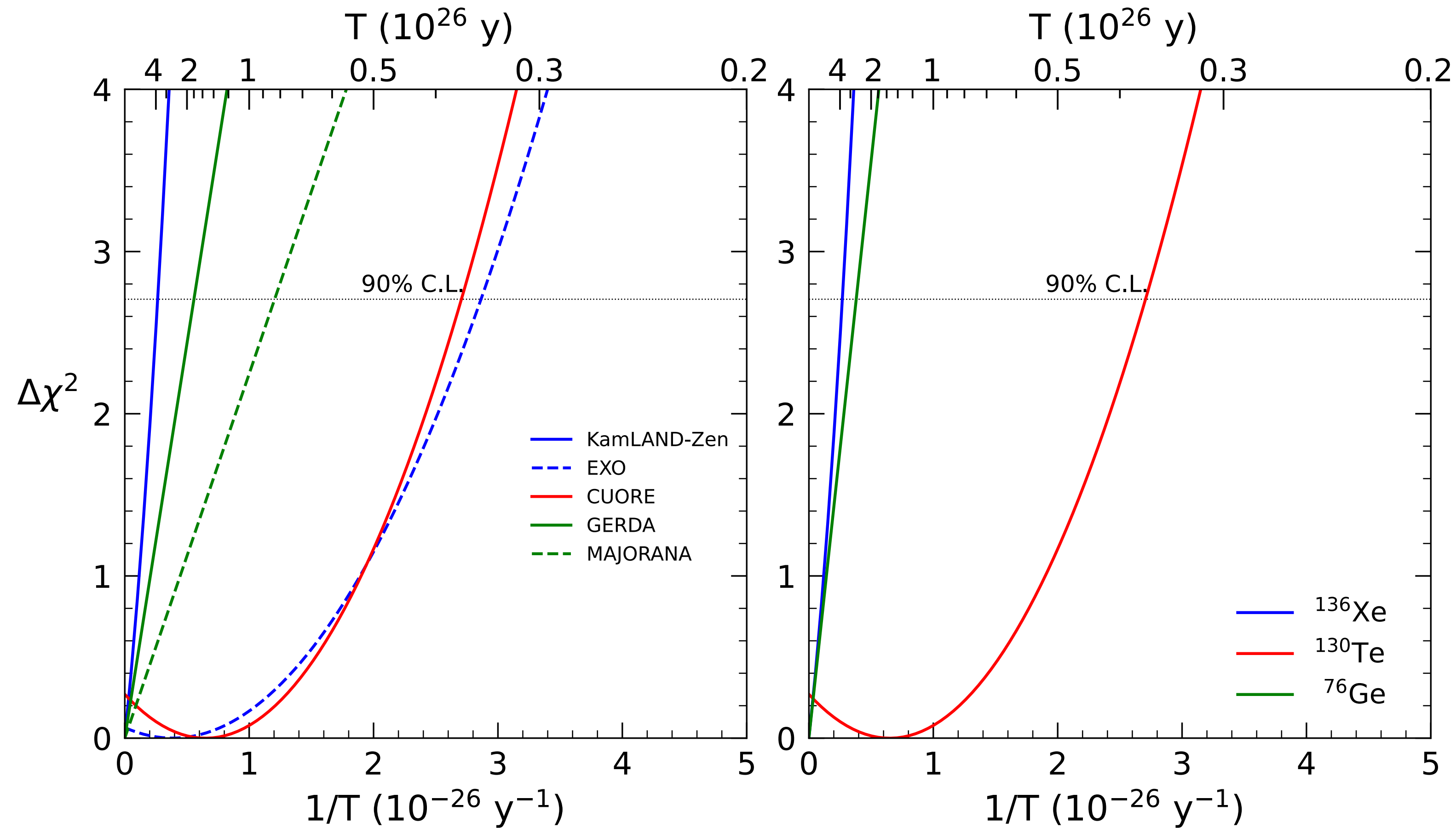


# Neutrinoless Double Beta Decay results

# Neutrinoless Double Beta Decay results



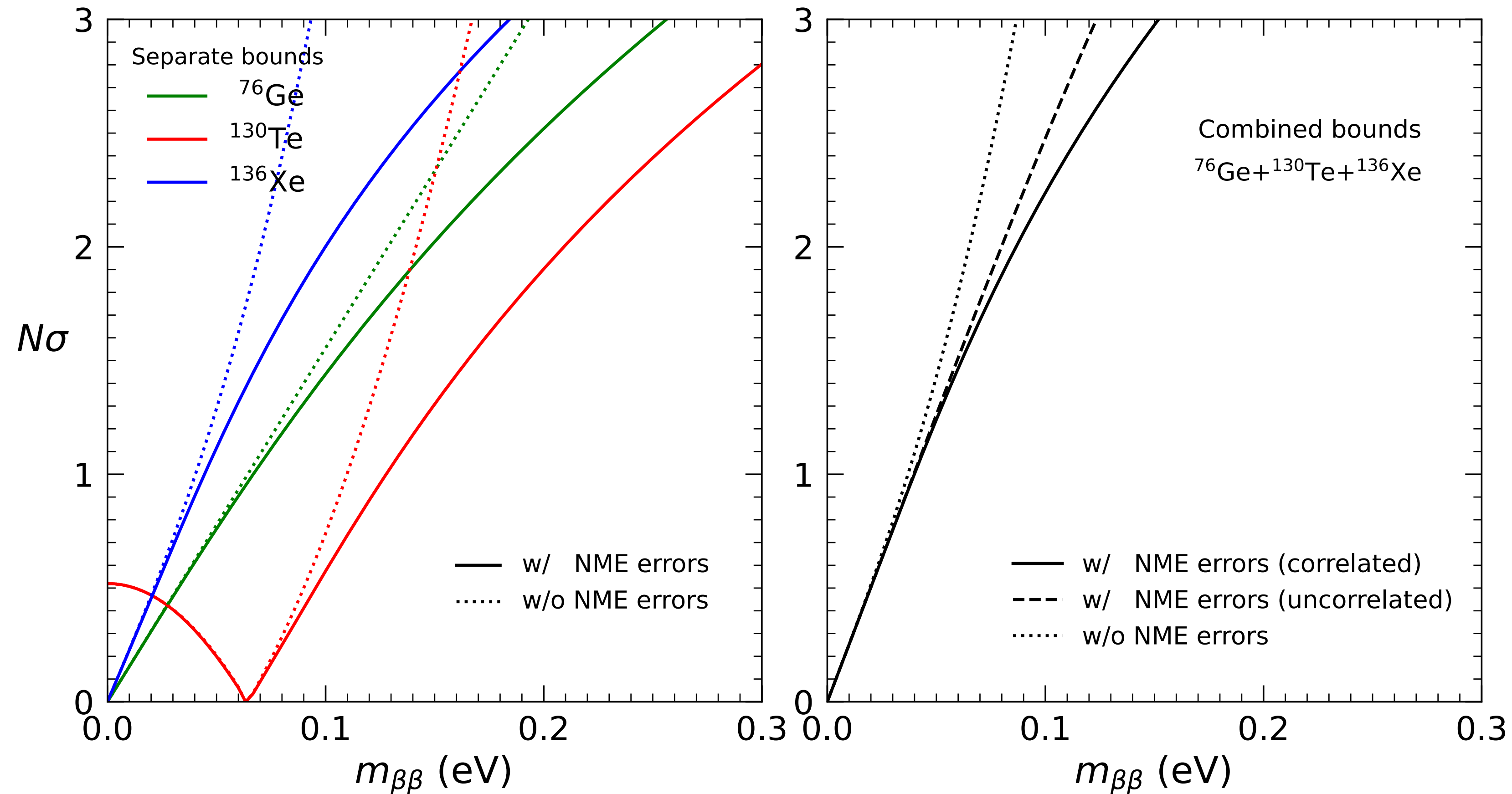
# Neutrinoless Double Beta Decay results



Translating bounds on the half-life  $T$  to bounds on  $m_{\beta\beta}$  requires the knowledge of the nuclear matrix element  $M$  (NME) for the decay at issue since

$$\frac{1}{T} = \text{phase space} \times |M|^2 \times m_{\beta\beta}^2$$

# Neutrinoless Double Beta Decay results



Translating bounds on the half-life  $T$  to bounds on  $m_{\beta\beta}$  requires the knowledge of the nuclear matrix element  $M$  (NME) for the decay at issue since

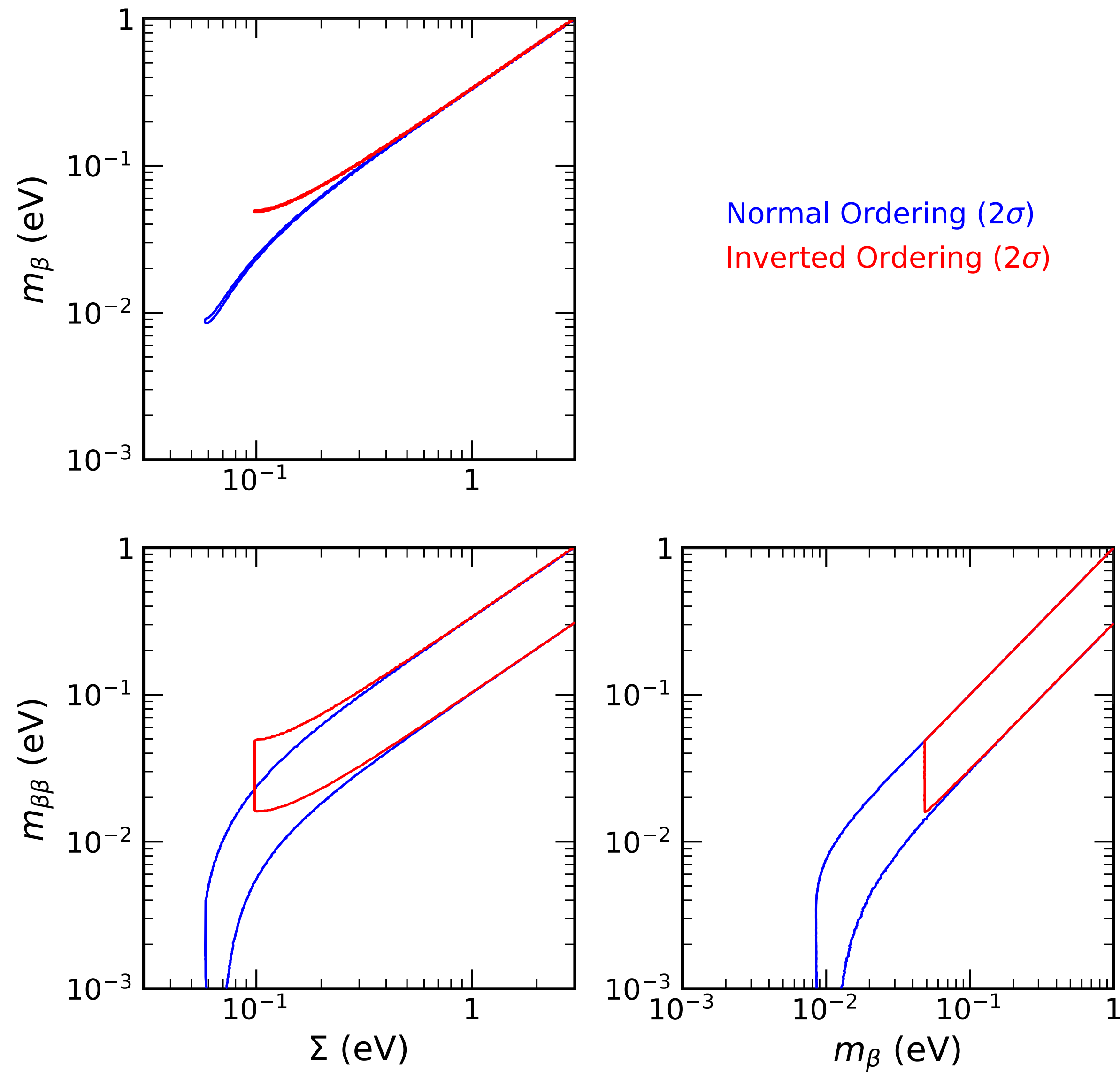
$$\frac{1}{T} = \text{phase space} \times |M|^2 \times m_{\beta\beta}^2$$



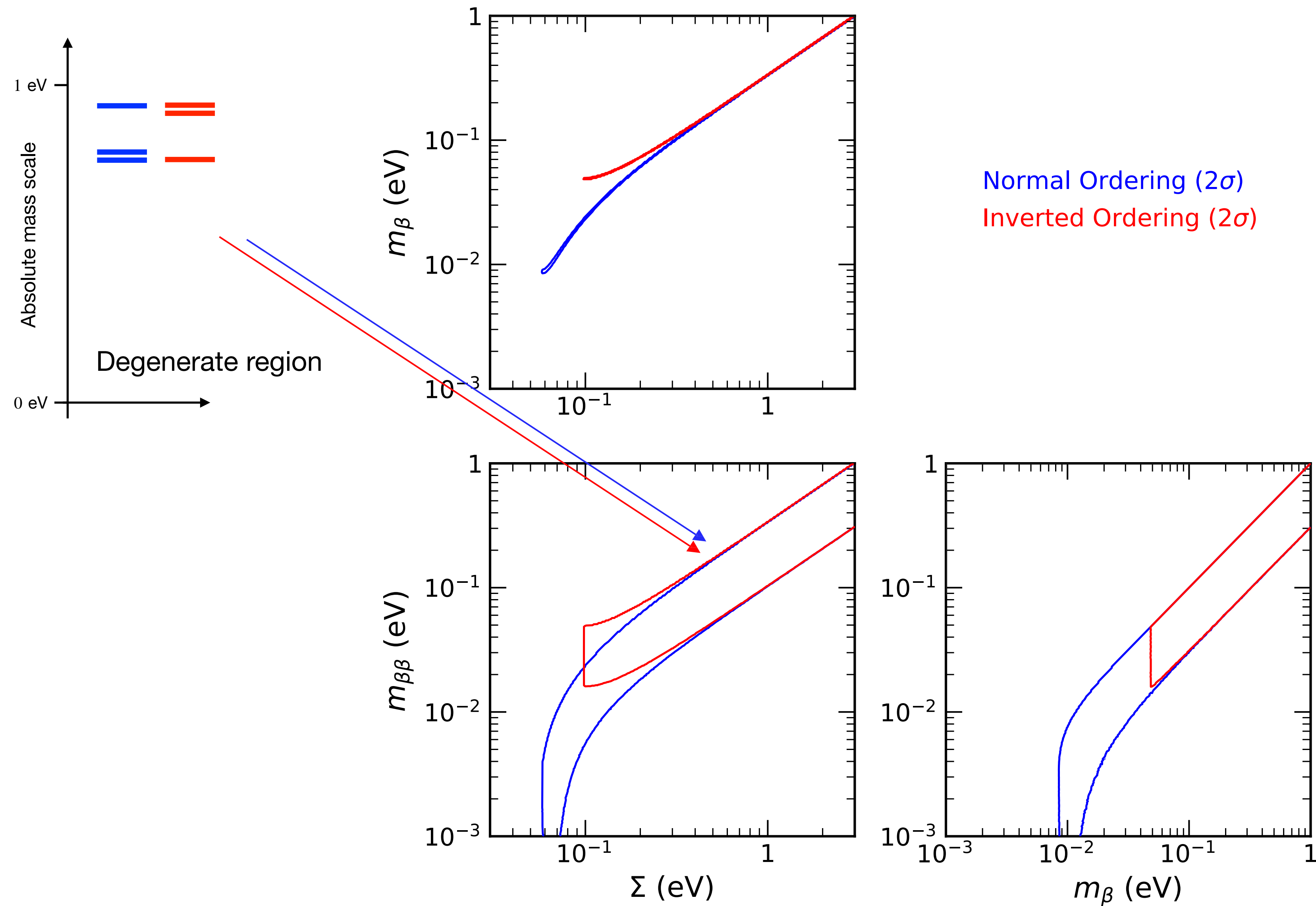


Regions allowed by oscillations on  $(\Sigma, m_\beta, m_{\beta\beta})$

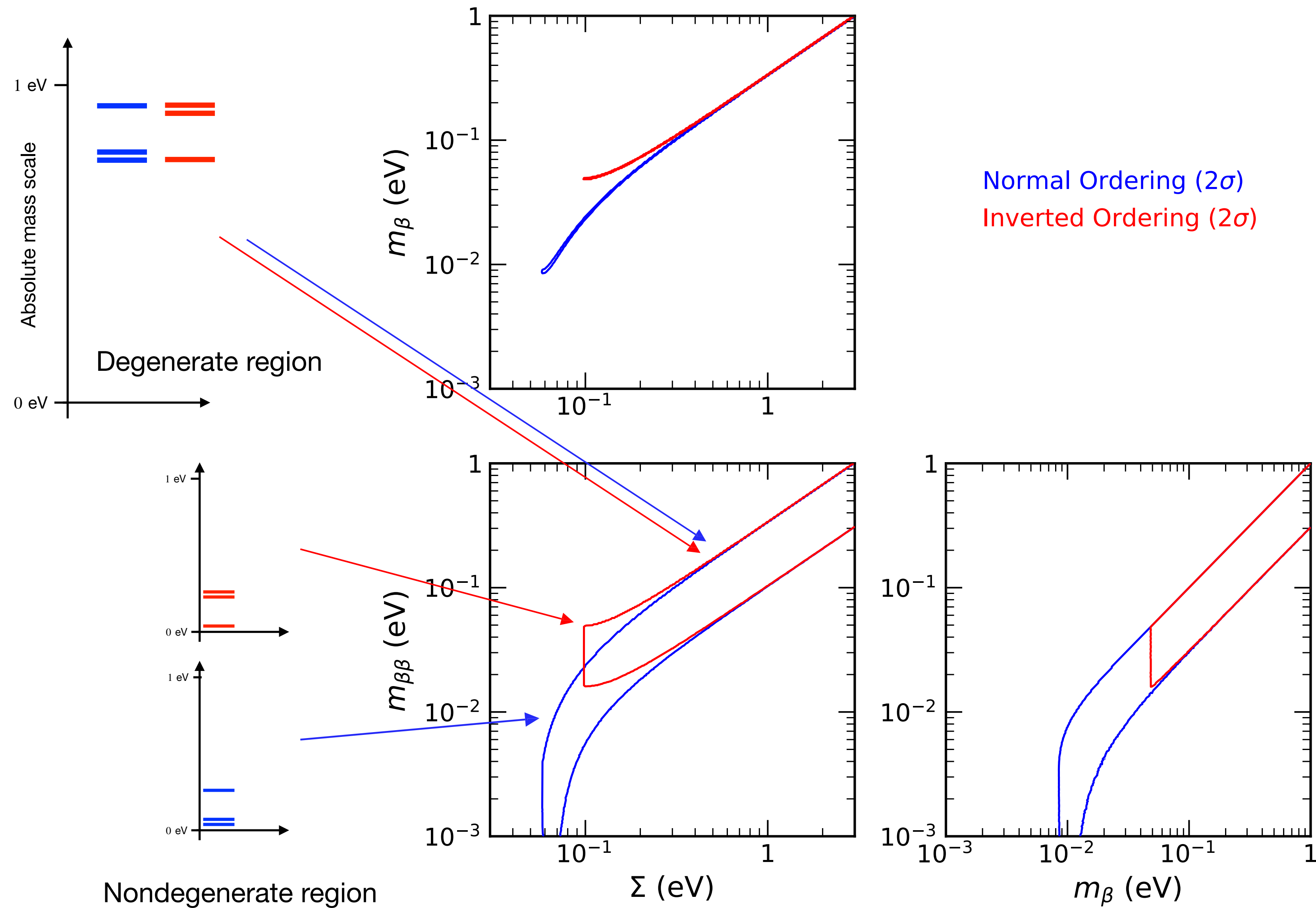
# Regions allowed by oscillations on $(\Sigma, m_\beta, m_{\beta\beta})$



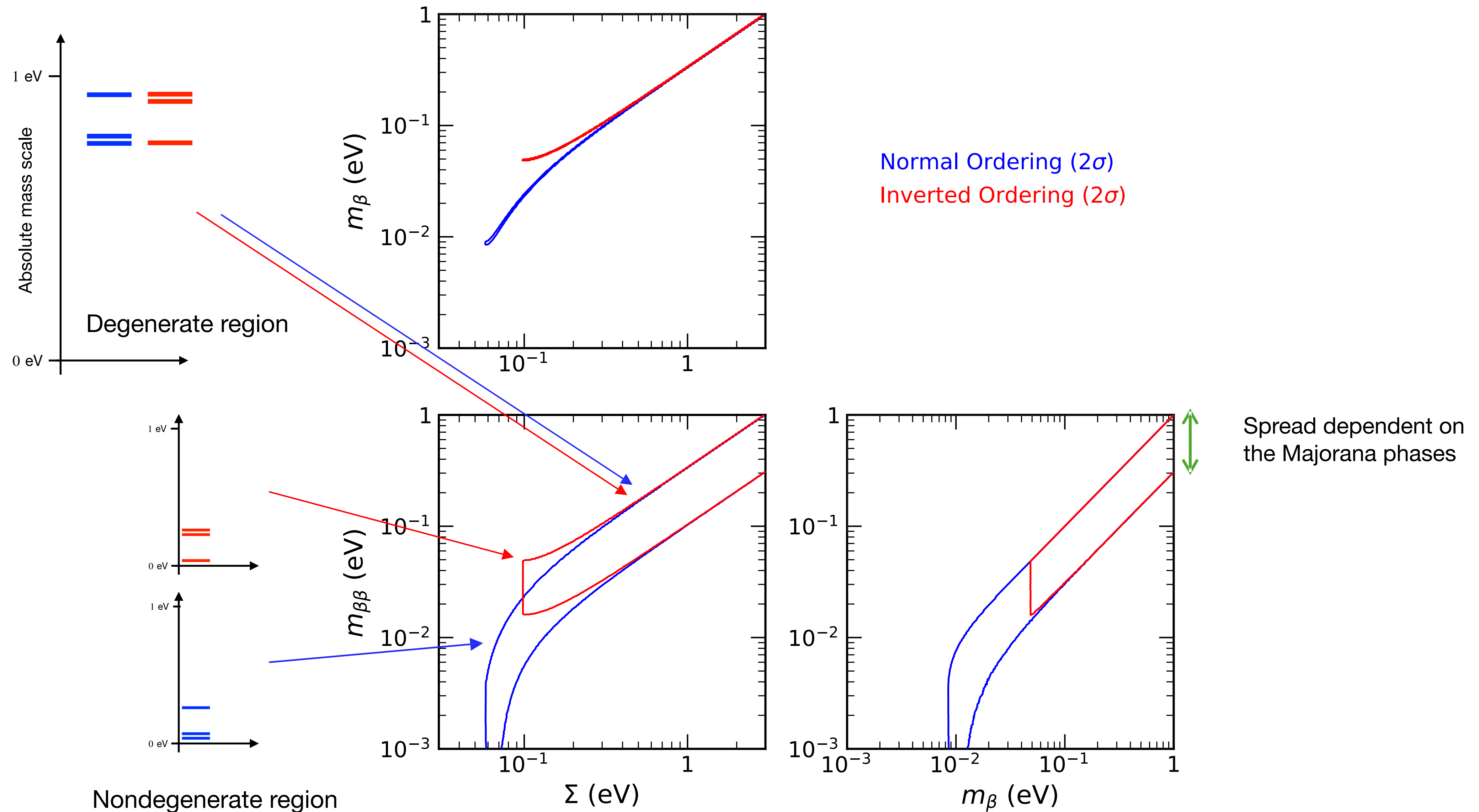
# Regions allowed by oscillations on $(\Sigma, m_\beta, m_{\beta\beta})$



# Regions allowed by oscillations on $(\Sigma, m_\beta, m_{\beta\beta})$

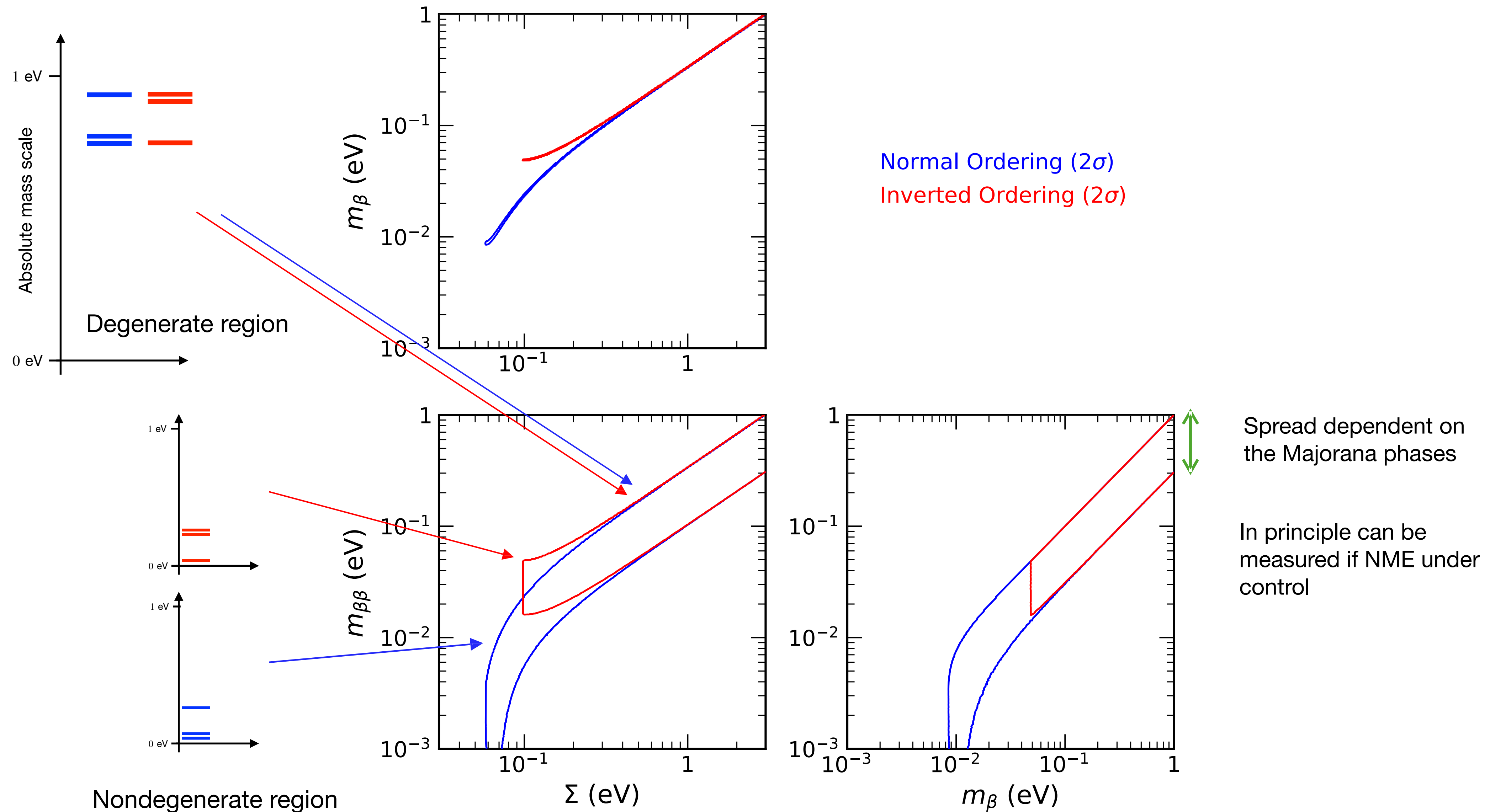


# Regions allowed by oscillations on $(\Sigma, m_\beta, m_{\beta\beta})$

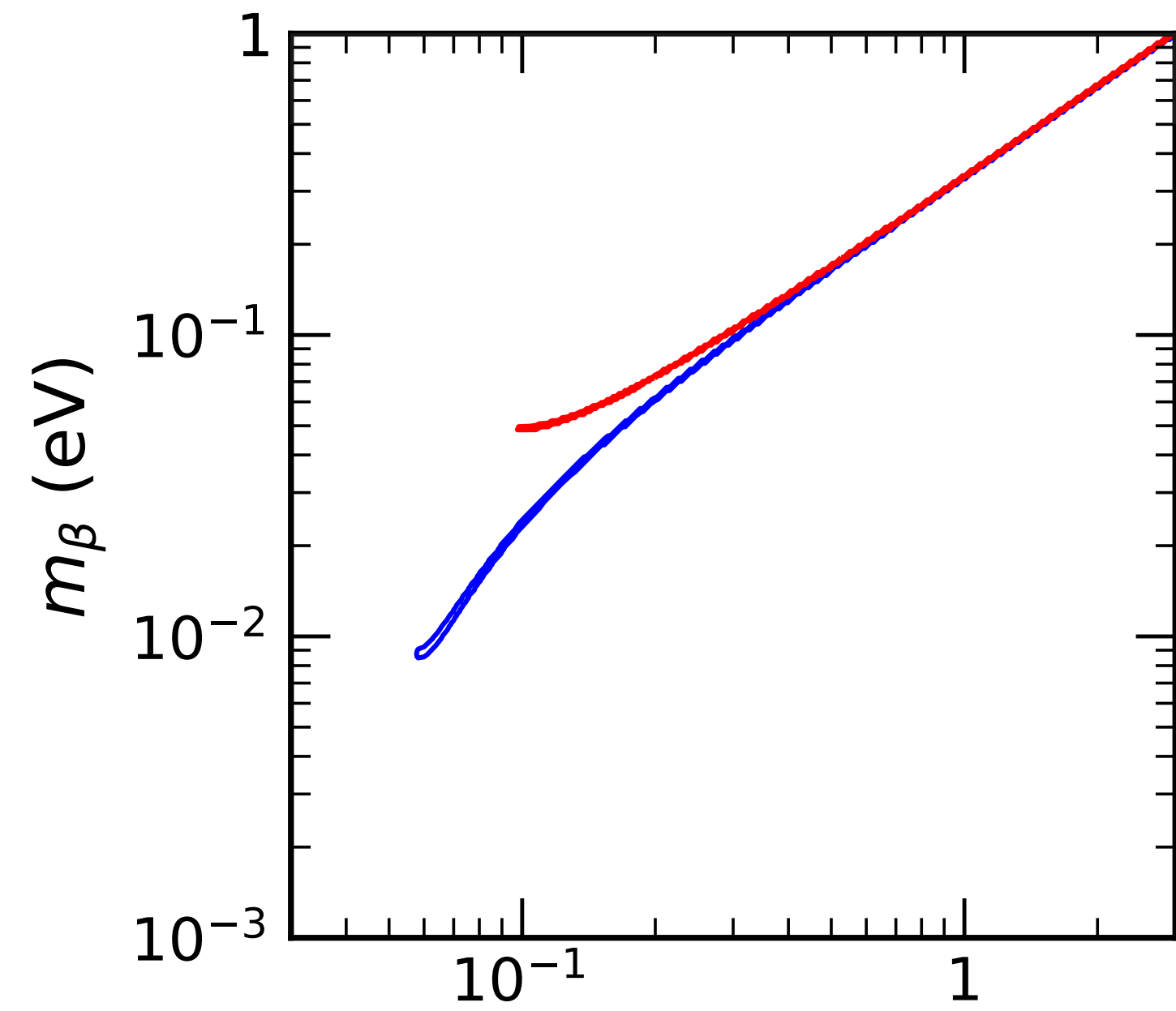




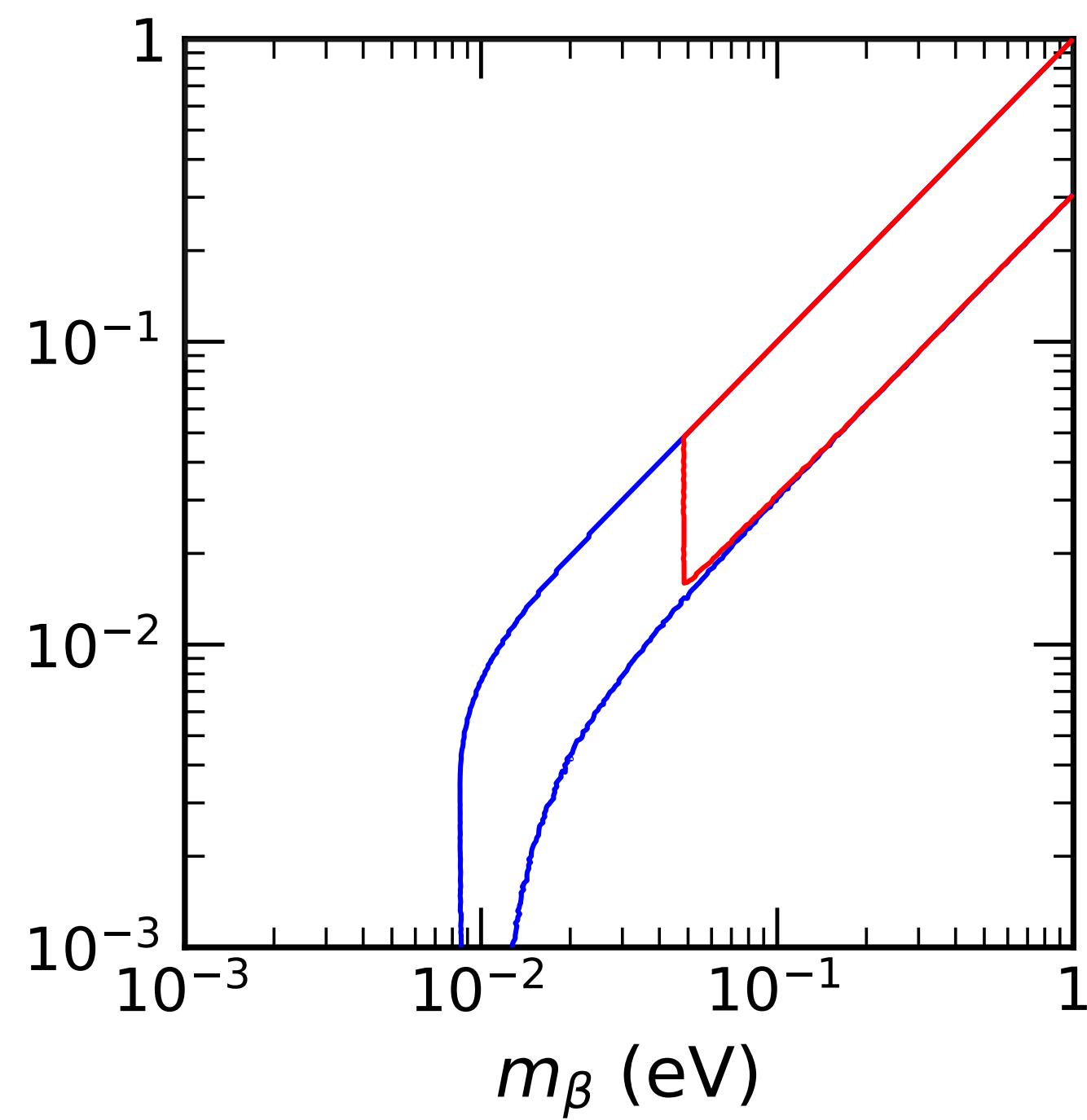
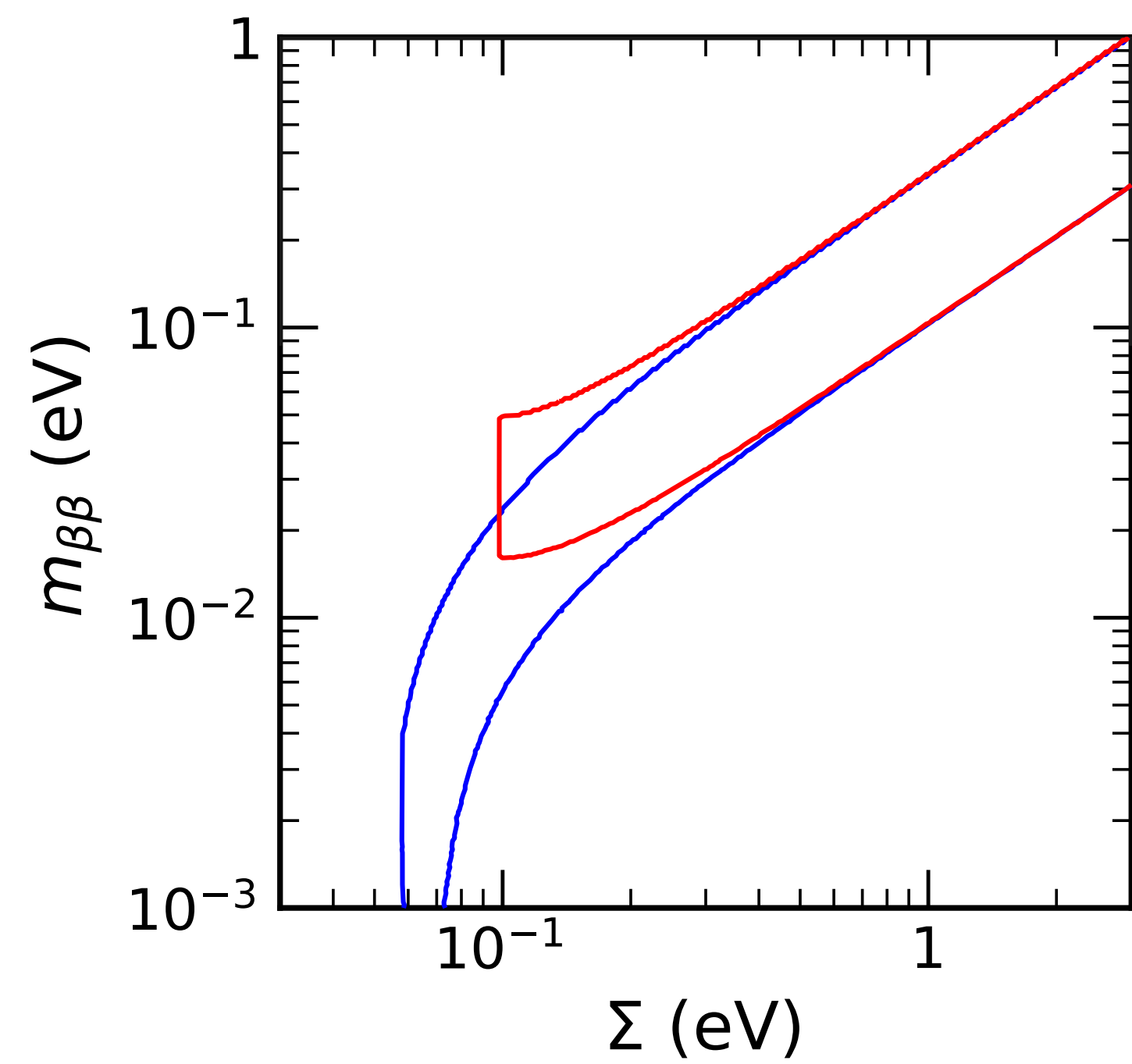
# Regions allowed by oscillations on $(\Sigma, m_\beta, m_{\beta\beta})$

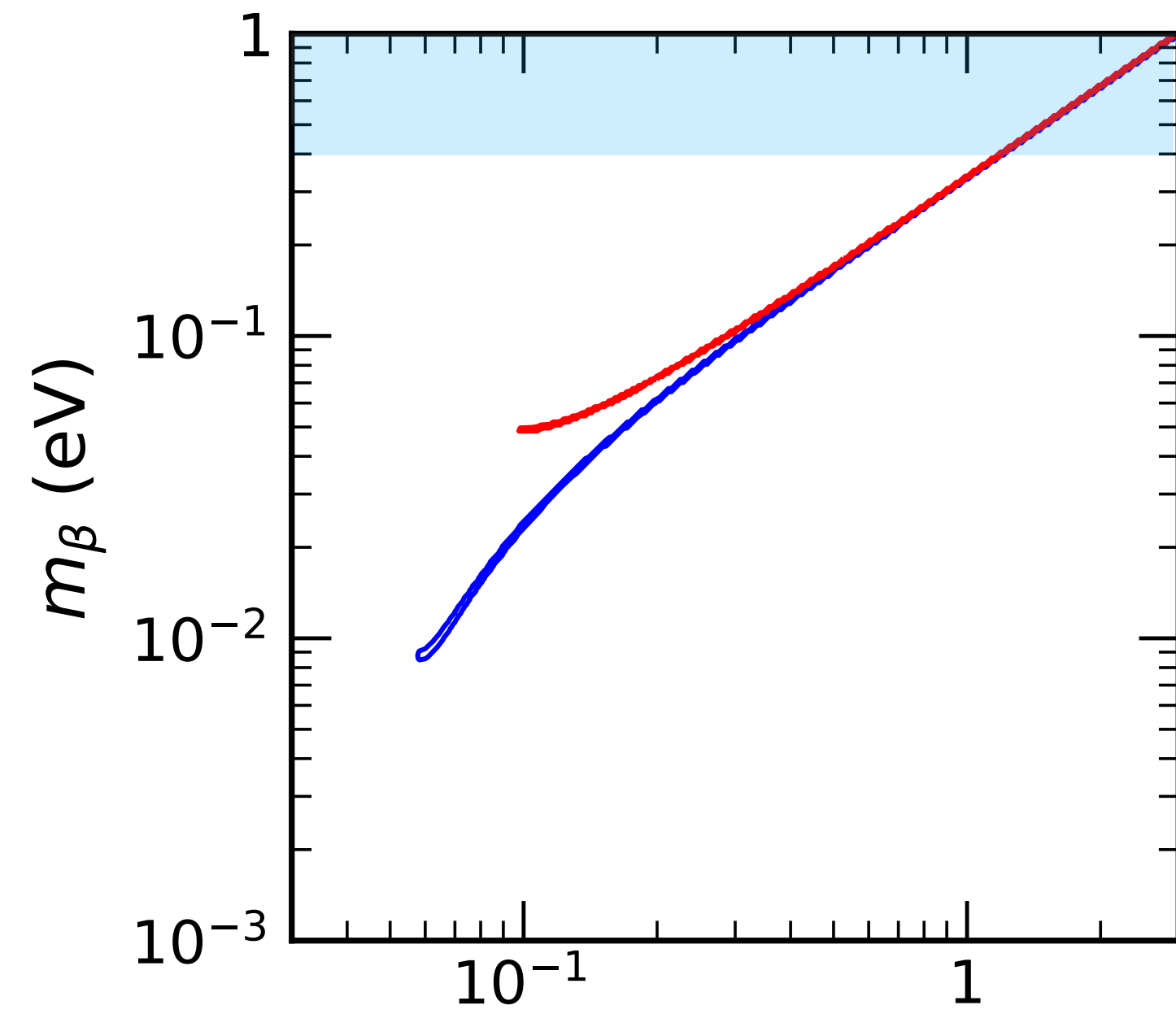






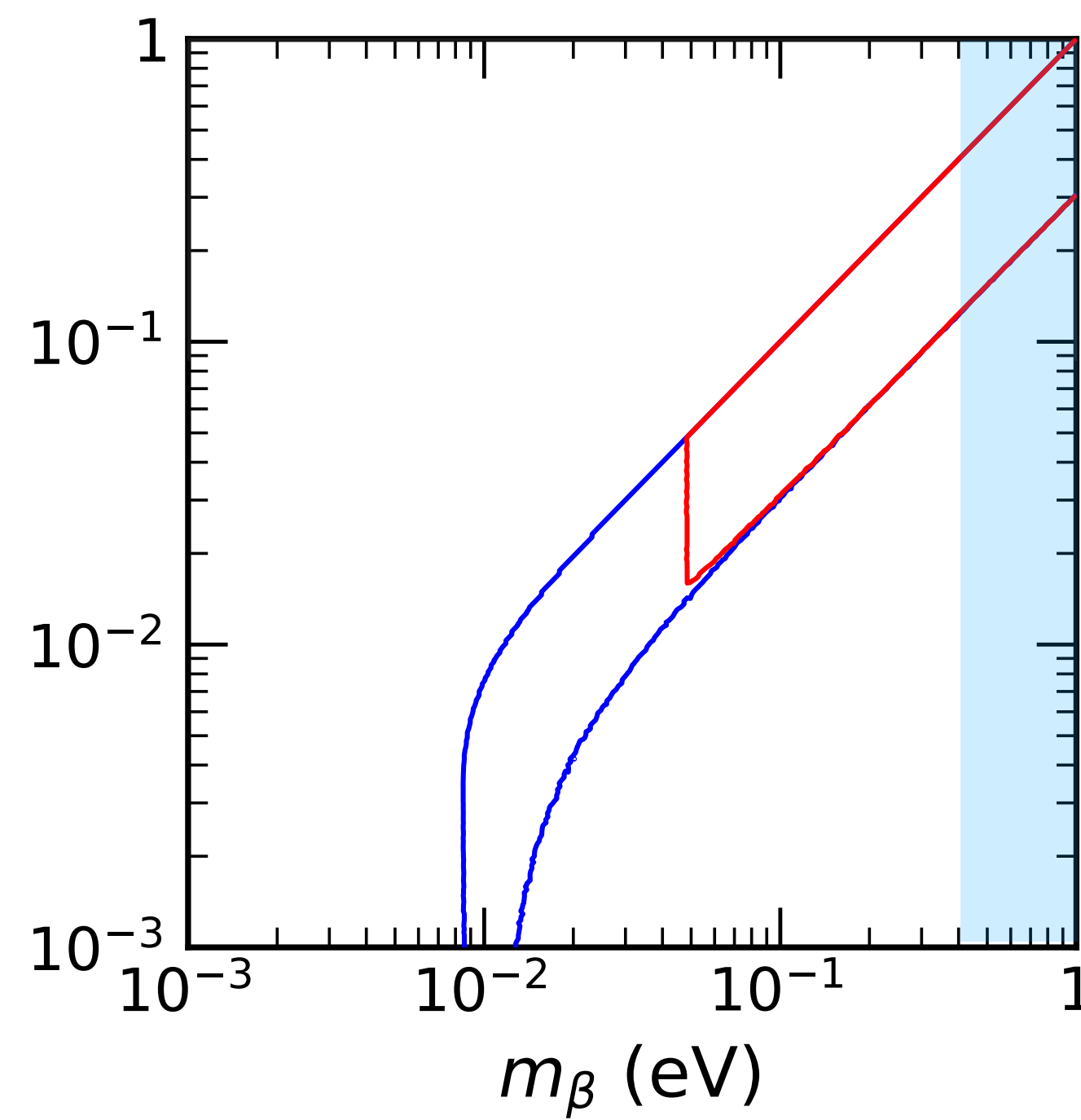
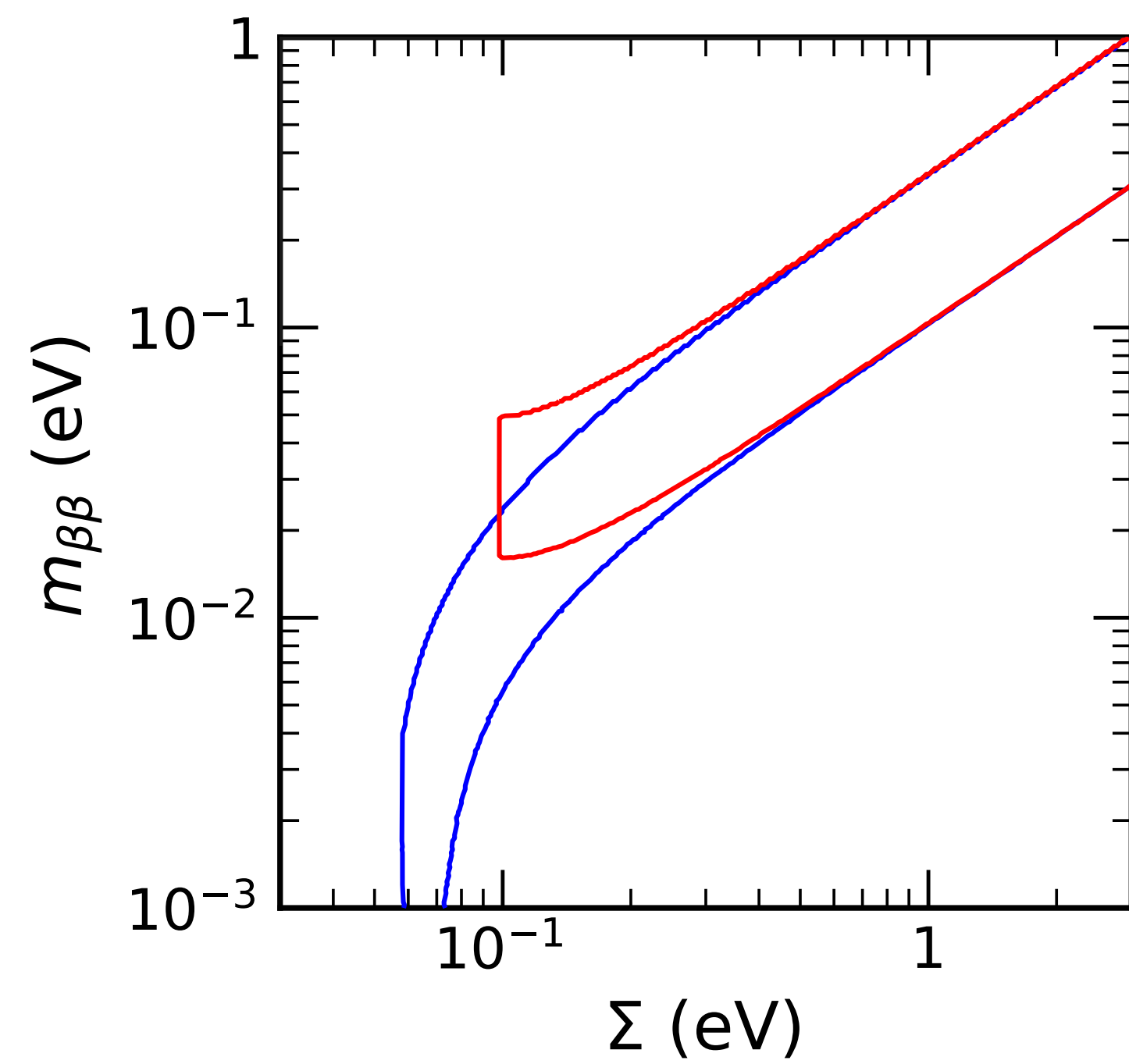
Normal Ordering ( $2\sigma$ )  
Inverted Ordering ( $2\sigma$ )

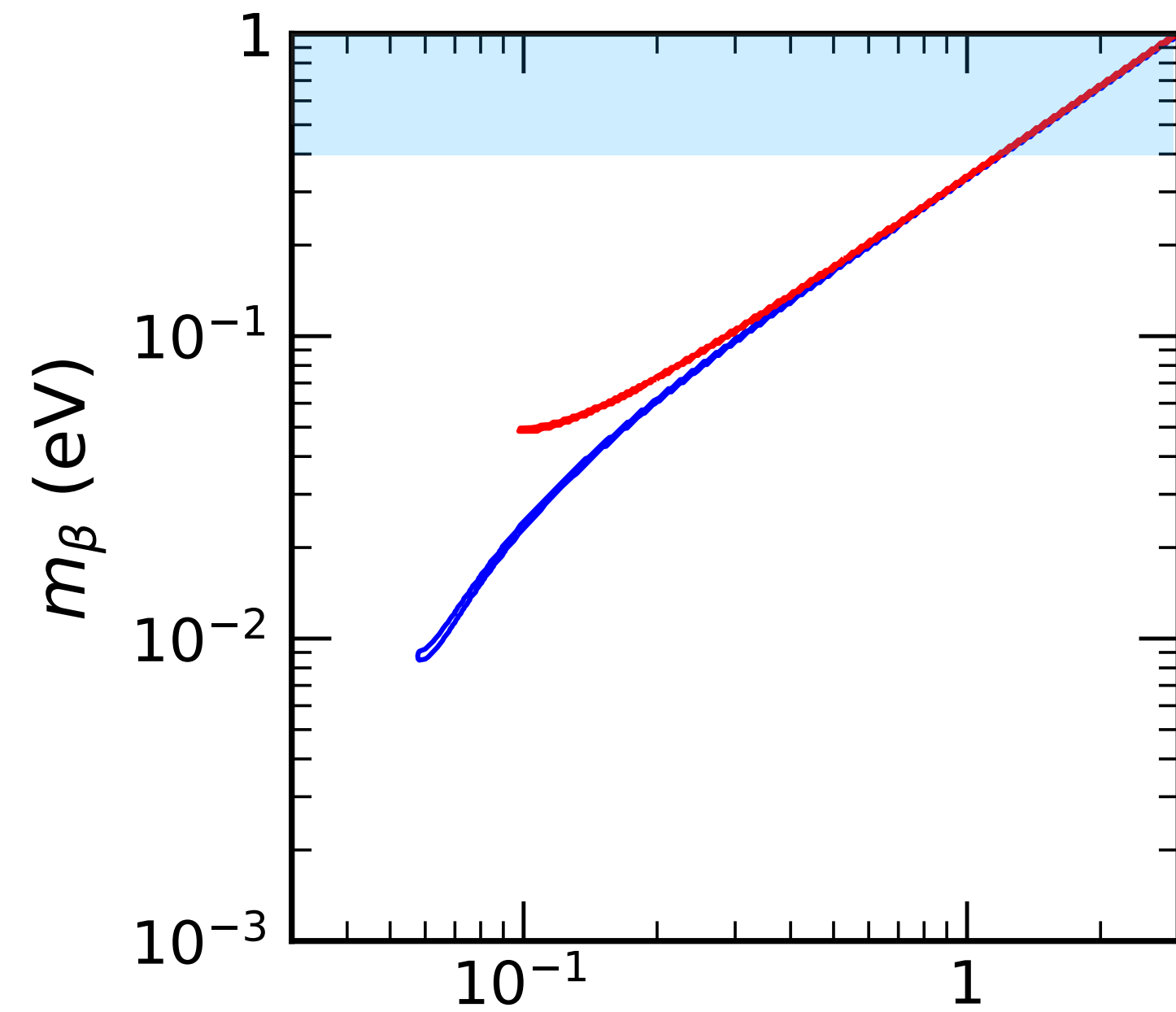




Normal Ordering ( $2\sigma$ )  
Inverted Ordering ( $2\sigma$ )

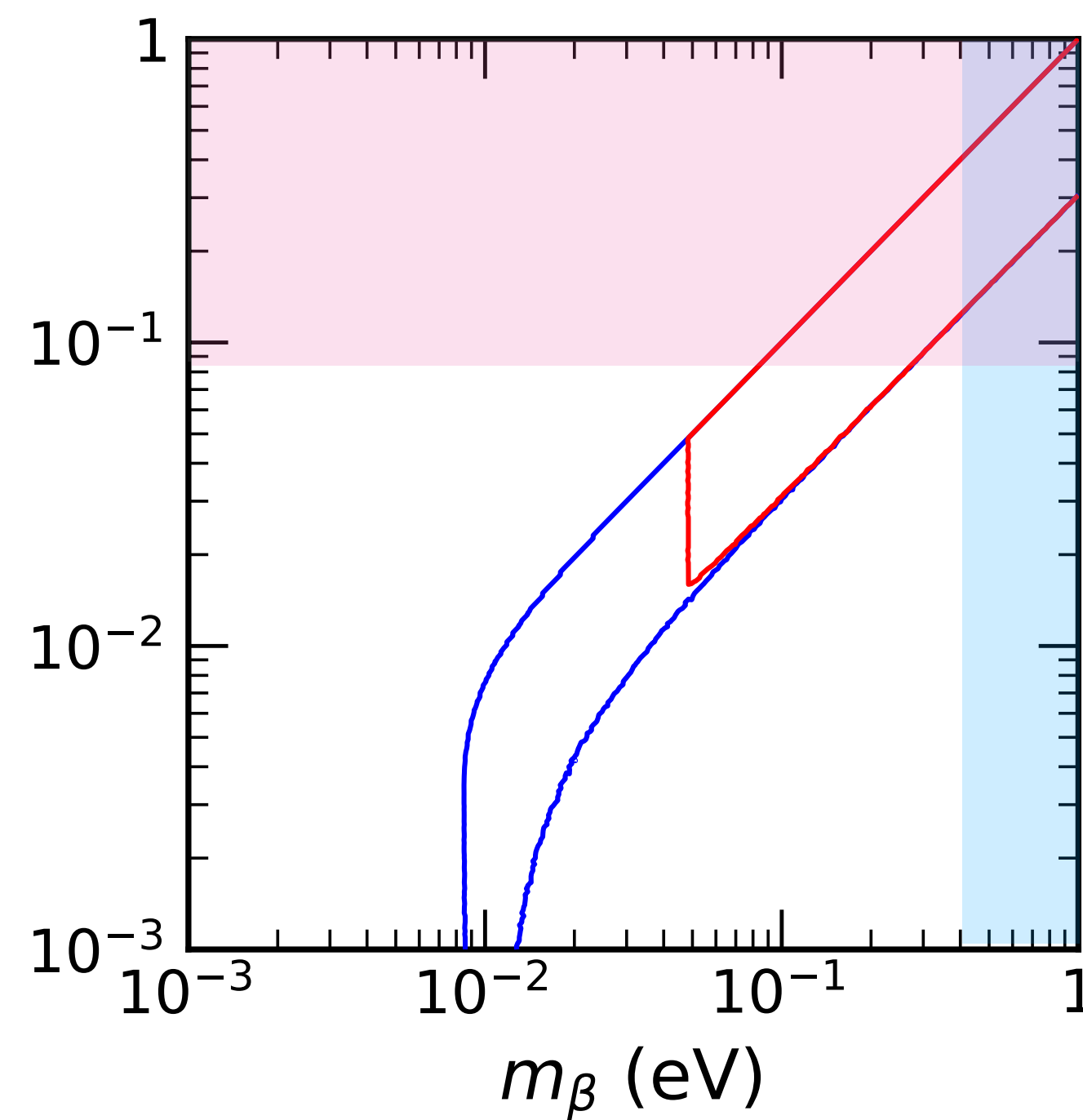
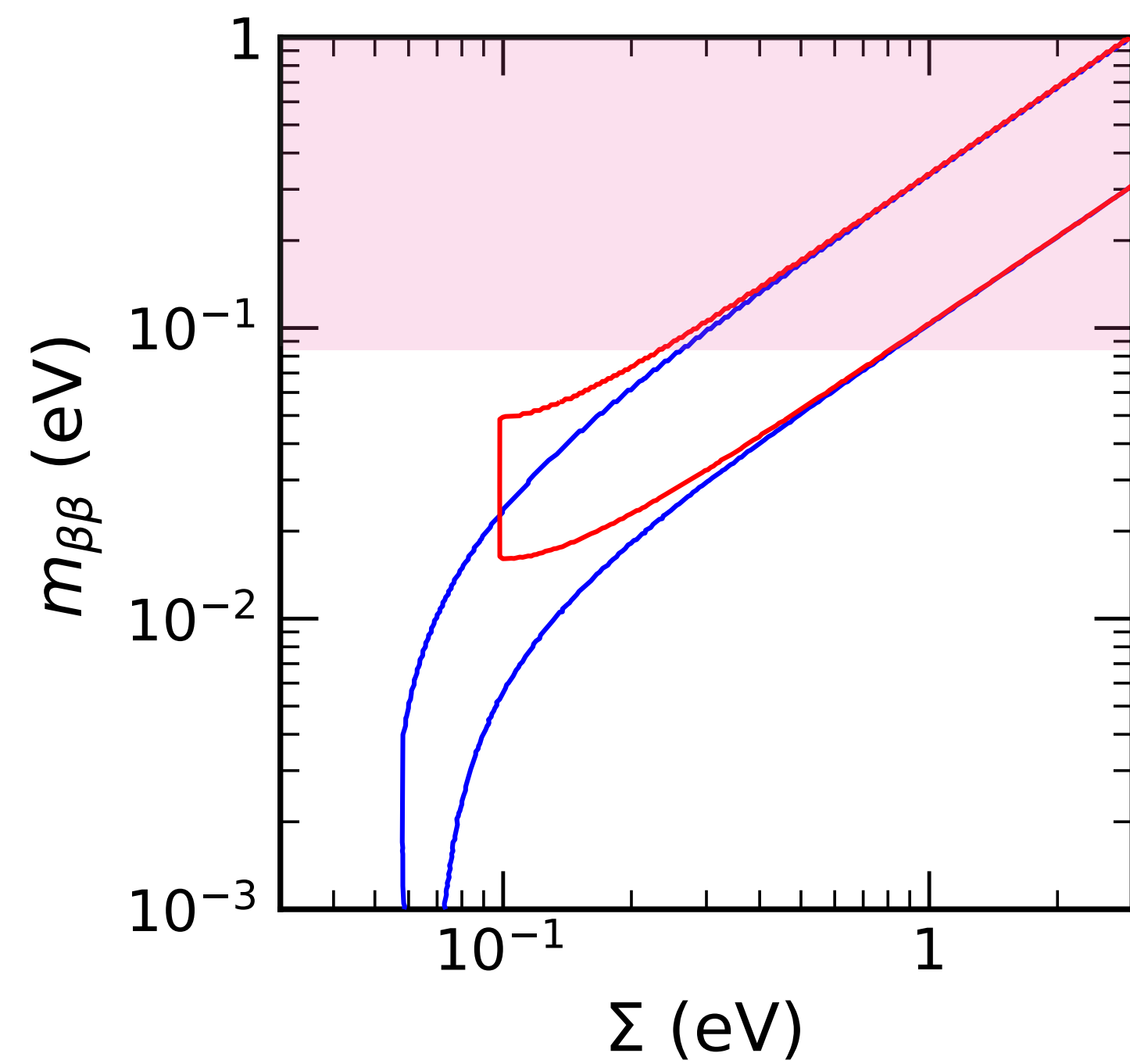
$\beta$  decay - KATRIN



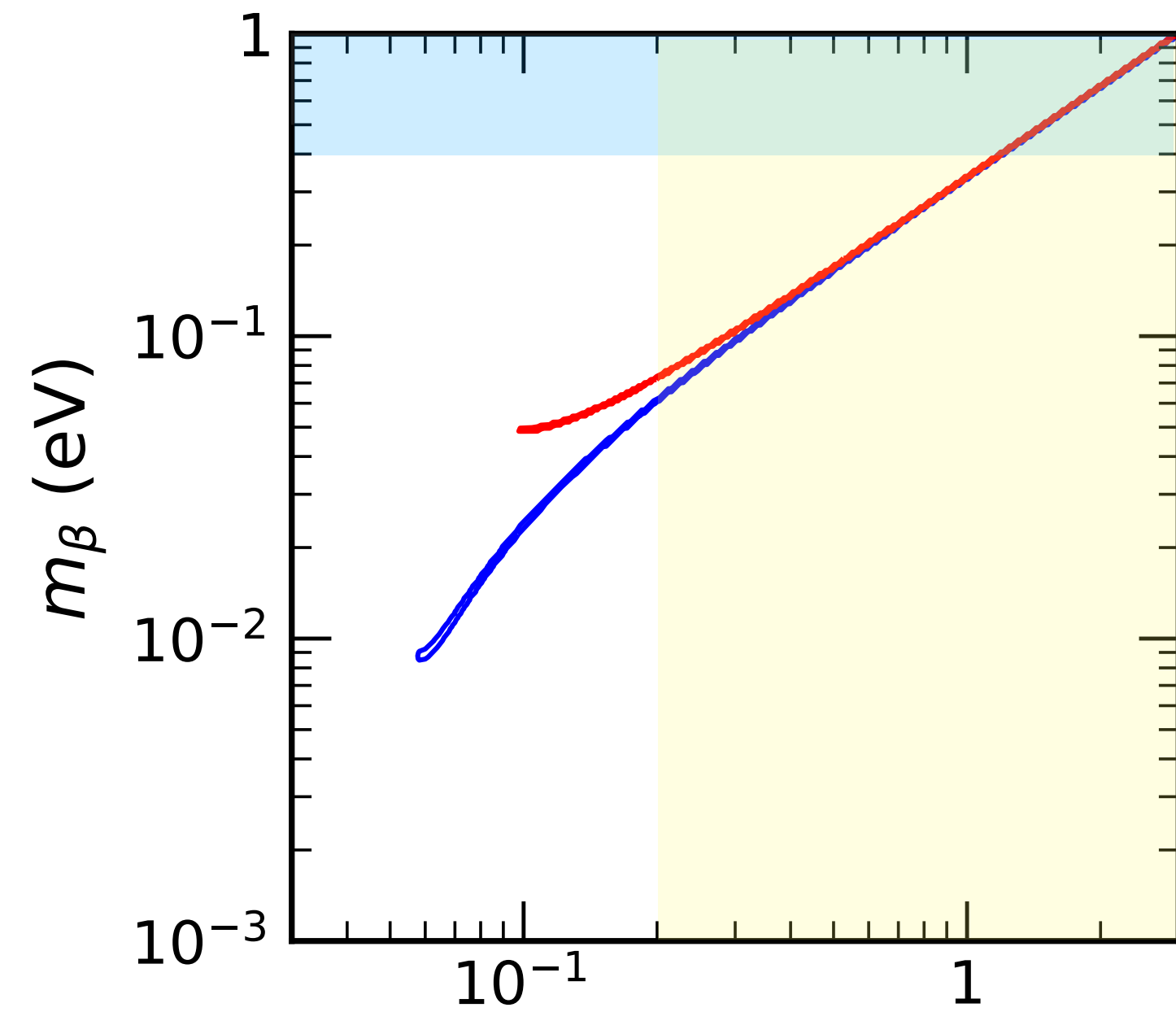


Normal Ordering ( $2\sigma$ )  
Inverted Ordering ( $2\sigma$ )

$\beta$  decay - KATRIN

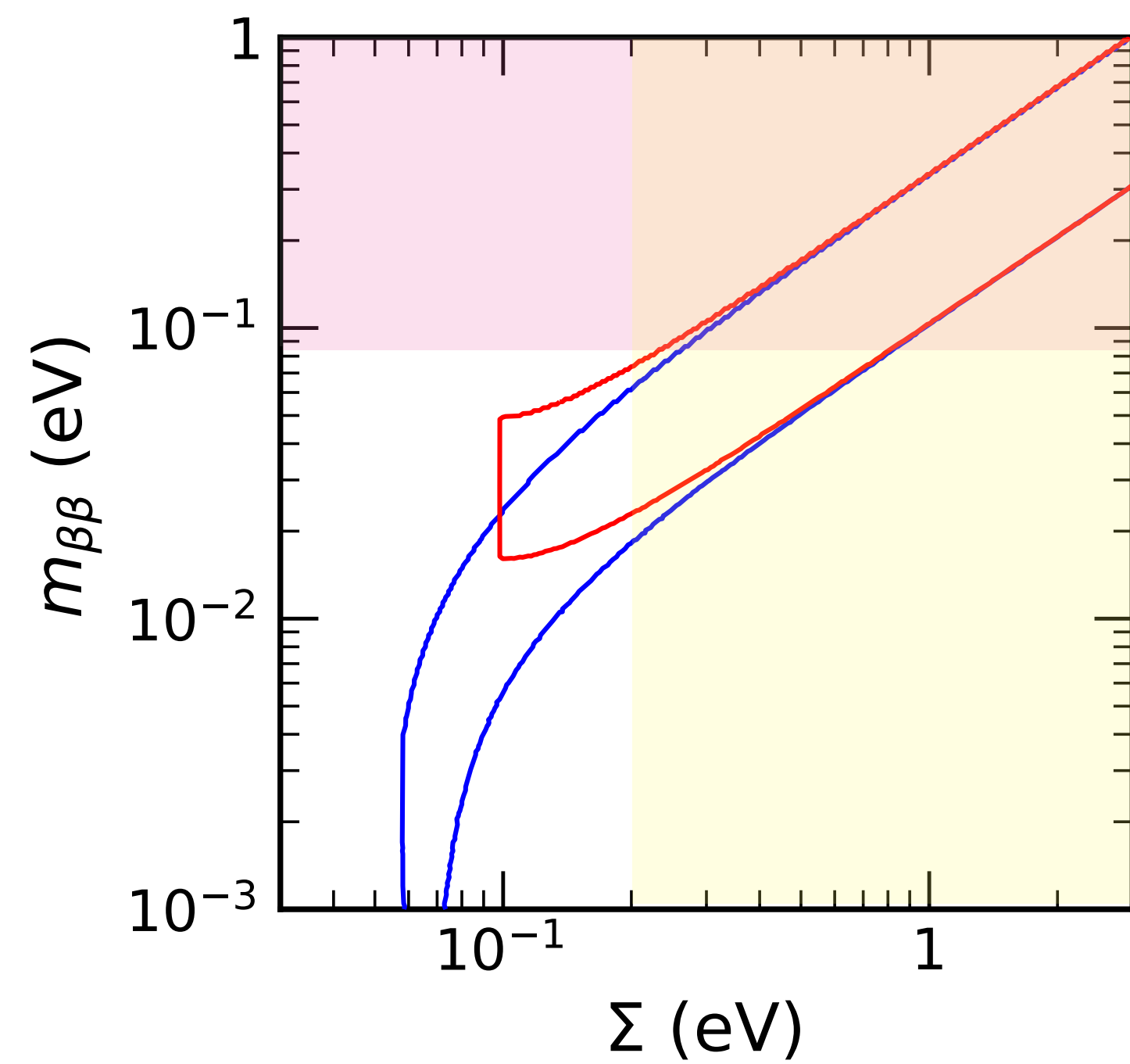


$0\nu\beta\beta$  decay  
KamLAND-Zen, EXO, CUORE, GERDA

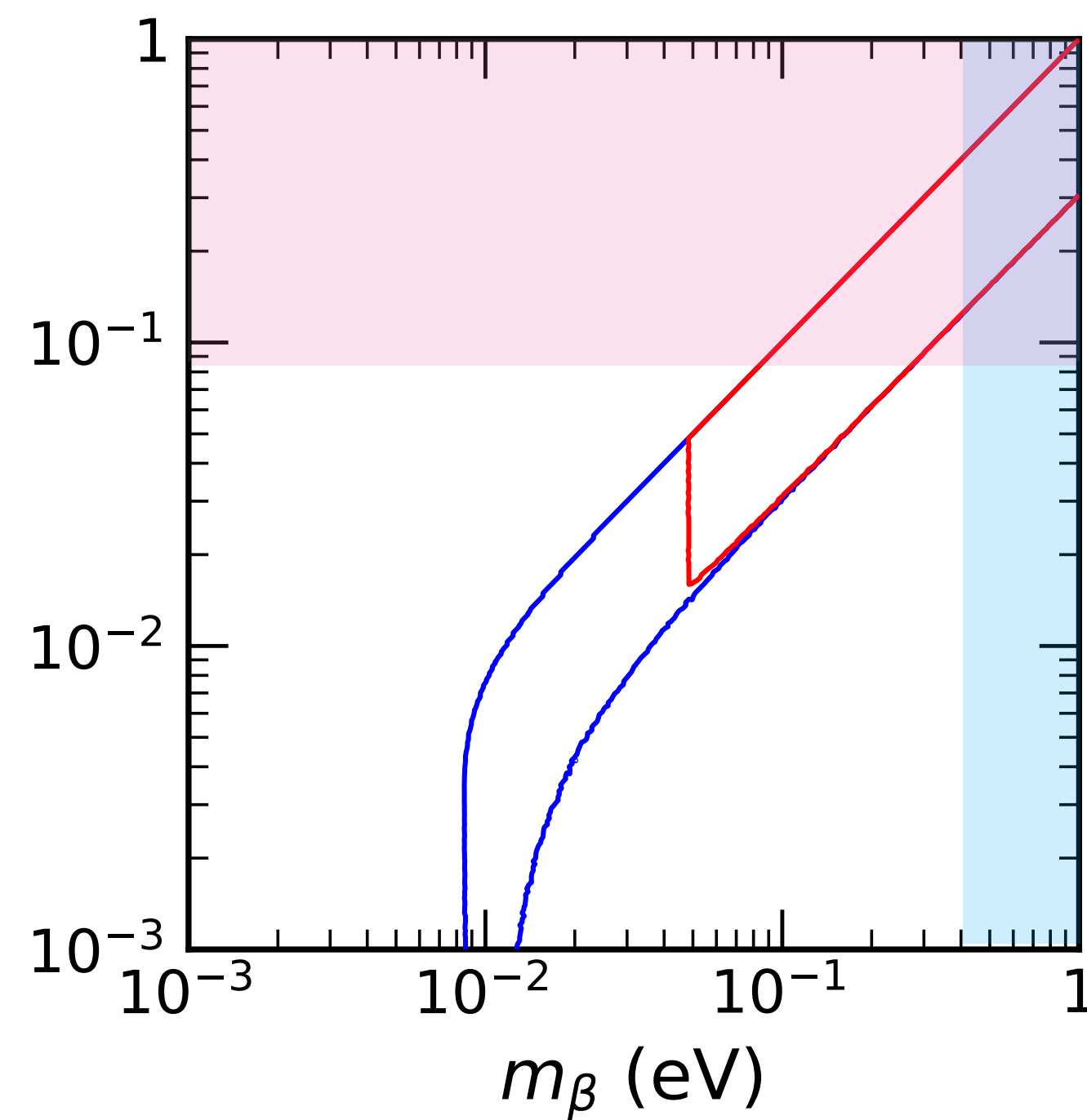


Normal Ordering (2 $\sigma$ )  
Inverted Ordering (2 $\sigma$ )

$\beta$  decay - KATRIN

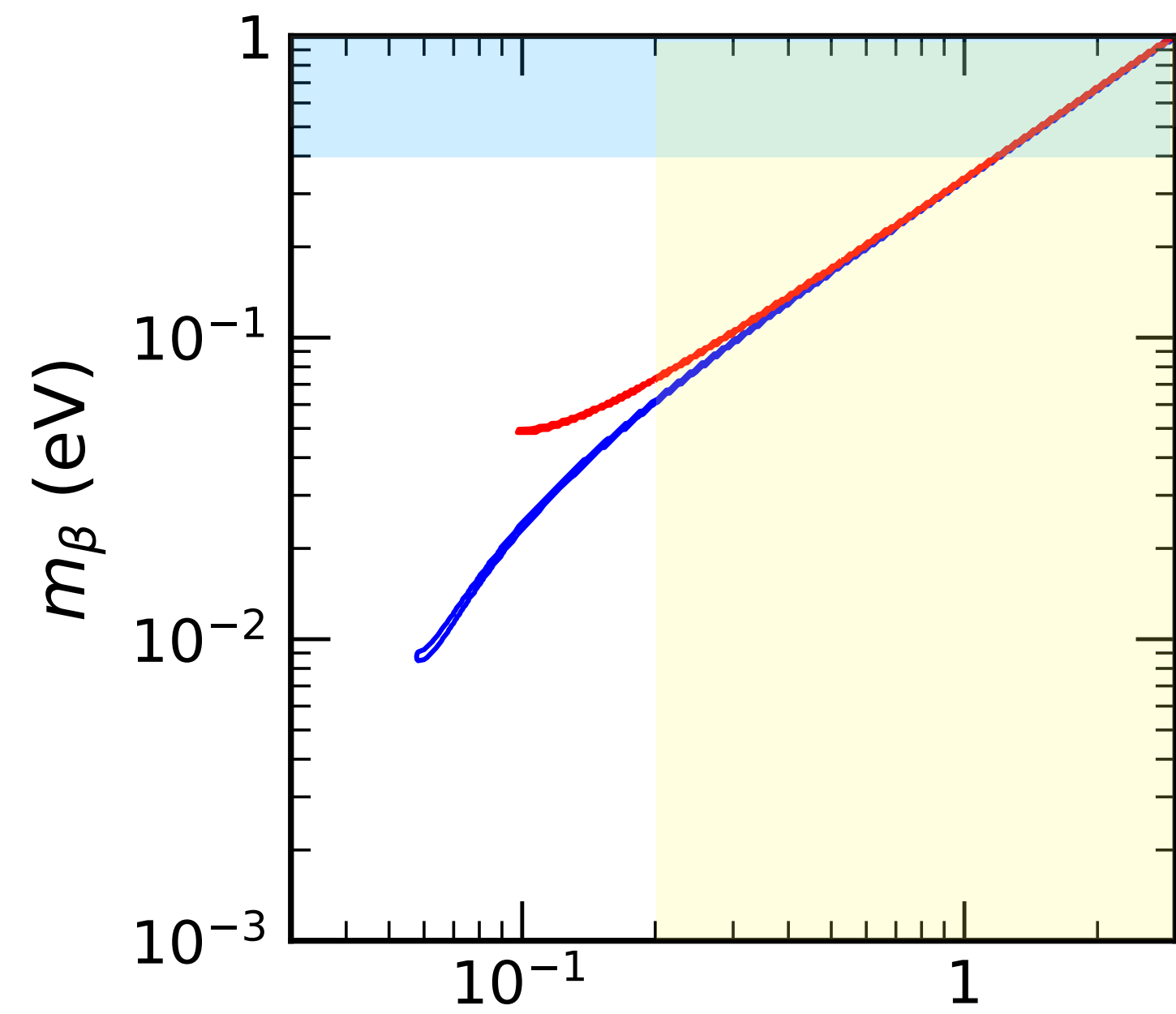


$0\nu\beta\beta$  decay  
KamLAND-Zen, EXO, CUORE, GERDA

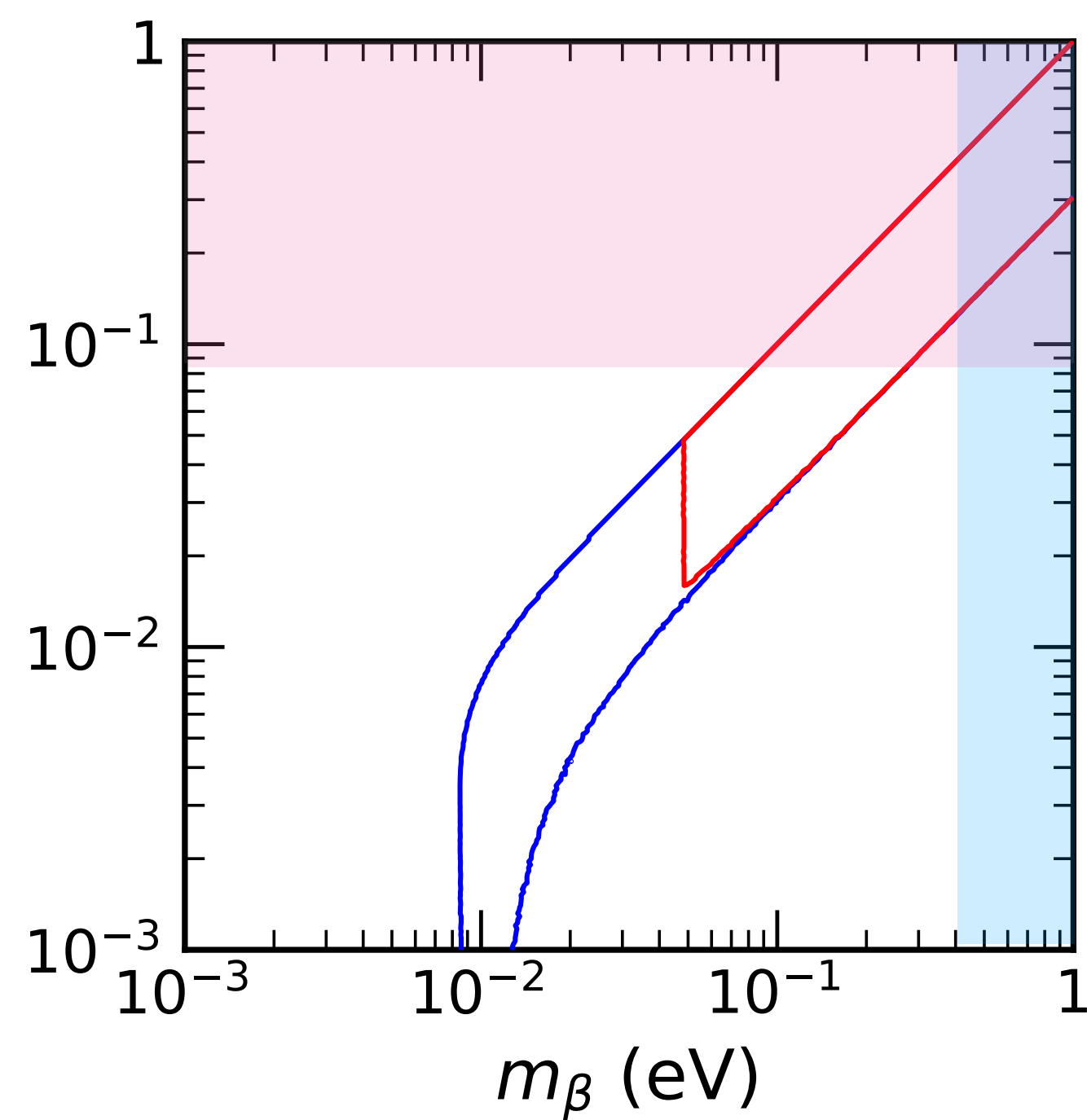
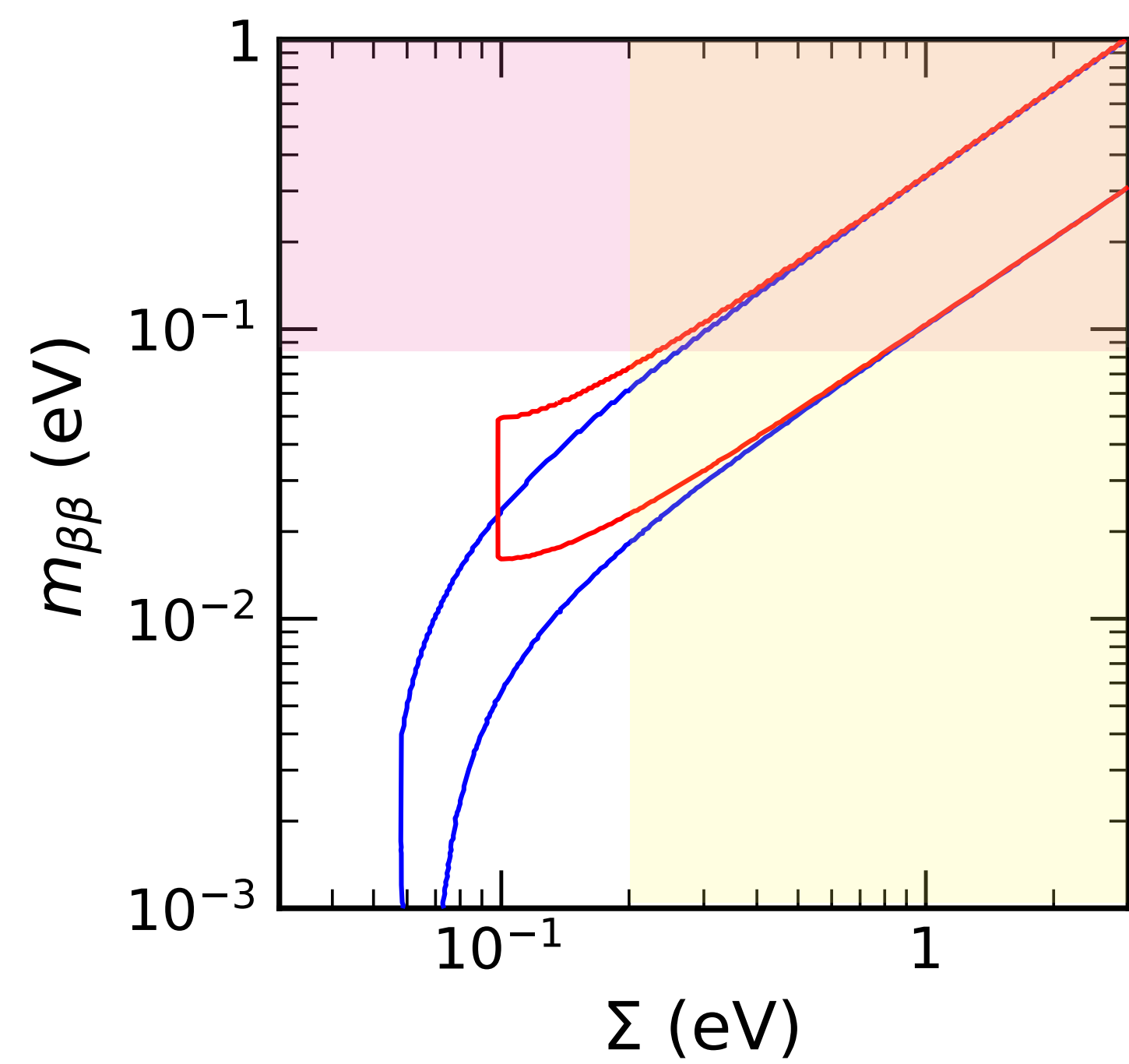


Astrophysics and Cosmology  
CMB, BAO, lensing, ...

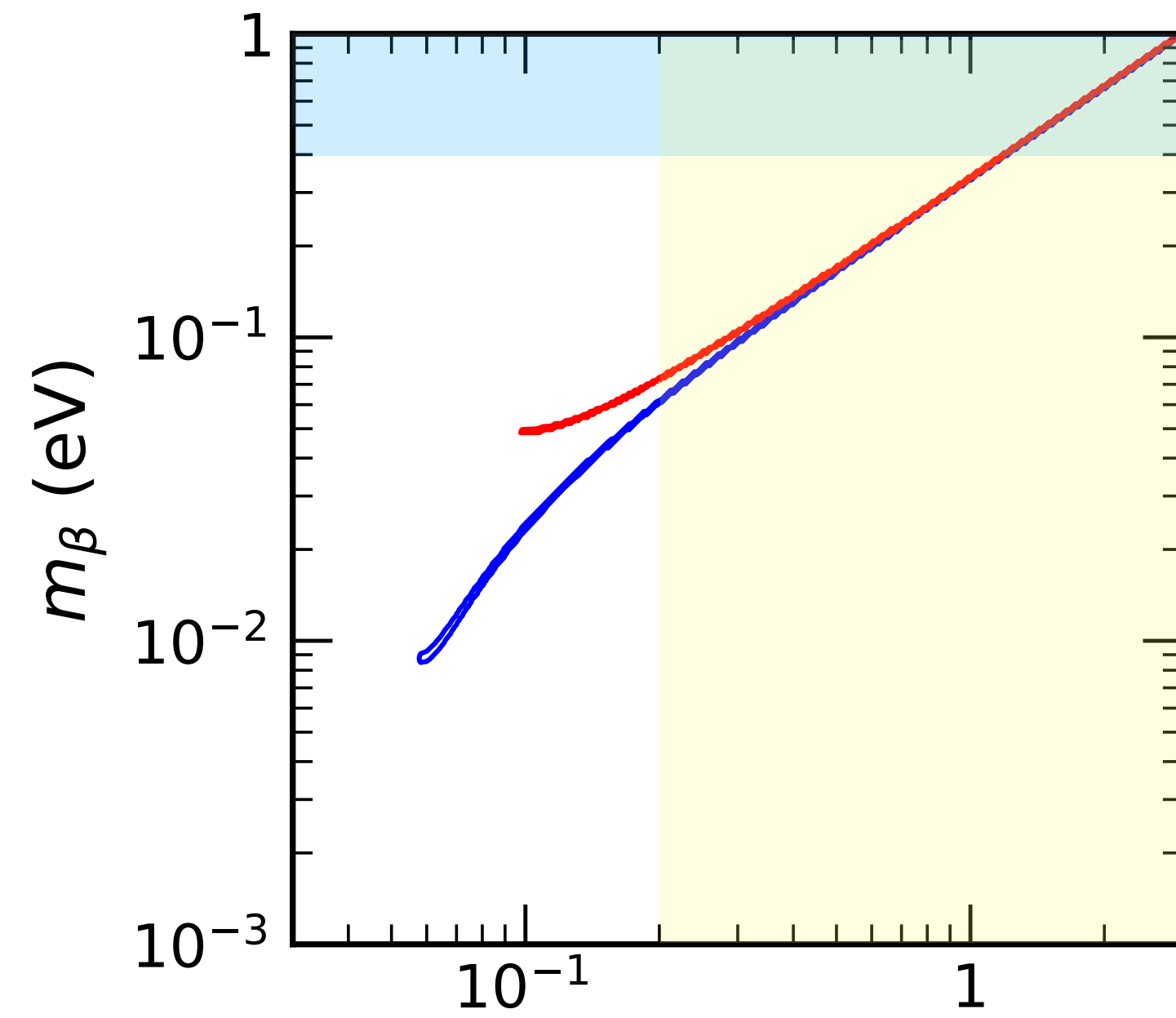




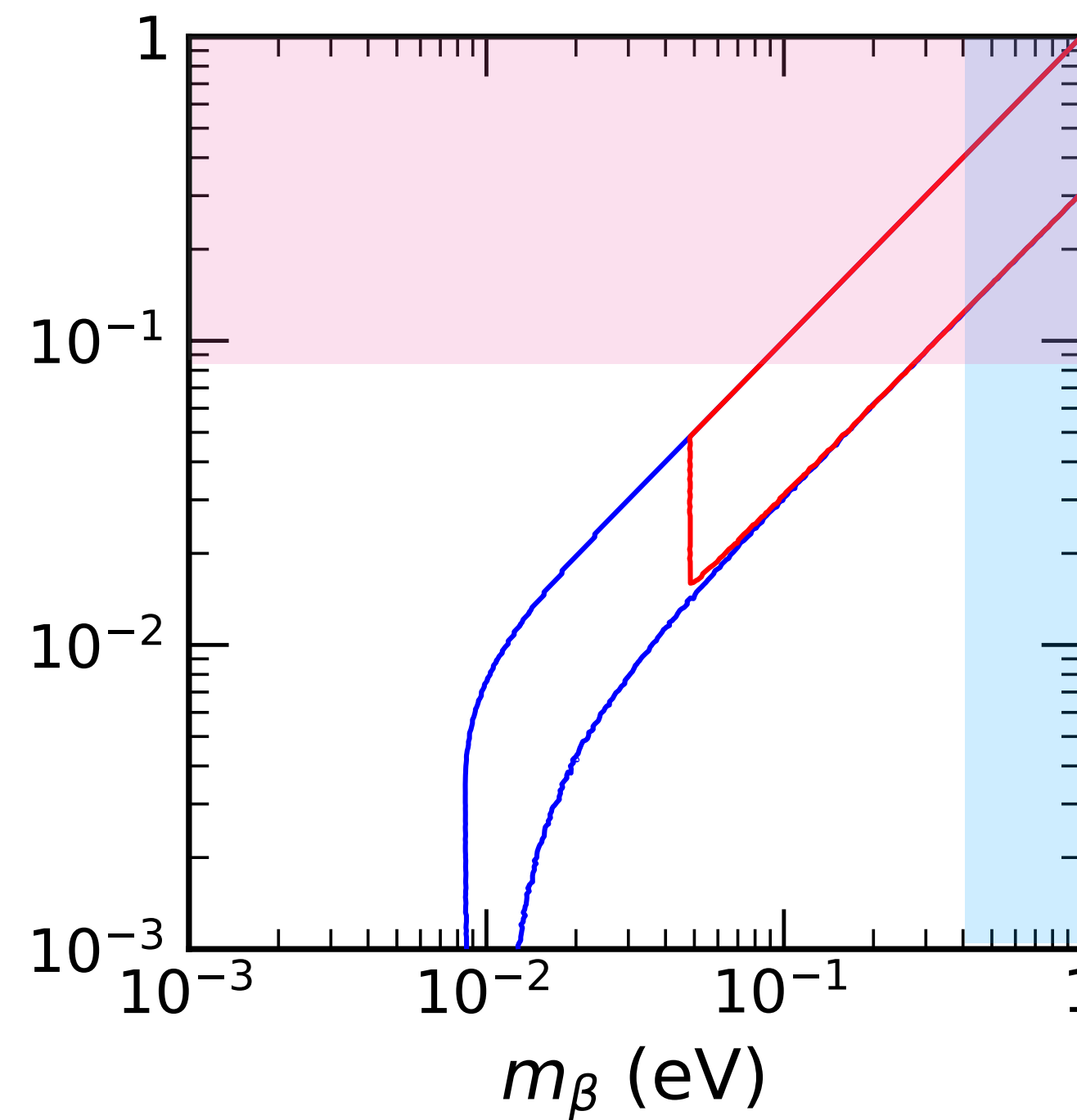
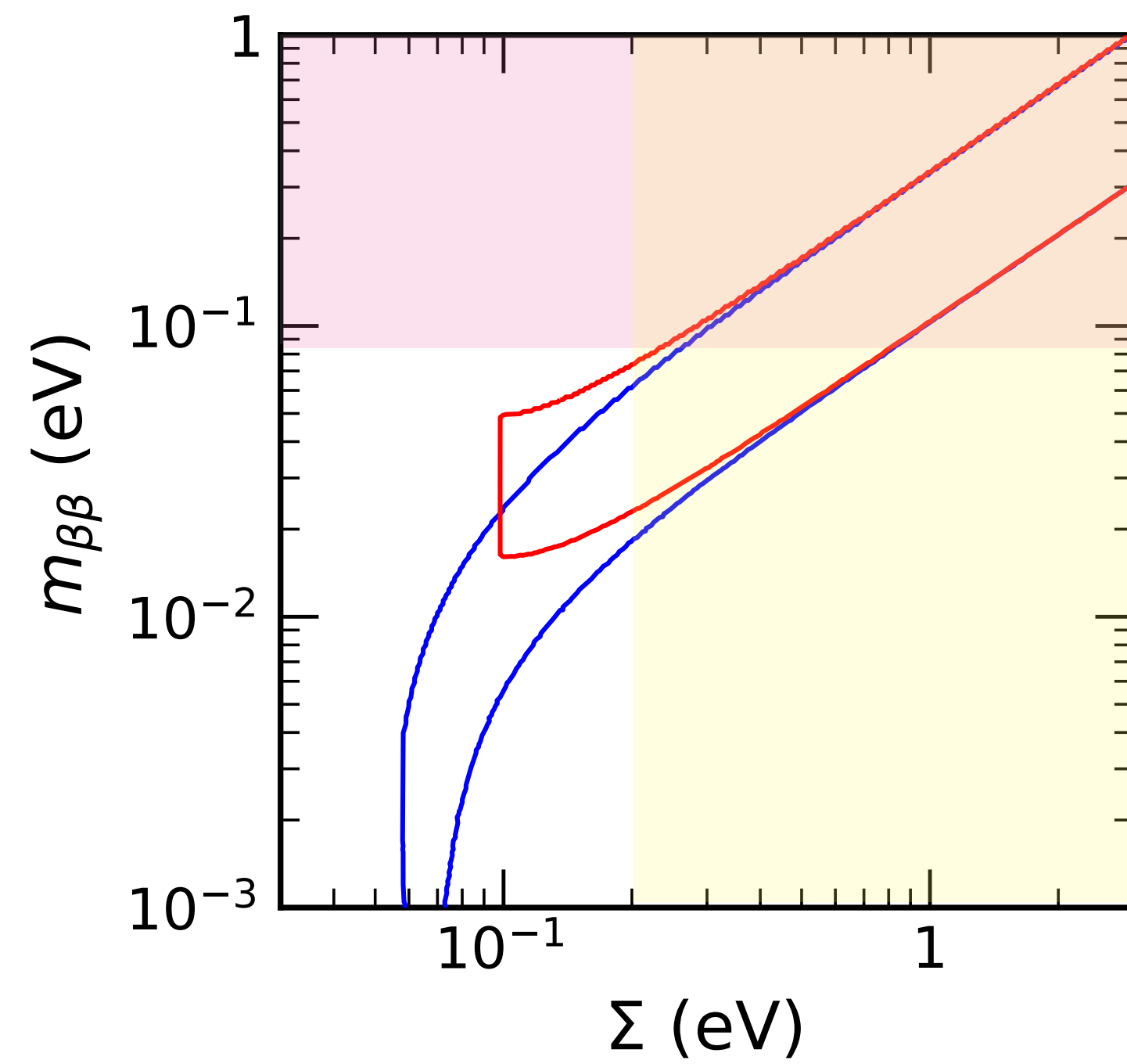
Normal Ordering (2 $\sigma$ )  
Inverted Ordering (2 $\sigma$ )



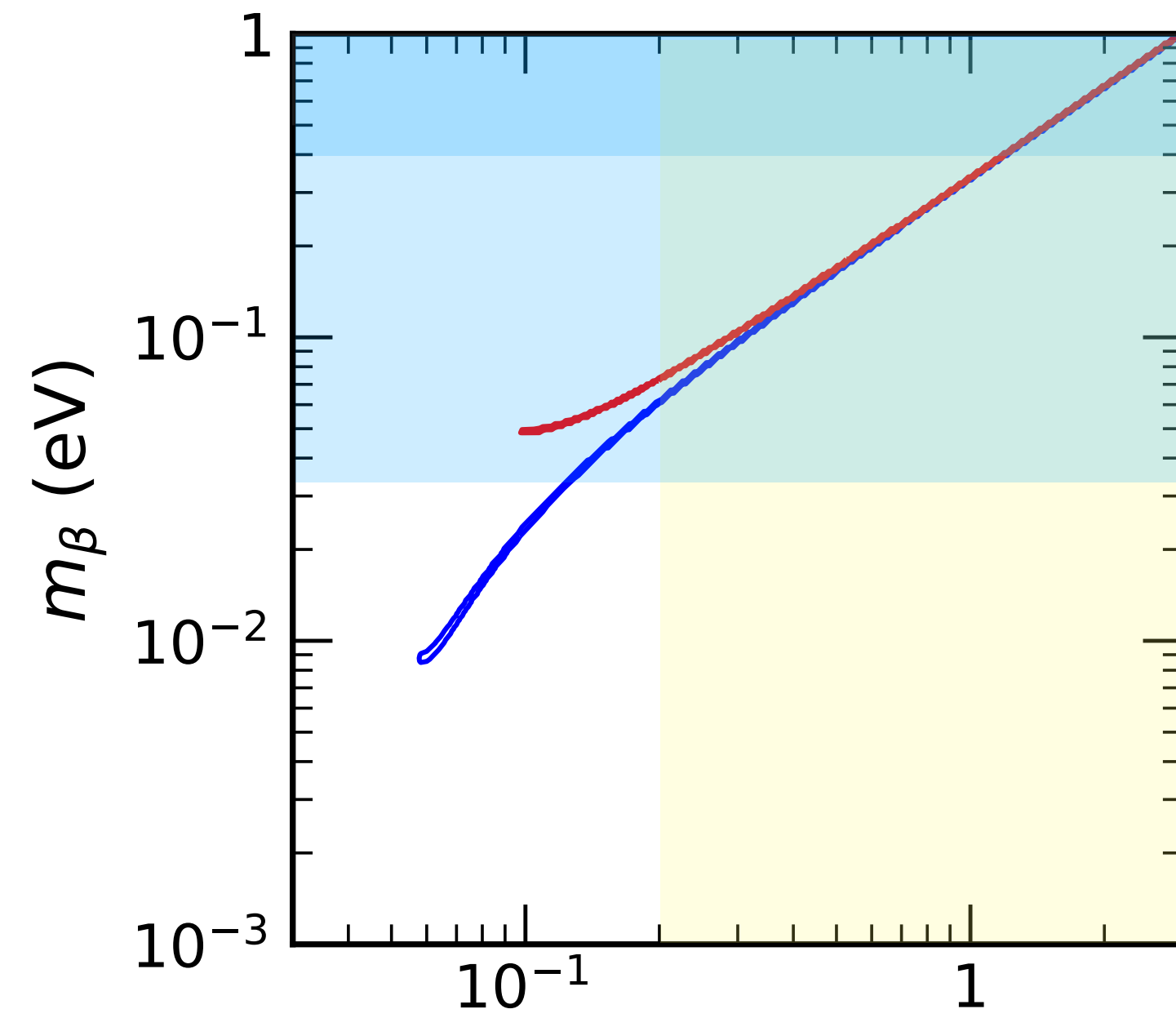
# Next-generation Projects ( $\gtrsim 10$ years)



Normal Ordering ( $2\sigma$ )  
Inverted Ordering ( $2\sigma$ )

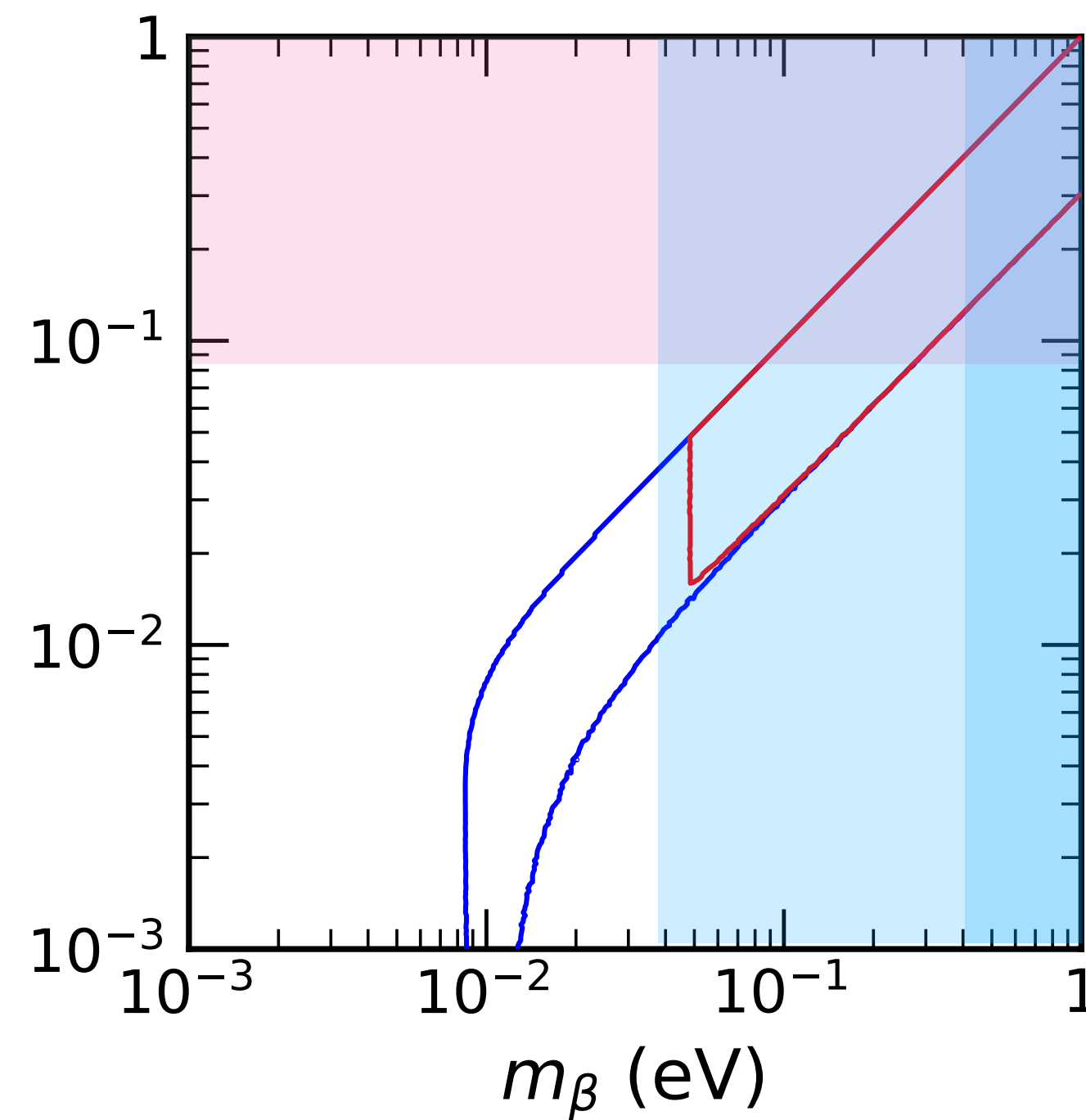
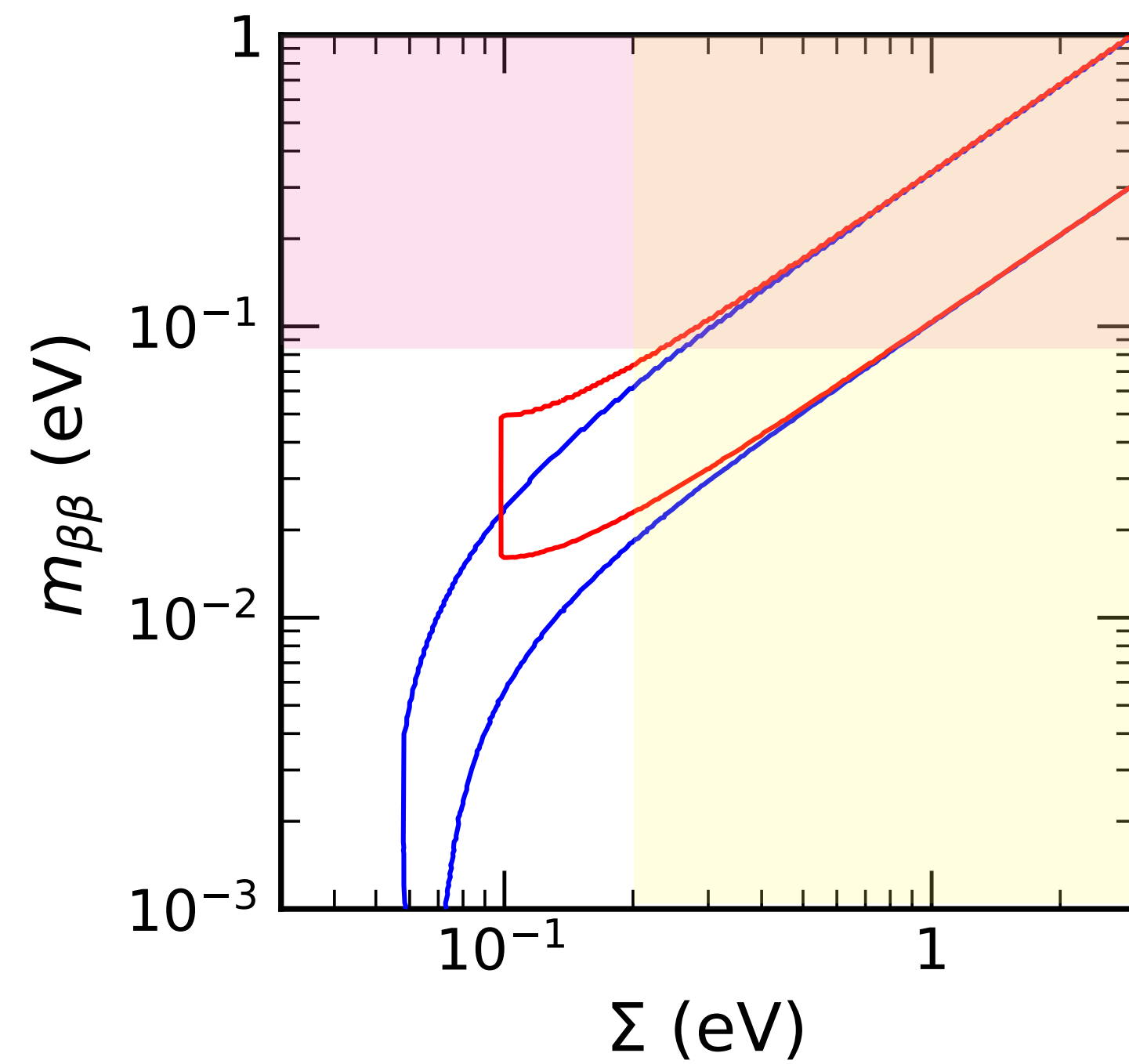


# Next-generation Projects ( $\gtrsim 10$ years)

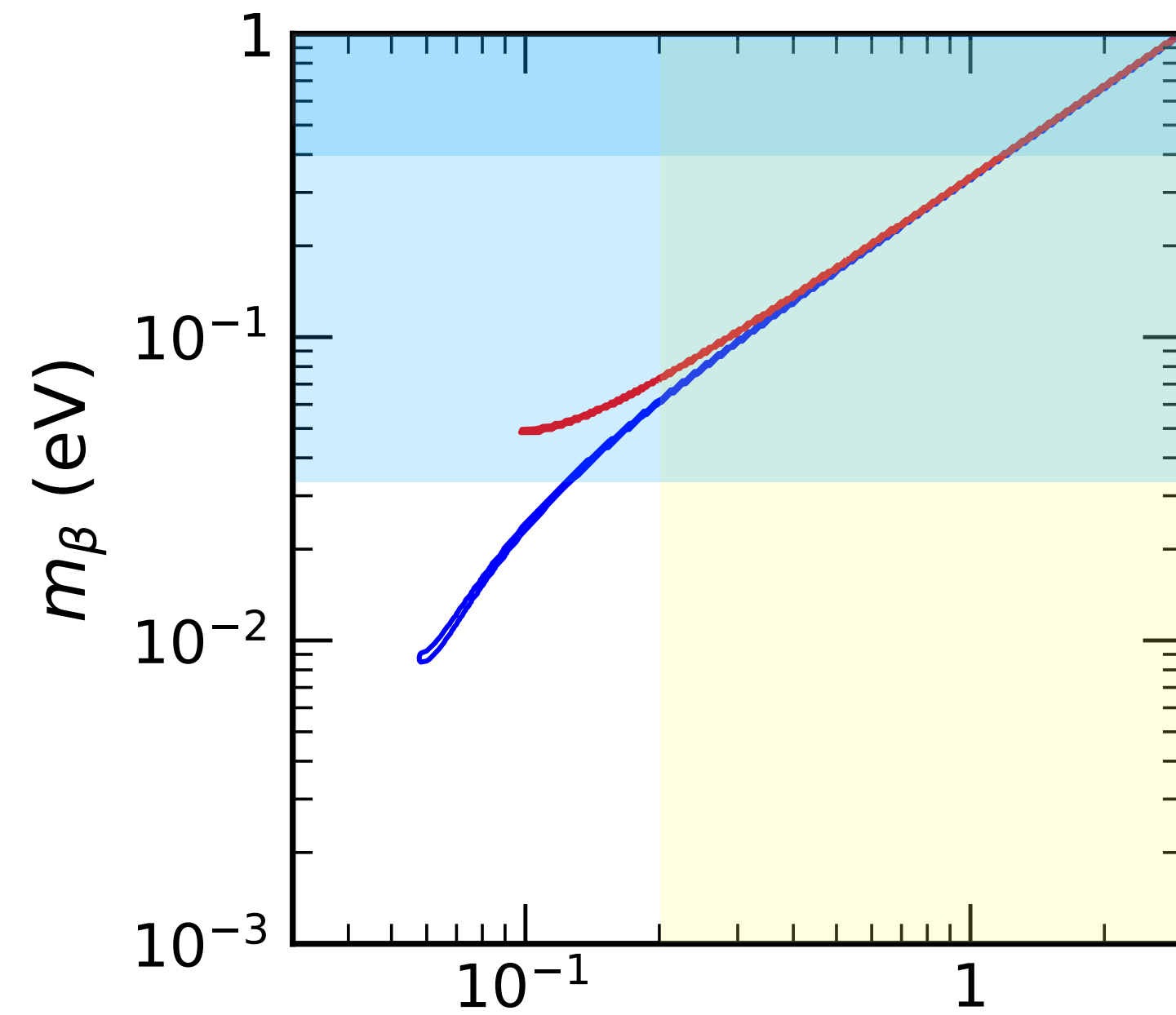


Normal Ordering ( $2\sigma$ )  
Inverted Ordering ( $2\sigma$ )

$\beta$  decay - Project-8, ...

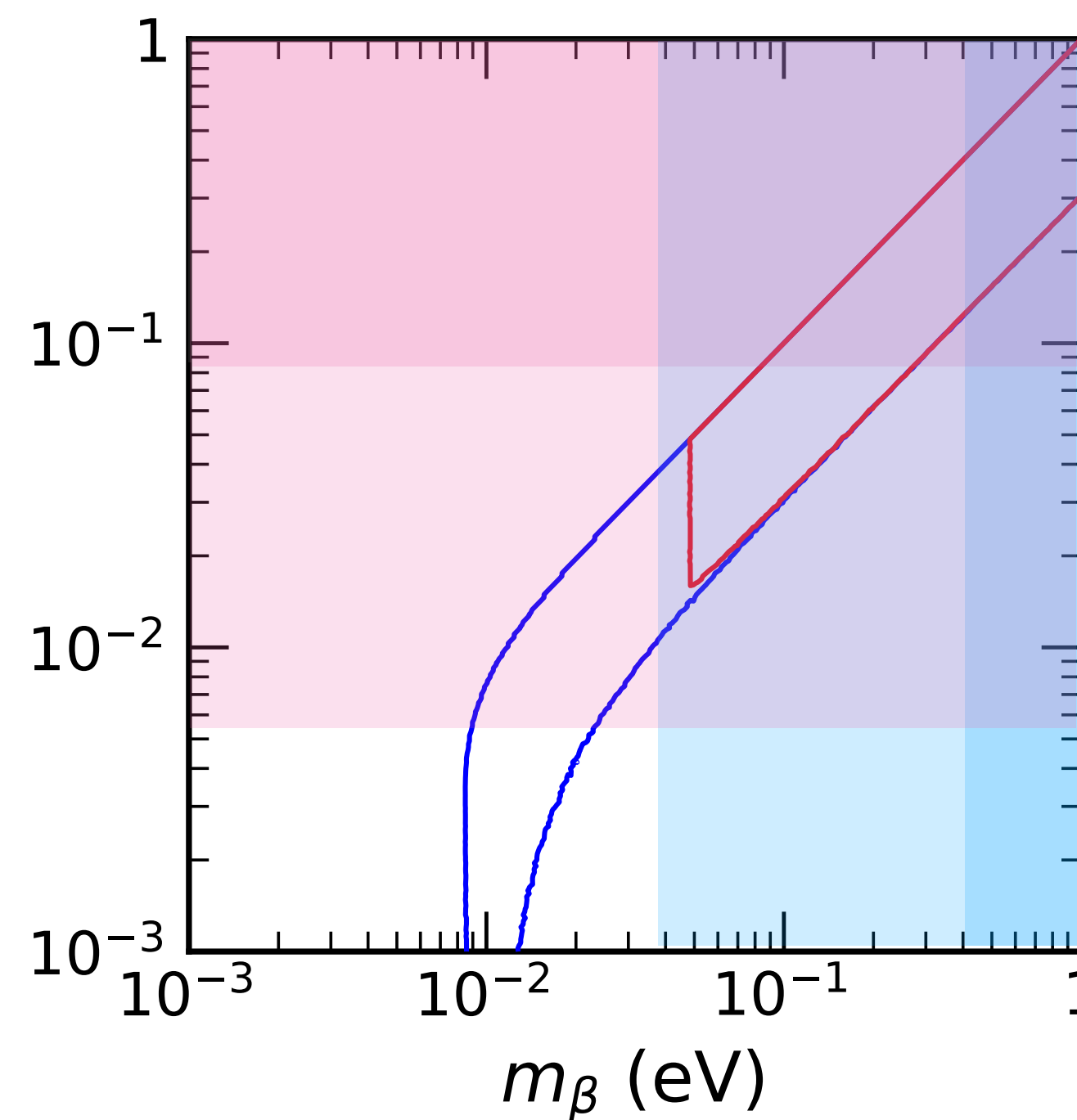
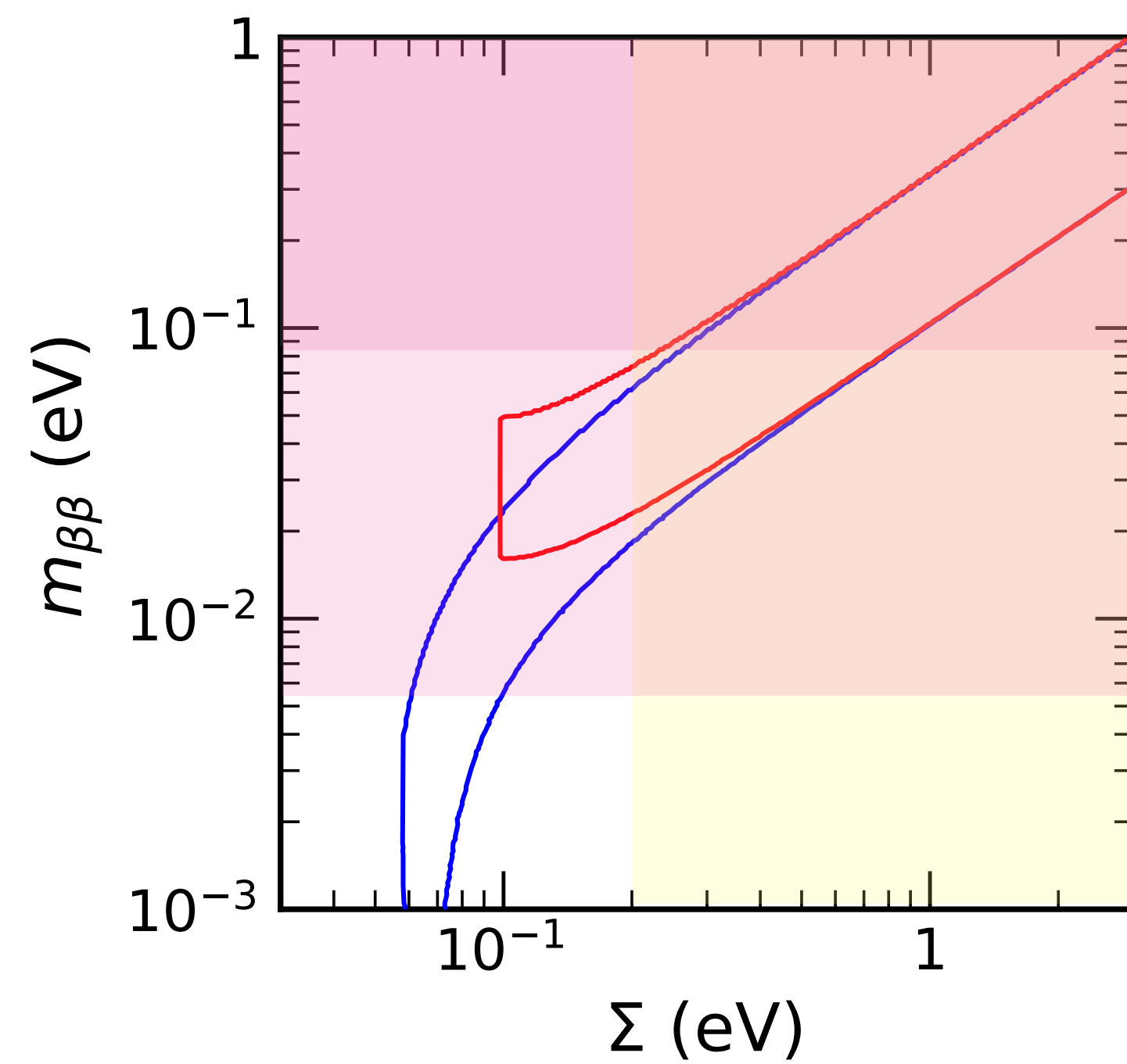


# Next-generation Projects ( $\gtrsim 10$ years)



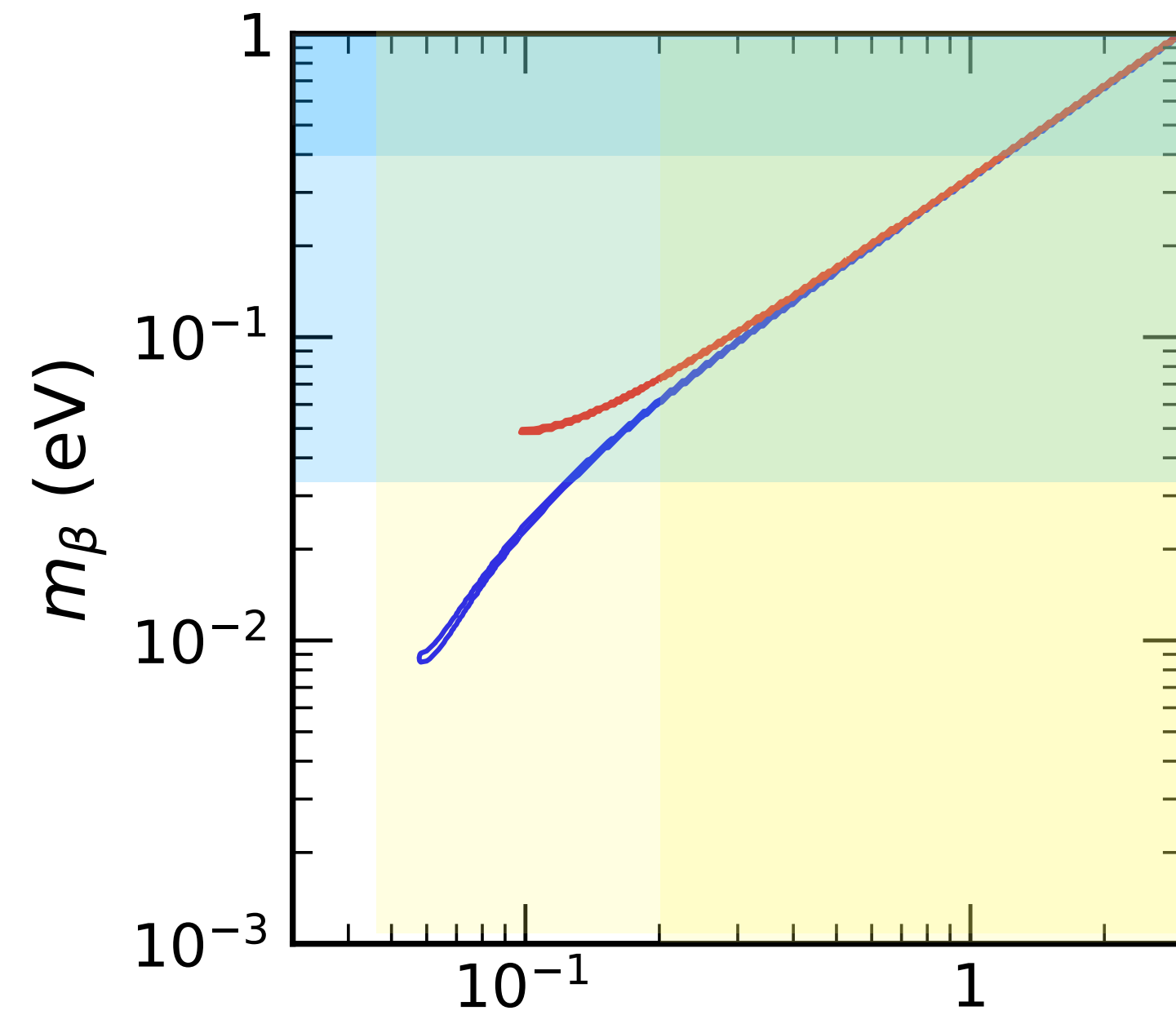
Normal Ordering ( $2\sigma$ )  
Inverted Ordering ( $2\sigma$ )

$\beta$  decay - Project-8, ...



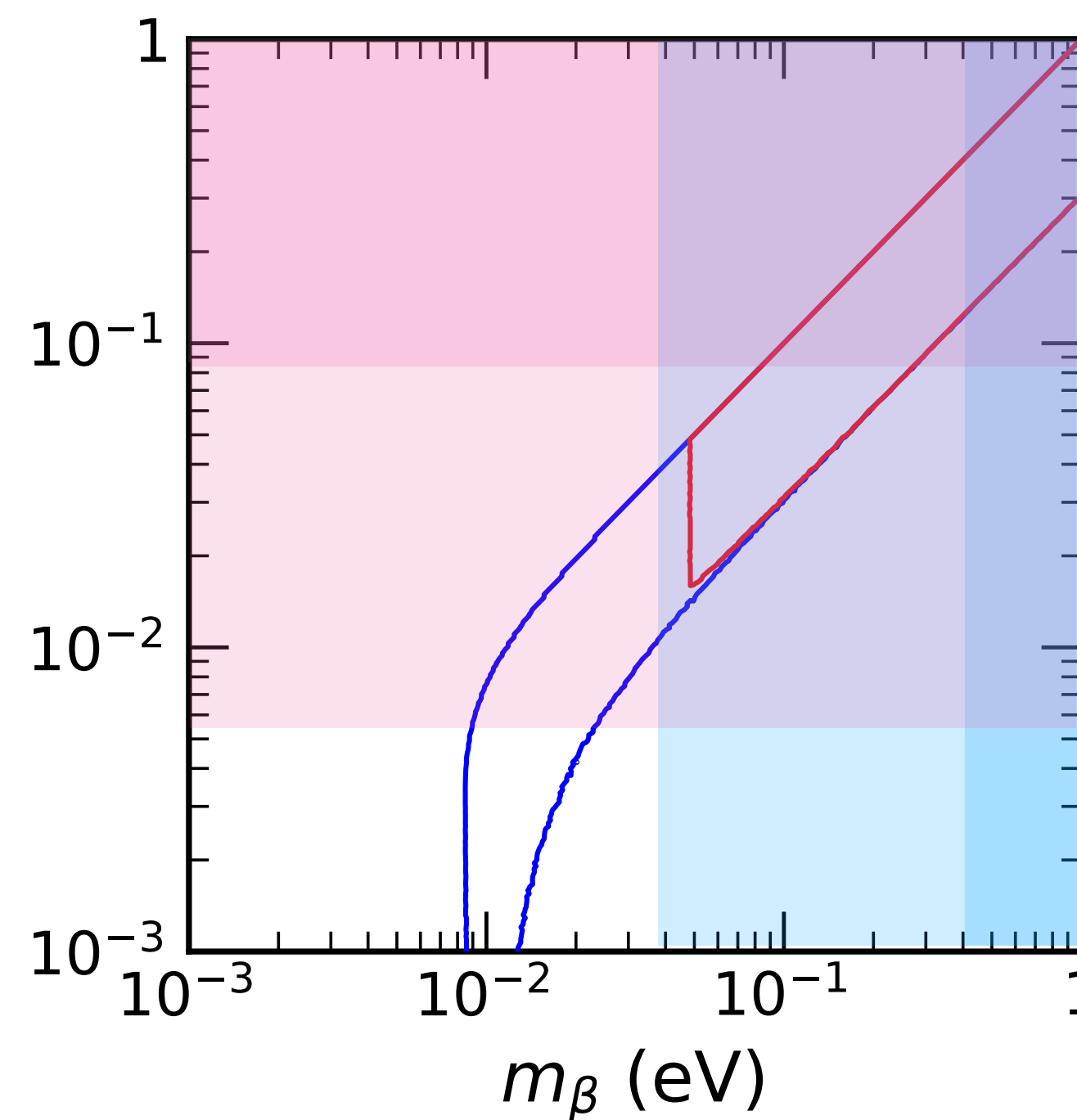
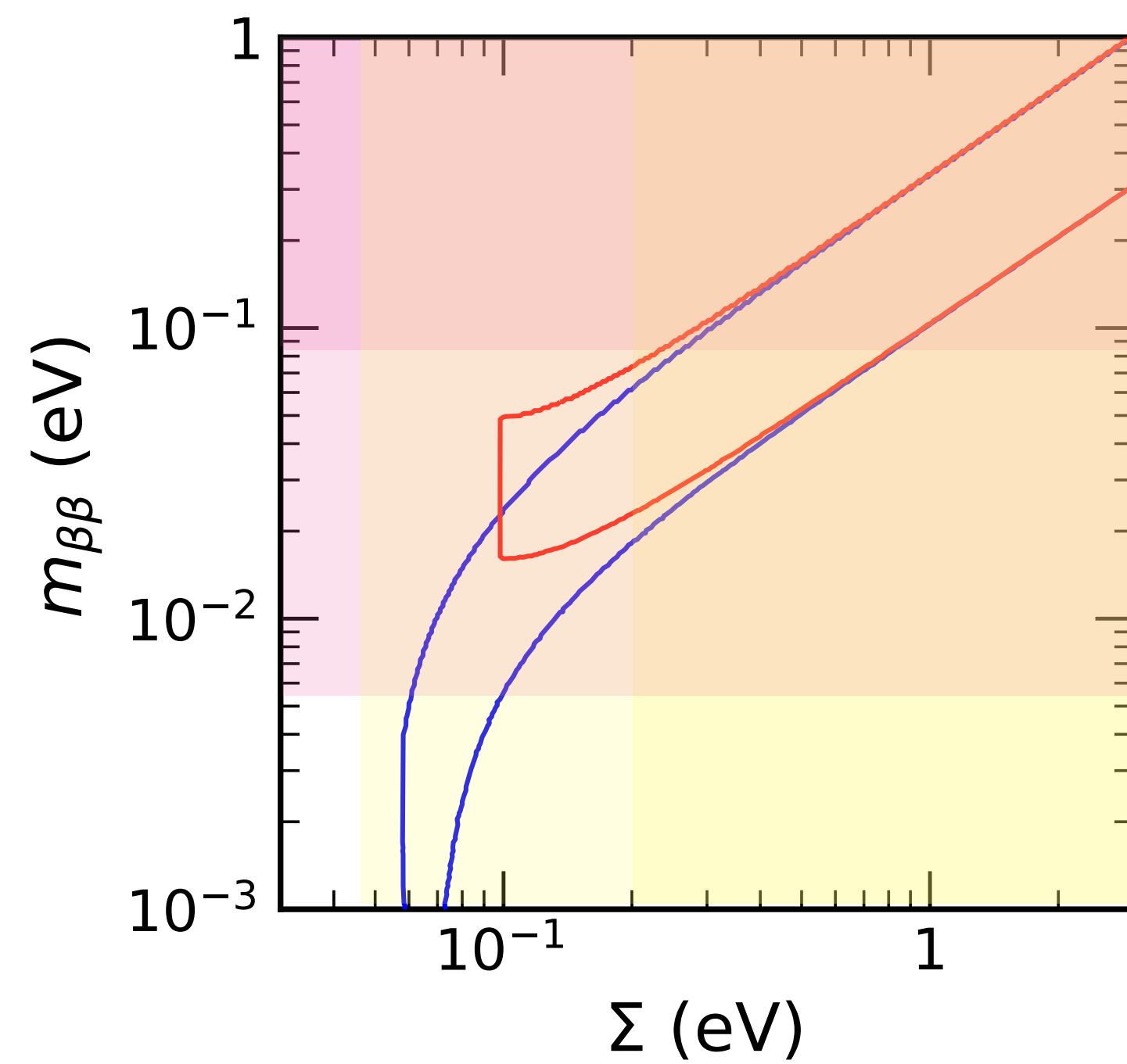
$0\nu\beta\beta$  decay  
LEGEND, NEXO, CUPID, ...

# Next-generation Projects ( $\gtrsim 10$ years)



Normal Ordering ( $2\sigma$ )  
Inverted Ordering ( $2\sigma$ )

$\beta$  decay - Project-8, ...



$0\nu\beta\beta$  decay  
LEGEND, NEXO, CUPID, ...

Astrophysics and Cosmology  
EUCLID, ...



Mass ordering can be determined by exploiting interference between oscillations driven by  $\Delta m^2$  and those driven by a second term  $Q$  whose sign is known



Mass ordering can be determined by exploiting interference between oscillations driven by  $\Delta m^2$  and those driven by a second term  $Q$  whose sign is known

$$Q \propto \delta m^2$$

medium-baseline reactors

Mass ordering can be determined by exploiting interference between oscillations driven by  $\Delta m^2$  and those driven by a second term  $Q$  whose sign is known

$$Q \propto \delta m^2$$

medium-baseline reactors

$$Q \propto G_F E N_e$$

matter effects in accelerator/atmospheric  $\nu$

Mass ordering can be determined by exploiting interference between oscillations driven by  $\Delta m^2$  and those driven by a second term  $Q$  whose sign is known

$$Q \propto \delta m^2$$

medium-baseline reactors

$$Q \propto G_F E N_e$$

matter effects in accelerator/atmospheric  $\nu$

$$Q \propto G_F E N_\nu$$

self-interaction effects in supernovae

Mass ordering can be determined by exploiting interference between oscillations driven by  $\Delta m^2$  and those driven by a second term  $Q$  whose sign is known

$$Q \propto \delta m^2$$

medium-baseline reactors

$$Q \propto G_F E N_e$$

matter effects in accelerator/atmospheric  $\nu$

$$Q \propto G_F E N_\nu$$

self-interaction effects in supernovae

Synergy across  $|\Delta m^2|$  determinations from reactor, accelerator, and atmospheric data: measurements converge in the true ordering and separate in the wrong one

Mass ordering can be determined by exploiting interference between oscillations driven by  $\Delta m^2$  and those driven by a second term  $Q$  whose sign is known

$$Q \propto \delta m^2$$

medium-baseline reactors

$$Q \propto G_F E N_e$$

matter effects in accelerator/atmospheric  $\nu$

$$Q \propto G_F E N_\nu$$

self-interaction effects in supernovae

Synergy across  $|\Delta m^2|$  determinations from reactor, accelerator, and atmospheric data: measurements converge in the true ordering and separate in the wrong one

In particular, JUNO will be sensitive to

$$\Delta m_{ee}^2 = |\Delta m^2| + \frac{1}{2}\alpha(\cos^2 \theta_{12} - \sin^2 \theta_{12})\delta m^2$$

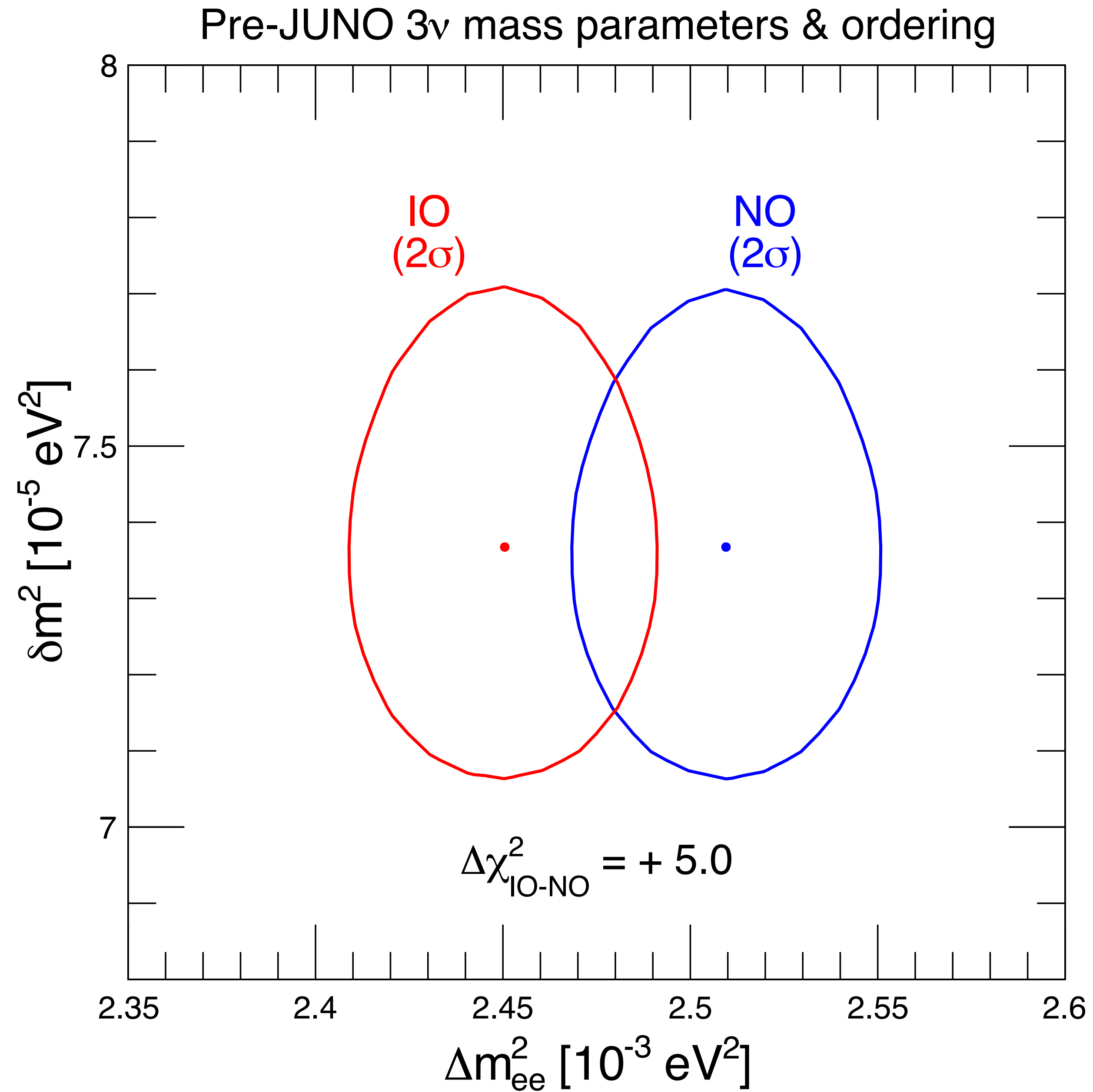
NO:  $\alpha = +1$

IO:  $\alpha = -1$



# Present knowledge about the two JUNO oscillation frequencies

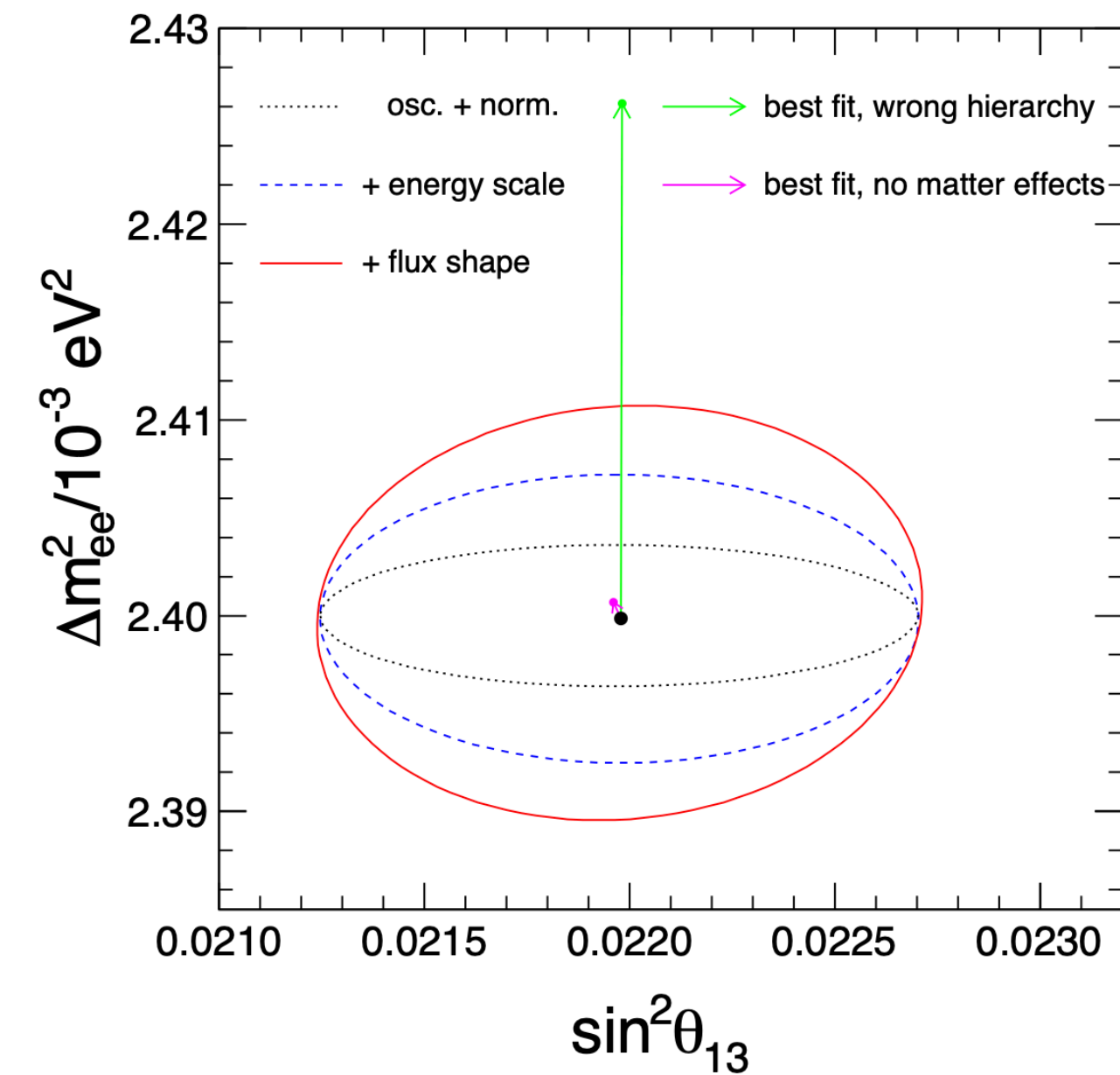
# Present knowledge about the two JUNO oscillation frequencies



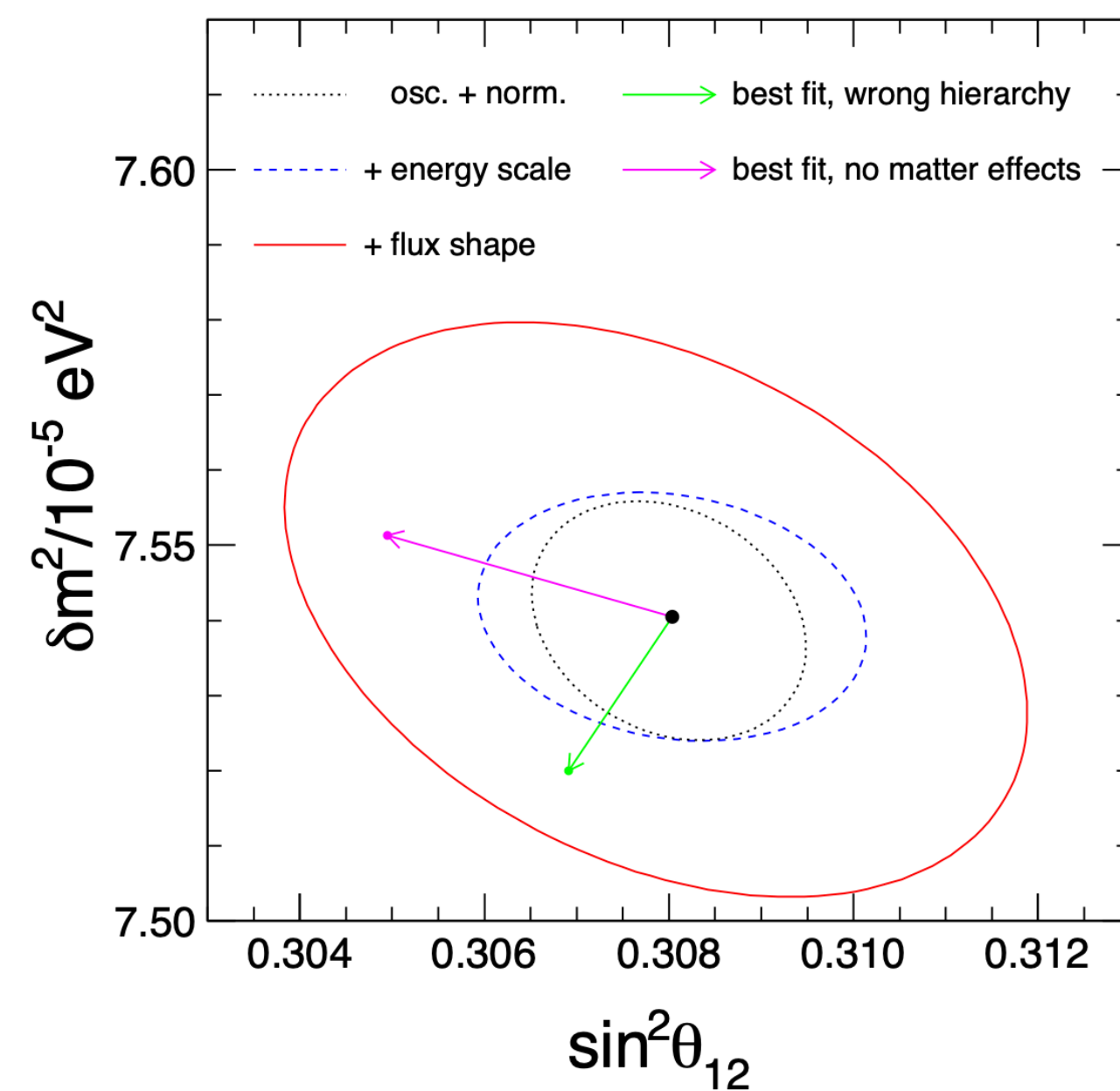


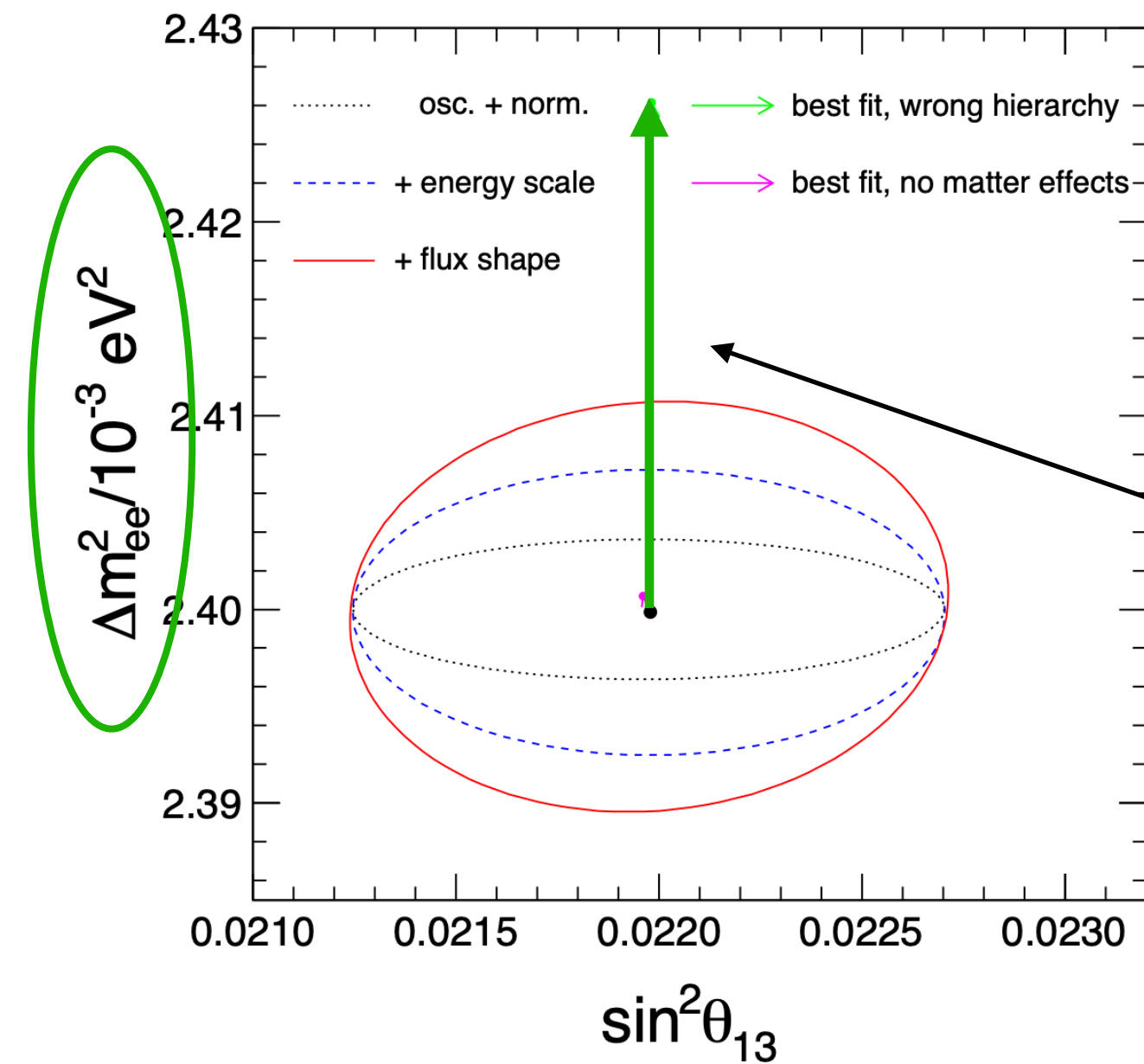


JUNO measurements will lead to  
slightly displaced best fits for the  
two frequencies in NO and IO



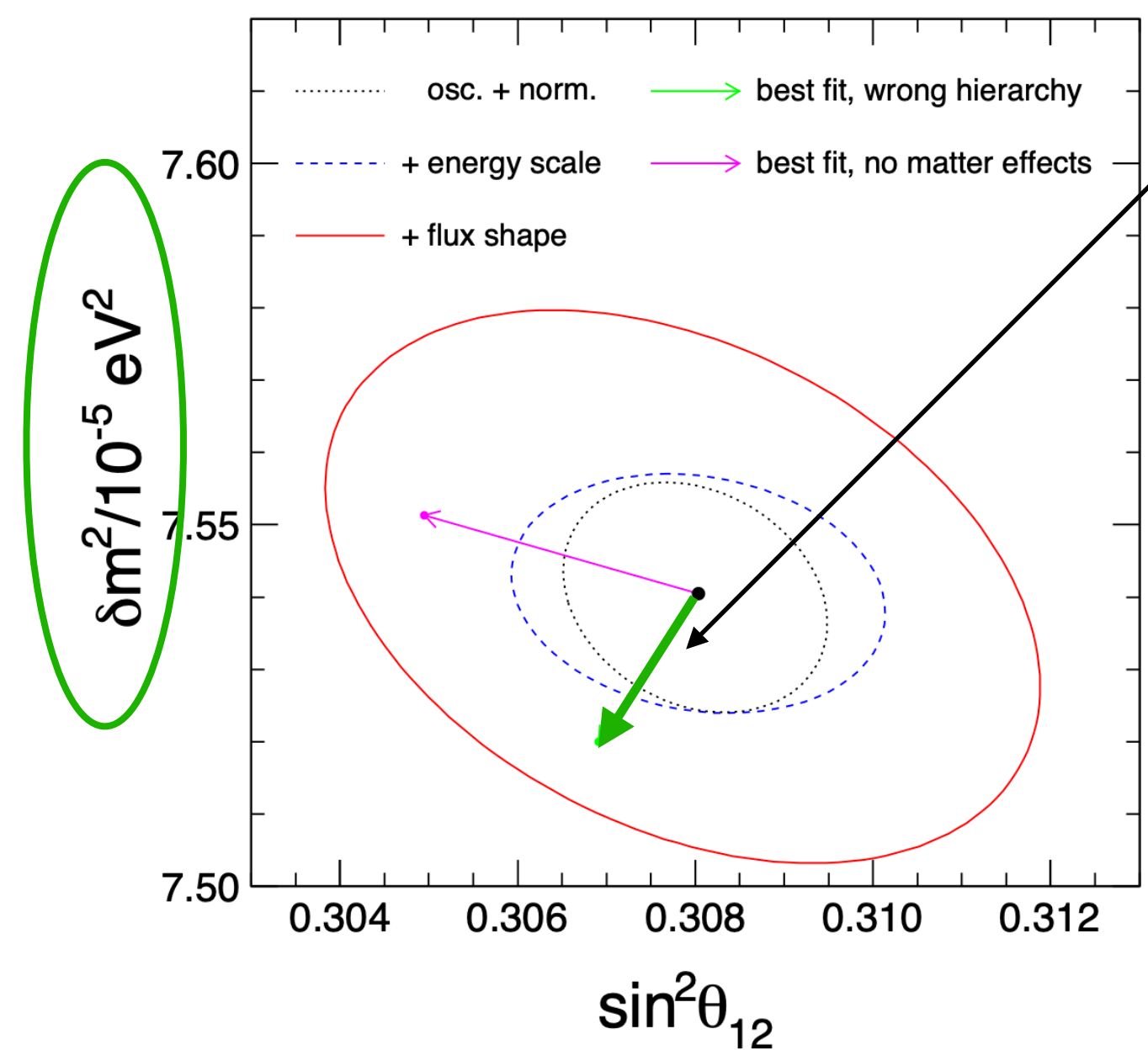
JUNO measurements will lead to  
 slightly displaced best fits for the  
 two frequencies in NO and IO

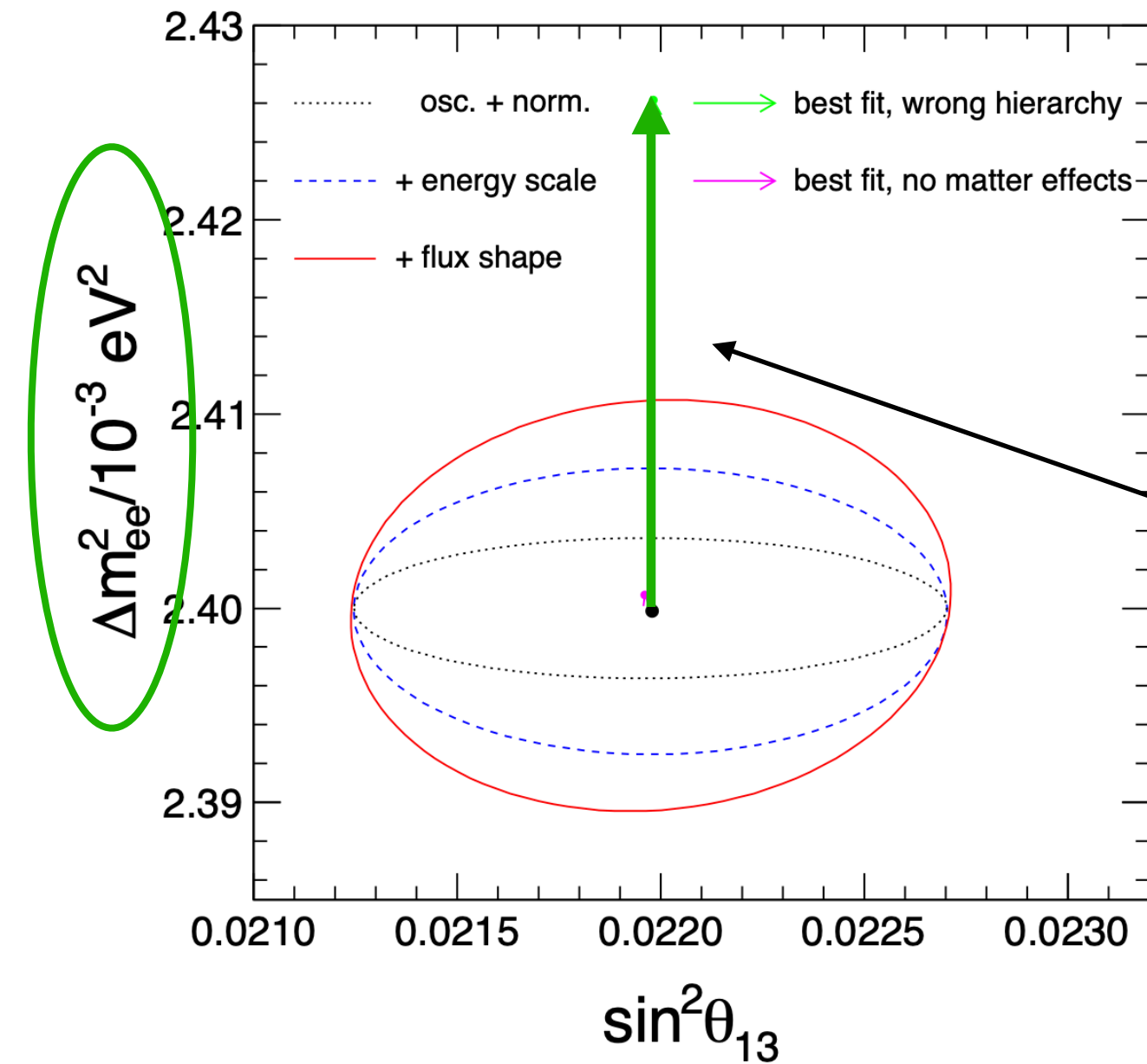




Green arrows: shifts of best fits when passing from NO to IO assumption

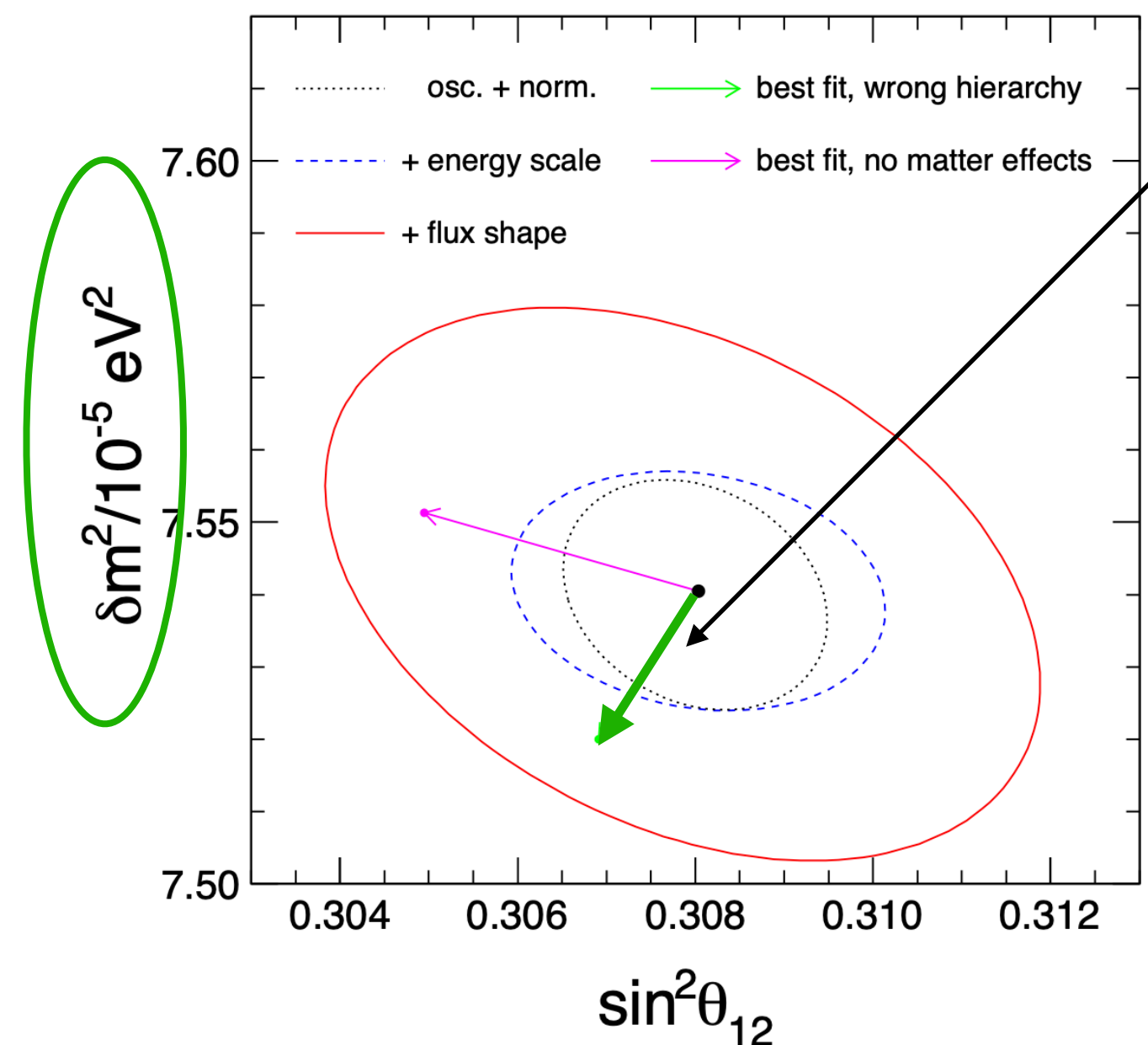
JUNO measurements will lead to slightly displaced best fits for the two frequencies in NO and IO





Green arrows: shifts of best fits when passing from NO to IO assumption

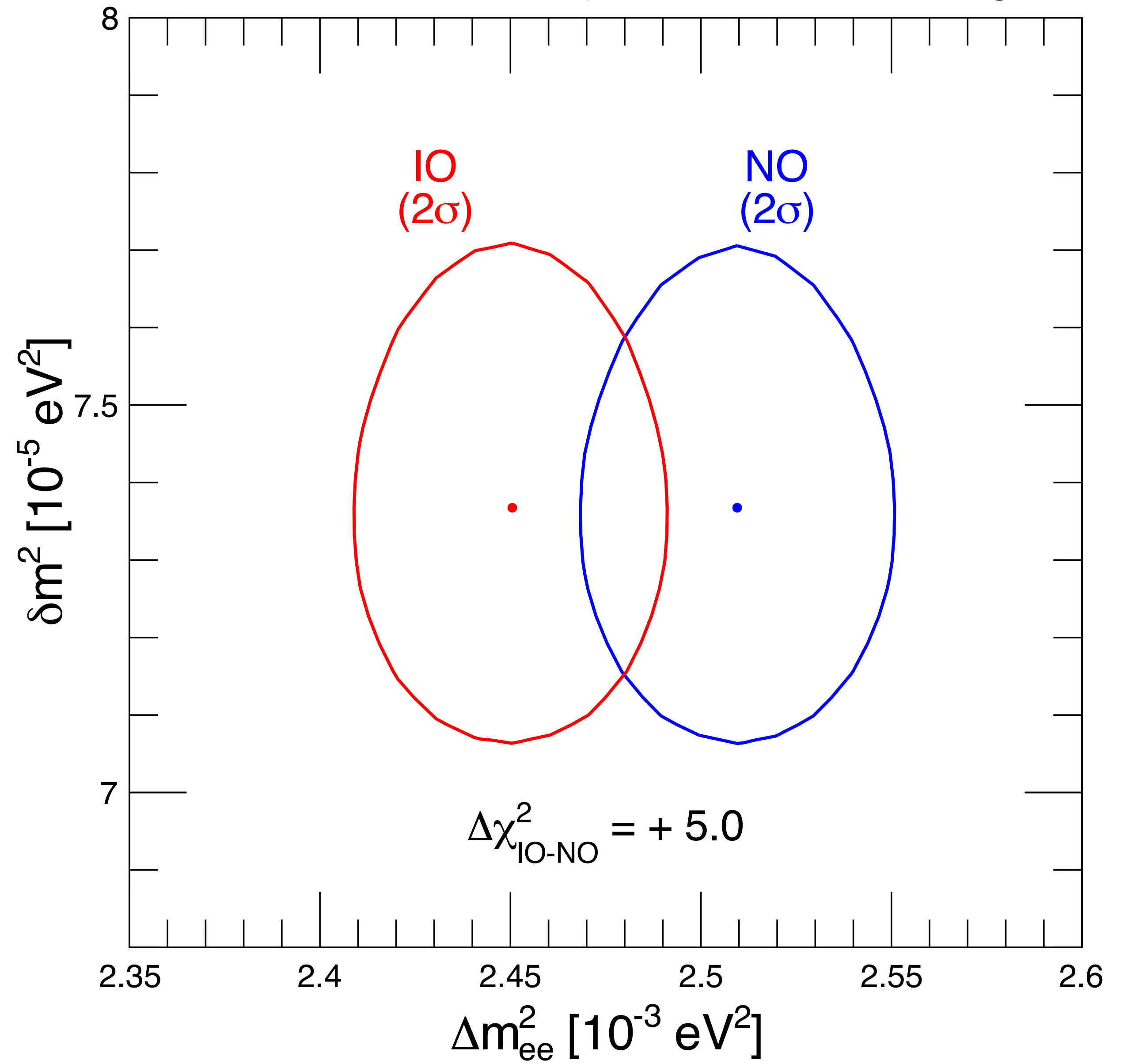
JUNO measurements will lead to slightly displaced best fits for the two frequencies in NO and IO



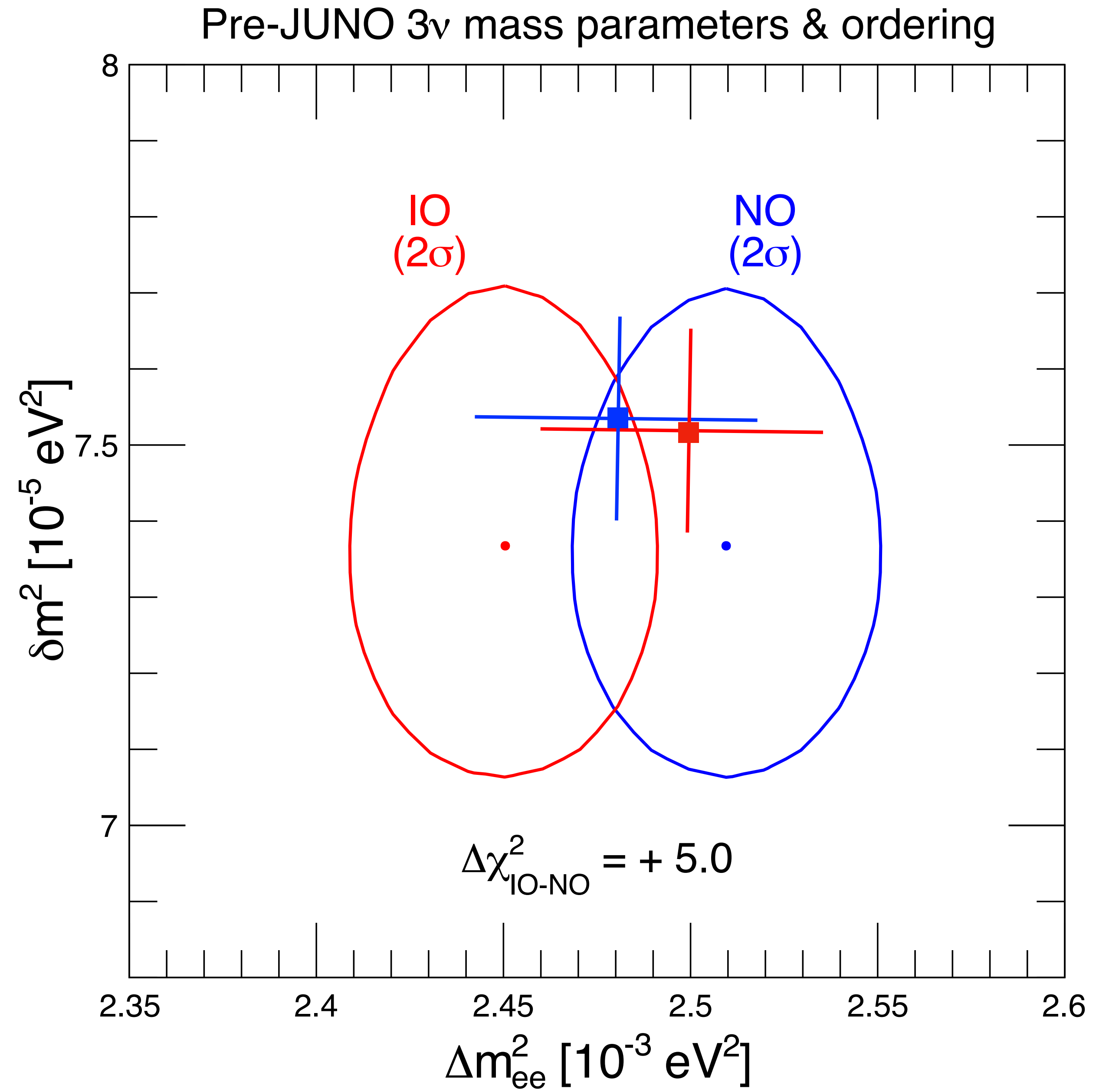
Shift of  $\Delta m^2$  discussed in many papers (see Parke et al. for instance)  
Specific values depend little on fit details.



# Pre-JUNO 3ν mass parameters & ordering



Typical relative displacement between  
JUNO bestfit point in **NO** and **IO**





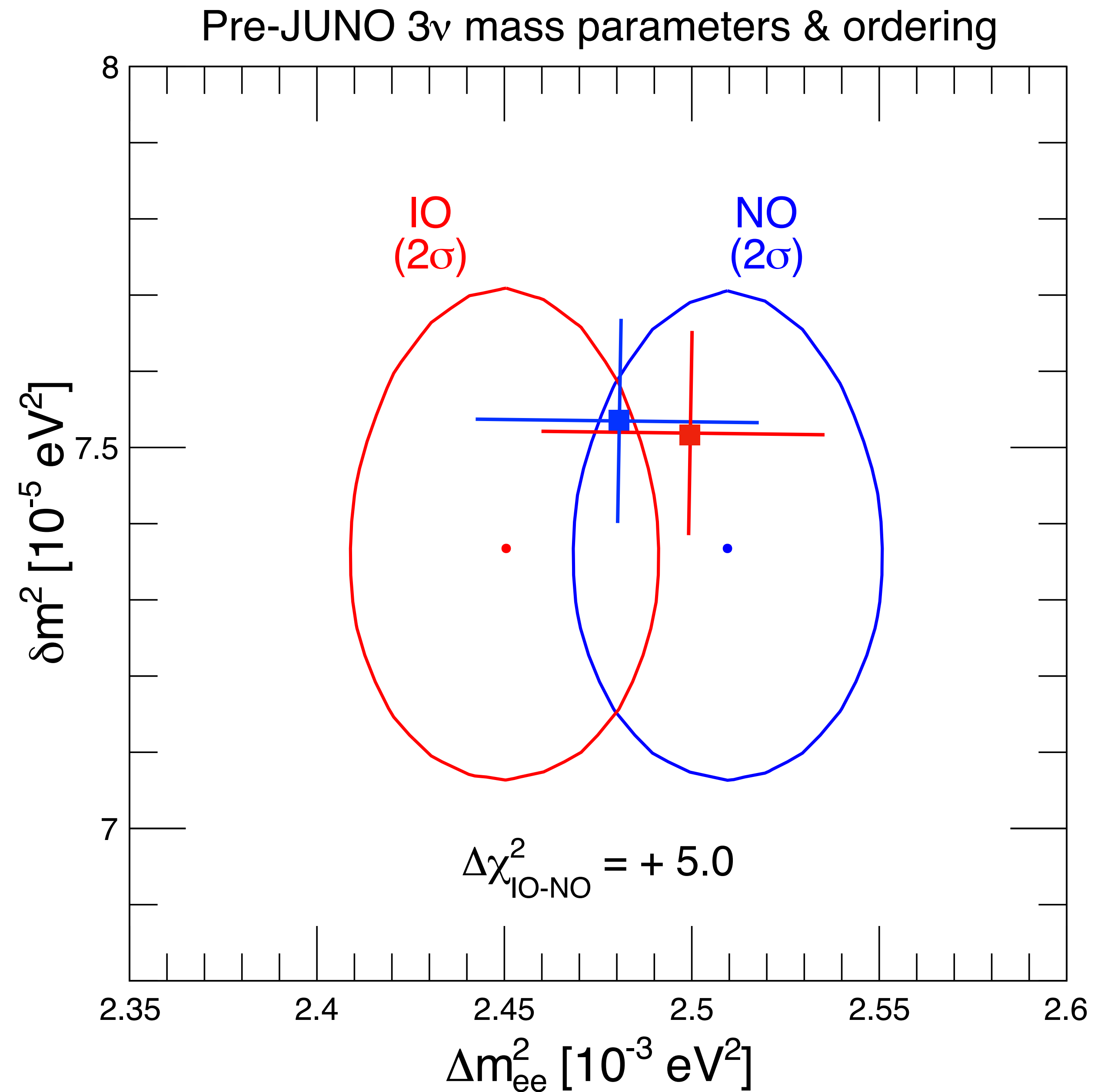
Typical relative displacement between  
JUNO bestfit point in **NO** and **IO**

$$\delta m^2 (2\sigma) \sim 0.15 \times 10^{-5} \text{ eV}^2$$

$$\Delta m_{ee}^2 (2\sigma) \sim 0.04 \times 10^{-3} \text{ eV}^2$$

Sub-percent precision measurement of neutrino  
oscillation parameters with JUNO<sup>\*</sup>

To cite this article: Angel Abusleme *et al* 2022 *Chinese Phys. C* 46 123001



Typical relative displacement between JUNO bestfit point in **NO** and **IO**

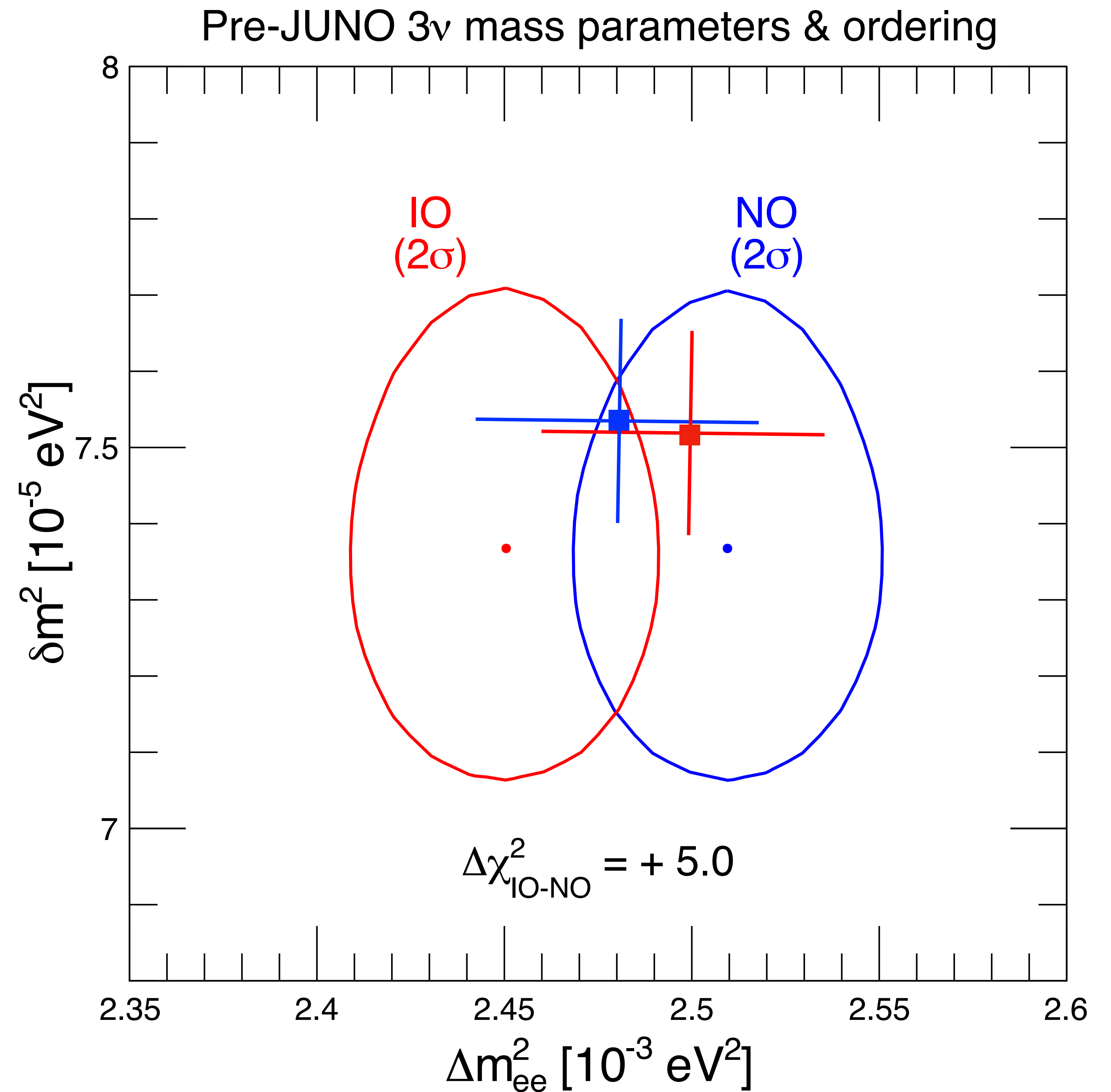
$$\delta m^2 (2\sigma) \sim 0.15 \times 10^{-5} \text{ eV}^2$$

$$\Delta m_{ee}^2 (2\sigma) \sim 0.04 \times 10^{-3} \text{ eV}^2$$

Sub-percent precision measurement of neutrino oscillation parameters with JUNO<sup>\*</sup>

To cite this article: Angel Abusleme *et al* 2022 *Chinese Phys. C* 46 123001

Relative shift between JUNO best-fit points is “opposite” to pre-JUNO best-fits → Adds to synergy



Typical relative displacement between JUNO bestfit point in **NO** and **IO**

$$\delta m^2 (2\sigma) \sim 0.15 \times 10^{-5} \text{ eV}^2$$

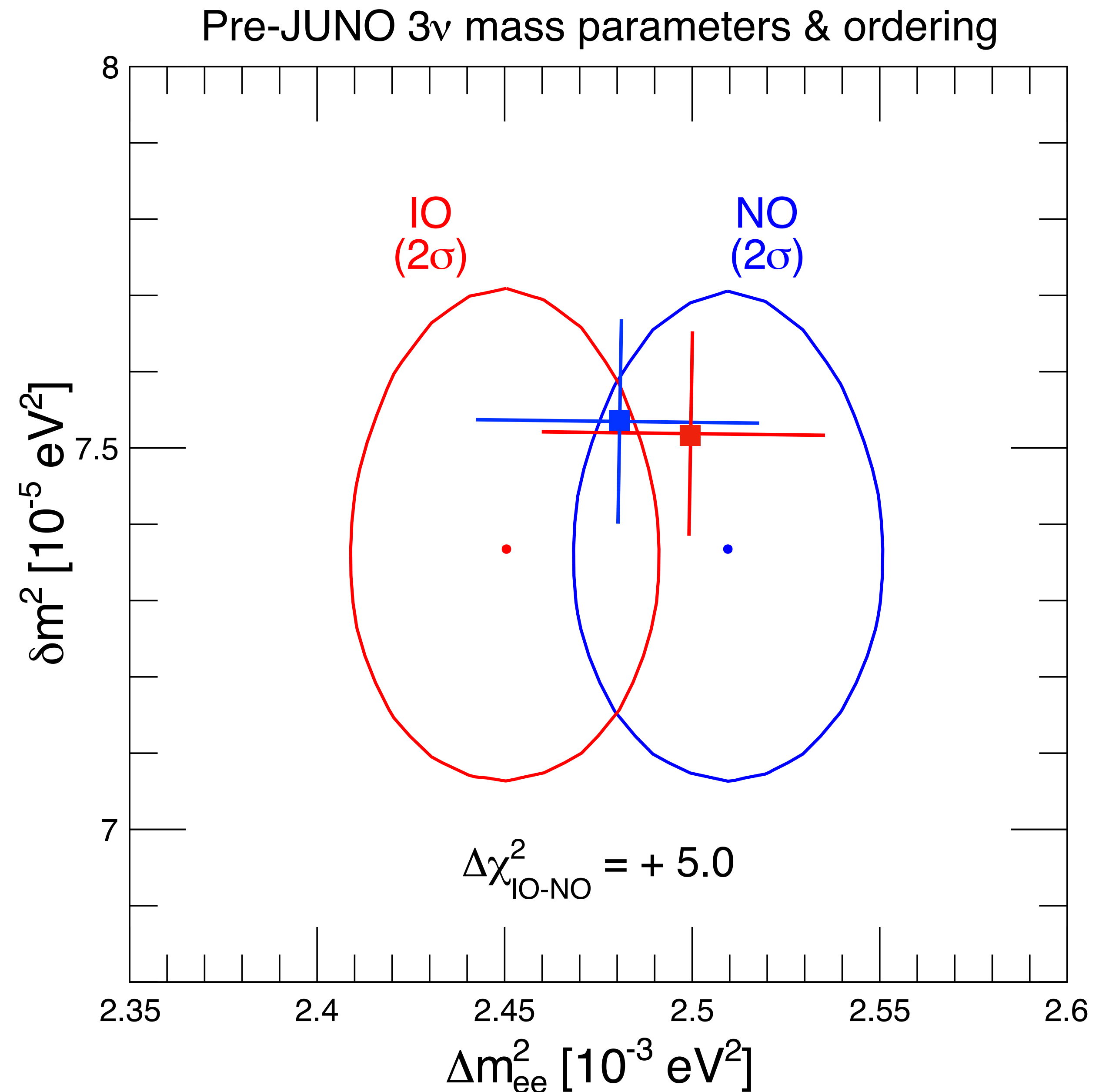
$$\Delta m_{ee}^2 (2\sigma) \sim 0.04 \times 10^{-3} \text{ eV}^2$$

Sub-percent precision measurement of neutrino oscillation parameters with JUNO<sub>\*</sub>

To cite this article: Angel Abusleme *et al* 2022 *Chinese Phys. C* 46 123001

Relative shift between JUNO best-fit points is “opposite” to pre-JUNO best-fits → Adds to synergy

While statistical fluctuations can initially mask the distinction between NO and IO, with higher exposure the true difference is expected to become evident

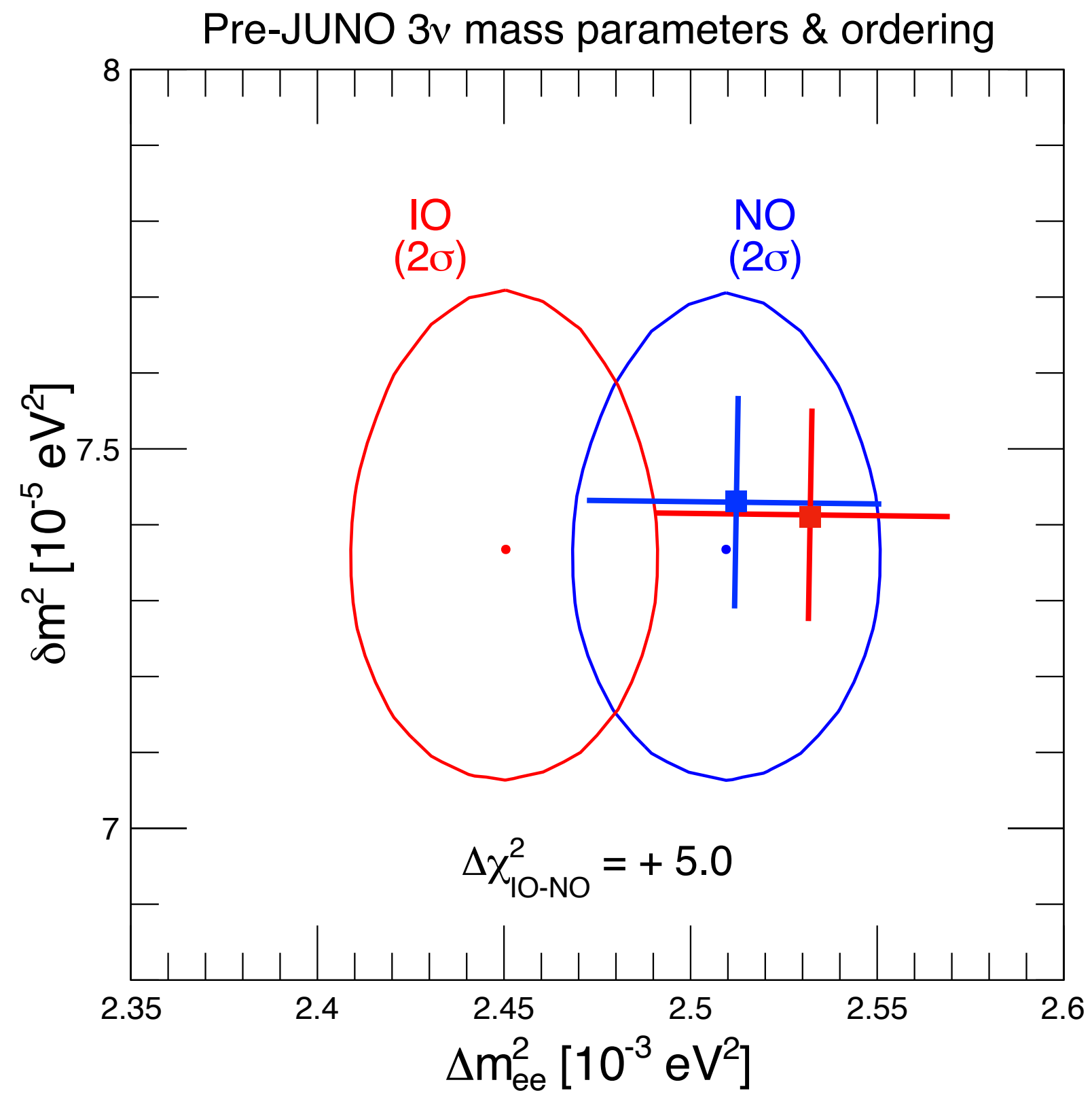




## Examples of possible JUNO first data compared with pre-JUNO data

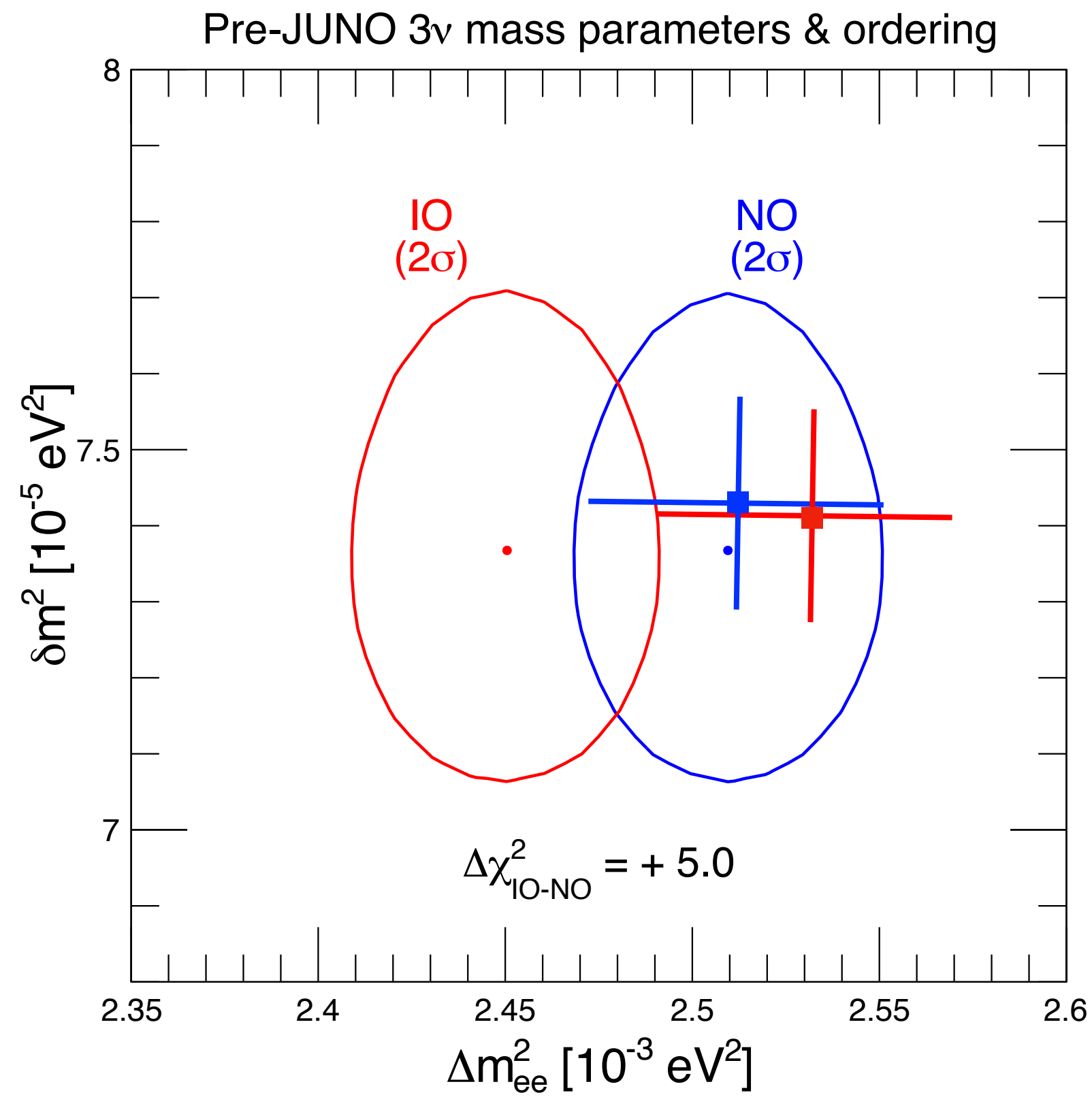
# Examples of possible JUNO first data compared with pre-JUNO data

## A synergy favoring NO

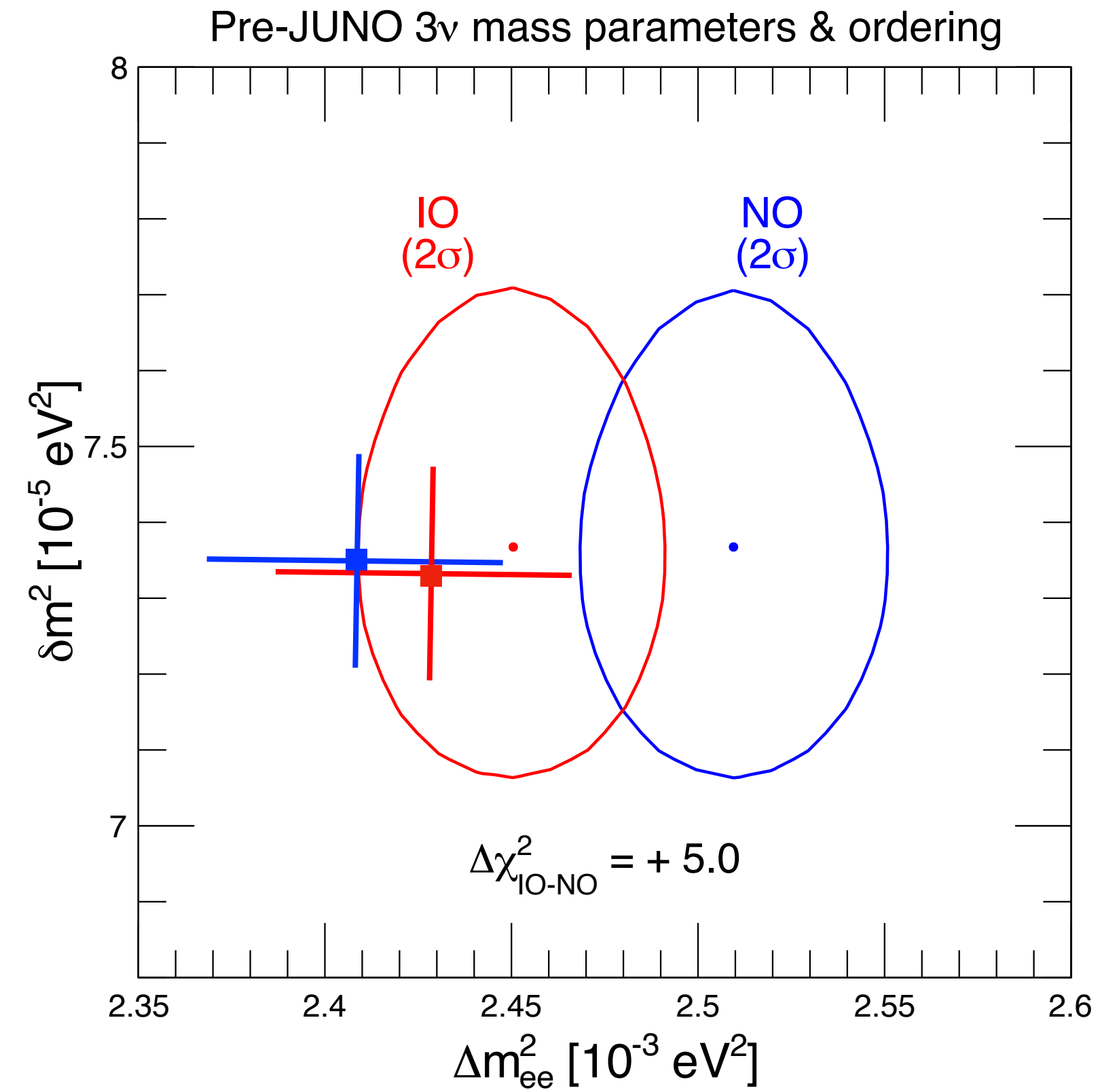


# Examples of possible JUNO first data compared with pre-JUNO data

## Asynergy favoring NO

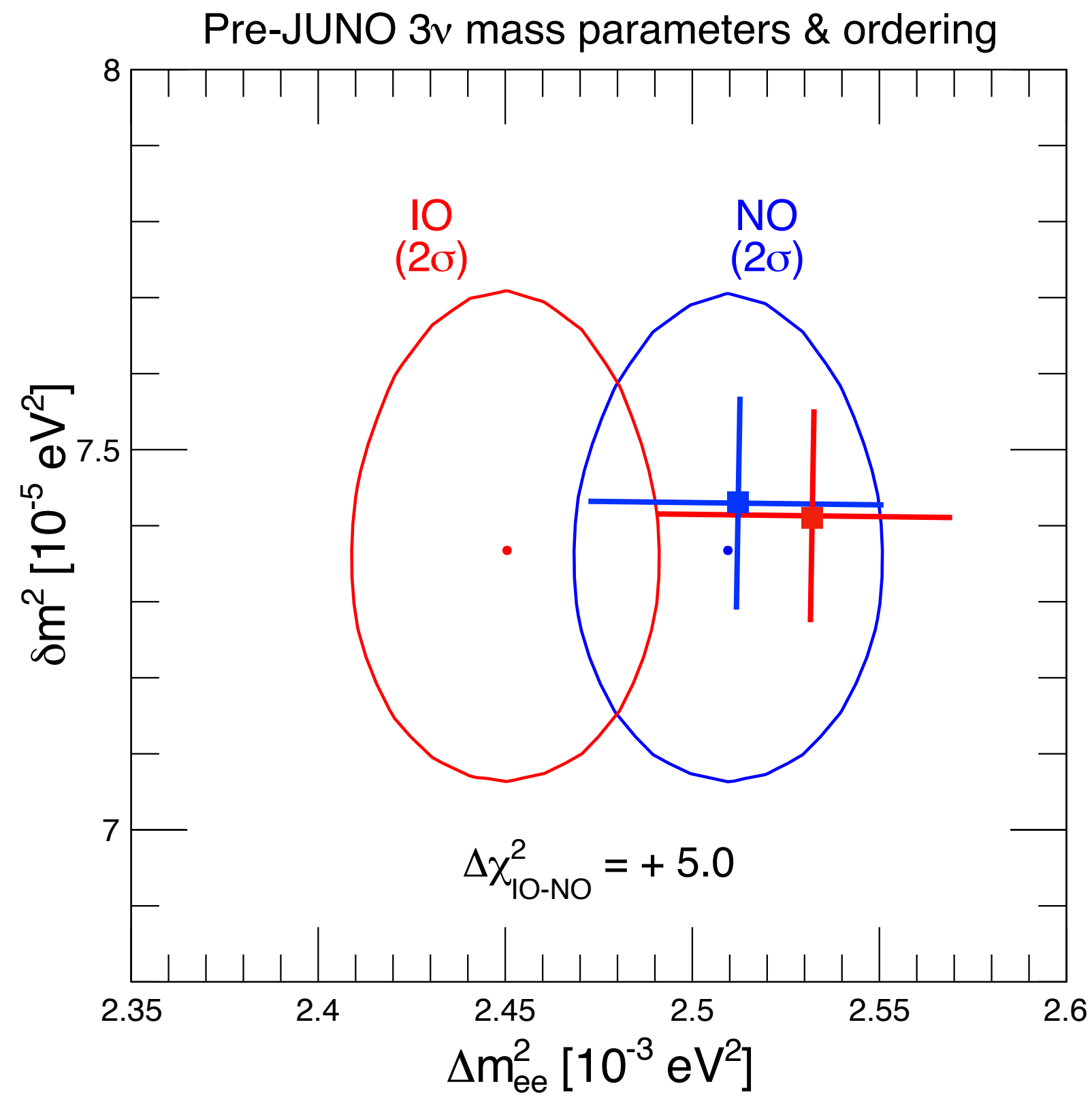


## Asynergy favoring IO

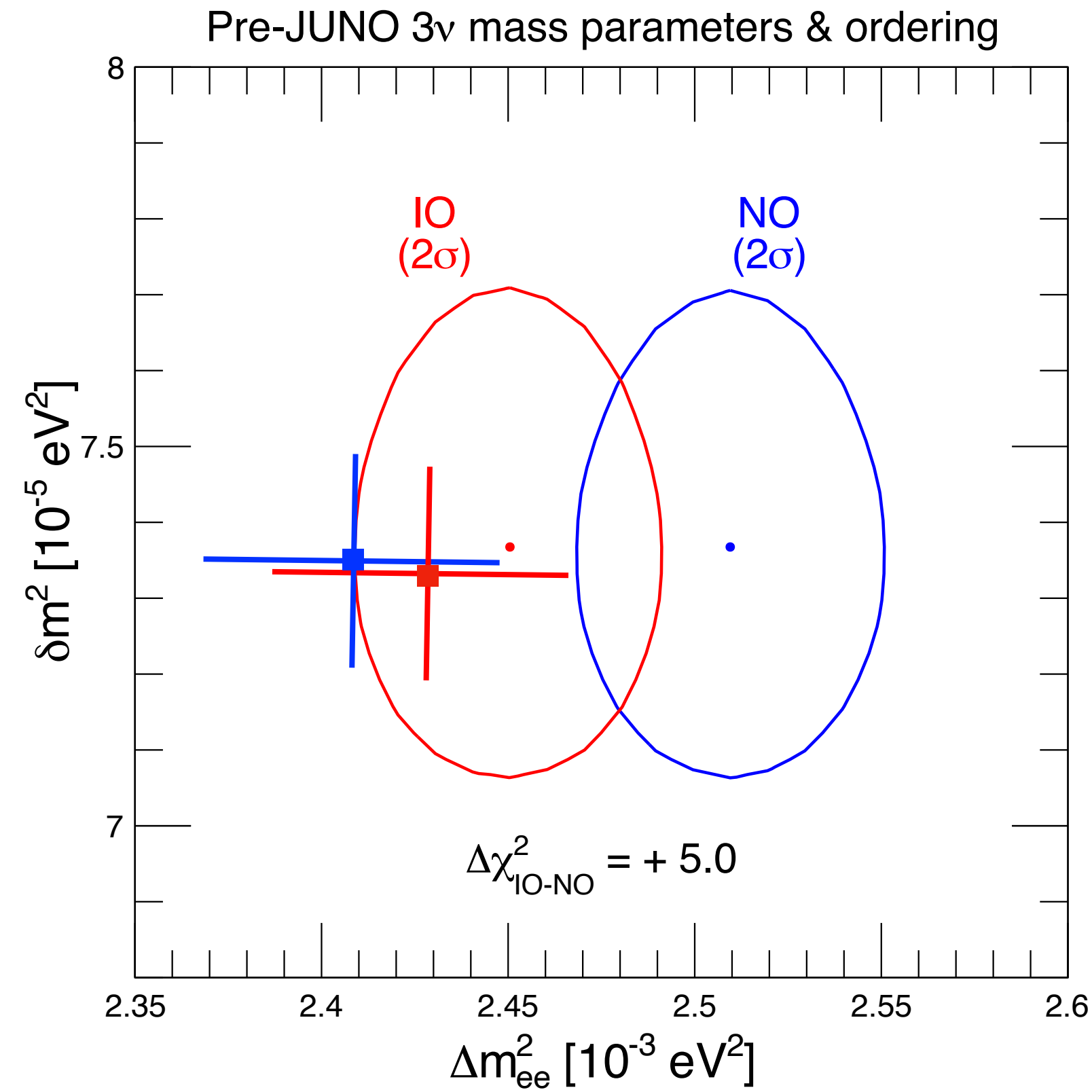


# Examples of possible JUNO first data compared with pre-JUNO data

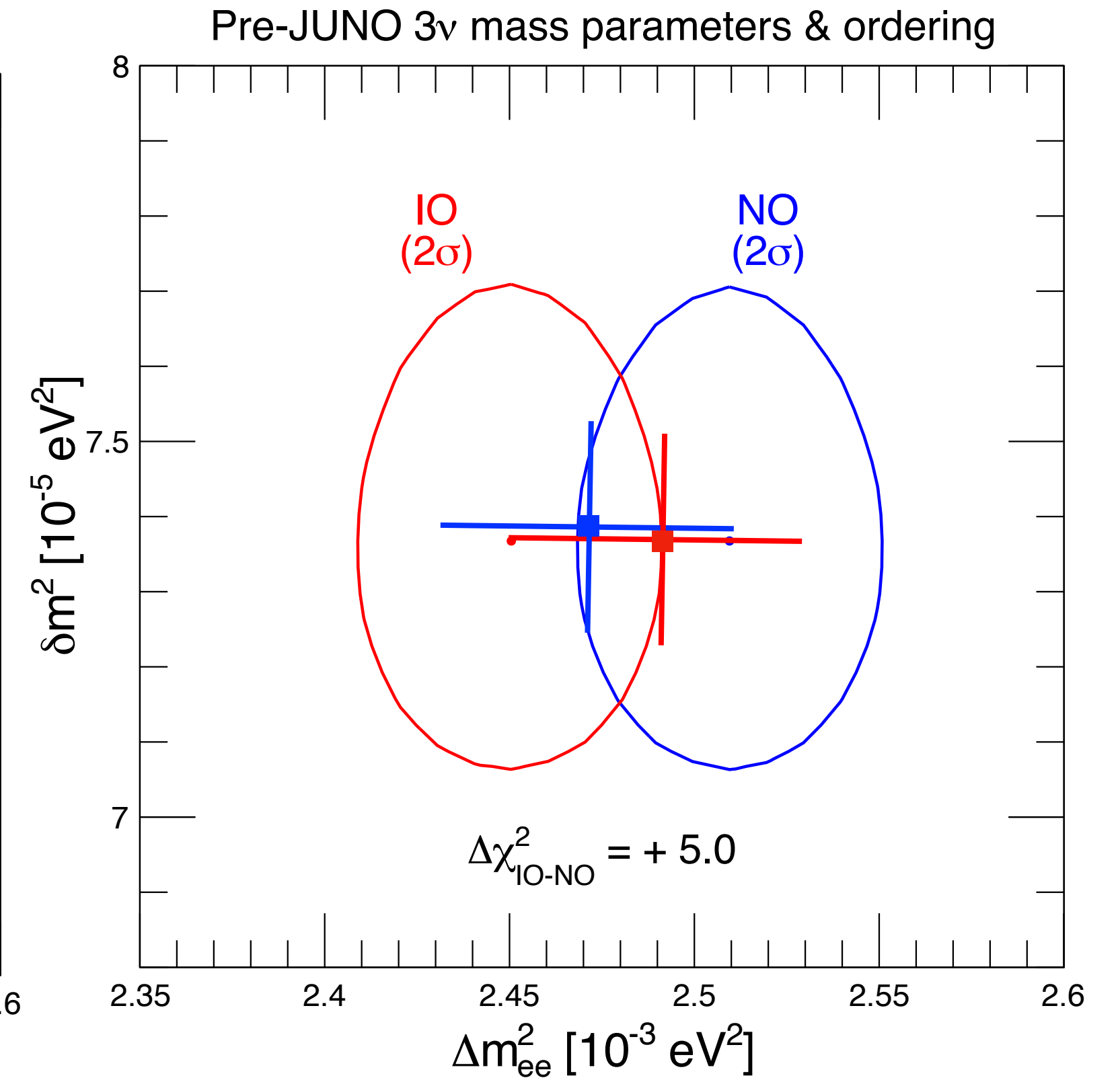
## Asynergy favoring NO



## Asynergy favoring IO



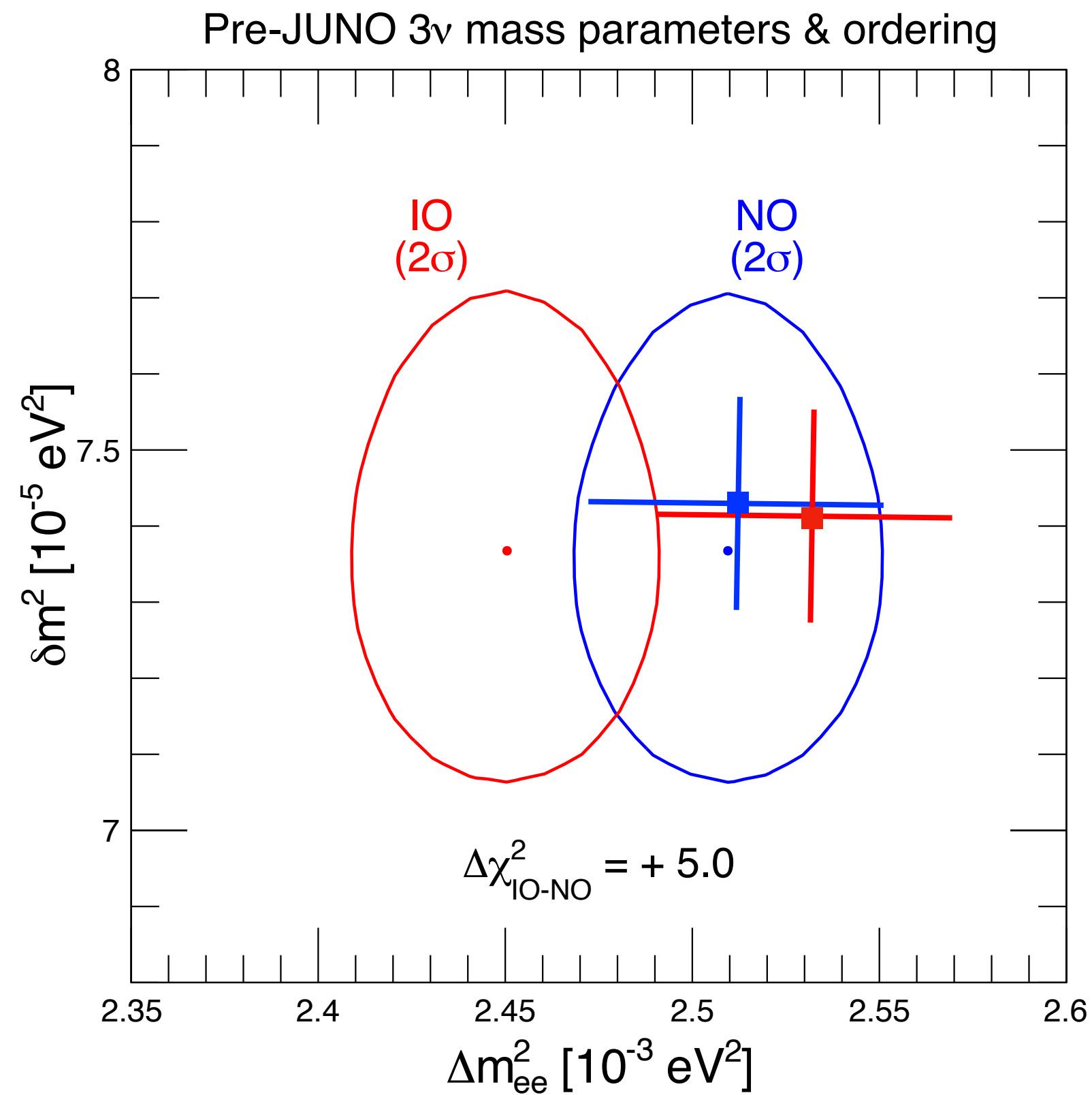
## An undecided NO/IO



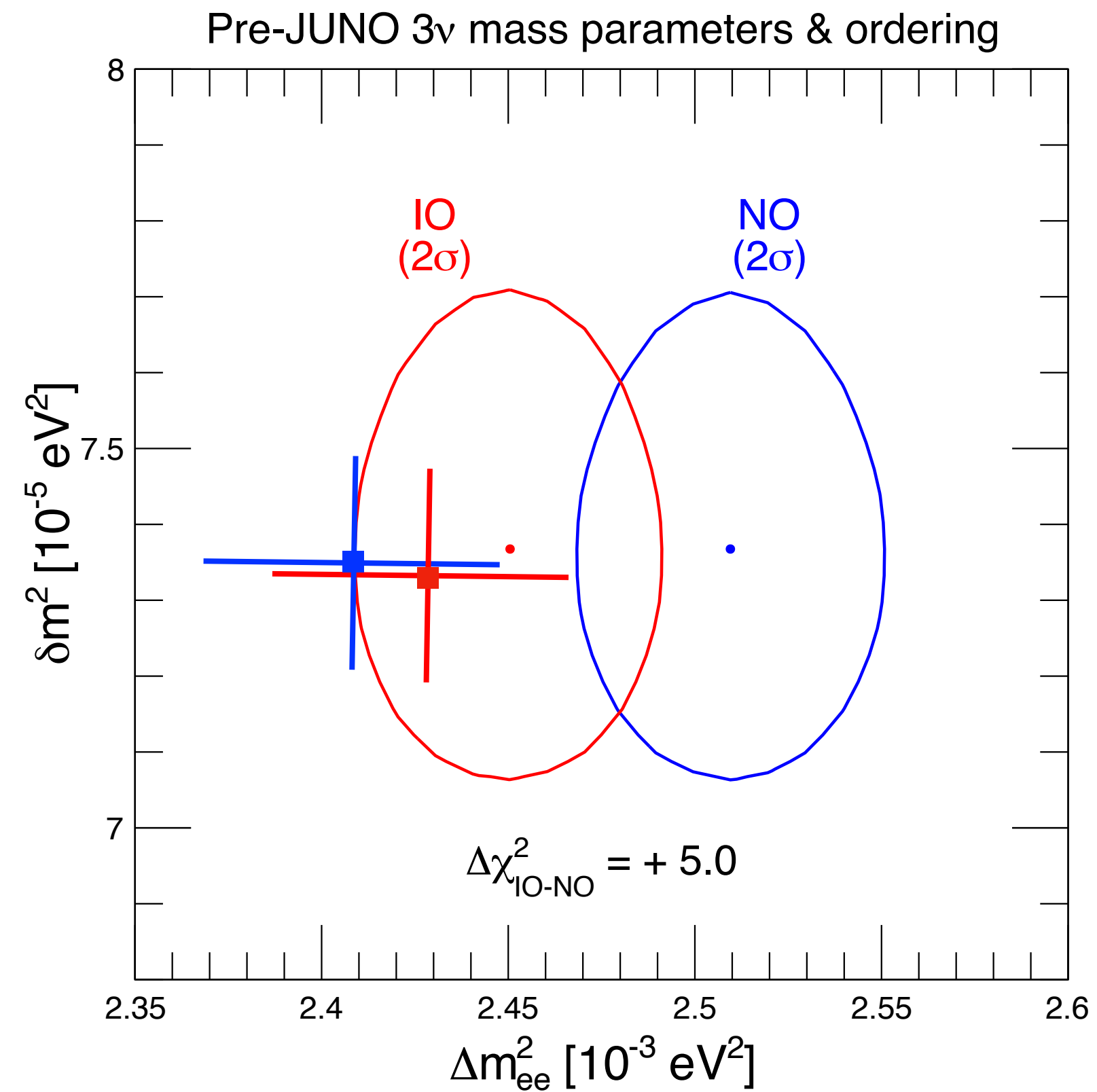


# Examples of possible JUNO first data compared with pre-JUNO data

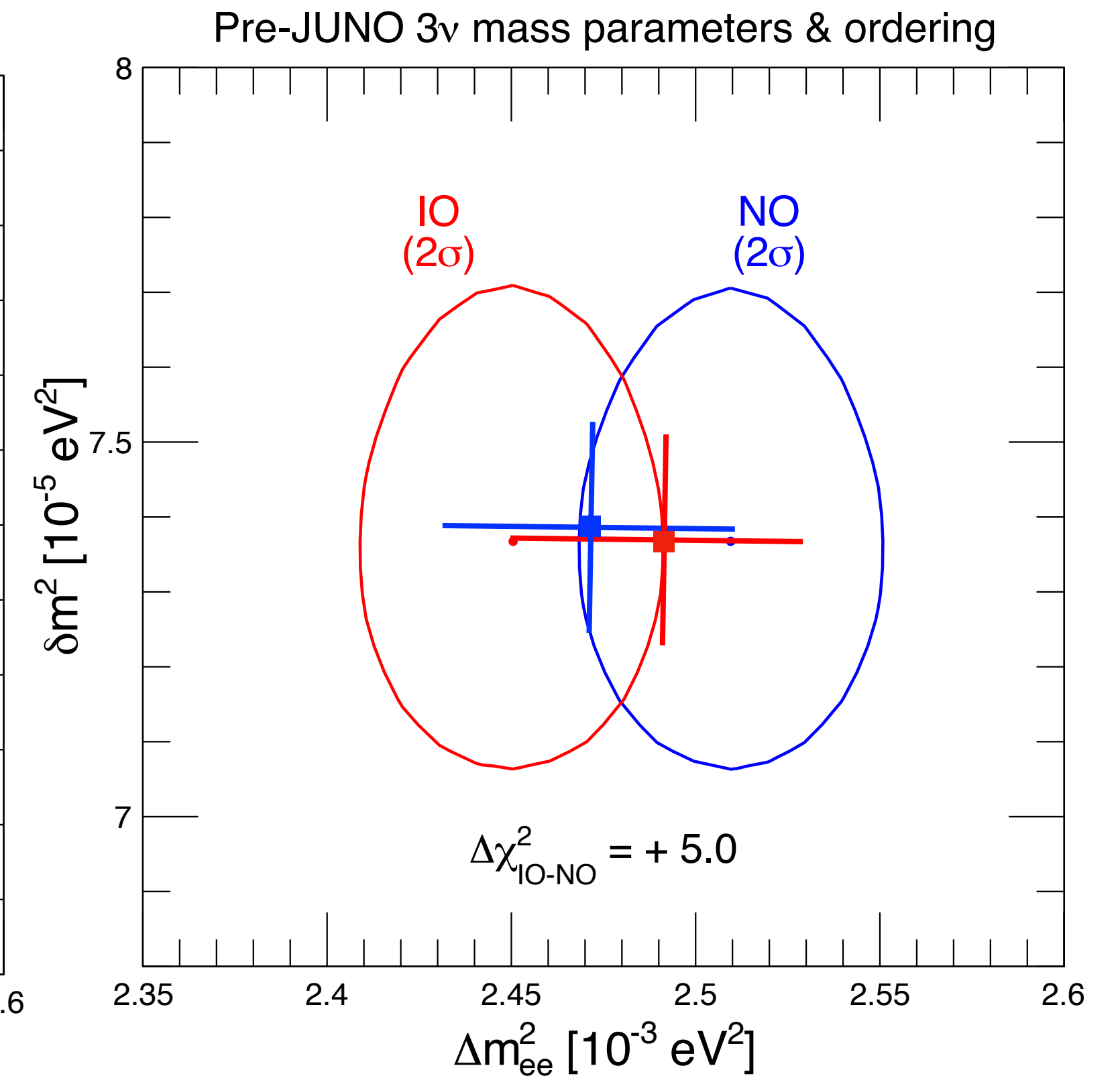
## Asynergy favoring NO



## Asynergy favoring IO



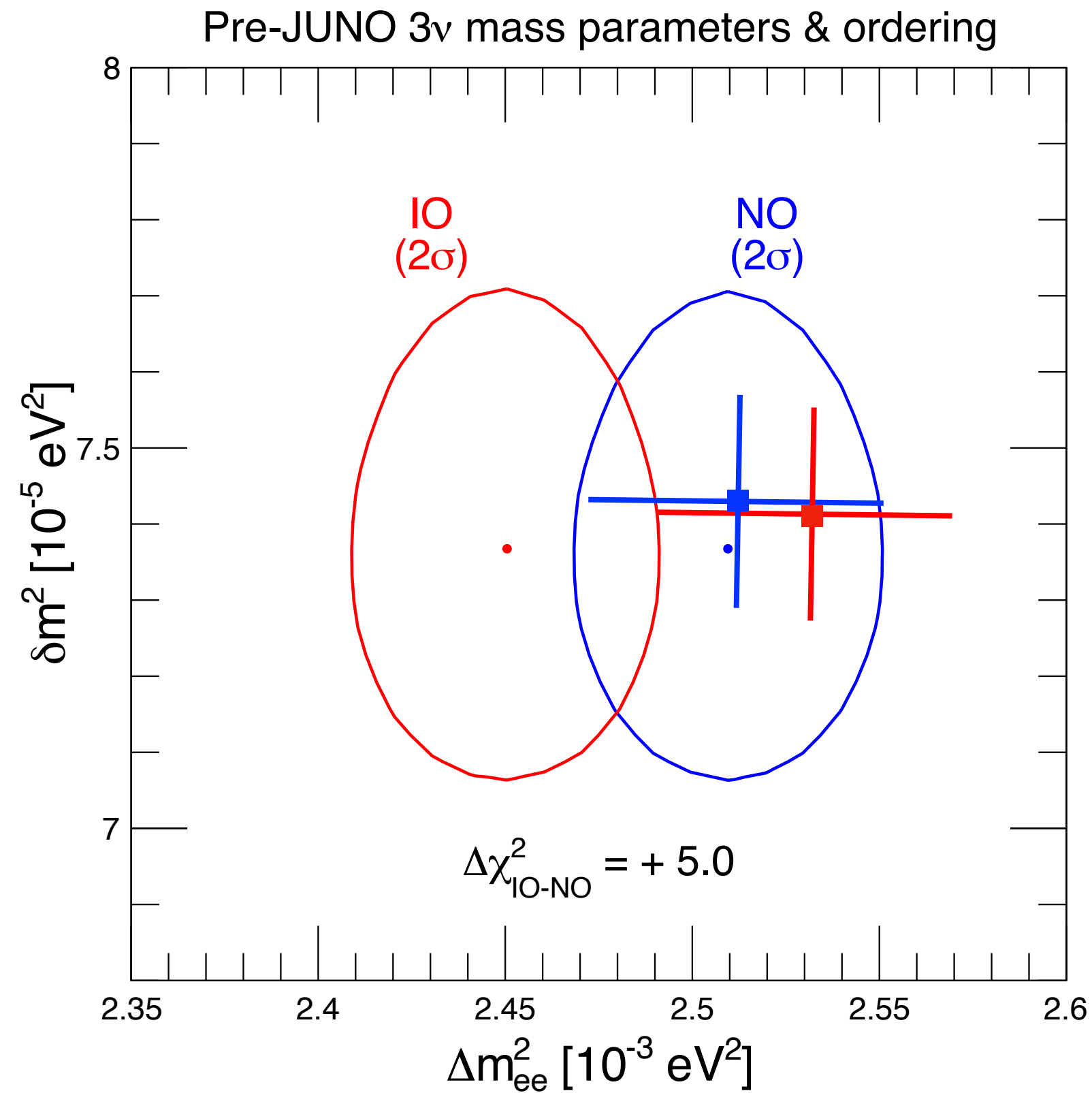
## An undecided NO/IO



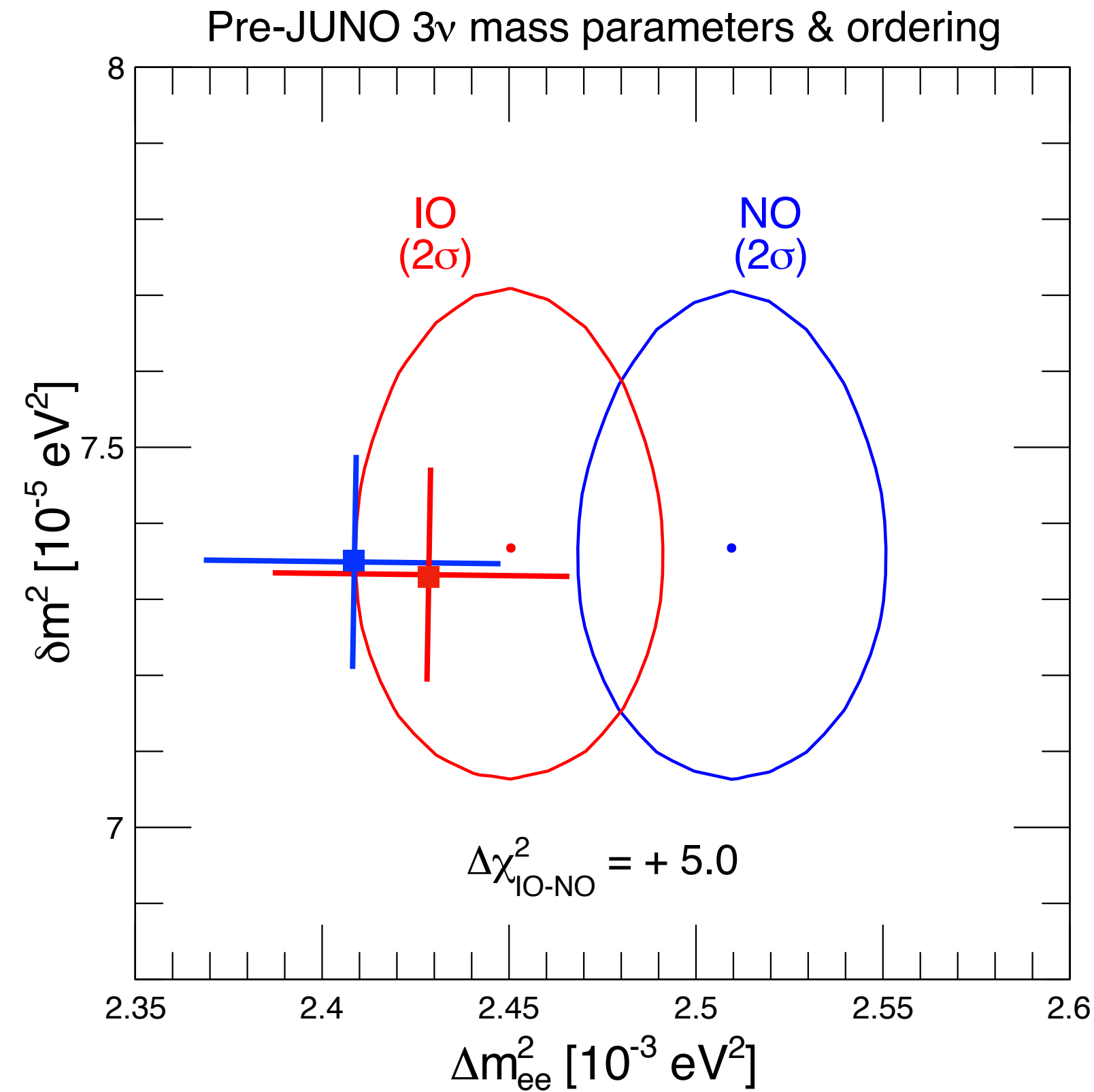
Global analyses helpful to understand correlated impact on other parameters

# Examples of possible JUNO first data compared with pre-JUNO data

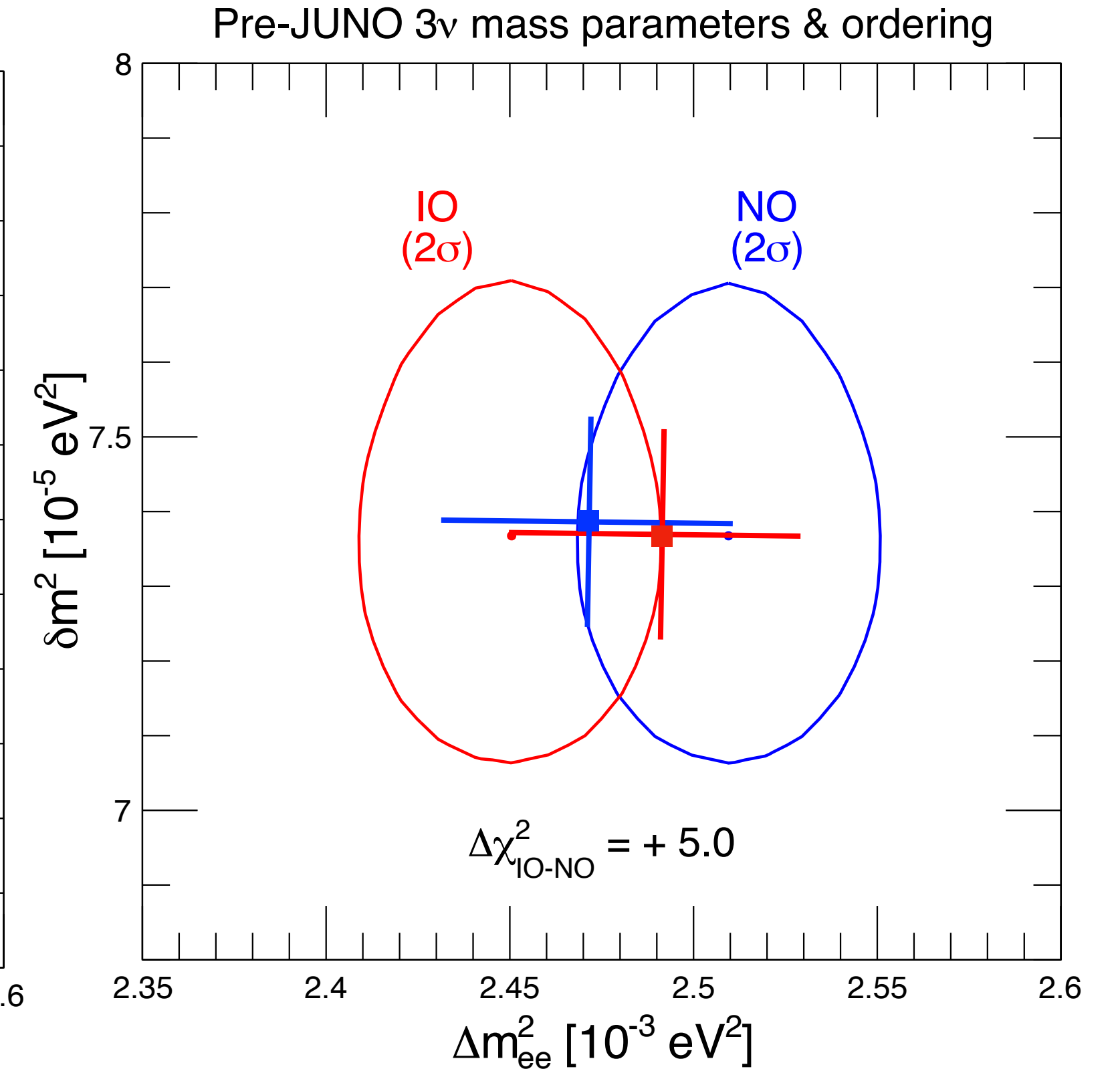
## Asynergy favoring NO



## Asynergy favoring IO



## An undecided NO/IO



Global analyses helpful to understand correlated impact on other parameters

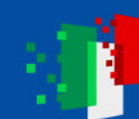
It will be instructive to locate the first JUNO data in this plane and eventually compare them with JUNO-alone NO/IO findings for convergence.



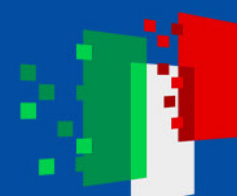
**Finanziato  
dall'Unione europea**  
NextGenerationEU



**Ministero  
dell'Università  
e della Ricerca**



**Italiadomani**  
PIANO NAZIONALE  
DI RIPRESA E RESILIENZA



# Summary and Perspectives

# Summary and Perspectives

- **Precision reached**

Percent accuracy on mixing angles and mass splittings - with  $|\Delta m^2|$  at **subpercent accuracy (0.8%)**

# Summary and Perspectives

- **Precision reached**

Percent accuracy on mixing angles and mass splittings - with  $|\Delta m^2|$  at **subpercent accuracy (0.8%)**

- **Oscillation unknowns**

Mass ordering,  $\theta_{23}$  octant, CP phase  $\delta_{CP}$  - weak hint for NO



# Summary and Perspectives

- **Precision reached**

Percent accuracy on mixing angles and mass splittings - with  $|\Delta m^2|$  at **subpercent accuracy (0.8%)**

- **Oscillation unknowns**

Mass ordering,  $\theta_{23}$  octant, CP phase  $\delta_{CP}$  - weak hint for NO

- **Absolute masses**

Current limits:  $m_\beta \leq 0.50$  eV,  $m_{\beta\beta} \leq 0.086$  eV,  $\Sigma \leq 0.2$  eV (cosmology uncertain)

# Summary and Perspectives

- **Precision reached**

Percent accuracy on mixing angles and mass splittings - with  $|\Delta m^2|$  at **subpercent accuracy (0.8%)**

- **Oscillation unknowns**

Mass ordering,  $\theta_{23}$  octant, CP phase  $\delta_{CP}$  - weak hint for NO

- **Absolute masses**

Current limits:  $m_\beta \leq 0.50$  eV,  $m_{\beta\beta} \leq 0.086$  eV,  $\Sigma \leq 0.2$  eV (cosmology uncertain)

- **Next steps**

JUNO  $\rightarrow (\delta m^2, \Delta m_{ee}^2)$  with subpercent precision and test mass ordering



# Summary and Perspectives

- **Precision reached**

Percent accuracy on mixing angles and mass splittings - with  $|\Delta m^2|$  at **subpercent accuracy (0.8%)**

- **Oscillation unknowns**

Mass ordering,  $\theta_{23}$  octant, CP phase  $\delta_{CP}$  - weak hint for NO

- **Absolute masses**

Current limits:  $m_\beta \leq 0.50$  eV,  $m_{\beta\beta} \leq 0.086$  eV,  $\Sigma \leq 0.2$  eV (cosmology uncertain)

- **Next steps**

JUNO  $\rightarrow (\delta m^2, \Delta m_{ee}^2)$  with subpercent precision and test mass ordering

- **Outlook**

The 3v framework is at a turning point: future synergies (or tensions) across oscillation,  $\beta$ -decay,  $0\nu\beta\beta$ , cosmology will be decisive